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Ex. 1.9 | Ch. 1 | P16 | Ch/oe | 2025.7.14

nf. (a) 1D-化波函数 $\int_{-\infty}^{+\infty} |\psi(x, t)|^2 dx = A^2 \int_{-\infty}^{+\infty} e^{-2amx^2/\hbar} dx$ (波函数归一化)

$$= 2A^2 \int_0^{+\infty} e^{-2amx^2/\hbar} dx \quad (\text{套用高斯积分公式}) = A^2 \sqrt{\frac{\pi\hbar}{2am}} = 1$$

$$A^2 = \frac{\sqrt{2am}}{\pi\hbar} \quad A = (\frac{2am}{\pi\hbar})^{\frac{1}{4}}$$

$$(b) \psi(x, t) = A e^{-amx^2/\hbar} e^{-i\omega t}$$

$$\frac{\partial \psi}{\partial t} = A e^{-amx^2/\hbar} \cdot (-i\omega) e^{-i\omega t} = -i\omega \psi(x, t)$$

$$\frac{\partial \psi}{\partial x} = A e^{-i\omega t} e^{-amx^2/\hbar} \cdot (-\frac{am}{\hbar}) \cdot 2x = -\frac{2amx}{\hbar} \psi(x, t)$$

$$\frac{\partial^2 \psi}{\partial x^2} = \frac{\partial}{\partial x} \left(\frac{\partial \psi}{\partial x} \right) = \left(-\frac{2am^2}{\hbar^2} \right) e^{-i\omega t} \cdot e^{-amx^2/\hbar} \left[1 - \frac{2amx^2}{\hbar^2} \right]$$

$$= \left(\frac{4a^2m^2x^2}{\hbar^2} - \frac{2am}{\hbar} \right) \psi(x, t)$$

$$\text{代入薛定谔方程 } i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2 \omega^2 \psi}{2m} + V\psi$$

$$i\hbar \cdot (-i\omega) \psi = -\hbar^2 \cdot \left(\frac{4a^2m^2x^2}{\hbar^2} - \frac{2am}{\hbar} \right) \psi + V\psi$$

$$\text{解得 } V = 2a^2mx^2$$

$$(c) \langle x \rangle = \int_{-\infty}^{+\infty} \psi^* x \psi dx = A^2 \int_{-\infty}^{+\infty} x e^{-2amx^2/\hbar} dx \quad (\text{奇}) = 0$$

$$\langle x^2 \rangle = \int_{-\infty}^{+\infty} \psi^* x^2 \psi dx = A^2 \int_{-\infty}^{+\infty} x^2 e^{-2amx^2/\hbar} dx \quad (\text{偶})$$

$$= 2A^2 \int_0^{+\infty} x^2 e^{-2amx^2/\hbar} dx \quad (\text{高斯积分}) = A^2 \hbar \sqrt{\frac{\pi\hbar}{4am}} = \frac{\hbar}{4am}$$

$$\langle p \rangle = \int_{-\infty}^{+\infty} \psi^* (-i\hbar \frac{\partial}{\partial x}) \psi dx = \int_{-\infty}^{+\infty} \psi^* (-i\hbar) \left(-\frac{2amx}{\hbar} \right) \psi dx$$

$$= \int_{-\infty}^{+\infty} -\frac{2am\hbar^2}{\hbar} x e^{-2amx^2/\hbar} dx = 0 \quad (\text{奇})$$

$$\langle p^2 \rangle = \int_{-\infty}^{+\infty} \psi^* (-i\hbar) \frac{\partial^2 \psi}{\partial x^2} dx = \int_{-\infty}^{+\infty} (2am - \frac{4a^2m^2x^2}{\hbar}) \psi^* \psi dx$$

$$= \int_{-\infty}^{+\infty} (2am - \frac{4a^2m^2x^2}{\hbar}) A^2 e^{-2amx^2/\hbar} dx \quad (\text{偶, 高斯积分})$$

$$= \sqrt{am\pi\hbar} \frac{\hbar A^2}{\sqrt{2}} = am$$

$$(d) \sigma_x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{am}$$

$$\sigma_p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \sqrt{am}$$

$$\sigma_x \cdot \sigma_p = \sqrt{am} \cdot \sqrt{am} = \frac{\hbar}{2}$$

满足海森伯不确定原理 $\sigma_x \cdot \sigma_p \geq \frac{\hbar}{2}$