Benchmarking Simulations against the ThermoML Database: Neat Liquid Densities and Static Dielectrics

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Useful atomistic simulations in the condensed phase require accurate depictions of solvent. While accurate experimental measurements of fundamental physical properties offer a straightforward approach for evaluating forcefield quality, the bulk of this information has been tied up in non-machine-readable formats that pose many risks to the automated evaluation of forcefield accuracy. Here we examine the feasibility of benchmarking atomistic forcefields against the NIST ThermoML database of physicochemical measurements, which aggregates thousands of experimental measurements in a portable machine-readable self-annotating format. We present a detailed benchmark of the generalized Amber small molecule forcefield (GAFF) using the AM1-BCC charge model against measurements (specifically liquid densities and static dielectric constants at ambient pressure) extracted from ThermoML, and discuss the extent of available data for neat liquids. We show that empirical polarizability models correct systematic biases inherent in predicting dielectric constants with fixed-charged forcefields. [IDC: Will need to be convinced!]

Keywords: molecular mechanics forcefields; forcefield parameterization; forcefield accuracy; forcefield validation; mass density; static dielectric constant

I. INTRODUCTION

Recent advances in hardware and molecular dynamics software have provided routine access to atomistic simulations at the 100 ns timescale and beyond. Leveraging these advances in combination with GPU clusters, distributed computing, or custom hardware has brought the microsecond and milliseconds within reach. These dramatic advances in sampling, however, have revealed forcefields as a critical barrier for truly predictive simulation.

Protein and water forcefields have been the subject of numerous benchmarks [] and enhancements [], with key outcomes including the ability to fold fast-folding proteins, improved fidelity of water thermodynamic properties, and improved prediction of NMR observables. Although small molecule forcefields have also been the subject of benchmarks [] and improvements [], such work has focused on small perturbations to specific functional groups. For example, a recent study found that modified hydroxyl nonbonded parameters led to improved prediction of static dielectrics and hydration free energies. There are also outstanding questions of generalizability of parameters. Will changes to a specific chemical moiety be compatible with seemingly un-

³⁴ related improvements? Addressing these questions re-³⁵ quires agreement on shared benchmarks that can be eas-³⁶ ily replicated with proposed forcefield enhancements.

A key barrier in forcefield development is that many 38 experimental datasets are heterogeneous, paywalled, 39 and unavailable in machine-readable formats (although 40 notable counterexamples exist, e.g. RCSB [1], Free-41 Solv [2] and BMRB [3]). While this inconvenience 42 is relatively minor for benchmarking a single target 43 (e.g. water), it becomes prohibitive for studies span-44 ning chemical space. To ameliorate problems of data 45 archival, the NIST Thermodynamics Research Center 46 has developed a IUPAC standard XML-based format— 47 ThermoML [4]—for storing physicochemical measure-48 ments, uncertainties, and metadata. Experimental re-49 searchers publishing measurements in several journals 50 (J. Chem. Eng. Data, J. Chem. Therm., Fluid Phase 51 Equil., Therm. Acta, and Int. J. Therm.) are now guided 52 through a data archival process that involves sanity 53 checks and eventual archival at the TRC (http://trc. 54 nist.gov/ThermoML.html).

Here we examine the ThermoML archive as a potential source for neat liquid density and static dielectric measurements, with the goal of developing a standard benchmark for validating these properties in fixed-charge forcefields of drug-like molecules. These two observables provide sensitive tests of forcefield accuracy that are nonetheless straightforward to calculate. Using the ThermoML data, we evaluate the generalized Amber small molecule forcefield (GAFF) [5] with the AM1-BCC charge model [6, 7] and identify systematic biases that might be improved upon.

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Filter	Mass Density	Static Dielectric
0. Single Component	130074	1649
1. Druglike Elements	120410	1649
2. Heavy Atoms	67897	1567
3. Temperature	36827	962
4. Pressure	13598	461
Liquid state	13573	461
6. Aggregate T, P	3573	432
7. Density+Dielectric	245	245

TABLE I. Number of ThermoML measurements matching sequentially applied filters.

II. RESULTS

A. Neat Liquid Measurements in ThermoML

We performed a number of queries to summarize the ThermoML content relevant for benchmarking organic molecule forcefields. Our aim is to explore neat liquid data with functional groups relevant to drug-like molecules. We therefore applied the following sequence of filters: has either density or static dielectric measurements, contains a single component, contains only druglike elements (H, N, C, O, S, P, F, Cl, Br), has low heavy atom count (< 10), has ambient temperature [K] (270 < $T \leq 330$), has ambient pressure [kPA] ($100 \leq P \leq 102$), ₇₈ and has density greater than 300 kg m^{-3} (a proxy for liquid state). After applying these filters, we also round all pressures within this range to exactly one atmosphere. We also round temperatures to one decimal place. These approximations are motived by common data entry errors; for example, an experiment performed at water's freezing point at ambient pressure might be entered as either 101.325 kPA or 100 kPA, with a temperature of either 273 K or 273.15 K. After the application of these filters (Table I), we are left with 245 conditions for which 88 both density and dielectric data are available. The func-89 tional groups present are summarized in Table I.

90 B. Benchmarking GAFF against ThermoML: Mass Density

Mass density has been widely used as a critical ingredient for parameterizing and testing forcefields, particularly the Lennard Jones parameters [9, 10]. We therefore used the present ThermoML compilation as a benchmark of the GAFF AM1-BCC forcefield (Fig. 1). Overall, the densities show reasonable accuracy (R2 + errobars), consistent with previous studies [11] reporting agreement of XYZ on a different benchmark set.

C. Benchmarking GAFF against ThermoML: Static Dielectric

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As a measure of the electronic medium, the static dilog electric constant of neat liquids provides a critical bench-

Functional Group	Counts
1,2-aminoalcohol	4
1,2-diol	3
alkene	3
aromatic compound	1
carbonic acid diester	2
carboxylic acid ester	4 7
dialkyl ether	7
heterocyclic compound	3
ketone	3 2
lactone	1
primary alcohol	19
primary aliphatic amine (alkylamine)	2
primary amine	2
secondary alcohol	4
secondary aliphatic amine (dialkylamine)	2
secondary aliphatic/aromatic amine (alkylarylamine)	1
secondary amine	3
sulfone	1
sulfoxide	1
tertiary aliphatic amine (trialkylamine)	3
tertiary amine	3

TABLE II. Functional group counts present in the dataset. The number of unique compounds is 44. Functional group classification was performed using checkmol version 0.5 [8].

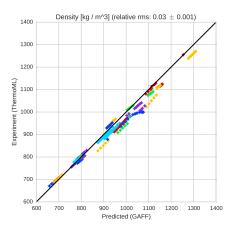


FIG. 1. Comparison of liquid densities between experiment and simulation. Liquid density measurements at ambient (1 atm) pressure extracted from ThermoML are compared against simulated densities at ambient pressure using the GAFF/AM1-BCC small molecule fixed-charge forcefield. Color groupings represent identical chemical species. Simulation error bars represent one standard error of the mean, with the number of effective (uncorrelated) samples estimated using pymbar. Experimental error bars indicate the standard deviation between independently reported measurements, when available, or author-reported standard deviations in ThermoML entries; for some measurements, neither uncertainty estimate is available. See Section B for further discussion of error.

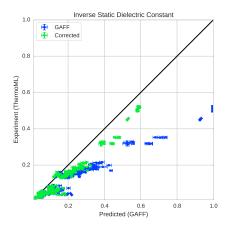
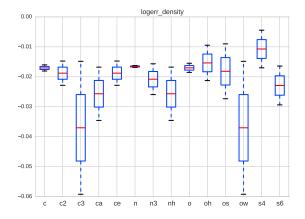
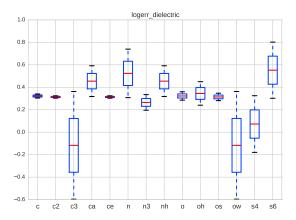


FIG. 2. Measured (ThermoML) versus predicted (GAFF) inverse static dielectrics (a). Color groupings represent identical chemical formulas. Simulation error bars represent one standard error of the mean estimated via block averaging with block sizes of 200 ps [15]. Experimental error bars indicate the larger of standard deviation between independently reported measurements and the authors reported standard deviations; for some measurements, neither uncertainty estimate is available. See section S2 for further discussion of error. The inverse dielectric $\frac{1}{\epsilon}$ is plotted instead of ϵ because $\frac{1}{\epsilon}$ is directly proportional to energy in continuum dielectric models: e.g. $U(r) = \frac{1}{4\pi} \frac{q_{12}}{r} \propto \frac{1}{\epsilon}.$

103 mark that is somewhat orthogonal to density and ther-104 modynamic quantities. We therefore compare simulations against the measurements in our ThermoML compilation. Overall, we find the dielectric constants to be qualitatively reasonable, but with clear deviations from experiment. In particular, GAFF AM1-BCC systematically underestimates the dielectric constants for nonpolar organics, with GAFF predictions of $\epsilon \approx 1.0 \pm 0.05$ being substantially smaller than the measured $\epsilon \approx 2$. Because this deviation likely stems from the lack of electronic polarization, we added a simple empirical correction for polarization [12], which leads to better agreement with experiment. A similar polarization correction was used in the development of the TIP4P-EW water model [13]; however, the need is much greater for the nonpolar organics, as the missing polarizability is the dominant contribution to the static dielectric constant. In the case of water, the Sales polarizability model predicts a dielectric correction of 0.52, while 0.79 was used for the TIP4P-EW model. For comparison, we also applied the same empirical correction to the VirtualChemistry dataset [11, 14] and saw similarly improved agreement with experiment for both the GAFF and OPLS 126 forcefields (Fig. 8).





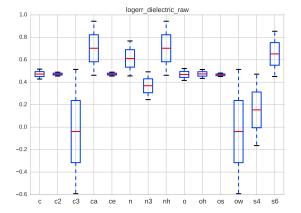


FIG. 3. TODO: MAKE A DECENT ANALYSIS THAT BREAKS THINGS DOWN ERRORS BY FUNCTIONAL GROUP!

III. DISCUSSION

A. Examining discrepancies by functional group

SEE FIG. 3

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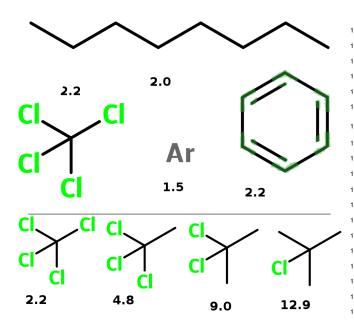


FIG. 4. (a). Measured static dielectric constants of various nonpolar or symmetric molecules [?]; fixed-charge forcefields give $\epsilon \approx 1$ for each species. (b). A congeneric series of chloromethanes show dielectric constants between 2 and 10. TODO: REDO THIS FIGURE, fixing Benzene greenness and 179 add explicit protons on CClH3

B. Fitting Forcefields to Dielectric Constants

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Recent forcefield development has seen a resurgence of papers fitting dielectric constants as primary data [16, 17]. However, a number of authors have pointed out potential challenges in constructing self-consistent fixed-charge forcefields [18, 19].

Interestingly, a recent work by Dill [18] pointed out 137 that, for CCl₄, reasonable choices of point charges are incapable of recapitulating the observed dielectric of = 2.2, instead producing dielectric constants in the range of $1.0 \le \epsilon \le 1.05$. This behavior is quite general: a fixed charge monopole force field predicts $\epsilon \approx 1$ for several nonpolar or symmetric molecules, but the measured dielectric constants are instead $\epsilon \approx 2$ (Fig. 4). While this behavior is well-known and results from missing physics of polarizability, we suspect it may have several unanticipated consequences.

Suppose, for example, that one attempts to di-198 148 rectly fit the static dielectric constants of CCl₄, CHCl₃, 149 CH₂Cl₂, and CH₃Cl. In moving from the tetrahedrally- 199 an appropriate choice of point charges. However, the 202 https://github.com/choderalab/ThermoPyL) 156 of polarizability. We hypothesize that this inconsistency 206 (http://media.iupac.org/namespaces/ 158 where symmetric molecules (e.g. benzene, CCl₄) have 208 data. This schema was converted into a Python object

159 qualitatively different properties than closely related 160 asymmetric molecules (e.g. toluene, CHCl₃). As a possible real-world example, we imagine that the missing po-162 larizability could be important in accurate transfer free energies invovling low-dielectric solvents. Using the 164 Onsager model for the transfer free energy of a dipole $_{165}$ $(\Delta G=-rac{\mu^2}{a^3}rac{\epsilon-1}{2\epsilon+1})$ gives an error of $\Delta\Delta G=\Delta G(\epsilon=0)$ $_{166} \,\, 2.2) - \Delta G(\epsilon = 1)$ of 2 kcal / mol for the transfer of water ($a=1.93~\text{Å}~\mu=2.2\text{D}$) into a low dielectric medium 168 such as tetrachloromethane or benzene.

As a temporary fix, we suggest using empirical polar-170 ization corrections before directly comparing measured 171 static dielectric constants to fixed-charge models particularly when examining low-dielectric solvents. Separating the contributions of fixed charges and polar-174 ization may also lead to the development of improved models of electrostatics that account for the missing po-176 larization physics; some such models have been pro-177 posed recently [19–21].

C. ThermoML as a Data Source

The present work has focused on the neat liquid den-180 sity and dielectric measurements present in ThermoML [4, 22, 23] as a target for molecular dynamics forcefield validation. While densities and dielectric constants have been widely used in forcefield work, several aspects of ThermoML make it a unique resource for the forcefield community. First, the aggregation, support, and dissemination of ThermoML is supported by NIST, whose mission makes these tasks a long-term priority. Second, ThermoML is actively growing, through partnerships 189 with journals such as J. Chem. Thermo-new experimen-190 tal measurements published in these journals are criti-191 cally examined by the TRC and included in the archive. 192 Finally, the files in ThermoML are machine readable via 193 a formal XML schema, allowing facile access to thou-194 sands of measurements. In the future, we hope to ex-195 amine additional measurement classes, including both 196 mixture and two-phase data.

IV. METHODS

ThermoML Processing

ThermoML XML files were obtained from the 150 symmetric CCl₄ to CHCl₃, it suddenly becomes possi- 200 the NIST TRC. To explore their content, we cre-151 ble to achieve the observed dielectric constant of 4.8 by 201 ated a python (version 2.7.9) tool (ThermoPyl: model for CHCl₃ uses fixed point charges to account for 203 that munges the XML content into a spreadsheetboth the net dipole moment and the (electronic) polariz- 204 like format accessible via the Pandas (version 155 ability, whereas the CCl₄ model contains no treatment 205 0.15.2) library. First, we obtained the XML schema 157 in parameterization may lead to strange mismatches, 207 ThermoML.xsd) defining the layout of the 209 via PyXB 1.2.4 (http://pyxb.sourceforge.net/). 256 ford University), Robert McGibbon (Stanford Univer-210 Finally, this schema and Pandas was used to extract the 257 sity), Jason Swails (Rutgers University), David L. Mob-211 data and apply the data filters described above.

B. Simulation

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Boxes of 1000 molecules were constructed using 213 AM1-BCC [6, 7] charges were gen-214 PackMol [24]. 215 erated using OpenEye Toolkit 2014-6-6 [25], using 216 the oequacpac.OEAssignPartialCharges module with the OECharges_AM1BCCSym. The selected con-218 former was then processed using antechamber in Am-219 berTools 14 [26]. The resulting AMBER files were con- 266 verted to OpenMM [27] ffxml forcefield XML files. Sim- 267 dards and Technology is not subject to copyright in the 221 ulation code used libraries gaff2xml 0.6, TrustButVer- 268 United States. Products or companies named here are 222 ify 0.1, OpenMM 6.2 [27], and MDTraj 1.2 [28]. [TODO: 269 cited only in the interest of complete technical descrip-

Molecular dynamics simulations were performed us- 272 be found to serve as well. 225 226 ing OpenMM 6.2 [27] using a Langevin integrator (with collision rate 1 ps^{-1}) and a 1 fs timestep; interestingly, we found that a 2 fs timestep led to insufficient accuracy in equilibrium densities (Table III). [JDC: Cite Langevin 230 integrator used in OpenMM.] Pressure coupling at 1 atmosphere was achieved with a Monte Carlo barostat uti-232 lizing molecular scaling and automated step size adjust-233 ment during equilibration, applied every 25 steps. Particle mesh Ewald [29] was used with a long-range cutoff of 0.95 nm and an long-range isotropic dispersion cor-236 rection. [JDC: Can we report the automatically-selected PME parameters?] Simulations were continued until density standard errors were less than 2×10^{-4} g / mL, 239 as estimated using the equilibration detection module in pymbar 2.1 [30]. Trajectory analysis was performed us-241 ing OpenMM [27] and MDTraj [28]. Density data was 242 output every 250 fs, while trajectory data was stored ev-243 ery 10 ps.

V. CONCLUSIONS

- ThermoML is a potentially useful resource for the forcefield community
- We have curated a subset of ThermoML for neat liquids with druglike atoms, with thousands of densities and hundreds of dielectrics
- Empirical polarization models correct a systematic bias in comparing fixed-charge forcefields to static dielectric constants

VI. ACKNOWLEDGEMENTS

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258 ley (University of California, Irvine), Christopher 259 I. Bayly (OpenEye Software), Michael R. Shirts (Uni-260 versity of Virginia), and members of Chodera lab for 261 helpful discussions. Support for JMB was provided 262 by the Tri-Institutional Training Program in Computa-263 tional Biology and Medicine (via NIH training grant 264 1T32GM083937).

VII. DISCLAIMERS

This contribution of the National Institute of Stan-Provide a script to install all of these versions via 270 tion, and neither constitute nor imply endorsement by 271 NIST or by the U.S. government. Other products may mu n neff sigma stderr error 0.5 0.903701 145510 20357.973571 0.007362 0.000052 0.000000 0.00 1.0 0.903114 159515 21988.457281 0.007415 0.000050 -0.000588 -0.00 2.0 0.901811 108346 15964.072327 0.007494 0.000059 -0.001891 -0.00

TABLE III. To probe the systematic error from finite time-step integration, we examined the timestep dependence of butyl acrylate density. The number of effective samples was estimated using pymbar's statistical inefficiency routine [30]. To approximate the timestep bias, we compare the density expectation ($\langle \rho \rangle$) to values calculated with a 0.5fs timestep. We find a 2fs timestep leads to systematic biases in the density on the order of 0.2%, while 1fs reduces the systematic bias to less than 0.1%—we therefore selected a 1fs timestep for the present work, where we aimed to achieve three digits of accuracy in density predictions.

Appendix A: Supplementary Information

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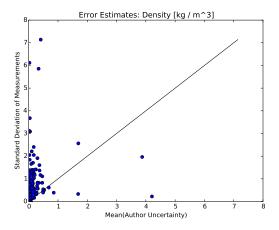
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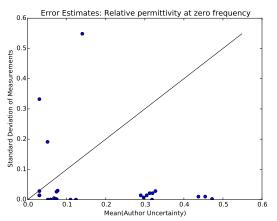
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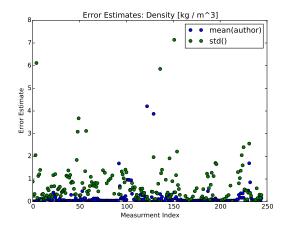
All information below this point will eventually be pulled into a separate SI. This will happen closer to submission, as the formatting may be journal-specific. The references may be split in two as well, depending on journal.

- Table: Timestep-dependence of density
- Figure: Error analysis for ThermoML dataset
- Table (CSV File): ThermoML Dataset used in present analysis.

Appendix B: Assessment of experimental error in ThermoML measurements







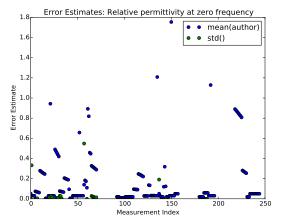


FIG. 5. Assessment of experimental error in ThermoML data. To assess the experimental error in our benchmark set, we consider two orthogonal error measurements. The first is the mean of the uncertainties reported by the measurement au-

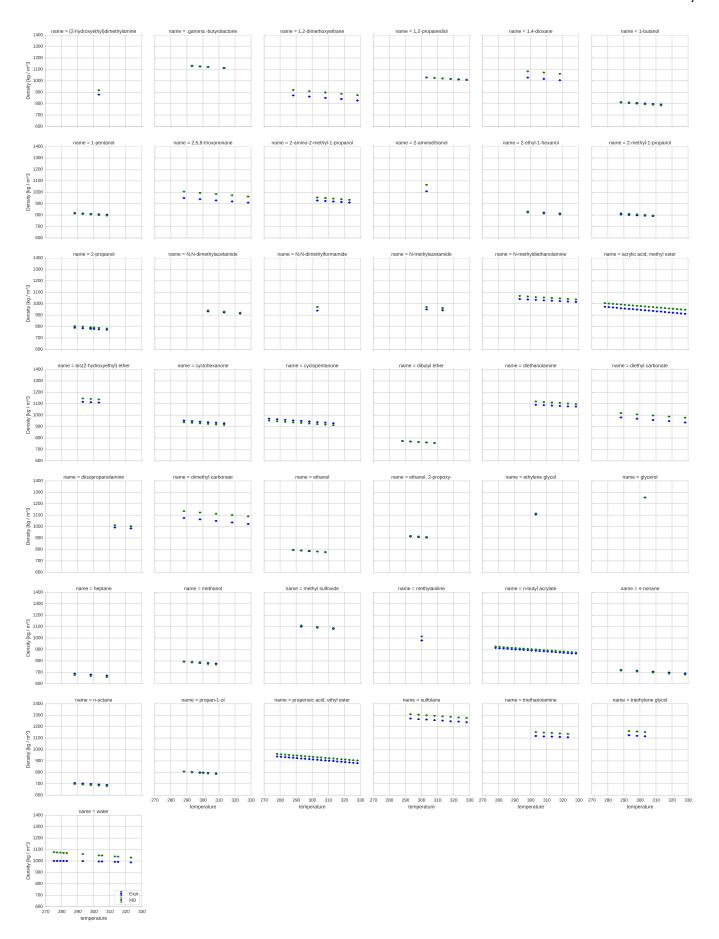


FIG. 6. Comparison of simulated and experimental densities for all compounds. Measured (blue) and simulated (green) densities are shown in units of kg/m^3 .

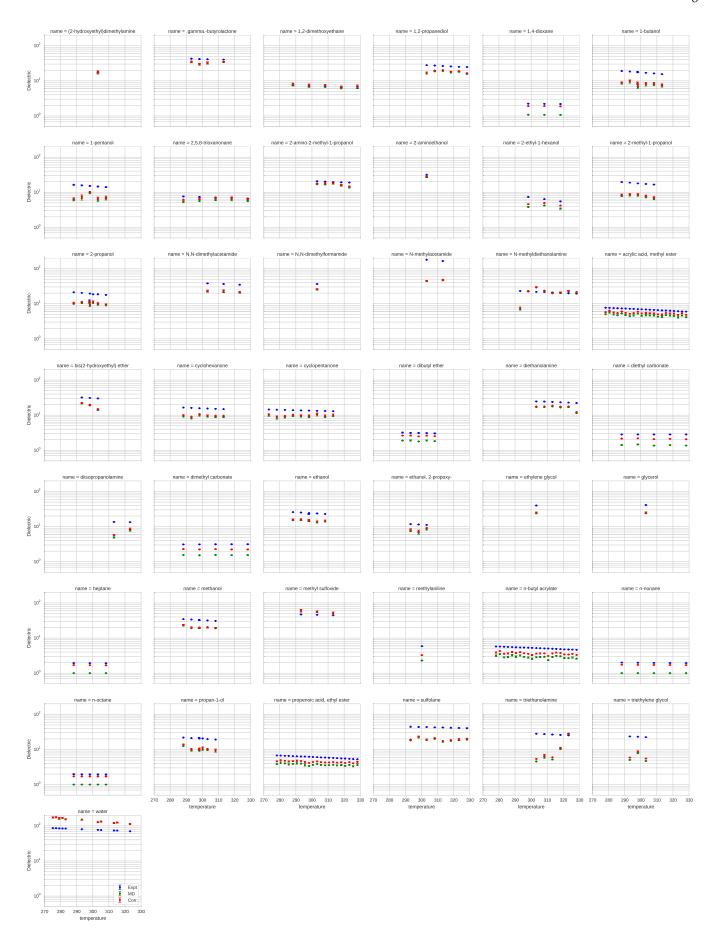
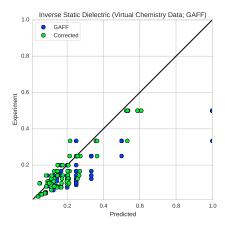


FIG. 7. Comparison of simulated and experimental static dielectric constants for all compounds. Measured (blue), simulated (green), and polarizability-corrected simulated (red) static dielectric constants are shown for all compounds. Note that dielectric constants, rather than inverse dielectric constants, are plotted here.



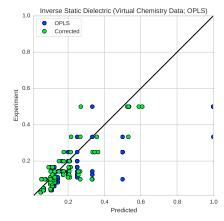


FIG. 8. Comparison of measured and simulated dielectric constants from the virtualchemistry dataset with and without polarizability correction. Measured (blue), MD (green), and MD + polarizability-corrected (red) dielectrics for the Virtual Chemistry dataset [11, 14].

- [1] H. M. Berman, J. Westbrook, Z. Feng, G. Gilliland, T. N. 328 [19] I. V. Leontyev and A. A. Stuchebrukhov, The Journal of Bhat, H. Weissig, I. N. Shindyalov, and P. E. Bourne, Nu- 329 cleic Acids Res. 28, 235 (2000).
- D. L. Mobley, Experimental and calculated small molecule 331 hydration free energies, Retrieved from: http://www. 332 escholarship.org/uc/item/6sd403pz, uC Irvine: 333 [21] Department of Pharmaceutical Sciences, UCI.
- [3] E. Ulrich, H. Akutsu, J. Doreleijers, Y. Harano, Y. Ioanni- 335 Nucleic Acids Res. 36, D402 (2008).
- [4] M. Frenkel, R. D. Chiroco, V. Diky, Q. Dong, K. N. 338 Marsh, J. H. Dymond, W. A. Wakeham, S. E. Stein, 339 E. Königsberger, and A. R. Goodwin, Pure and applied 340 [23] R. D. Chirico, M. Frenkel, V. V. Diky, K. N. Marsh, and chemistry 78, 541 (2006).
- [5] J. Wang, R. M. Wolf, J. W. Caldwell, P. A. Kollman, and 342 D. A. Case, J. Comput. Chem. 25, 1157 (2004). 300
 - [6] A. Jakalian, B. L. Bush, D. B. Jack, and C. I. Bayly, J. Com- 344 put. Chem. 21, 132 (2000).
- [7] A. Jakalian, D. B. Jack, and C. I. Bayly, J. Comput. Chem. 346 303 23, 1623 (2002). 304
 - [8] N. Haider, Molecules 15, 5079 (2010).

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- [9] W. L. Jorgensen, J. Chandrasekhar, J. D. Madura, R. W. 349 306 Impey, and M. L. Klein, The Journal of chemical physics 350 307 **79**, 926 (1983). 308
- [10] W. L. Jorgensen, J. D. Madura, and C. J. Swenson, Journal 309 of the American Chemical Society 106, 6638 (1984). 310
- [11] C. Caleman, P. J. van Maaren, M. Hong, J. S. Hub, L. T. 354 311 Costa, and D. van der Spoel, Journal of chemical theory 355 312 and computation 8, 61 (2011). 313
- [12] R. Bosque and J. Sales, Journal of chemical information 357 [29] 314 and computer sciences 42, 1154 (2002). 315
- 13] H. Horn, W. Swope, J. Pitera, J. Madura, T. Dick, G. Hura, 359 316 and T. Head-Gordon, J. Chem. Phys. 120, 9665 (2004). 317
- 14] D. van der Spoel, P. J. van Maaren, and C. Caleman, Bioin- 361 [31] Mettler toledo density meters, [Online; accessed 15-Jan-318 formatics 28, 752 (2012). 319
- [15] H. Flyvbjerg and H. G. Petersen, J. Chem. Phys. 91, 461 363 320 (1989).321
- [16] L.-P. Wang, T. J. Martínez, and V. S. Pande, The Journal of 365 322 Physical Chemistry Letters (2014). 323
- C. J. Fennell, K. L. Wymer, and D. L. Mobley, The Journal 367 324 of Physical Chemistry B (2014). 325
- C. J. Fennell, L. Li, and K. A. Dill, The Journal of Physical 369 326 Chemistry B 116, 6936 (2012). 327

- chemical physics 141, 014103 (2014).
- 330 [20] J.-F. Truchon, A. Nicholl's, J. A. Grant, R. I. Iftimie, B. Roux, and C. I. Bayly, Journal of computational chemistry 31, 811 (2010).
- J.-F. Truchon, A. Nicholls, B. Roux, R. I. Iftimie, and C. I. Bayly, Journal of chemical theory and computation 5, 1785 (2009).
- dis, J. Lin, M. Livny, S. Mading, D. Maziuk, and Z. Miller, 336 [22] M. Frenkel, R. D. Chirico, V. V. Diky, Q. Dong, S. Frenkel, P. R. Franchois, D. L. Embry, T. L. Teague, K. N. Marsh, and R. C. Wilhoit, Journal of Chemical & Engineering Data 48, 2 (2003).
 - R. C. Wilhoit, Journal of Chemical & Engineering Data 48, 1344 (2003).
 - 343 [24] L. Martínez, R. Andrade, E. G. Birgin, and J. M. Martínez, Journal of computational chemistry 30, 2157 (2009).
 - 345 [25] Openeye toolkits 2014, URL http://www.eyesopen. com.
 - 347 [26] D. Case, V. Babin, J. Berryman, R. Betz, O. Cai, D. Cerutti, T. Cheatham III, T. Darden, R. Duke, H. Gohlke, et al., 348 University of California, San Francisco (2014).
 - P. Eastman, M. S. Friedrichs, J. D. Chodera, R. J. Radmer, C. M. Bruns, J. P. Ku, K. A. Beauchamp, T. J. Lane, L.-P. Wang, D. Shukla, et al., J. Chem. Theory Comput. 9, 461 (2012).353
 - R. T. McGibbon, K. A. Beauchamp, C. R. Schwantes, L.-P. Wang, C. X. Hernández, M. P. Harrigan, T. J. Lane, J. M. Swails, and V. S. Pande, bioRxiv p. 008896 (2014). 356
 - T. Darden, D. York, and L. Pedersen, J. Chem. Phys. 98, 358 10089 (1993).
 - M. R. Shirts and J. D. Chodera, J. Chem. Phys. 129, 124105 360 (2008).
 - 362 2015], URL. http://us.mt.com/us/en/home/ products/Laboratory_Analytics_Browse/ Density_Family_Browse_main/DE_Benchtop. tabs.models-and-specs.html.
 - 366 [32] R. D. Chirico, M. Frenkel, J. W. Magee, V. Diky, C. D. Muzny, A. F. Kazakov, K. Kroenlein, I. Abdulagatov, G. R. Hardin, and W. E. Acree Jr, Journal of Chemical & Engineering Data 58, 2699 (2013).