

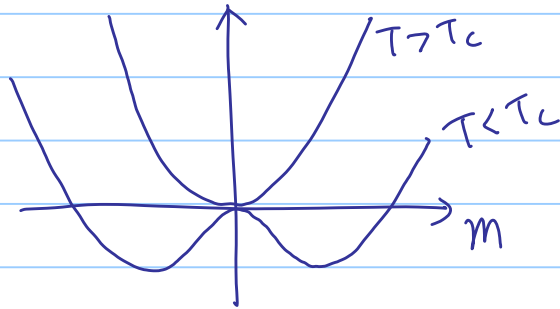
MFT - Continued

Note Title

13/02/2012

Bragg-Williams form of Ising model

$$f = \frac{F}{N} = \frac{1}{2} (T - T_c) m^2 + \frac{T}{12} m^4 + \dots$$



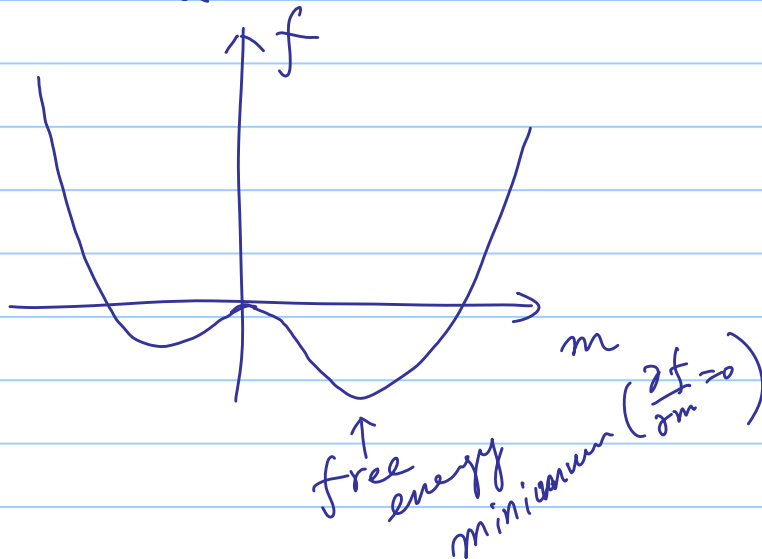
$$m \sim (T_c - T)^\beta$$

$$\beta = \frac{1}{2}$$

Add magnetic field (external field) : H

→ free energy term: $-Hm$

$$f = \frac{1}{2} (T - T_c) m^2 + \frac{T}{12} m^4 - Hm + \dots$$



$$\frac{\partial f}{\partial m} = (T - T_c) m + \frac{T}{3} m^3 - H = 0$$

eq. state

At $T = T_c$

$$m^3 = \frac{3H}{T}$$

$$m \propto H^{1/3} \text{ (@ } T_c \text{)}$$

$$m \propto H^{1/\delta} \quad \delta = 3$$

mean field

Susceptibility

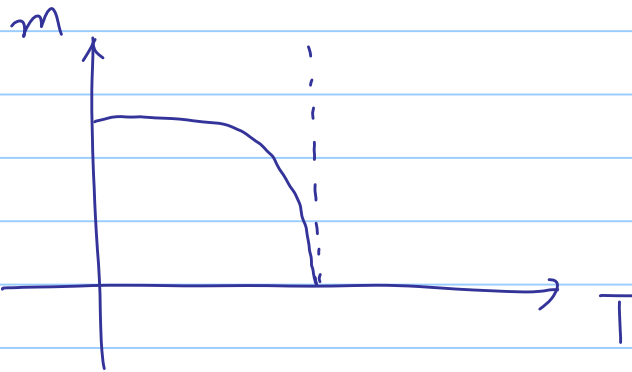
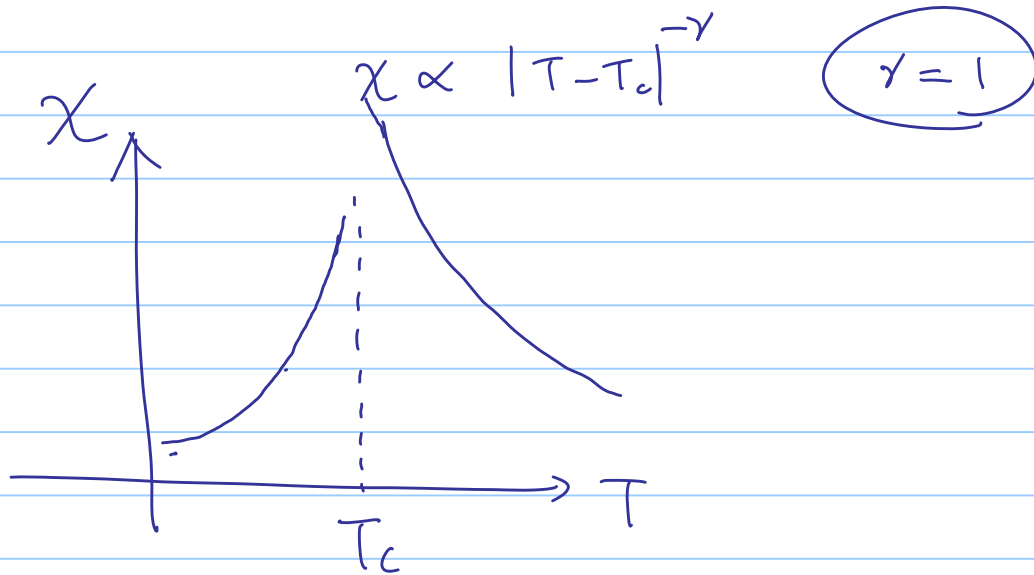
$$\chi = \left. \frac{\partial m}{\partial H} \right|_{H \rightarrow 0}$$

$$\frac{\partial}{\partial H} \left(\frac{\partial f}{\partial m} \right) = (T - T_c) \underbrace{\frac{\partial m}{\partial H}}_{\chi} + T m^2 \underbrace{\frac{\partial m}{\partial H}}_{\chi} - 1 = 0$$

$$\chi = \frac{1}{T - T_c + T_c m^2} \quad (T \approx T_c)$$

For $T > T_c$: $m = 0$ $\chi = \frac{1}{T - T_c} = (T - T_c)^{-\gamma}$

$T < T_c$: $m^2 = \frac{3(T_c - T)}{T_c} \Rightarrow \chi = \frac{1}{2(T_c - T)} = \frac{1}{2}(T_c - T)^{-\gamma}$



Fluctuations are
suppressed by order
parameter saturation

More critical exponents: ν , η

Landau theory \rightarrow Free energy can be expanded as a power series in the order parameter

* Only those terms compatible with the symmetry of the system

example: $f = \frac{1}{2}(T - T_c)m^2 + \frac{T}{12}m^4 + \dots$

Typically, $f = a_0 + a_1 \phi + a_2 \phi^2 + a_3 \phi^3 + \dots$

ϕ : order parameter.

$(T - T_c)$ term

$$t \equiv \frac{T - T_c}{T_c}$$

$$a_2 \phi^2 = \tilde{a}_2 t \phi^2$$

Back to Ising model (for correlation fn.)

Non-uniform fluctuations

Allow magnetization variation as a fn. of position
(block spin) $m \rightarrow m(x)$ (1D example)

$$\mathcal{H} = -J \sum_x m(x) m(x+a)$$

$$m(x) m(x+a) = \left\{ - \left[\frac{m(x+a) - m(x)}{a} \right]^2 \cdot a^2 + m(x)^2 + m(x+a)^2 \right\} \frac{1}{2}$$

In Continuum limit $a \rightarrow 0$

$$\mathcal{H} = -J \sum_x m(x)^2 + \underbrace{\frac{J a^2}{2} \sum_x \left(\frac{\partial m}{\partial x} \right)^2}_{\text{Additional term taking into account of fluctuations}} + \text{higher order terms.}$$

With this, Landau MFT (Ising model example)

$$f = a_0 + a_2 \sum_x m(x)^2 + g \sum_x \left(\frac{\partial m}{\partial x} \right)^2$$

Suppose $m(x)$ varies slowly, and small

$$f = a_0 + a_2 \int m(x)^2 dx + g \int \left(\frac{\partial m}{\partial x} \right)^2 dx$$

$$\text{F.T.} \left[\begin{array}{l} m(x) = \frac{1}{2\pi} \int dq e^{iqx} m(q) \\ m(q) = \int dx e^{-iqx} m(x) \end{array} \right] \uparrow$$

$$f - a_0 = \frac{1}{2\pi} \int dq (a_2 + q^2) m(q) m(-q)$$

Orstein-Zernike (OZ) assumption.

Independent q contribution to f

→ Each q : $(a_2 + q^2) \langle m(q) m(-q) \rangle$

should equally contribute to f

$$(a_2 + q^2) \langle m(q) m(-q) \rangle = kT$$

↓

$$S(q)$$

What is $\langle m(q) m(-q) \rangle$?

Recall $C(x) = \langle m(x) m(0) \rangle - \langle m(0) \rangle^2$ $T > T_c$

$$S(q) = \langle m(q) m(-q) \rangle$$

$$= \frac{kT}{a_2 + q^2}$$

inv. F.T.

$$C(x) = e^{-x/\xi}$$

$$\xi = \sqrt{\frac{q}{a_2}}$$

$$a_2 = \tilde{a}_2 (T - T_c)$$

$$\xi \propto |T - T_c|^{-1/2} = |T - T_c|^{-\nu}$$

$$\nu = \frac{1}{2}$$