

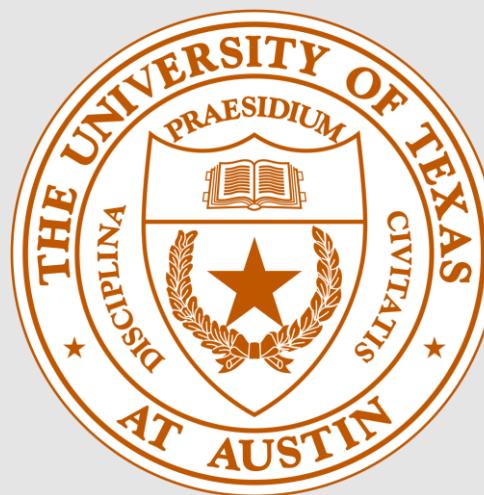


Reciprocity, passivity, and causality in fully coupled acousto-electrodynamic media

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185th Meeting of the Acoustical Society of America
Sydney, Australia
Tuesday, December 5th, 2023

Outline

Background

1. Piezoelectricity
2. Electrodynamic and acoustic biaxiality
3. Acousto-electrodynamic constitutive relations

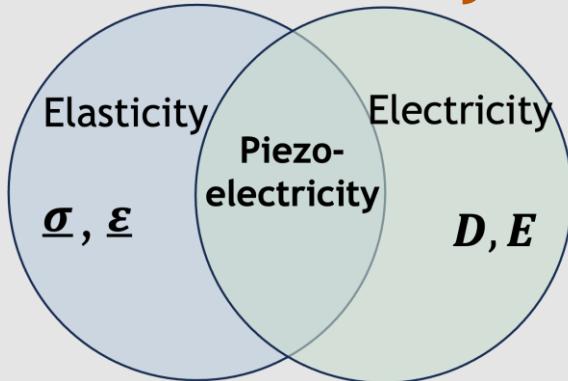
Derivations for acousto-electrodynamic media

1. Reciprocity
2. Passivity
3. Causality

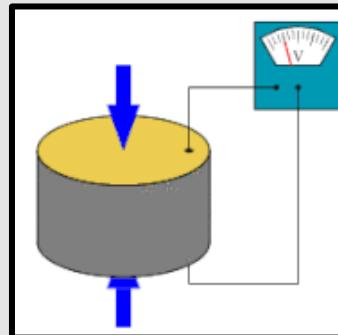
Closing

1. Verification of source-driven constitutive relations
2. Interpretation

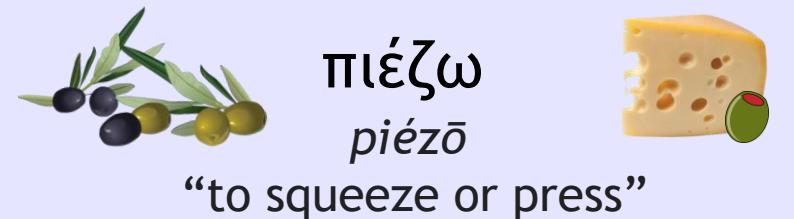
Piezoelectricity



1880



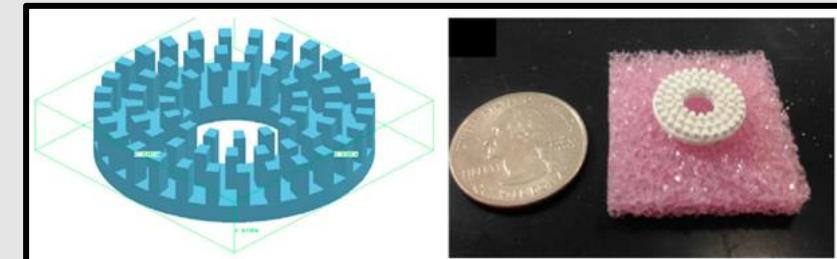
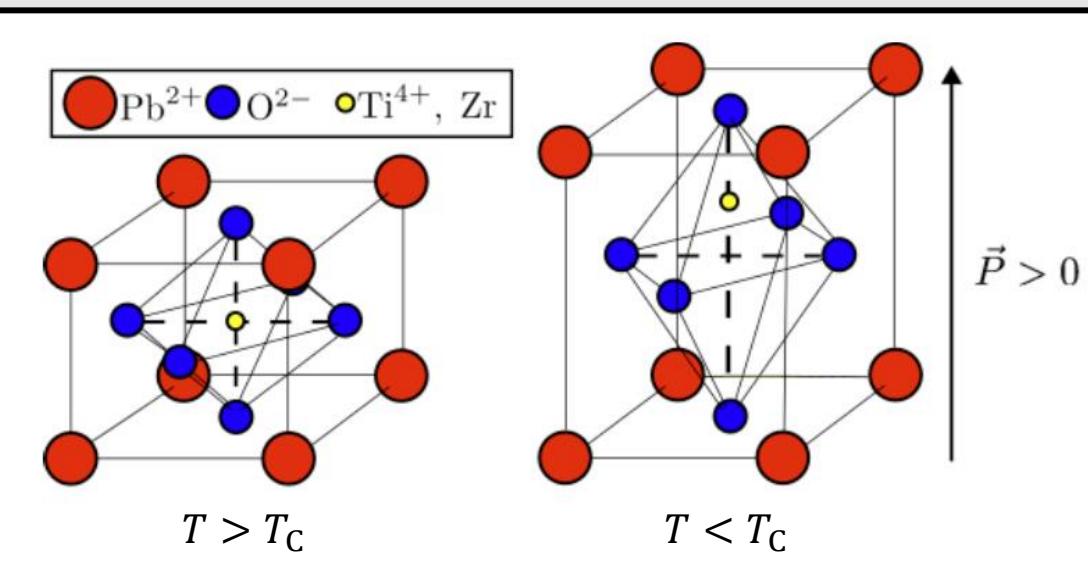
Etymology



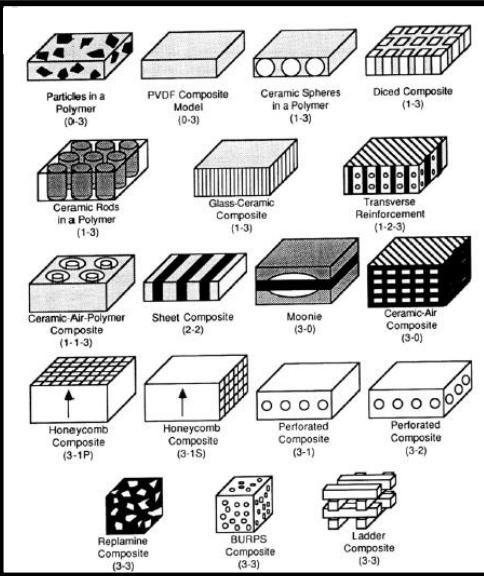
Piezoelectric constitutive relations

$$\text{strain} \rightarrow \underline{\varepsilon} = \mathcal{S} : \underline{\sigma} + \tilde{d}^T \cdot \underline{E}$$

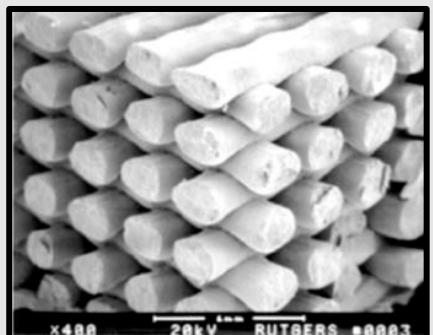
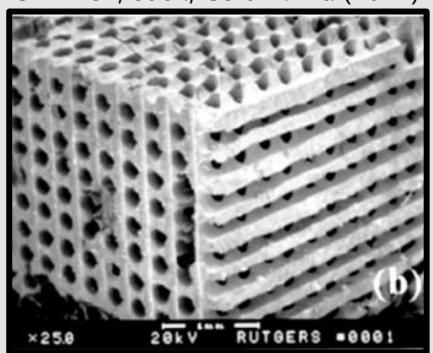
$$\begin{matrix} \text{compliance} & \text{stress} & \text{electric field} \\ \downarrow & \nearrow & \leftarrow \\ \text{Electric displacement} & \underline{D} = \tilde{d} : \underline{\sigma} + \underline{\epsilon}^T \cdot \underline{E} & \text{permittivity} \\ \uparrow \text{coupling} & & \end{matrix}$$



Smirnov, et al., Ceram. Int. (2021)



Smirnov, et al., Ceram. Int. (2021)



Electrodynamic and acoustic bianisotropy

$$D = \epsilon E$$

$$B = \mu H$$

Isotropic



Hubble deep field, NASA

$$D = \underline{\epsilon} E$$

$$B = \underline{\mu} H$$

Anisotropic

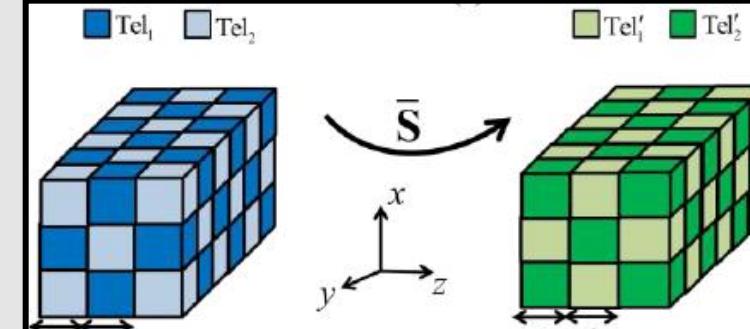


Birefringence, Wikipedia

$$D = \epsilon E + \zeta H$$

$$B = \xi E + \mu H$$

Biisotropic

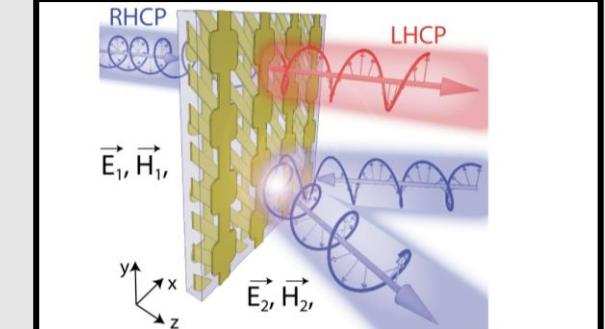


Prudencio et al., IEEE (2014)

$$D = \underline{\epsilon} E + \underline{\zeta} H$$

$$B = \underline{\xi} E + \underline{\mu} H$$

Bianisotropic



C. Pfeiffer et al., PRL (2014)

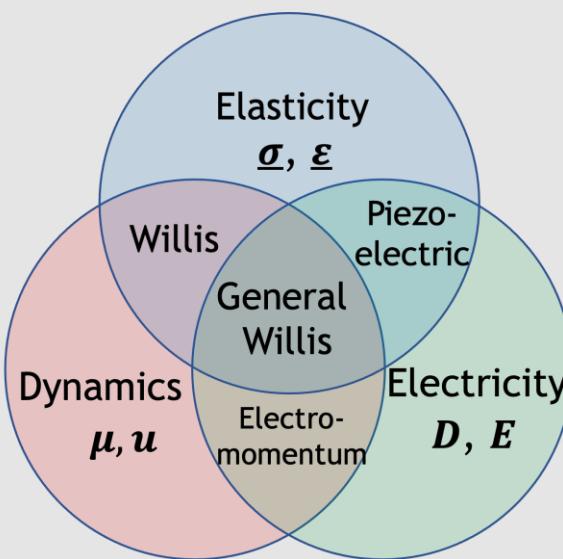
$$\varepsilon = -\beta p$$

$$\mu = \rho u$$

Isotropic



Meyerson Symphony Center

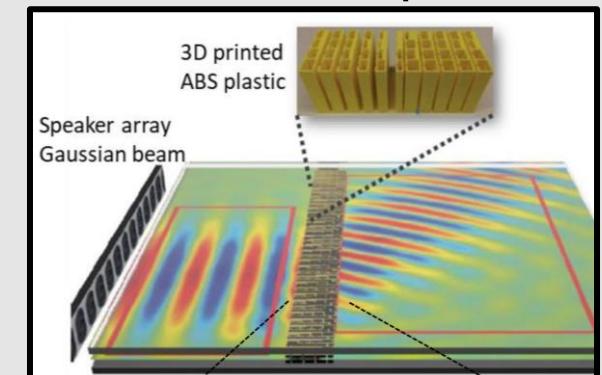


Pernas-Salomón and Shmuel, J. Mech. Phys. Solids 134, 103770 (2020)

$$\varepsilon = -\beta p + \gamma \cdot u$$

$$\mu = -\eta p + \underline{\rho} u$$

Bianisotropic



J. Li et al., Nat. Comm. (2018)

Acousto-electrodynamic constitutive relations

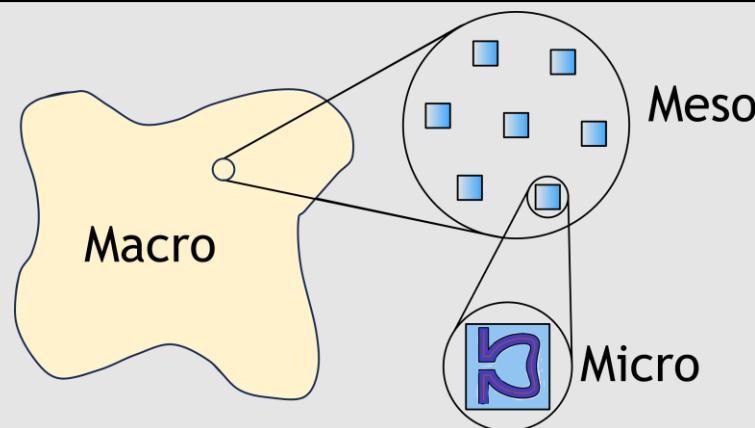
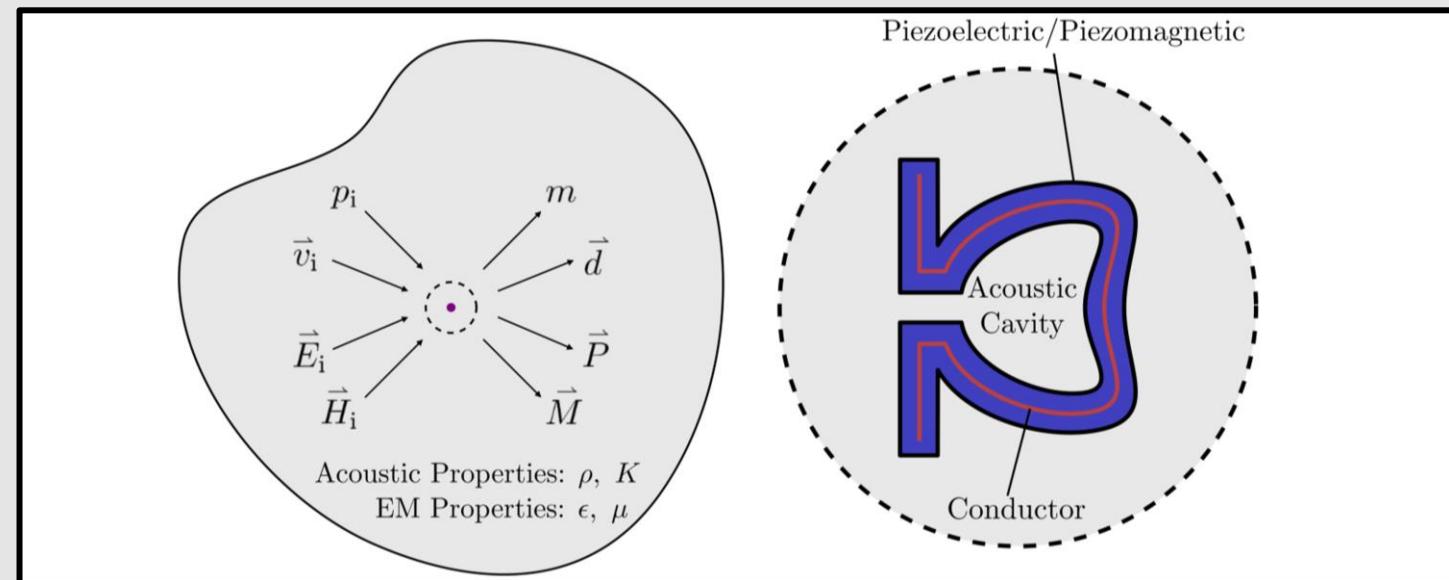
$$\begin{pmatrix} \varepsilon \\ \mu \\ D \\ B \end{pmatrix} = \begin{pmatrix} -\beta & \gamma^T & \underline{\mathbf{b}}^T & \underline{\mathbf{d}}^T \\ -\eta & \underline{\rho} & \underline{\mathbf{v}} & \underline{\mathbf{m}} \\ c & \underline{\mathbf{w}} & \underline{\epsilon} & \underline{\xi} \\ e & \underline{\mathbf{n}} & \underline{\zeta} & \underline{\mu} \end{pmatrix} \begin{pmatrix} p \\ u \\ E \\ H \end{pmatrix}$$

"[O]ne can think about metasurfaces simultaneously controlling and transforming waves of different nature, which can lead to the creation of multiphysics bianisotropic metasurfaces."

—V. S. Asadchy et al., Nanophotonics 7, 6 (2018)

Components of constitutive block matrix

1. acoustic bianisotropy
2. electromagnetic bianisotropy
3. piezoelectricity
4. piezomagnetism
5. electromomentum
6. magnetomomentum



Acousto-electrodynamic constitutive relations

$$\begin{pmatrix} \varepsilon \\ \mu \\ D \\ B \end{pmatrix} = \begin{pmatrix} -\beta & \gamma^T & b^T & d^T \\ -\eta & \underline{\rho} & \underline{v} & \underline{m} \\ c & \underline{w} & \underline{\epsilon} & \underline{\xi} \\ e & \underline{n} & \underline{\zeta} & \underline{\mu} \end{pmatrix} \begin{pmatrix} p \\ u \\ E \\ H \end{pmatrix}$$

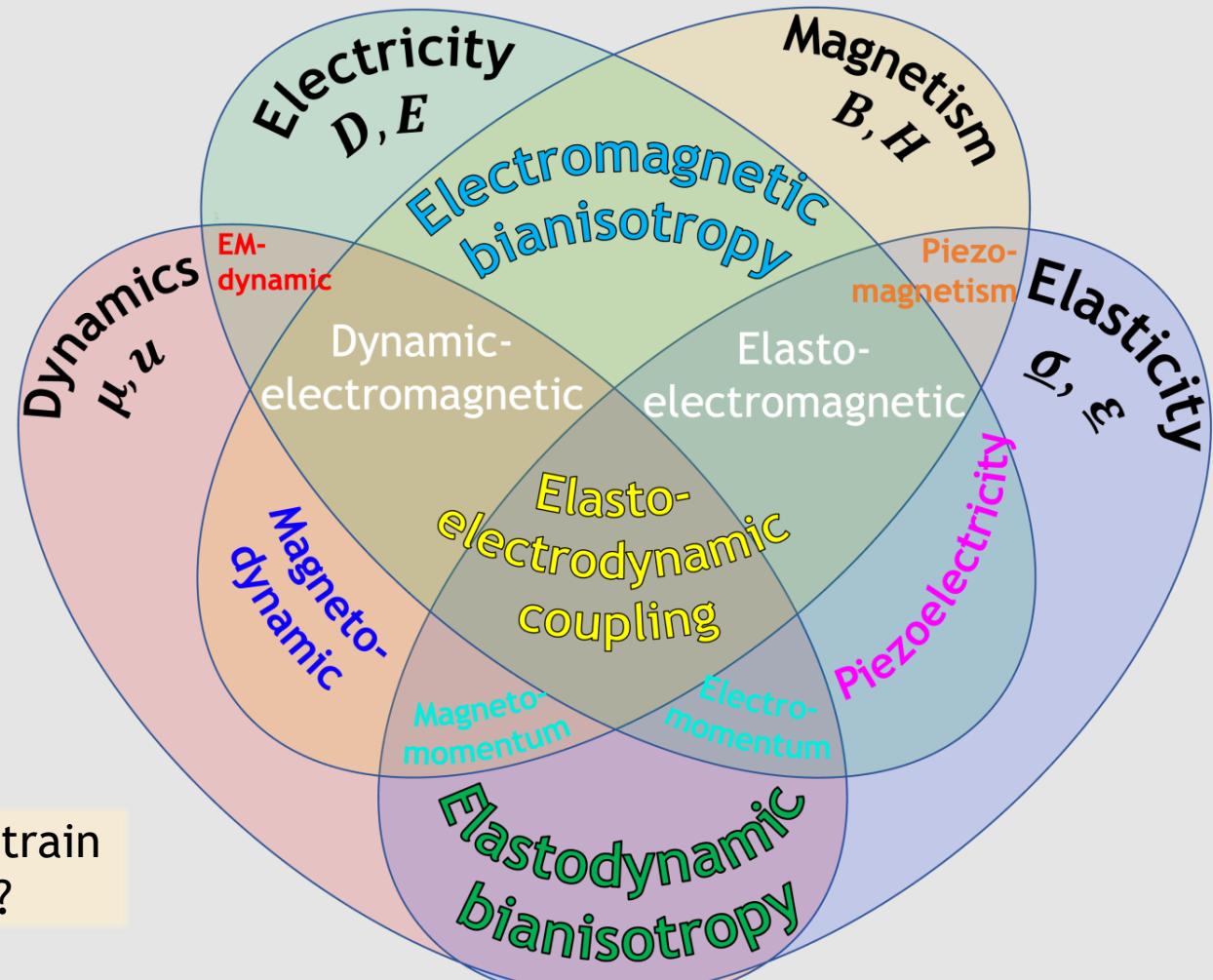
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How do reciprocity, passivity, and causality constrain the acousto-electrodynamic constitutive relations?



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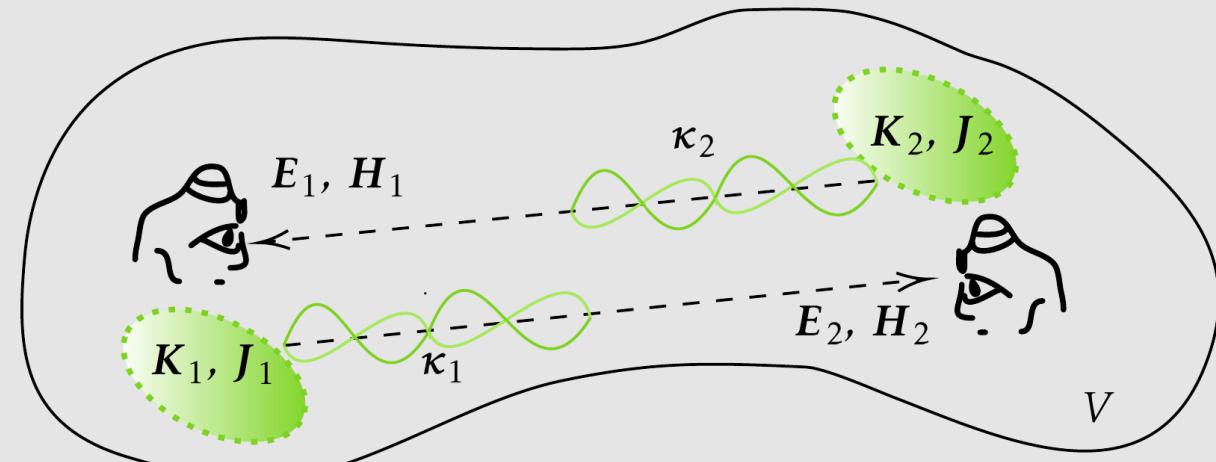
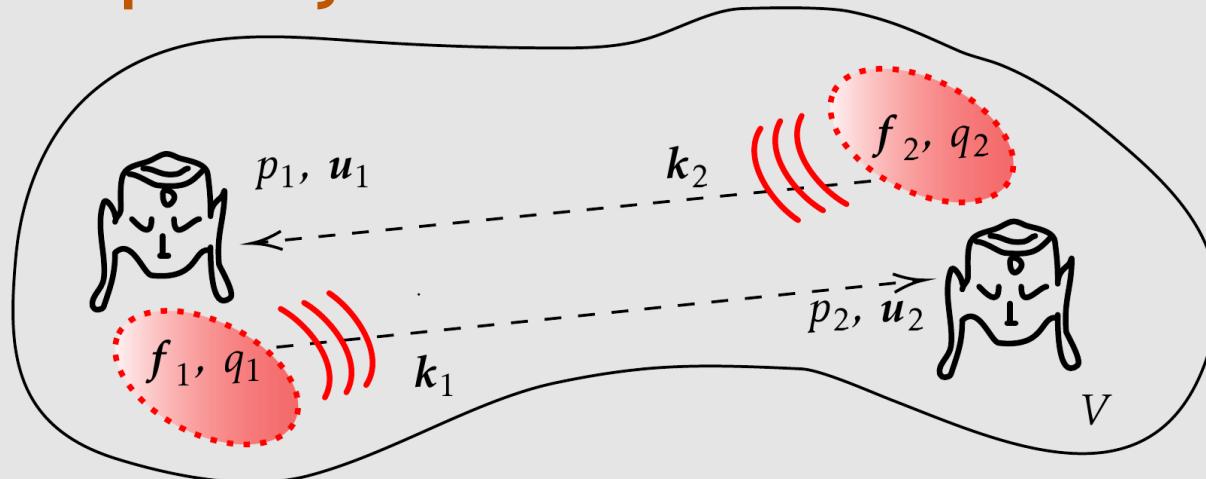
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Reciprocity



Governing equations for 1

$$\begin{aligned}\nabla p_1 &= i\omega \boldsymbol{\mu}_1 + \mathbf{f}_1 \\ \nabla \cdot \mathbf{u}_1 &= -i\omega \varepsilon_1 + q_1 \\ \nabla \times \mathbf{E}_1 &= i\omega \mathbf{B}_1 - \mathbf{K}_1 \\ \nabla \times \mathbf{H}_1 &= -i\omega \mathbf{D}_1 + \mathbf{J}_1\end{aligned}$$

$$\begin{aligned}\mathbf{u}_2 \cdot \nabla p_1 &= i\omega \mathbf{u}_2 \cdot \boldsymbol{\mu}_1 + \mathbf{u}_2 \cdot \mathbf{f}_1 \\ p_2 \nabla \cdot \mathbf{u}_1 &= -i\omega p_2 \varepsilon_1 + p_2 q_1 \\ \mathbf{H}_2 \cdot (\nabla \times \mathbf{E}_1) &= i\omega \mathbf{H}_2 \cdot \mathbf{B}_1 - \mathbf{H}_2 \cdot \mathbf{K}_1 \\ \mathbf{E}_2 \cdot (\nabla \times \mathbf{H}_1) &= -i\omega \mathbf{E}_2 \cdot \mathbf{D}_1 + \mathbf{E}_2 \cdot \mathbf{J}_1\end{aligned}$$

+

Governing equations for 2

$$\begin{aligned}\nabla p_2 &= i\omega \boldsymbol{\mu}_2 + \mathbf{f}_2 \\ \nabla \cdot \mathbf{u}_2 &= -i\omega \varepsilon_2 + q_2 \\ \nabla \times \mathbf{E}_2 &= i\omega \mathbf{B}_2 - \mathbf{K}_2 \\ \nabla \times \mathbf{H}_2 &= -i\omega \mathbf{D}_2 + \mathbf{J}_2.\end{aligned}$$

$$\begin{cases} \nabla \cdot (p_1 \mathbf{u}_2) = i\omega (\mathbf{u}_2 \cdot \boldsymbol{\mu}_1 - p_1 \varepsilon_2) + \mathbf{u}_2 \cdot \mathbf{f}_1 + p_1 q_2 \\ \nabla \cdot (p_2 \mathbf{u}_1) = i\omega (\mathbf{u}_1 \cdot \boldsymbol{\mu}_2 - p_2 \varepsilon_1) + \mathbf{u}_1 \cdot \mathbf{f}_2 + p_2 q_1 \\ \nabla \cdot (\mathbf{E}_1 \times \mathbf{H}_2) = i\omega (\mathbf{H}_2 \cdot \mathbf{B}_1 + \mathbf{E}_1 \cdot \mathbf{D}_2) - \mathbf{E}_1 \cdot \mathbf{J}_2 - \mathbf{H}_2 \cdot \mathbf{K}_1 \\ \nabla \cdot (\mathbf{E}_2 \times \mathbf{H}_1) = i\omega (\mathbf{H}_1 \cdot \mathbf{B}_2 + \mathbf{E}_2 \cdot \mathbf{D}_1) - \mathbf{E}_2 \cdot \mathbf{J}_1 - \mathbf{H}_1 \cdot \mathbf{K}_2 \end{cases}$$

Add, integrate over all space, and apply divergence theorem

$$\begin{aligned}\nabla \cdot (p_1 \mathbf{u}_2 - p_2 \mathbf{u}_1) &= \mathbf{u}_2 \cdot \mathbf{f}_1 - \mathbf{u}_1 \cdot \mathbf{f}_2 + p_1 q_2 - p_2 q_1 \\ &\quad + i\omega (\mathbf{u}_2 \cdot \boldsymbol{\mu}_1 - \mathbf{u}_1 \cdot \boldsymbol{\mu}_2 - p_1 \varepsilon_2 + p_2 \varepsilon_1)\end{aligned}$$

$$\begin{aligned}\nabla \cdot (\mathbf{E}_1 \times \mathbf{H}_2 - \mathbf{E}_2 \times \mathbf{H}_1) &= \mathbf{E}_2 \cdot \mathbf{J}_1 - \mathbf{E}_1 \cdot \mathbf{J}_2 + \mathbf{H}_1 \cdot \mathbf{K}_2 - \mathbf{H}_2 \cdot \mathbf{K}_1 \\ &\quad + i\omega (\mathbf{H}_2 \cdot \mathbf{B}_1 - \mathbf{H}_1 \cdot \mathbf{B}_2 + \mathbf{E}_1 \cdot \mathbf{D}_2 - \mathbf{E}_2 \cdot \mathbf{D}_1)\end{aligned}$$

Reciprocity

$$\int_A (p_1 \mathbf{u}_2 - p_2 \mathbf{u}_1 + \mathbf{E}_1 \times \mathbf{H}_2 - \mathbf{E}_2 \times \mathbf{H}_1) \cdot \hat{\mathbf{n}} dA = 0 \text{ by the Sommerfeld radiation condition}$$

$$= \int_V (\mathbf{u}_2 \cdot \mathbf{f}_1 - \mathbf{u}_1 \cdot \mathbf{f}_2 + p_1 q_2 - p_2 q_1 + \mathbf{E}_2 \cdot \mathbf{J}_1 - \mathbf{E}_1 \cdot \mathbf{J}_2 + \mathbf{H}_1 \cdot \mathbf{K}_2 - \mathbf{H}_2 \cdot \mathbf{K}_1) dV = 0 \text{ by reciprocity}$$

$$- i\omega \int_V (\mathbf{u}_1 \cdot \underline{\mu}_2 - \mathbf{u}_2 \cdot \underline{\mu}_1 + p_1 \varepsilon_2 - p_2 \varepsilon_1 + \mathbf{H}_1 \cdot \mathbf{B}_2 - \mathbf{H}_2 \cdot \mathbf{B}_1 - \mathbf{E}_1 \cdot \mathbf{D}_2 + \mathbf{E}_2 \cdot \mathbf{D}_1) dV \stackrel{!}{=} 0$$

Acousto-electrodynamic constitutive relations

$$\begin{pmatrix} \varepsilon \\ \mu \\ D \\ B \end{pmatrix} = \begin{pmatrix} -\beta & \gamma^\top & \mathbf{b}^\top & \mathbf{d}^\top \\ -\eta & \underline{\rho} & \underline{\mathbf{v}} & \underline{\mathbf{m}} \\ c & \underline{\mathbf{w}} & \underline{\boldsymbol{\epsilon}} & \underline{\boldsymbol{\xi}} \\ e & \underline{\mathbf{n}} & \underline{\zeta} & \underline{\mu} \end{pmatrix} \begin{pmatrix} p \\ \mathbf{u} \\ E \\ H \end{pmatrix}$$

$$\gamma_1 = \gamma^e + \gamma^o$$

$$\gamma_2 = \gamma^e - \gamma^o$$

$$\eta_1 = \eta^e + \eta^o$$

$$\eta_2 = \eta^e - \eta^o$$

$$\underline{\zeta}_1 = \underline{\zeta}^e + \underline{\zeta}^o$$

$$\underline{\zeta}_2 = \underline{\zeta}^e - \underline{\zeta}^o$$

etc....

$$\begin{aligned} \underline{\rho} &= \underline{\rho}^\top && \leftarrow \text{dynamic density } \in \text{Sym}^1 \\ \underline{\mu} &= \underline{\mu}^\top && \leftarrow \text{permeability } \in \text{Sym}^2 \\ \underline{\boldsymbol{\epsilon}} &= \underline{\boldsymbol{\epsilon}}^\top && \leftarrow \text{permittivity } \in \text{Sym}^2 \end{aligned}$$

$$\begin{aligned} \gamma &= \eta^o - \eta^e && \leftarrow \text{acoustic bianisotropy vector}^1 \\ \underline{\boldsymbol{\xi}} &= (\underline{\zeta}^o)^\top - (\underline{\zeta}^e)^\top && \leftarrow \text{electromagnetic bianisotropy tensor}^2 \\ \underline{\mathbf{v}} &= (\underline{\mathbf{w}}^o)^\top - (\underline{\mathbf{w}}^e)^\top && \leftarrow \text{electromomentum coupling tensor}^3 \\ \underline{\mathbf{m}} &= -(\underline{\mathbf{n}}^o)^\top + (\underline{\mathbf{n}}^e)^\top && \leftarrow \text{magnetomomentum coupling tensor} \\ \mathbf{b} &= \mathbf{c}^o - \mathbf{c}^e && = "d," \text{ conventional piezoelectric constant in strain-charge form}^3 \\ \mathbf{d} &= -\mathbf{e}^o + \mathbf{e}^e && \leftarrow \text{piezomagnetic coupling vector} \end{aligned}$$

All constitutive properties are generally *dispersive*, i.e., effective properties can change dramatically if sub-wavelength resonances exist in the microstructure.

¹Sieck et al., PRB 96, 104303 (2017)

²Kong, Proc. IEEE 60, 9 (1972)

³Pernas-Salomón and Shmuel, J. Mech. Phys. Solids 134, 103770 (2020)

Passivity

Lossless, passive condition

$$\Re(\nabla \cdot I + \nabla \cdot S) = 0$$

Time-avg. Poynting vectors

$$I = p\mathbf{u}^*/2, S = \mathbf{E} \times \mathbf{H}^*/2$$

Time-harmonic governing eqs.

$$\nabla p = i\omega \boldsymbol{\mu} \quad \leftarrow \text{momentum}$$

$$\nabla \cdot \mathbf{u} = -i\omega \varepsilon \quad \leftarrow \text{continuity}$$

$$\nabla \times \mathbf{E} = i\omega \mathbf{B} \quad \leftarrow \text{Faraday}$$

$$\nabla \times \mathbf{H} = -i\omega \mathbf{D} \quad \leftarrow \text{Ampere}$$

Acousto-electrodynamic constitutive relations

$$\begin{pmatrix} \varepsilon \\ \boldsymbol{\mu} \\ \mathbf{D} \\ \mathbf{B} \end{pmatrix} = \begin{pmatrix} -\beta & \boldsymbol{\gamma}^\top & \mathbf{b}^\top & \mathbf{d}^\top \\ -\eta & \underline{\rho} & \underline{\mathbf{v}} & \underline{\mathbf{m}} \\ c & \underline{\mathbf{w}} & \underline{\epsilon} & \underline{\xi} \\ e & \underline{\mathbf{n}} & \underline{\zeta} & \underline{\mu} \end{pmatrix} \begin{pmatrix} p \\ \mathbf{u} \\ \mathbf{E} \\ \mathbf{H} \end{pmatrix}$$

substitute in constitutive relations

$$\frac{i\omega}{4}(\boldsymbol{\mu} \cdot \mathbf{u}^* - \boldsymbol{\mu}^* \cdot \mathbf{u} + \varepsilon^* p - \varepsilon p^* + \mathbf{B} \cdot \mathbf{H}^* - \mathbf{B}^* \cdot \mathbf{H} - \mathbf{D}^* \cdot \mathbf{E} + \mathbf{D} \cdot \mathbf{E}^*) = 0$$

$$\operatorname{Im} \beta = 0$$

$$\underline{\rho}^\dagger = \underline{\rho}$$

$$\eta = \gamma^*$$

$$\underline{\mu} = \underline{\mu}^\dagger$$

$$\underline{\epsilon} = \underline{\epsilon}^\dagger$$

$$\underline{\zeta} = \underline{\xi}^\dagger$$

← compressibility $\in \Re^1$

← dynamic density¹

← acous. bianisotropy¹

← permeability²

← permittivity²

← EM bianisotropy²

$$\underline{\mathbf{w}}^\dagger = \underline{\mathbf{v}}$$

$$\underline{\mathbf{n}}^\dagger = \underline{\mathbf{m}}$$

$$c = -\mathbf{b}^*$$

$$\mathbf{d} = -\mathbf{e}^*$$

← electromomentum³

← magnetomomentum

← piezoelectricity³

← piezomagnetism

Define $\tilde{p} = -p = \operatorname{tr}(\boldsymbol{\alpha})/3$

$$\begin{pmatrix} \varepsilon \\ \boldsymbol{\mu} \\ \mathbf{D} \\ \mathbf{B} \end{pmatrix} = \begin{pmatrix} \beta & \boldsymbol{\gamma}^\top & \mathbf{b}^\top & \mathbf{d}^\top \\ \gamma^* & \underline{\rho} & \underline{\mathbf{v}} & \underline{\mathbf{m}} \\ b^* & \underline{\mathbf{v}}^\dagger & \underline{\epsilon} & \underline{\xi} \\ d^* & \underline{\mathbf{m}}^\dagger & \underline{\zeta}^\dagger & \underline{\mu} \end{pmatrix} \begin{pmatrix} \tilde{p} \\ \mathbf{u} \\ \mathbf{E} \\ \mathbf{H} \end{pmatrix}$$

What about a reciprocal and passive acousto-electrodynamic medium?

¹Sieck et al., PRB **96**, 104303 (2017)

²Kong, Proc. IEEE **60**, 9 (1972)

³Pernas-Salomón and Shmuel, J. Mech. Phys. Solids **134**, 103770 (2020)

Reciprocity and passivity

Hermitian tensors

$$\underline{\rho} = \underline{\rho}^\dagger$$

$$\underline{\epsilon} = \underline{\epsilon}^\dagger$$

$$\underline{\mu} = \underline{\mu}^\dagger$$

Lossless, passive, linear media

Real symmetric tensors

$$\underline{\rho} = \underline{\rho}^\top$$

$$\underline{\epsilon} = \underline{\epsilon}^\top$$

$$\underline{\mu} = \underline{\mu}^\top$$

Lossless reciprocal media

To satisfy reciprocity and passivity, $\underline{\rho}$, $\underline{\epsilon}$, and $\underline{\mu}$ must be real and symmetric.

$\Im(\cdot) = 0$ for lossless case

$$\Im(\underline{\rho}) \geq 0$$

$$\Im(\beta) \geq 0$$

$$\Im(\underline{\epsilon}) \geq 0$$

$$\Im(\underline{\mu}) \geq 0$$

Results from reciprocity

$$\underline{\rho} = \underline{\rho}^\top$$

$$\underline{\mu} = \underline{\mu}^\top$$

$$\underline{\epsilon} = \underline{\epsilon}^\top$$

$$\gamma = \eta^o - \eta^e$$

$$\underline{\xi} = (\underline{\zeta}^o)^\top - (\underline{\zeta}^e)^\top$$

$$\underline{v} = (\underline{w}^o)^\top - (\underline{w}^e)^\top$$

$$\underline{m} = -(\underline{n}^o)^\top + (\underline{n}^e)^\top$$

$$\underline{b} = \underline{c}^o - \underline{c}^e$$

$$\underline{d} = -\underline{e}^o + \underline{e}^e$$

Passive condition: $\Re(\nabla \cdot I + \nabla \cdot S) \geq 0$ (equality for lossless case)

$$\frac{i\omega}{4}(\mu \cdot \underline{u}^* - \mu^* \cdot \underline{u} + \epsilon^* p - \epsilon p^* + \underline{B} \cdot \underline{H}^* - \underline{B}^* \cdot \underline{H} - \underline{D}^* \cdot \underline{E} + \underline{D} \cdot \underline{E}^*) \geq 0$$

- Substitute results from reciprocity into lossless, passive condition
- Group like terms and insert definitions following Sieck et al.¹
- Note that $\Re iz = -\Im z$ for complex scalar, vector, or tensor z
- Since fields are arbitrary, require $\Im(\text{all } \chi) = 0$ to satisfy equality

Definitions following Sieck¹

$$\underline{\eta}^o = \underline{\chi}_{pu}^o, \quad \underline{\eta}^e = i\underline{\chi}_{pu}^e$$

$$\underline{\zeta}^o = \underline{\chi}_{EH}^o, \quad \underline{\zeta}^e = i\underline{\chi}_{EH}^e$$

$$\underline{n}^o = i\underline{\chi}_{uH}^o, \quad \underline{n}^e = \underline{\chi}_{uH}^e$$

$$\underline{c}^o = i\underline{\chi}_{pE}^o, \quad \underline{c}^e = \underline{\chi}_{pE}^e$$

$$\underline{w}^o = \underline{\chi}_{uE}^o, \quad \underline{w}^e = i\underline{\chi}_{uE}^e$$

$$\underline{e}^o = \underline{\chi}_{pH}^o, \quad \underline{e}^e = i\underline{\chi}_{pH}^e$$

$\Im(\cdot) = 0$ for lossless case

$$\Im(\text{all } \chi) \geq 0$$

Causality

- Effect cannot precede cause
- Guaranteed for passive media
- Describes how β , ρ , ϵ , and μ depend on ω

$A(\omega)$ represents any constitutive quantity

$$A(\omega) = A'(\omega) + iA''(\omega)$$

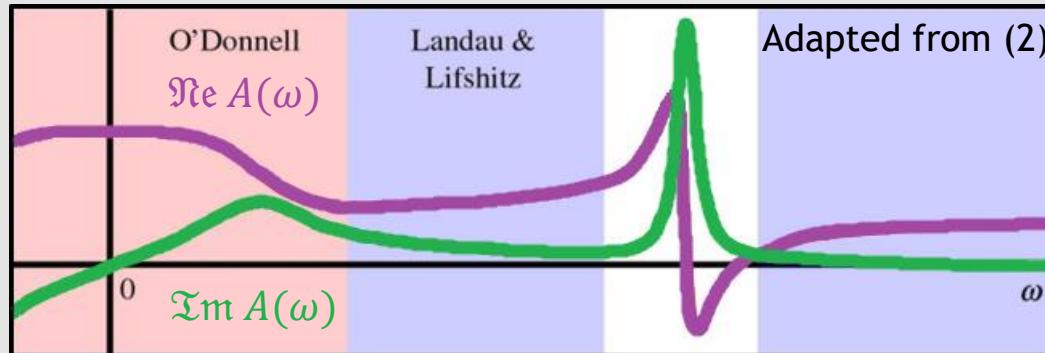
(A can be a scalar, vector, or tensor)

$A'(\infty) = \text{frozen state}$
 $(\omega \gg 2\pi/\tau)$

Kramers-Kronig relates $A''(\omega)$ to $A'(\omega) - A'(\infty)$ ^{1,2}

$$A'(\omega) - A'(\infty) = \frac{2}{\pi} \mathcal{P} \int_0^\infty \frac{\lambda A''(\lambda)}{\lambda^2 - \omega^2} d\lambda$$

$$A''(\omega) = -\frac{2}{\pi} \mathcal{P} \int_0^\infty \frac{\omega A'(\lambda)}{\lambda^2 - \omega^2} d\lambda$$



$\lambda = \omega$ dominates $\lambda = \omega$ is negligible

Relaxation dominates³

$$A''(\omega) \simeq -\frac{\pi}{2} \omega \frac{\partial A'}{\partial \omega}$$

Loss is negligible:¹ $A''(\omega) \simeq 0$

$$\frac{\partial A'}{\partial \omega} = \frac{4}{\pi} \int_0^\infty \frac{\lambda \omega}{(\lambda^2 - \omega^2)^2} A''(\lambda) d\lambda.$$

Recall that passivity requires that $A''(\omega) \geq 0$.

$$\frac{\partial}{\partial \omega} \operatorname{Re} A(\omega) < 0$$

$$\frac{\partial}{\partial \omega} \operatorname{Re} A(\omega) > 0$$

Relevant case away from bandgaps and resonances

- Storage and loss moduli related by Kramers-Kronig relations

- Measured/predicted dispersion must obey

$$\frac{\partial \beta}{\partial \omega} \geq 0, \quad \frac{\partial \rho}{\partial \omega} \geq 0, \quad \frac{\partial \epsilon}{\partial \omega} \geq 0, \quad \frac{\partial \mu}{\partial \omega} \geq 0.$$

¹ L. D. Landau and E. M. Lifshitz. Electrodynamics of continuous media (1960)

² M. B. Muhlestein et al., P. Roy Soc. A - Math Phy. (2016)

³ M. O'Donnell et al., J. Acoust. Soc. Am. (1981)

R. Christensen, Theory of viscoelasticity: An Introduction (2012)

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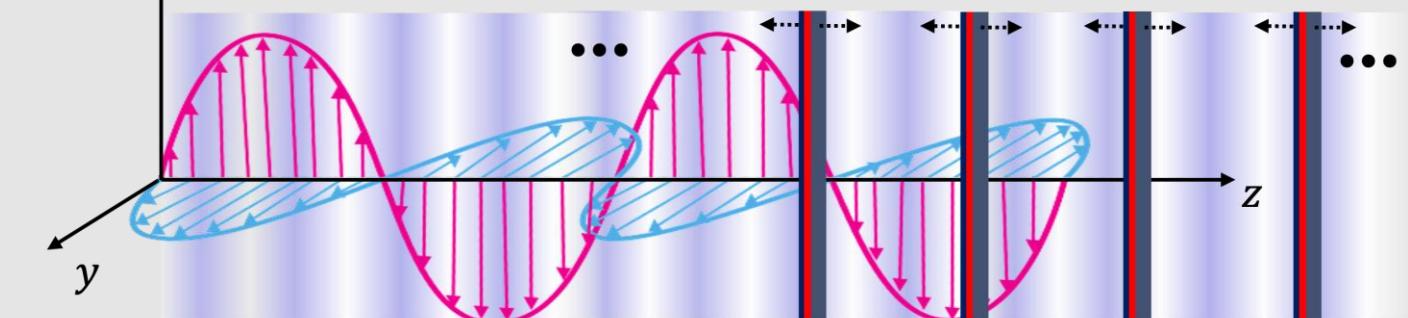
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Verification of source-driven homogenization theory

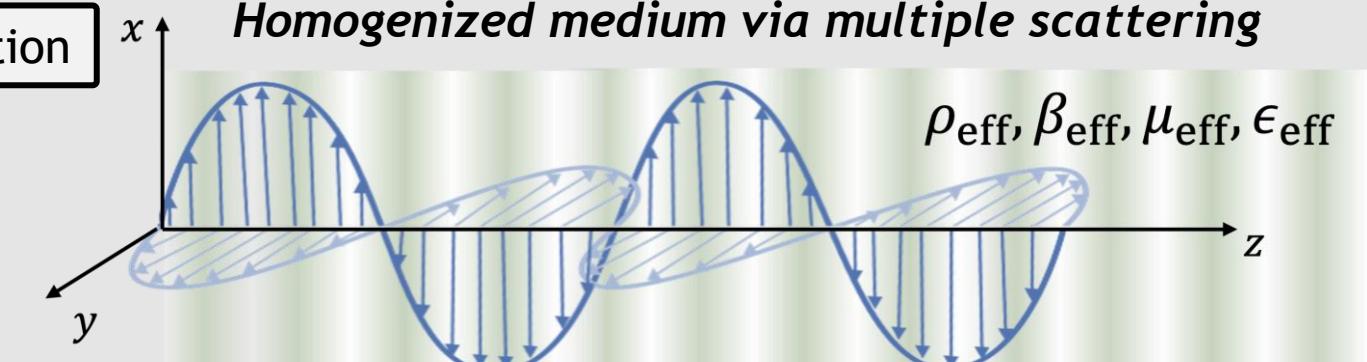
periodic passive identical
 1D reciprocal $ka, \kappa a \ll 1$

Micromechanical behaviour due to asymmetric scatterers



Homogenization

Homogenized medium via multiple scattering



Effective constitutive relations
obey the symmetries mandated
by reciprocity and passivity

$$\begin{pmatrix} \varepsilon_{\text{eff}} \\ \mu_{v,\text{eff}} \\ D_{\text{eff}} \\ B_{\text{eff}} \end{pmatrix} = \begin{pmatrix} -\beta_{\text{eff}} & \chi_{pu}^o - i\chi_{pu}^e & i\chi_{pE}^o - \chi_{pE}^e & -(\chi_{pH}^o - i\chi_{pH}^e) \\ -(\chi_{pu}^o + i\chi_{pu}^e) & \rho_{\text{eff}} & \chi_{uE}^o - i\chi_{uE}^e & -i\chi_{uH}^o + \chi_{uH}^e \\ i\chi_{pE}^o + \chi_{pE}^e & \chi_{uE}^o + i\chi_{uE}^e & \epsilon_{\text{eff}} & \chi_{EH}^o \\ \chi_{pH}^o + i\chi_{pH}^e & i\chi_{uH}^o + \chi_{uH}^e & \chi_{EH}^o & \mu_{\text{eff}} \end{pmatrix} \begin{pmatrix} p_{\text{eff}} \\ u_{\text{eff}} \\ E_{\text{eff}} \\ H_{\text{eff}} \end{pmatrix}$$

Thanks for listening!

Summary

- Derived constraints on acousto-electrodynamic coupled media imposed by reciprocity, passivity, and causality
- Recovered relations for dynamic piezoelectricity, acoustic bianisotropy, and electrodynamic bianisotropy
- Showed existing source-driven homogenization theory for coupled acousto-electrodynamic media obeys reciprocity and passivity

Student award competition:



Acknowledgments

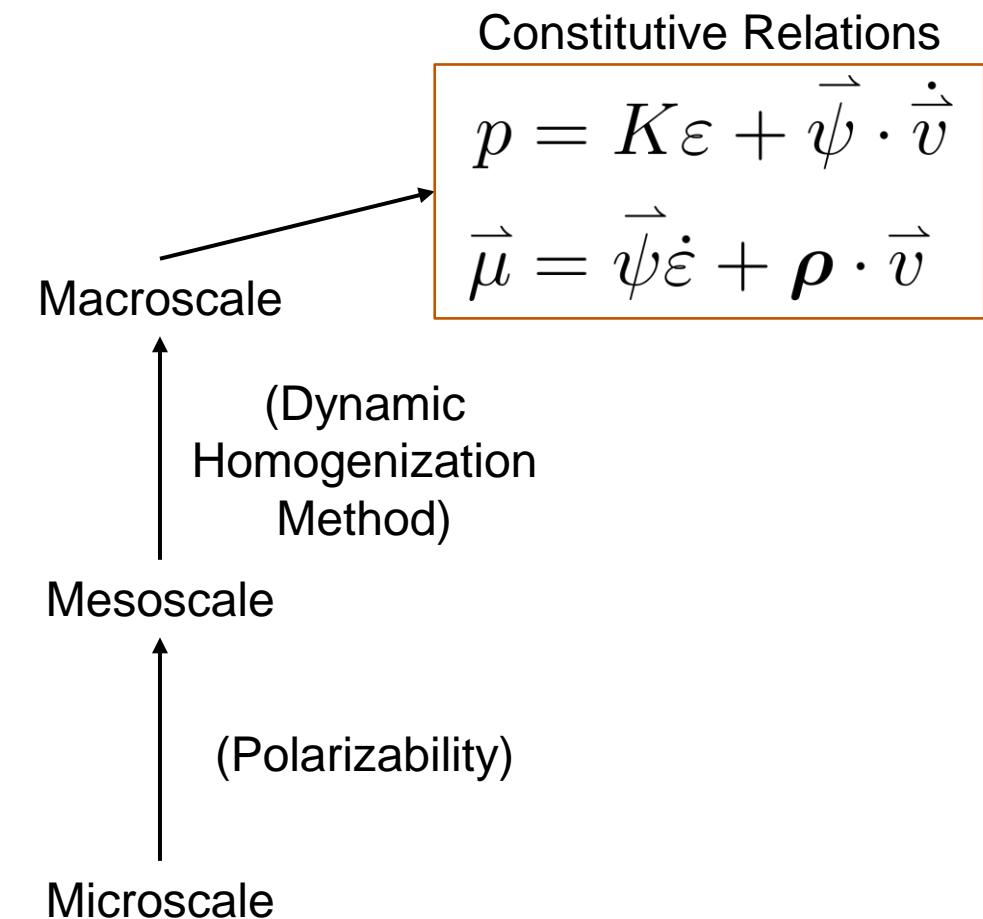
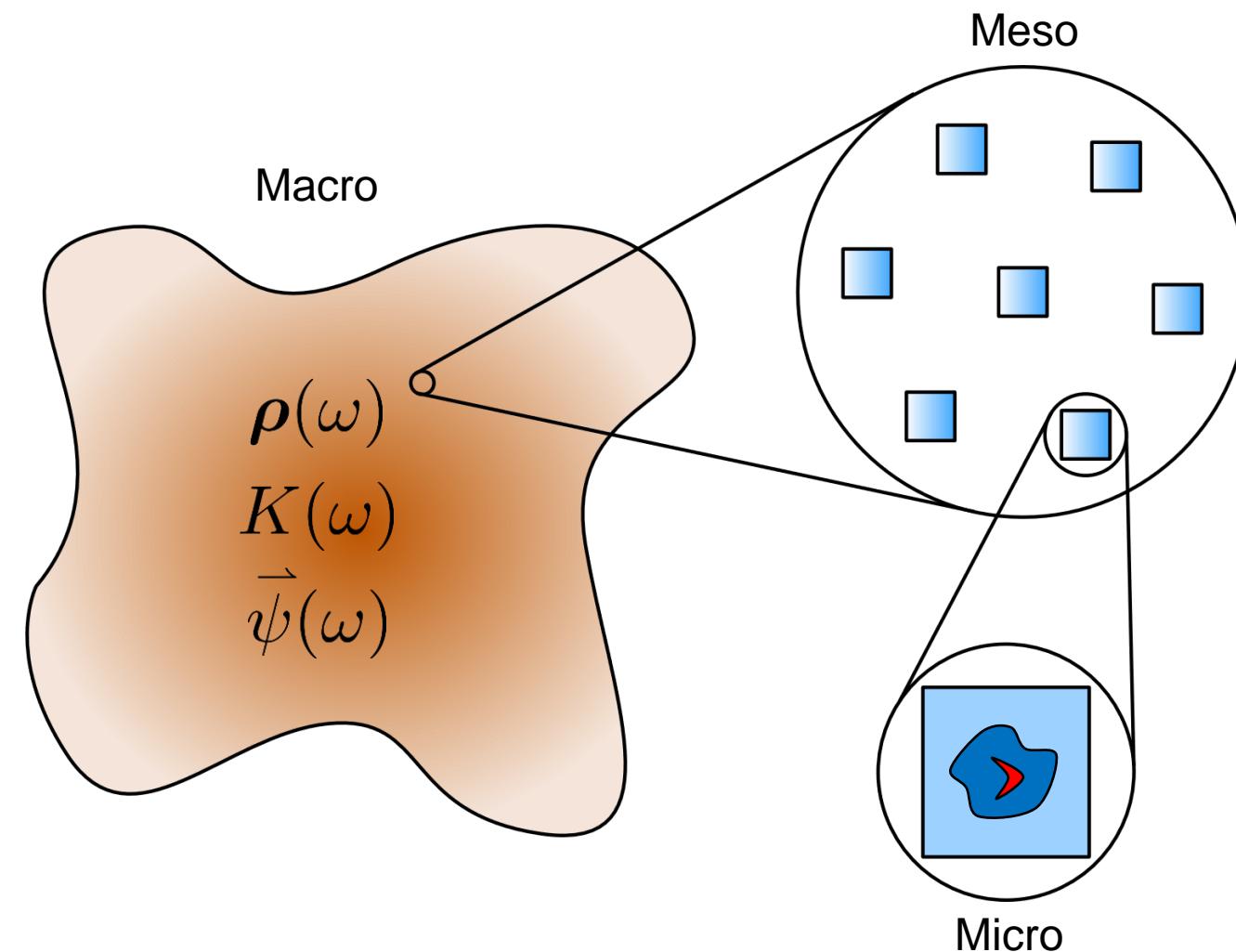
ARL:UT Chester M. McKinney Graduate Fellowship in Acoustics
Defense Advanced Research Projects Agency (DARPA)



ACOUSTICS 23
Sydney

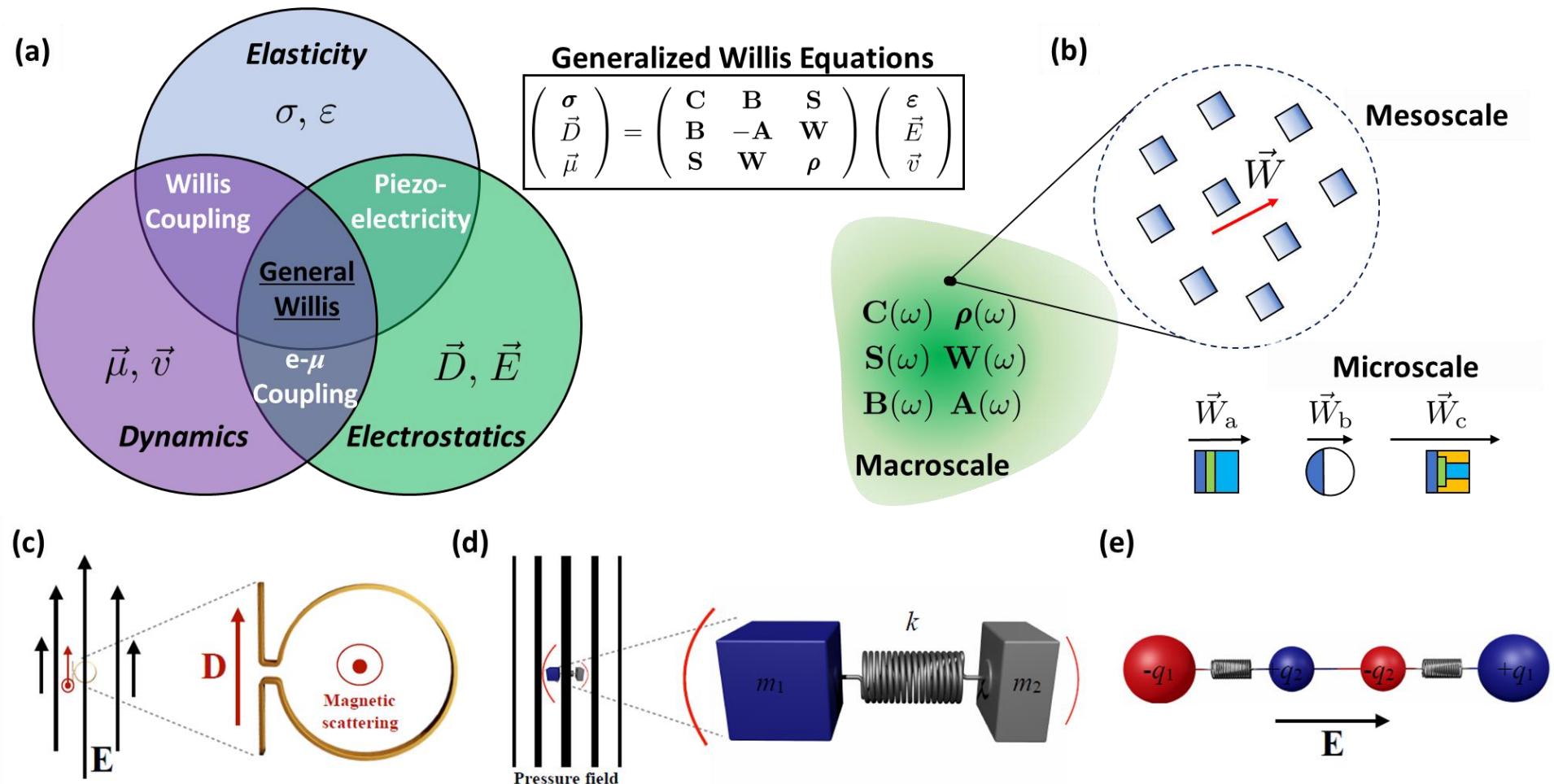


Extra slides



Sieck, C.F., Alu, A., Haberman, M.R., Phys. Rev. B, **96**, 104303, (2017).

Electro-momentum Coupling

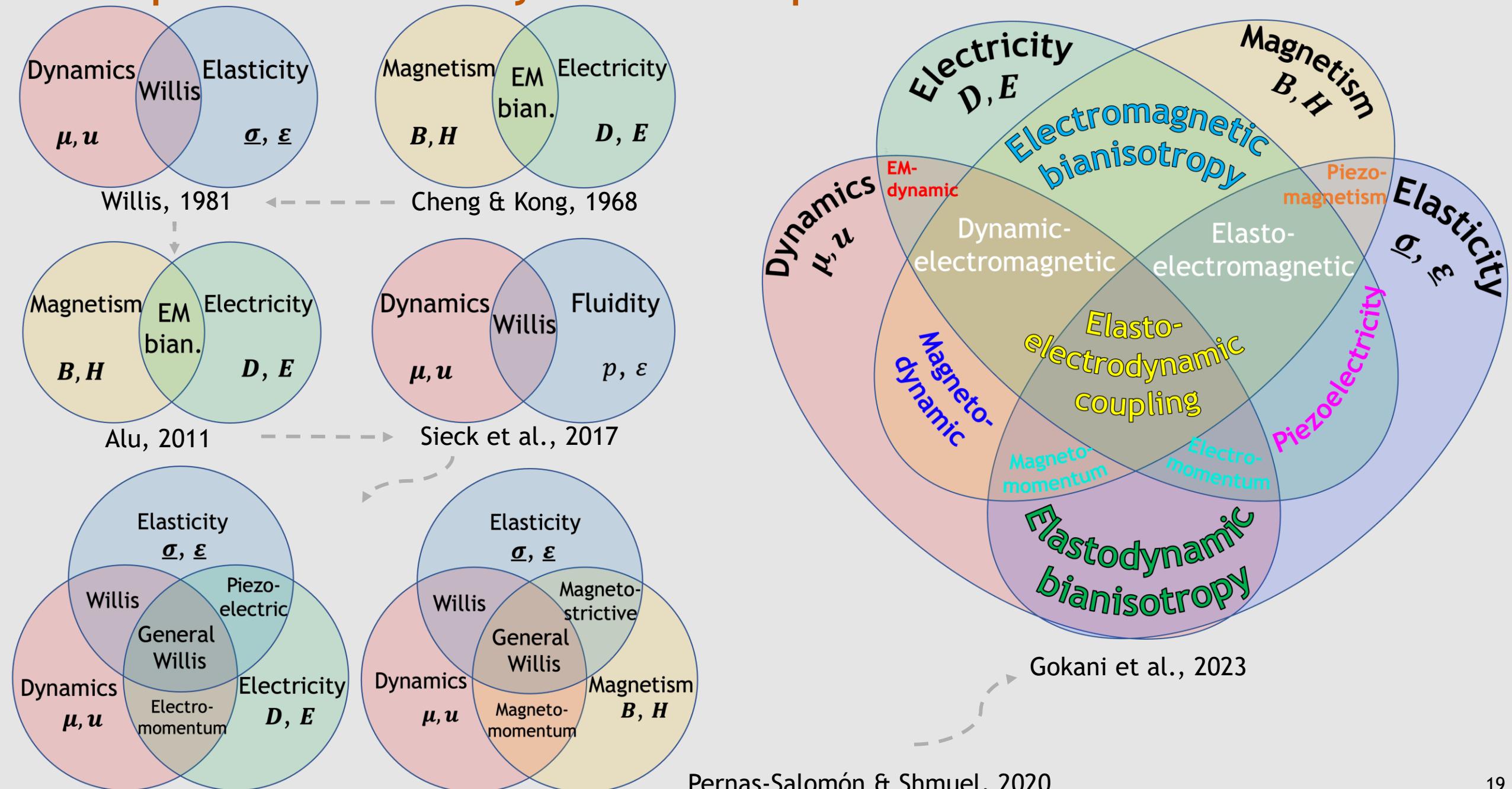


Pernas-Salomon and Shmuel, J. Mech. Phys. Solids 134, 103770 (2020)

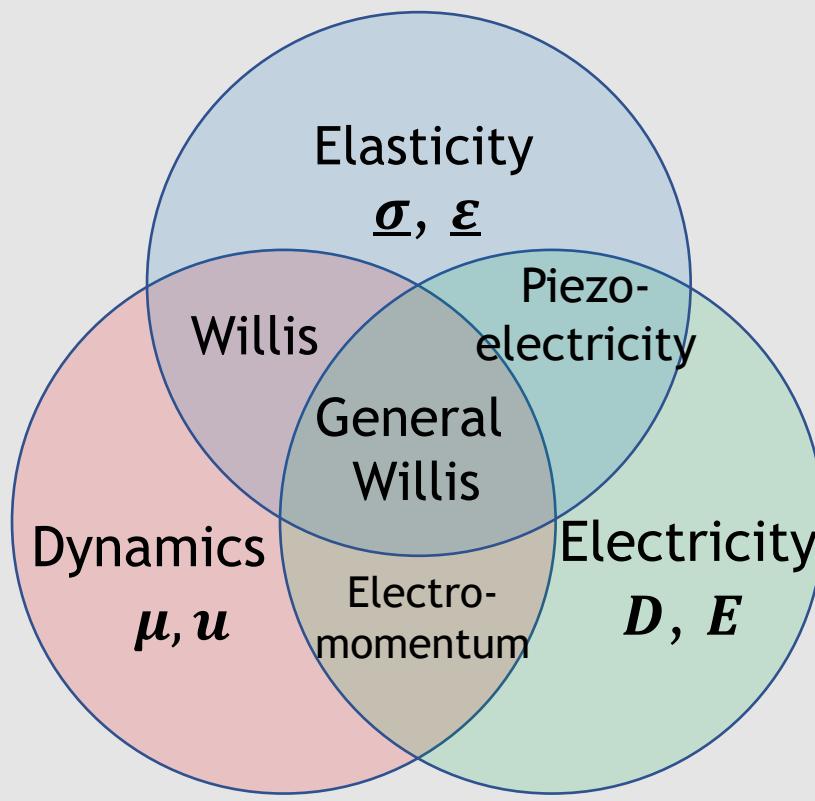
Pernas-Salomon and Shmuel, Phys. Rev. Appl. 14, 064005 (2020)

Pernas-Salomon, Haberman, Norris and Shmuel, Wave Motion 106, 102797 (2021)

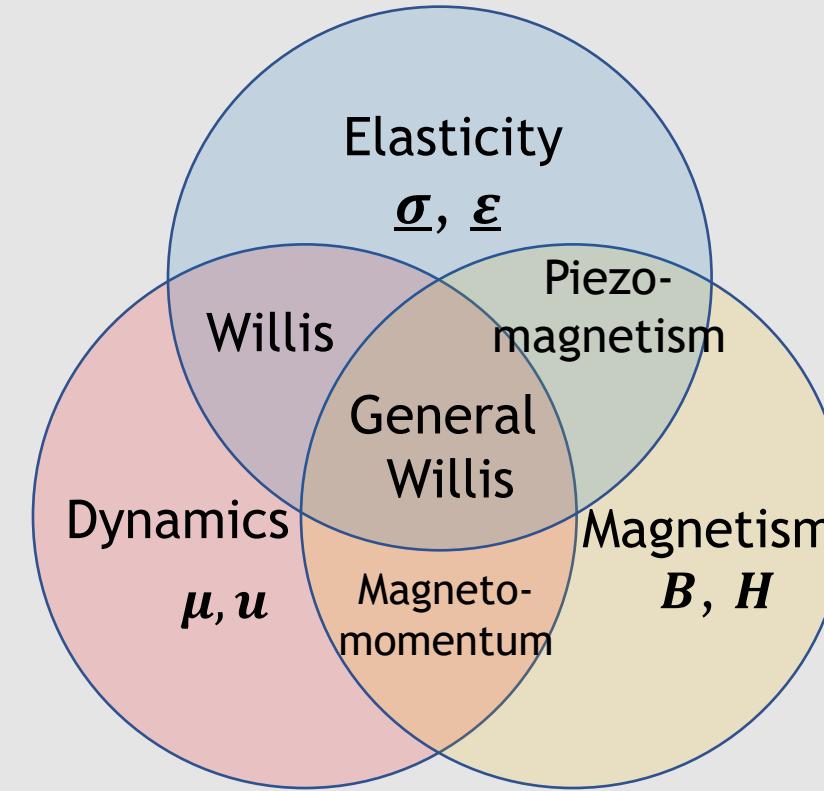
Developments in the study of bianisotropic media



Electromomentum coupling usually does not consider electrodynamics

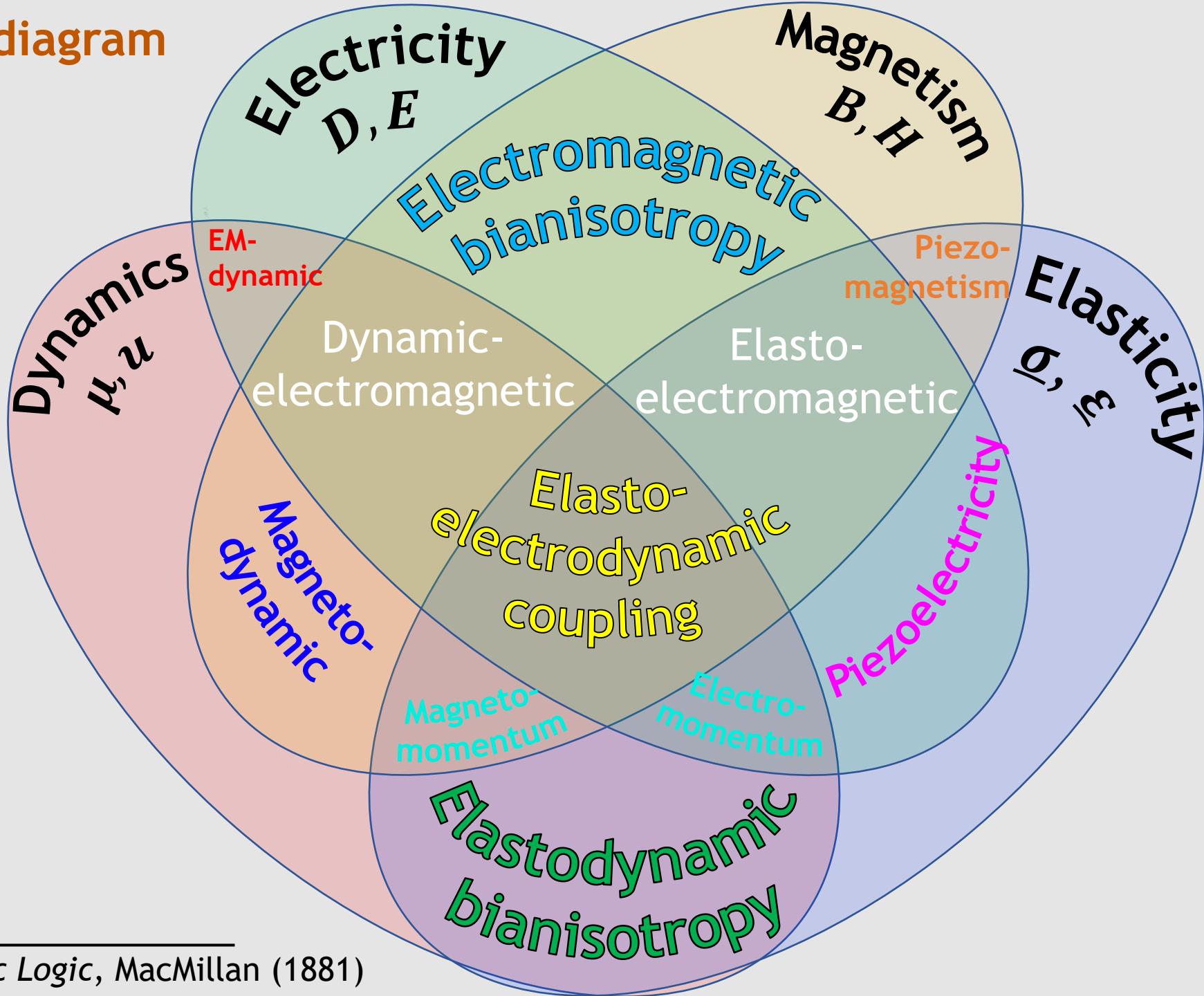


quasielectrostatic



quasimagnetostatic

4-way Venn diagram



Reciprocity

$$\int_A (p_1 \mathbf{u}_2 - p_2 \mathbf{u}_1 + \mathbf{E}_1 \times \mathbf{H}_2 - \mathbf{E}_2 \times \mathbf{H}_1) \cdot \hat{\mathbf{n}} dA = 0 \text{ by the Sommerfeld radiation condition}$$

$$= \int_V (\mathbf{u}_2 \cdot \mathbf{f}_1 - \mathbf{u}_1 \cdot \mathbf{f}_2 + p_1 q_2 - p_2 q_1 + \mathbf{E}_2 \cdot \mathbf{J}_1 - \mathbf{E}_1 \cdot \mathbf{J}_2 + \mathbf{H}_1 \cdot \mathbf{K}_2 - \mathbf{H}_2 \cdot \mathbf{K}_1) dV = 0 \text{ by reciprocity}$$

$$- i\omega \int_V (\underbrace{\mathbf{u}_1 \cdot \boldsymbol{\mu}_2 - \mathbf{u}_2 \cdot \boldsymbol{\mu}_1 + p_1 \boldsymbol{\varepsilon}_2 - p_2 \boldsymbol{\varepsilon}_1 + \mathbf{H}_1 \cdot \mathbf{B}_2 - \mathbf{H}_2 \cdot \mathbf{B}_1 - \mathbf{E}_1 \cdot \mathbf{D}_2 + \mathbf{E}_2 \cdot \mathbf{D}_1}_{\equiv 0}) dV$$

Acousto-electrodynamic constitutive relations

$$\begin{pmatrix} \boldsymbol{\varepsilon} \\ \boldsymbol{\mu} \\ \mathbf{D} \\ \mathbf{B} \end{pmatrix} = \begin{pmatrix} -\beta & \boldsymbol{\gamma}^\top & \mathbf{b}^\top & \mathbf{d}^\top \\ -\boldsymbol{\eta} & \underline{\rho} & \underline{\mathbf{v}} & \underline{\mathbf{m}} \\ \mathbf{c} & \underline{\mathbf{w}} & \underline{\boldsymbol{\epsilon}} & \underline{\boldsymbol{\xi}} \\ \mathbf{e} & \underline{\mathbf{n}} & \underline{\boldsymbol{\zeta}} & \underline{\boldsymbol{\mu}} \end{pmatrix} \begin{pmatrix} \mathbf{p} \\ \mathbf{u} \\ \mathbf{E} \\ \mathbf{H} \end{pmatrix}$$

$$\gamma_1 = \gamma^e + \gamma^o$$

$$\gamma_2 = \gamma^e - \gamma^o$$

$$\eta_1 = \eta^e + \eta^o$$

$$\eta_2 = \eta^e - \eta^o$$

$$\underline{\boldsymbol{\zeta}}_1 = \underline{\boldsymbol{\zeta}}^e + \underline{\boldsymbol{\zeta}}^o$$

$$\underline{\boldsymbol{\zeta}}_2 = \underline{\boldsymbol{\zeta}}^e - \underline{\boldsymbol{\zeta}}^o$$

etc....

$$\underline{\boldsymbol{\rho}} = \underline{\boldsymbol{\rho}}^\top, \quad \underline{\boldsymbol{\mu}} = \underline{\boldsymbol{\mu}}^\top, \quad \underline{\boldsymbol{\epsilon}} = \underline{\boldsymbol{\epsilon}}^\top$$

$$\begin{aligned} & \mathbf{u}_1 \cdot [(\underline{\boldsymbol{\rho}} - \underline{\boldsymbol{\rho}}^\top) \mathbf{u}_2] + \mathbf{H}_1 \cdot [(\underline{\boldsymbol{\mu}} - \underline{\boldsymbol{\mu}}^\top) \mathbf{H}_2] + \mathbf{E}_2 \cdot [(\underline{\boldsymbol{\epsilon}} - \underline{\boldsymbol{\epsilon}}^\top) \mathbf{E}_1] \\ & + (\boldsymbol{\eta}^e + \boldsymbol{\gamma}^e) \cdot (p_1 \mathbf{u}_2 - p_2 \mathbf{u}_1) + (\boldsymbol{\eta}^o - \boldsymbol{\gamma}^o) \cdot (p_1 \mathbf{u}_2 + p_2 \mathbf{u}_1) \\ & - \mathbf{E}_1 \cdot [(\underline{\boldsymbol{\zeta}}^e)^\top + \underline{\boldsymbol{\xi}}^e] \mathbf{H}_2 - \mathbf{E}_1 \cdot [(\underline{\boldsymbol{\zeta}}^o)^\top - \underline{\boldsymbol{\xi}}^o] \mathbf{H}_2 \\ & + \mathbf{E}_2 \cdot [(\underline{\boldsymbol{\zeta}}^e)^\top + \underline{\boldsymbol{\xi}}^e] \mathbf{H}_1 - \mathbf{E}_2 \cdot [(\underline{\boldsymbol{\zeta}}^o)^\top - \underline{\boldsymbol{\xi}}^o] \mathbf{H}_1 \\ & + \mathbf{u}_1 \cdot [\underline{\mathbf{v}}^e + (\underline{\mathbf{w}}^e)^\top] \mathbf{E}_2 - \mathbf{u}_1 \cdot [\underline{\mathbf{v}}^o - (\underline{\mathbf{w}}^o)^\top] \mathbf{E}_2 \\ & - \mathbf{u}_2 \cdot [(\underline{\mathbf{w}}^e)^\top + \underline{\mathbf{v}}^e] \mathbf{E}_1 + \mathbf{u}_2 \cdot [(\underline{\mathbf{w}}^o)^\top - \underline{\mathbf{v}}^o] \mathbf{E}_1 \\ & + \mathbf{u}_1 \cdot [\underline{\mathbf{m}}^e - (\underline{\mathbf{n}}^e)^\top] \mathbf{H}_2 - \mathbf{u}_1 \cdot [\underline{\mathbf{m}}^o + (\underline{\mathbf{n}}^o)^\top] \mathbf{H}_2 \\ & + \mathbf{u}_2 \cdot [(\underline{\mathbf{n}}^e)^\top - \underline{\mathbf{m}}^e] \mathbf{H}_1 - \mathbf{u}_2 \cdot [(\underline{\mathbf{n}}^o)^\top + \underline{\mathbf{m}}^o] \mathbf{H}_1 \\ & + p_1 (\mathbf{b}^e + \mathbf{c}^e) \cdot \mathbf{E}_2 - p_1 (\mathbf{b}^o - \mathbf{c}^o) \cdot \mathbf{E}_2 \\ & - p_2 (\mathbf{c}^e + \mathbf{b}^e) \cdot \mathbf{E}_1 + p_2 (\mathbf{c}^o - \mathbf{b}^o) \cdot \mathbf{E}_1 \\ & + p_2 (\mathbf{e}^e - \mathbf{d}^e) \cdot \mathbf{H}_1 - p_2 (\mathbf{e}^o + \mathbf{d}^o) \cdot \mathbf{H}_1 \\ & + p_1 (\mathbf{d}^e - \mathbf{e}^e) \cdot \mathbf{H}_2 - p_1 (\mathbf{d}^o + \mathbf{e}^o) \cdot \mathbf{H}_2 = 0 \end{aligned}$$

$$\begin{aligned} \boldsymbol{\gamma} &= \boldsymbol{\eta}^o - \boldsymbol{\eta}^e \\ \underline{\boldsymbol{\xi}} &= (\underline{\boldsymbol{\zeta}}^o)^\top - (\underline{\boldsymbol{\zeta}}^e)^\top \\ \underline{\mathbf{v}} &= (\underline{\mathbf{w}}^o)^\top - (\underline{\mathbf{w}}^e)^\top \\ \underline{\mathbf{m}} &= (\underline{\mathbf{n}}^e)^\top - (\underline{\mathbf{n}}^o)^\top \\ \mathbf{b} &= \mathbf{c}^o - \mathbf{c}^e \\ \mathbf{d} &= \mathbf{e}^e - \mathbf{e}^o \end{aligned}$$

} new results

Passivity

Lossless, passive condition

$$\Re(\nabla \cdot I + \nabla \cdot S) = 0$$

Time-avg. Poynting vectors

$$I = p\mathbf{u}^*/2, S = \mathbf{E} \times \mathbf{H}^*/2$$

Time-harmonic governing eqs.

$$\nabla p = i\omega \boldsymbol{\mu} \leftarrow \text{momentum}$$

$$\nabla \cdot \mathbf{u} = -i\omega \varepsilon \leftarrow \text{continuity}$$

$$\nabla \times \mathbf{E} = i\omega \mathbf{B} \leftarrow \text{Faraday}$$

$$\nabla \times \mathbf{H} = -i\omega \mathbf{D} \leftarrow \text{Ampere}$$

Acousto-electrodynamic constitutive relations

$$\begin{pmatrix} \varepsilon \\ \boldsymbol{\mu} \\ \mathbf{D} \\ \mathbf{B} \end{pmatrix} = \begin{pmatrix} -\beta & \gamma^\top & \mathbf{b}^\top & \mathbf{d}^\top \\ -\eta & \underline{\rho} & \underline{\mathbf{v}} & \underline{\mathbf{m}} \\ c & \underline{\mathbf{w}} & \underline{\epsilon} & \underline{\xi} \\ e & \underline{\mathbf{n}} & \underline{\zeta} & \underline{\mu} \end{pmatrix} \begin{pmatrix} p \\ \mathbf{u} \\ \mathbf{E} \\ \mathbf{H} \end{pmatrix}$$

$$\frac{i\omega}{4}(\boldsymbol{\mu} \cdot \mathbf{u}^* - \boldsymbol{\mu}^* \cdot \mathbf{u} + \varepsilon^* p - \varepsilon p^* + \mathbf{B} \cdot \mathbf{H}^* - \mathbf{B}^* \cdot \mathbf{H} - \mathbf{D}^* \cdot \mathbf{E} + \mathbf{D} \cdot \mathbf{E}^*) = 0$$

$$\begin{aligned} \frac{i\omega}{4} &[-\mathbf{u}^* \cdot (\underline{\rho}^\dagger - \underline{\rho}) \mathbf{u} + 2i\operatorname{Im}(\beta)|p|^2 - \\ &(\eta - \gamma^*) \cdot p\mathbf{u}^* + (\eta^* - \gamma) \cdot p^*\mathbf{u} \\ &+ \mathbf{E}^* \cdot (\underline{\epsilon} - \underline{\epsilon}^\dagger) \mathbf{E} + \mathbf{H}^* \cdot (\underline{\mu} - \underline{\mu}^\dagger) \mathbf{H} \\ &- \mathbf{E}^* \cdot (\underline{\zeta}^\dagger - \underline{\xi}) \mathbf{H} + \mathbf{H}^* \cdot (\underline{\zeta} - \underline{\xi}^\dagger) \mathbf{E} + \mathbf{u}^* \cdot (\underline{\mathbf{v}} - \underline{\mathbf{w}}^\dagger) \mathbf{E} \\ &+ \mathbf{u}^* \cdot (\underline{\mathbf{m}} - \underline{\mathbf{n}}^\dagger) \mathbf{H} + \mathbf{E}^* \cdot (\underline{\mathbf{w}} - \underline{\mathbf{v}}^\dagger) \mathbf{u} + \mathbf{H}^* \cdot (\underline{\mathbf{n}} - \underline{\mathbf{m}}^\dagger) \mathbf{u} \\ &+ p(\mathbf{b}^* + \mathbf{c}) \cdot \mathbf{E}^* + p(\mathbf{d}^* + \mathbf{e}) \cdot \mathbf{H}^* \\ &- p^*(\mathbf{b} + \mathbf{c}^*) \cdot \mathbf{E} - p^*(\mathbf{d} + \mathbf{e}^*) \cdot \mathbf{H}] = 0 \end{aligned}$$

$$\begin{aligned} \operatorname{Im}\beta &= 0 \\ \underline{\rho}^\dagger &= \underline{\rho} \\ \eta &= \gamma^* \\ \underline{\mu} &= \underline{\mu}^\dagger \\ \underline{\epsilon} &= \underline{\epsilon}^\dagger \\ \underline{\zeta} &= \underline{\xi}^\dagger \end{aligned}$$

$$\begin{aligned} \underline{\mathbf{w}}^\dagger &= \underline{\mathbf{v}} \\ \underline{\mathbf{n}}^\dagger &= \underline{\mathbf{m}} \\ \mathbf{c} &= -\mathbf{b}^* \\ \mathbf{d} &= -\mathbf{e}^* \end{aligned}$$

} new results

What about a reciprocal and passive acousto-electrodynamic medium?

Reciprocity and passivity

Results from reciprocity

$$\underline{\rho} = \underline{\rho}^T$$

$$\underline{\mu} = \underline{\mu}^T$$

$$\underline{\epsilon} = \underline{\epsilon}^T$$

$$\gamma = \eta^o - \eta^e$$

$$\underline{\xi} = (\underline{\zeta}^o)^T - (\underline{\zeta}^e)^T$$

$$\underline{v} = (\underline{w}^o)^T - (\underline{w}^e)^T$$

$$\underline{m} = (\underline{n}^e)^T - (\underline{n}^o)^T$$

$$\underline{b} = \underline{c}^o - \underline{c}^e$$

$$\underline{d} = \underline{e}^e - \underline{e}^o$$

Equation for $\Re e(\nabla \cdot I + \nabla \cdot S) \geq 0$ for reciprocity and passivity:

$$\omega \left[\frac{1}{2} \underline{u}^* \cdot \operatorname{Im}(\underline{\rho}) \underline{u} + \frac{1}{2} \operatorname{Im}(\beta) |p|^2 - \operatorname{Im}(\chi_{pu}^o) \cdot \Re e(p \underline{u}^*) + \operatorname{Im}(\chi_{pu}^e) \cdot \operatorname{Im}(p \underline{u}^*) \right.$$

$$+ \frac{1}{2} \underline{E}^* \cdot \operatorname{Im}(\underline{\epsilon}) \underline{H} + \frac{1}{2} \underline{H}^* \cdot \operatorname{Im}(\underline{\mu}) \underline{H} - \operatorname{Im}(\chi_{EH}^o) : \Re e(\underline{H}^* \otimes \underline{E}) - \operatorname{Im}(\chi_{EH}^e) : \operatorname{Im}(\underline{H} \otimes \underline{E}^*)$$

$$+ \operatorname{Im}(\chi_{uE}^o) : \Re e(\underline{E} \otimes \underline{u}^*) + \operatorname{Im}(\chi_{uE}^e) : \operatorname{Im}(\underline{E} \otimes \underline{u}^*) - \operatorname{Im}(\chi_{uH}^o) : \operatorname{Im}(\underline{H}^* \otimes \underline{u}) + \operatorname{Im}(\chi_{uH}^e) : \Re e(\underline{H} \otimes \underline{u})$$

$$\left. - \operatorname{Im}(\chi_{pE}^o) \cdot \operatorname{Im}(p \underline{E}^*) + \operatorname{Im}(\chi_{pE}^e) \cdot \Re e(p \underline{E}^*) + \operatorname{Im}(\chi_{pH}^o) \cdot \Re e(p \underline{H}^*) - \operatorname{Im}(\chi_{pH}^e) \cdot \operatorname{Im}(p \underline{H}^*) \right] \geq 0.$$

$$\Re e(\nabla \cdot I + \nabla \cdot S) \geq 0$$

$$-\frac{i\omega}{4}[-\underline{u}^* \cdot (\underline{\rho}^\dagger - \underline{\rho}) \underline{u} + 2i\operatorname{Im}(\beta)|p|^2 - (\eta - \gamma^*) \cdot p \underline{u}^* + (\eta^* - \gamma) \cdot p^* \underline{u}$$

$$+ \underline{E}^* \cdot (\underline{\epsilon} - \underline{\epsilon}^\dagger) \underline{E} + \underline{H}^* \cdot (\underline{\mu} - \underline{\mu}^\dagger) \underline{H} - \underline{E}^* \cdot (\underline{\zeta}^\dagger - \underline{\xi}) \underline{H} + \underline{H}^* \cdot (\underline{\zeta} - \underline{\xi}^\dagger) \underline{E}$$

$$+ \underline{u}^* \cdot (\underline{v} - \underline{w}^\dagger) \underline{E} + \underline{u}^* \cdot (\underline{m} - \underline{n}^\dagger) \underline{H} + \underline{E}^* \cdot (\underline{w} - \underline{v}^\dagger) \underline{u} + \underline{H}^* \cdot (\underline{n} - \underline{m}^\dagger) \underline{u}$$

$$+ p(\underline{b}^* + \underline{c}) \cdot \underline{E}^* + p(\underline{d}^* + \underline{e}) \cdot \underline{H}^* - p^*(\underline{b} + \underline{c}^*) \cdot \underline{E} - p^*(\underline{d} + \underline{e}^*) \cdot \underline{H}] \geq 0.$$

Definitions following Sieck¹

$$\begin{aligned} \eta^o &= \chi_{pu}^o & \eta^e &= i\chi_{pu}^e \\ \underline{\zeta}^o &= \underline{\chi}_{EH}^o & \underline{\zeta}^e &= i\underline{\chi}_{EH}^e \\ \underline{n}^o &= i\underline{\chi}_{uH}^o, & \underline{n}^e &= \underline{\chi}_{uH}^e \\ \underline{c}^o &= i\underline{\chi}_{pE}^o, & \underline{c}^e &= \underline{\chi}_{pE}^e \\ \underline{w}^o &= \underline{\chi}_{uE}^o, & \underline{w}^e &= i\underline{\chi}_{uE}^e \\ \underline{e}^o &= \underline{\chi}_{pH}^o, & \underline{e}^e &= i\underline{\chi}_{pH}^e \end{aligned}$$

Lossless case

$$\operatorname{Im}(\underline{\rho}) = 0$$

$$\operatorname{Im}(\beta) = 0$$

$$\operatorname{Im}(\underline{\epsilon}) = 0$$

$$\operatorname{Im}(\underline{\mu}) = 0$$

and

$$\operatorname{Im}(\text{all } \chi) = 0$$

Reciprocity and passivity

Reciprocity

$$\underline{\rho} = \underline{\rho}^T$$

$$\underline{\mu} = \underline{\mu}^T$$

$$\underline{\epsilon} = \underline{\epsilon}^T$$

$$\gamma = \eta^o - \eta^e$$

$$\underline{\xi} = (\underline{\zeta}^o)^T - (\underline{\zeta}^e)^T$$

$$\underline{v} = (\underline{w}^o)^T - (\underline{w}^e)^T$$

$$\underline{m} = (\underline{n}^e)^T - (\underline{n}^o)^T$$

$$\underline{b} = c^o - c^e$$

$$\underline{d} = e^e - e^o$$

+

Passivity

$$\begin{aligned}
-\Re(\nabla \cdot I + \nabla \cdot S) = & -\frac{i\omega}{4}[-\underline{u}^* \cdot (\underline{\rho}^\dagger - \underline{\rho})\underline{u} + 2i\Im(\beta)|p|^2 - (\eta - \gamma^*) \cdot p\underline{u}^* + (\eta^* - \gamma) \cdot p^*\underline{u} \\
& + E^* \cdot (\underline{\epsilon} - \underline{\epsilon}^\dagger)E + H^* \cdot (\underline{\mu} - \underline{\mu}^\dagger)H - E^* \cdot (\underline{\zeta}^\dagger - \underline{\xi})H + H^* \cdot (\underline{\zeta} - \underline{\xi}^\dagger)E \\
& + \underline{u}^* \cdot (\underline{v} - \underline{w}^\dagger)E + \underline{u}^* \cdot (\underline{m} - \underline{n}^\dagger)H + E^* \cdot (\underline{w} - \underline{v}^\dagger)\underline{u} + H^* \cdot (\underline{n} - \underline{m}^\dagger)\underline{u} \\
& + p(\underline{b}^* + c) \cdot E^* + p(\underline{d}^* + e) \cdot H^* - p^*(\underline{b} + \underline{c}^*) \cdot E - p^*(\underline{d} + \underline{e}^*) \cdot H] \geq 0
\end{aligned}$$

=

All \Im parts vanish for lossless case

$$\begin{aligned}
-\Re(\nabla \cdot I + \nabla \cdot S) = & \omega \left[\frac{1}{2} \underline{u}^* \cdot \Im(\underline{\rho})\underline{u} + \frac{1}{2} \Im(\beta)|p|^2 - \Im(\chi_{pu}^o) \cdot \Re(p\underline{u}^*) + \Im(\chi_{pu}^e) \cdot \Im(p\underline{u}^*) \right. \\
& + \frac{1}{2} E^* \cdot \Im(\underline{\epsilon})H + \frac{1}{2} H^* \cdot \Im(\underline{\mu})H - \Im(\chi_{EH}^o) : \Re(H^* \otimes E) - \Im(\chi_{EH}^e) : \Im(H \otimes E^*) \\
& + \Im(\chi_{uE}^o) : \Re(E \otimes \underline{u}^*) + \Im(\chi_{uE}^e) : \Im(E \otimes \underline{u}^*) - \Im(\chi_{uH}^o) : \Im(H^* \otimes \underline{u}) + \Im(\chi_{uH}^e) : \Re(H \otimes \underline{u}) \\
& \left. - \Im(\chi_{pE}^o) \cdot \Im(pE^*) + \Im(\chi_{pE}^e) \cdot \Re(pE^*) + \Im(\chi_{pH}^o) \cdot \Re(pH^*) - \Im(\chi_{pH}^e) \cdot \Im(pH^*) \right] \geq 0
\end{aligned}$$

Verification of source-driven constitutive relations

Acousto-electrodynamic constitutive relations

$$\begin{pmatrix} \varepsilon \\ \mu \\ D \\ B \end{pmatrix} = \begin{pmatrix} -\beta & \gamma^\top & \underline{b}^\top & \underline{d}^\top \\ -\eta & \underline{\rho} & \underline{v} & \underline{m} \\ c & \underline{w} & \underline{\epsilon} & \underline{\xi} \\ e & \underline{n} & \underline{\zeta} & \underline{\mu} \end{pmatrix} \begin{pmatrix} p \\ u \\ E \\ H \end{pmatrix}$$



$$\begin{aligned} \text{Im } \beta &= 0 \\ \underline{\rho}^\dagger &= \underline{\rho} \\ \underline{\eta} &= \underline{\gamma}^* \\ \underline{\mu} &= \underline{\mu}^\dagger \\ \underline{\epsilon} &= \underline{\epsilon}^\dagger \\ \underline{\zeta} &= \underline{\xi}^\dagger \end{aligned}$$

$$\begin{aligned} \underline{w}^\dagger &= \underline{v} \\ \underline{n}^\dagger &= \underline{m} \\ c &= -\underline{b}^* \\ d &= -\underline{e}^* \end{aligned}$$

Matches result of passivity: theory conserves energy!

Box 1. Source-driven theory¹

$$\begin{aligned} \eta &= \eta^o + \eta^e = \chi_{pu}^o + i\chi_{pu}^e \\ \underline{\zeta} &= \underline{\zeta}^o + \underline{\zeta}^e = \underline{\chi}_{EH}^o + i\underline{\chi}_{EH}^e \\ \underline{w} &= \underline{w}^o + \underline{w}^e = \underline{\chi}_{uE}^o + i\underline{\chi}_{uE}^e \\ \underline{n} &= \underline{n}^o + \underline{n}^e = i\underline{\chi}_{uH}^o + \underline{\chi}_{uH}^e \\ c &= c^o + c^e = i\chi_{pE}^o + \chi_{pE}^e \\ e &= e^o + e^e = \chi_{pH}^o + i\chi_{pH}^e \end{aligned}$$

Box 2. By reciprocity

$$\begin{aligned} \gamma &= \eta^o - \eta^e = \chi_{pu}^o - i\chi_{pu}^e \\ \underline{\xi} &= (\underline{\zeta}^o - \underline{\zeta}^e)^\top = (\underline{\chi}_{EH}^o - i\underline{\chi}_{EH}^e)^\top \\ \underline{v} &= (\underline{w}^o)^\top - (\underline{w}^e)^\top = (\underline{\chi}_{uE}^o - i\underline{\chi}_{uE}^e)^\top \\ \underline{m} &= -(\underline{n}^o)^\top + (\underline{n}^e)^\top = (-i\underline{\chi}_{uH}^o + \underline{\chi}_{uH}^e)^\top \\ b &= c^o - c^e = i\chi_{pE}^o - \chi_{pE}^e \\ d &= -e^o + e^e = -(\chi_{pH}^o - i\chi_{pH}^e) \end{aligned}$$

Compare Box 1, 2

$$\begin{aligned} \eta &= \gamma^* \\ \underline{\zeta} &= \underline{\xi}^\dagger \\ \underline{w}^\dagger &= \underline{v} \\ d &= -e^* \\ \underline{n}^\dagger &= \underline{m} \\ c &= -b^* \end{aligned}$$

¹Gokani et al., *J. Acoust. Soc. Am.* 153, A120 (2023).

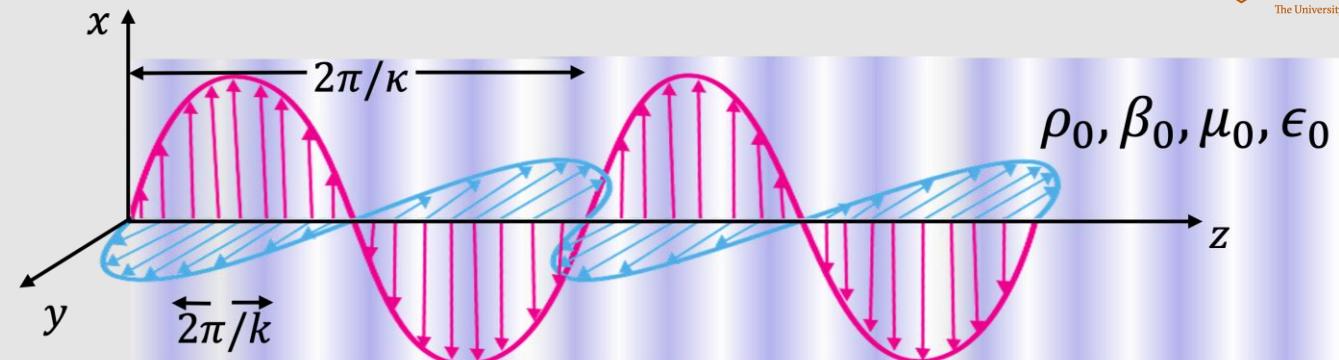
Even & odd couplings

$$ik\mathbf{p}_{\text{ext}} = i\omega\rho_0\mathbf{u}_{\text{ext}} + \mathbf{f}_{\text{ext}}$$

$$ik\mathbf{u}_{\text{ext}} = i\omega\beta_0\mathbf{p}_{\text{ext}} + \mathbf{q}_{\text{ext}}$$

$$i\kappa\mathbf{E}_{\text{ext}} = i\omega\mu_0\mathbf{H}_{\text{ext}} - \mathbf{K}_{\text{ext}}$$

$$i\kappa\mathbf{H}_{\text{ext}} = i\omega\epsilon_0\mathbf{E}_{\text{ext}} - \mathbf{J}_{\text{ext}}$$



$$\mathbf{p}_{\text{eff}} = \mathbf{p}_{\text{ext}} + \frac{1}{1 - k_0^2/k^2} \left(-\frac{k_0^2}{k^2} \beta_0^{-1} \bar{\mathbf{m}}_{\text{eff}} + \frac{k_0}{k} c_0 \bar{\mathbf{d}}_{\text{eff}} \right)$$

$$\mathbf{u}_{\text{eff}} = \mathbf{u}_{\text{ext}} + \frac{1}{1 - k_0^2/k^2} \left(-\frac{k_0}{k} c_0 \bar{\mathbf{m}}_{\text{eff}} + \frac{k_0^2}{k^2} \rho_0^{-1} \bar{\mathbf{d}}_{\text{eff}} \right)$$

$$\mathbf{E}_{\text{eff}} = \mathbf{E}_{\text{ext}} + \frac{1}{1 - \kappa_0^2/\kappa^2} \left(\frac{\kappa_0^2}{\kappa^2} \epsilon_0^{-1} \bar{\mathbf{P}}_{\text{eff}} + \frac{\kappa_0}{\kappa} C_0 \bar{\mathbf{M}}_{\text{eff}} \right)$$

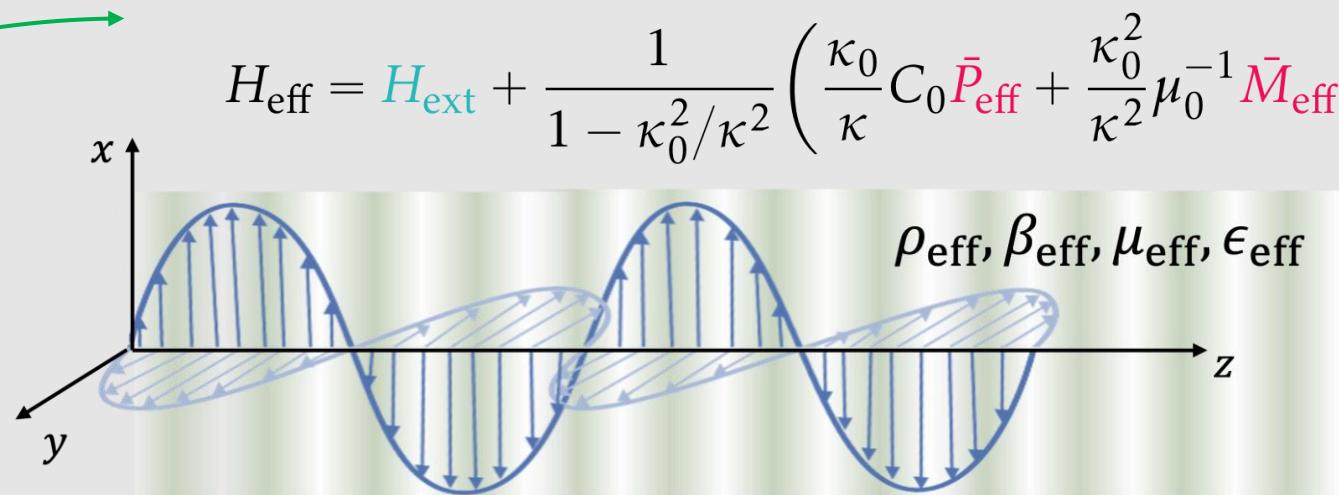
$$\mathbf{H}_{\text{eff}} = \mathbf{H}_{\text{ext}} + \frac{1}{1 - \kappa_0^2/\kappa^2} \left(\frac{\kappa_0}{\kappa} C_0 \bar{\mathbf{P}}_{\text{eff}} + \frac{\kappa_0^2}{\kappa^2} \mu_0^{-1} \bar{\mathbf{M}}_{\text{eff}} \right)$$

$$\varepsilon_{\text{eff}} = -\beta_0 \mathbf{p}_{\text{eff}} + \bar{\mathbf{m}}_{\text{eff}}$$

$$\mu_{\nu, \text{eff}} = \rho_0 \mathbf{u}_{\text{eff}} + \bar{\mathbf{d}}_{\text{eff}}$$

$$D_{\text{eff}} = \epsilon_0 \mathbf{E}_{\text{eff}} + \bar{\mathbf{P}}_{\text{eff}}$$

$$B_{\text{eff}} = \mu_0 \mathbf{H}_{\text{eff}} + \bar{\mathbf{M}}_{\text{eff}}$$



Interpretation of form of constitutive relations

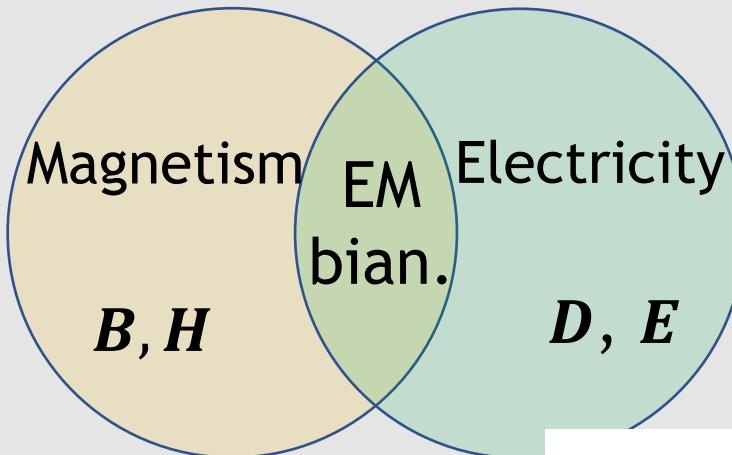
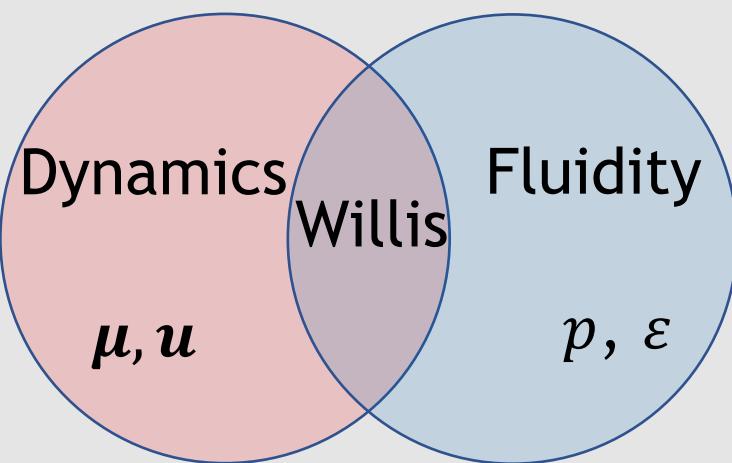
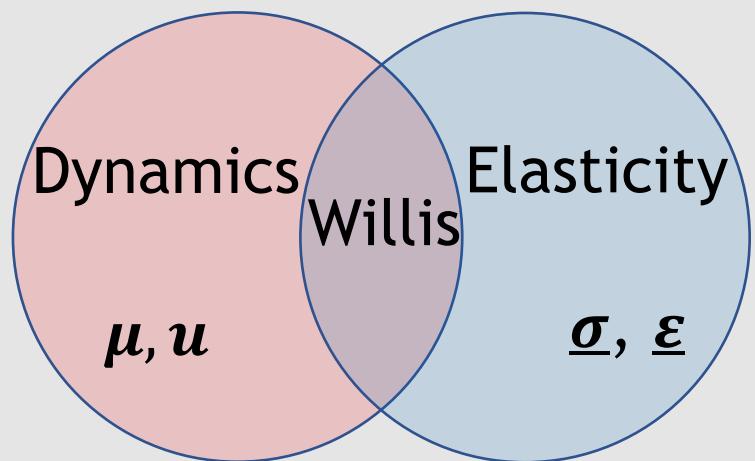
$$\begin{pmatrix} \varepsilon_{\text{eff}} \\ \mu_{v,\text{eff}} \\ D_{\text{eff}} \\ B_{\text{eff}} \end{pmatrix} = \begin{pmatrix} -\beta_{\text{eff}} & \chi_{pu}^{\text{o}} - i\chi_{pu}^{\text{e}} & i\chi_{pE}^{\text{o}} - \chi_{pE}^{\text{e}} & -(\chi_{pH}^{\text{o}} - i\chi_{pH}^{\text{e}}) \\ -(\chi_{pu}^{\text{o}} + i\chi_{pu}^{\text{e}}) & \rho_{\text{eff}} & \chi_{uE}^{\text{o}} - i\chi_{uE}^{\text{e}} & -i\chi_{uH}^{\text{o}} + \chi_{uH}^{\text{e}} \\ i\chi_{pE}^{\text{o}} + \chi_{pE}^{\text{e}} & \chi_{uE}^{\text{o}} + i\chi_{uE}^{\text{e}} & \epsilon_{\text{eff}} & \chi_{EH}^{\text{o}} \\ \chi_{pH}^{\text{o}} + i\chi_{pH}^{\text{e}} & i\chi_{uH}^{\text{o}} + \chi_{uH}^{\text{e}} & \chi_{EH}^{\text{o}} & \mu_{\text{eff}} \end{pmatrix} \begin{pmatrix} p_{\text{eff}} \\ u_{\text{eff}} \\ E_{\text{eff}} \\ H_{\text{eff}} \end{pmatrix}$$

Phase relations

- ε_{eff} in phase with E_{eff} ; D_{eff} in phase with p_{eff}
- $\mu_{v,\text{eff}}$ in phase with H_{eff} ; B_{eff} in phase with u_{eff}
- ε_{eff} and H_{eff} related by time derivative
- $\mu_{v,\text{eff}}$ and E_{eff} related by time derivative

Insights

Figure board



\hat{q} is a function of a complex variable: resort to numerical root finding

Fubini coefficients (–)

$$B_1 = 1 - \frac{1}{8}\sigma^2 + \mathcal{O}(\sigma^4)$$

$$B_2 = \frac{1}{2}\sigma + \mathcal{O}(\sigma^3)$$

$$B_3 = \frac{3}{8}\sigma^2 + \mathcal{O}(\sigma^4)$$

$$B_4 = \frac{1}{3}\sigma^3 + \mathcal{O}(\sigma^5)$$

$$B_5 = 0.3255\sigma^4 + \mathcal{O}(\sigma^6)$$

$$B_n = \frac{(n\sigma)^{n-1}}{2^{n-1}n!} + \mathcal{O}(\sigma^{n+1}).$$

