Benchmarking Atomistic Simulations against the ThermoML Data Archive: **Neat Liquid Densities and Static Dielectric Constants**

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Useful atomistic simulations in the condensed phase require accurate depictions of solvent. While experimental measurements of fundamental physical properties offer a straightforward approach for evaluating forcefield quality, the bulk of this information has been tied up in formats that are not machine-readable. These formats require substantial human effort to compile benchmark datasets which are prone to accumulation of human errors, hindering the development of reproducible benchmarks of forcefield accuracy. Here, we examine the feasibility of benchmarking atomistic forcefields against the NIST ThermoML data archive of physicochemical measurements, which aggregates thousands of experimental measurements in a portable, machine-readable, self-annotating format. As a proof of concept, we present a detailed benchmark of the generalized Amber small molecule forcefield (GAFF) using the AM1-BCC charge model against measurements (specifically liquid densities and static dielectric constants at ambient pressure) automatically extracted from the archive, and discuss the extent of available data for neat liquids. The results of this benchmark highlights a general problem with fixed-charge forcefields in the representation of liquids of low dielectric.

Keywords: molecular mechanics forcefields; forcefield parameterization; forcefield accuracy; forcefield validation; mass density; static dielectric constant

I. INTRODUCTION

Recent advances in hardware and software for molec-13 ular dynamics simulation now permits routine access to atomistic simulations at the 100 ns timescale and beyond [1]. Leveraging these advances in combination with consumer GPU clusters, distributed computing, or custom hardware has brought microsecond and millisecond simulation timescales within reach of many laboratories. These ¹⁹ dramatic advances in sampling, however, have revealed deficiencies in forcefields as a critical barrier to enabling truly predictive simulations of physical properties of biomolecular systems.

Protein and water forcefields have been the subject of numerous benchmarks [2] and enhancements [3-5], with key outcomes including the ability to fold fast-folding proteins [6-8], improved fidelity of water thermodynamic properties [9], and improved prediction of NMR observables. Although small molecule forcefields have also been the subject of benchmarks [10] and improvements [11], such 30 work has typically focused on small perturbations to spe-31 cific functional groups. For example, a recent study found

32 that modified hydroxyl nonbonded parameters led to im-33 proved prediction of static dielectric constants and hydra-34 tion free energies [11]. There are also outstanding ques-35 tions of generalizability of these targeted perturbations; it is uncertain whether changes to the parameters for a specific chemical moiety will be compatible with seemingly un-38 related improvements to other groups. Addressing these 39 questions requires establishing a community agreement on 40 shared benchmarks that can be easily replicated among 41 laboratories to test proposed forcefield enhancements and 42 expanded as the body of experimental data grows.

A key barrier to establishing reproducible and extensible 44 forcefield accuracy benchmarks is that many experimen-45 tal datasets are heterogeneous, paywalled, and unavail-46 able in machine-readable formats (although notable coun-47 terexamples exist, e.g. the RCSB [12], FreeSolv [13], and 48 the BMRB [14]). While this inconvenience is relatively mi-49 nor for benchmarking forcefield accuracy for a single tar-50 get (e.g. water), it becomes prohibitive for studies span-51 ning the relevant chemical space. To ameliorate prob-52 lems of data archival, the NIST Thermodynamics Research 53 Center (TRC) has developed a IUPAC standard XML-based 54 format—ThermoML [15]—for storing physicochemical mea-55 surements, uncertainties, and metadata. Experimental 56 researchers publishing measurements in several journals 57 (J. Chem. Eng. Data, J. Chem. Therm., Fluid Phase Equil., 58 Therm. Acta, and Int. J. Therm.) are guided through a data 59 archival process that involves sanity checks, conversion to a 50 standard machine-readable format, and archival at the TRC 61 (http://trc.nist.gov/ThermoML.html).

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Here, we examine the ThermoML archive as a potential source for providing the foundation for a reproducible, extensible accuracy benchmark of biomolecular forcefields. In particular, we concentrate on two important physical property measurements easily computable in many simulation codes—neat liquid density and static dielectric constant measurements—with the goal of developing a standard benchmark for validating these properties in fixedcharge forcefields of drug-like molecules and biopolymer residue analogues. These two properties provide sensitive tests of forcefield accuracy that are nonetheless straightforward to calculate. Using these data, we evaluate the generalized Amber small molecule forcefield (GAFF) [16] with the AM1-BCC charge model [17, 18] and identify systematic biases to aid further forcefield refinement.

II. RESULTS

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A. Extracting neat liquid measurements from the NIST TRC ThermoML Archive

We retrieved a copy of the ThermoML Archive from the NIST TRC (http://trc.nist.gov/ThermoML.html accessed 13 Sep 2014) and performed a number of sequential filtering steps to produce an extract of the ThermoML Archive relevant for benchmarking forcefields describing small organic molecules. [JDC: This is the date I had on the ThermoML.tar.gz archive in GitHub. We should check to make sure this is accurate.] As our aim is to explore neat liquid data with functional groups relevant to biopolymers and drug-like molecules, we applied the following ordered of filters, starting with all data containing density or static diselectric constants:

- 1. The measured solution contains only a single component (e.g. no binary mixtures)
- 2. The molecule contains only the druglike elements (defined here as H, N, C, O, S, P, F, Cl, Br)
 - 3. The molecule has \leq 10 heavy atoms
 - 4. The measurement was performed in a biophysically relevant temperature range [K] $(270 \le T \le 330)$
 - 5. The measurement was performed at ambient pressure [kPA] $(100 \le P \le 102)$
 - 6. Measured densities below 300 kg m⁻³ were discarded to eliminate gas-phase measurements
 - 7. The temperature and pressure were rounded to nearby values (as described below), averaging all measurements within each group of like conditions 137
 - 8. Only conditions (molecule, temperature, pressure) for which *both* density and dielectric constants were available were retained

	Number of measurements remaining			
Filter step	Mass density	Static dielectric		
1. single component	130074	1649		
2. only druglike elements	120410	1649		
3. ≤10 heavy atoms	67897	1567		
4. $(270 \le T \le 330)$ [K]	36827	962		
5. ambient pressure	13598	461		
6. liquid state	13573	461		
7. aggregate T, P	3573	432		
8. density and dielectric	245	245		

TABLE I. Successive filtration of the ThermoML Archive. A set of successive filters were applied to all measurements in the ThermoML Archive (accessed 13 Sep 2014) that contained either mass density or static dielectric constant measurements. Each column reports the number of measurements remaining after successive application of the corresponding filtration step.

The temperature and pressure rounding step was motivated by common data reporting variations; for example, an experiment performed at water's freezing point at ambient pressure might be entered as either 101.325 kPA or 100 kPA, with a temperature of either 273 K or 273.15 K. Therefore all pressures within the range [kPA] $(100 \le P \le 102)$ were rounded to exactly one atmosphere. Temperatures were rounded to one decimal place. The application of these filters (Table I) leaves 245 conditions—where a condition here indicates a (molecule, temperature, pressure) tuple—for which both density and dielectric data are available. The functional groups present in the resulting dataset are summarized in Table II.

[JDC: It might be useful to point the users to the scripts that were used to do this extraction. Also, can we automate the downloading of the complete up-to-date archive, perhaps with Kenneth's help in identifying the least intrusive way to do so?]

B. Benchmarking GAFF/AM1-BCC against the ThermoML Archive

C. Methods Summary

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We used OpenMM 6.2 to perform highly precise constantpressure simulations of mass density. The Langevin timestep was fixed at 1 fs to reduce systematic error. All simulations used the GAFF / AM1-BCC forcefield. Addition details are provided in Methods.

1. Mass density

Mass density has been widely used for parameterizing and testing forcefields, particularly the Lennard-Jones parameters representing dispersive and repulsive interactions [20, 21]. We therefore used the present ThermoML extract as a benchmark of the GAFF/AM1-BCC forcefield (Fig. 1). We estimate the mass density via $\rho = \langle \frac{M}{V} \rangle$.

Functional Group	Occurrences
1,2-aminoalcohol	4
1,2-diol	3
alkene	3
aromatic compound	1
carbonic acid diester	2
carboxylic acid ester	4
dialkyl ether	7
heterocyclic compound	3
ketone	2
lactone	1
primary alcohol	19
primary aliphatic amine (alkylamine)	2
primary amine	2
secondary alcohol	4
secondary aliphatic amine (dialkylamine)	2
secondary aliphatic/aromatic amine (alkylarylamine)	1
secondary amine	3
sulfone	1
sulfoxide	1
tertiary aliphatic amine (trialkylamine)	3
tertiary amine	3

TABLE II. Functional groups present in filtered dataset. The filtered ThermoML dataset contained 245 distinct (molecule, temperature, pressure) conditions, spanning 44 unique compounds. The functional groups represented in these compounds (as identified by the program checkmol v0.5 [19]) is summarized here.

Overall, the densities show reasonable accuracy, with a root-mean square (RMS) relative error over all measurements of 2.8±0.1% (with one standard error of the mean determined by bootstrapping over all measurements), especially encouraging given that this forcefield was not designed with the intention of modeling bulk liquid properties of organic molecules [16] This is reasonably consistent with ment of electronic polarization, we used a simple empirical benchmark set [10].

[JDC: Discuss outliers here. There must be more things we can say about densities. Some of the densities are quite good, while others seem poor, with systematic bias toward higher densities than experiment. We can also point out that densities at different temperatures for a given molecule seem to be biased in a consistent way.]

Static dielectric constant

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As a measure of the dielectric response, the static dielec-

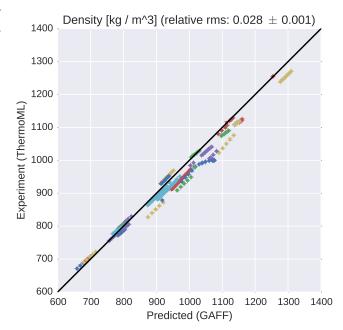


FIG. 1. Comparison of liquid densities between experiment and simulation. Liquid density measurements extracted from ThermoML are compared against densities predicted using the GAFF/AM1-BCC small molecule fixed-charge forcefield. Color groupings represent identical chemical species. Simulation error bars represent one standard error of the mean, with the number of effective (uncorrelated) samples estimated using pymbar. Experimental error bars indicate the standard deviation between independently reported measurements, when available, or authorreported standard deviations in ThermoML entries; for some measurements, neither uncertainty estimate is available. See Appendix B for further discussion of error.

previous studies reporting agreement of 4% on a different polarization model that computes the molecular electronic polarizability α a sum of elemental atomic polarizability contributions [22]. [JDC: I've commented out the equation because I don't think it is central to our point. Essentially, it is just an atom-based linear model for computing the molecular polarizability, and the parameters can be looked up elsewhere.] From the computed molecular electronic polarizability lpha, an additive correction to the simulation-derived 178 static dielectric constant accounting for the missing elec-179 tronic polarizability can be computed [9]

$$\Delta\epsilon = 4\pi N \frac{\alpha}{\langle V \rangle} \tag{1}$$

tric constant of neat liquids provides a critical benchmark 180 While a similar polarization correction was used in the deof the accuracy electrostatic treatment in forcefield mod- 181 velopment of the TIP4P-Ew water model, where it had a miels. We therefore compare simulations against the measure- 182 nor effect [9], missing polarizability is a dominant contriments in our ThermoML extract. Overall, we find the dielec- 183 bution to the static dielectric constant of nonpolar organic tric constants to be qualitatively reasonable, but with clear 184 molecules; in the case of water, the empirical atomic polardeviations from experiment. In particular, GAFF/AM1-BCC 185 izability model predicts a dielectric correction of 0.52, while systematically underestimates the dielectric constants for 186 0.79 was used for the TIP4P-Ew model. [JDC: What were the nonpolar organics, with the predictions of $\epsilon \approx 1.0 \pm 0.05$ being substantially smaller than the measured $\epsilon \approx 2$. Because 188 we also applied the same empirical correction to the Virthis deviation likely stems from the lack of an explicit treat- 189 tualChemistry dataset [10, 23] and saw similarly improved

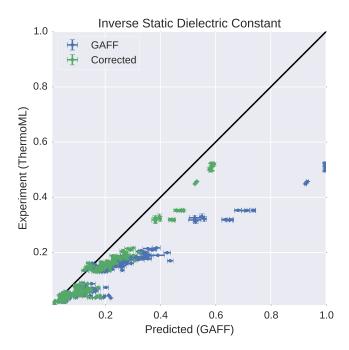


FIG. 2. Measured (ThermoML) versus predicted (GAFF/AM1-BCC) inverse static dielectrics (a). Simulation error bars represent one standard error of the mean estimated via block averaging with block sizes of 200 ps [24]. [JDC: Why are we using block averaging here? Why didn't we just use timeseries.py. We should not be using block averaging, especially without a justification that 200 ps is a reasonable block size for every specific system and condition Let's talk about this.] Experimental error bars indicate the larger of standard deviation between independently reported measurements and the authors reported standard deviations; for some measurements, neither uncertainty estimate is available. See Section B for further discussion of error. The inverse dielectric constant ϵ^{-1} is plotted instead of ϵ because ϵ^{-1} is directly proportional to the Coulomb interaction energy between point charges embedded in a dielectric material [e.g. $U(r) \propto$ $q_1q_2/r \propto \epsilon^{-1}$]. [JDC: We need to trim the whitespace of all sides of the figures that you are outputting in order for the figure to actually fill the column width. There must be some option to set that. See the figure.tight_layout() option in matplotlib, along with matplotlib.backends.backend_pdf.PdfPages.]

agreement with experiment for both the GAFF and OPLS forcefields (Fig. 7). [JDC: Not sure if we should keep the VirtualChemistry stuff here other than to compare our computed pipeline results with Virtual Chemistry as validation to show we didn't screw up our pipeline.]

DISCUSSION

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Fitting Forcefields to Dielectric Constants

Recent forcefield development has seen a resurgence of papers fitting dielectric constants during forcefield parameterization [11, 25]. However, a number of authors have pointed out potential challenges in constructing self- 226 lated asymmetric molecules (e.g. toluene and $\mathrm{CHCl_3}$). 201 consistent fixed-charge forcefields [26, 27].

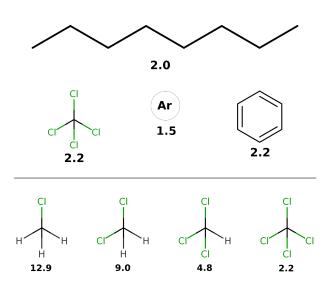


FIG. 3. Typical experimental static dielectric constants of some **nonpolar compounds.** (a). Measured static dielectric constants of various nonpolar or symmetric molecules [?]; fixed-charge forcefields give $\epsilon \approx 1$ for each species. (b). A congeneric series of chlorosubstituted methanes have static dielectric constants between 2 and 13. [Can we use a better citable source for these numbers instead of Wikipedia? Also, what temperatures/pressures are these measurements cited at? Maybe we can just say "near ambient"?] [JDC: We should not use PNG files for figure graphics—only vector graphics (when possible). Can you use a vector graphics PDF instead?] [JDC: Can we show both experimental and GAFF/AM1-BCC computed dielectric constants for some of these compounds?]

Interestingly, recent work by Dill and coworkers [26] ob-203 served that, for CCl₄, reasonable choices of point charges 204 are incapable of recapitulating the observed dielectric of $\epsilon = 2.2$, instead producing dielectric constants in the range of $1.0 < \epsilon < 1.05$. This behavior is quite general: fixed point charge forcefields will predict $\epsilon \approx 1$ for many non-208 polar or symmetric molecules, but the measured dielectric 209 constants are instead $\epsilon \approx 2$ (Fig. 3). While this behavior 210 is well-known and results from missing physics of polariz-211 ability, we suspect it may have several unanticipated consequences, which we discuss below.

Suppose, for example, that one attempts to fit forcefield parameters to match the static dielectric constants of CCl₄, CHCl₃, CH₂Cl₂, and CH₃Cl. In moving from the tetrahedrally-symmetric CCl₄ to the asymmetric CHCl₃, 217 it suddenly becomes possible to achieve the observed dielectric constant of 4.8 by an appropriate choice of point charges. However, the model for CHCl₃ uses fixed point 220 charges to account for both the permanent dipole moment and the electronic polarizability, whereas the CCl₄ model contains no treatment of polarizability. We hypothesize that this inconsistency in parameterization may lead to strange mismatches, where symmetric molecules (e.g. benzene and CCl₄) have qualitatively different properties than closely re-

How important is this effect? As a possible real-world

228 example, we imagine that the missing atomic polarizabil- 280 in the ThermoML Archive are portable and machine read-234 medium such as tetrachloromethane or benzene.

$$\Delta G = -\frac{\mu^2}{a^3} \frac{\epsilon - 1}{2\epsilon + 1} \tag{2}$$

Similarly, we calculated the mean polarization error for solvation free energies (gas to solvent transfer free energies) of druglike molecules in cyclohexane. For each molecule in the latest (Feb. 20) FreeSolv database [13, 28], we took the cavity radius a to be the half the maximum interatomic distance and calculated $\mu = \sum_i q_i r_i$ using the provided mol2 oordinates and AM1-BCC charges. This calculation predicts 292 (where the standard error is computed from bootstrapping 294 To explore the content of this archive, over measurements), suggesting that the missing atomic 295 ated a Python (version 2.7.9) polarizabilty unrepresentable by fixed point charge force- 296 https://github.com/choderalab/ThermoPyL) also conjecture that

izability, the inconsistency of fixed-charge models in low- 305 Section II A. dielectric media is sufficiently alarming to motivate further study of polarizable forcefields. In particular, continuum methods [29-31], point dipole methods [32, 33], and 306 Drude methods [34, 35] have been maturing rapidly. Finding the optimal balance of accuracy and performance remains an open question; however, the use of experimentallyparameterized direct polarization methods [36] may pro- 309 version?]. fixed charge forcefields.

ThermoML as a data source

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electric constants have already been widely used in force- 320 script to install all of these versions via conda.] field work, several aspects of ThermoML make it a unique 321 ₂₇₉ number of journal here also expanding?] Finally, the files ₃₃₀ every 25 steps. The particle mesh Ewald (PME) method [44]

ity could be important in accurate transfer free energies 281 able via a formal XML schema, allowing facile access to huninvolving low-dielectric solvents. The Onsager model for 282 dreds of thousands of measurements. Numerous additional the transfer free energy of a dipole (Eq. 2) gives an error of 283 physical properties contained in ThermoML—including ac- $\Delta\Delta G=\Delta G(\epsilon=2.2)-\Delta G(\epsilon=1)$ of -2 kcal / mol for the $_{284}$ tivity coefficients, diffusion constants, boiling point tempertransfer of water (a=1.93 Å, $\mu=2.2$ D) into a low dielectric 285 atures, critical pressures and densities, coefficients of ex-286 pansion, speed of sound measurements, viscosities, excess 287 molar enthalpies, heat capacities, and volumes—for neat (2) 288 phases and mixtures represent a rich dataset of high utility 289 for forcefield validation and parameterization.

METHODS

ThermoML Processing

A tarball archive of the ThermoML Archive was mean error of -0.91 ± 0.07 kcal/mol for the 643 molecules 293 obtained from the the NIST TRC on 13 Sep 2014. tool (ThermoPyL: fields could contribute substantially to errors in predicted 297 formats the XML content into a spreadsheet-like format transfer and solvation properties of druglike molecules. We 298 accessible via the Pandas (version 0.15.2) library. First, we 299 obtained the XML schema (http://media.iupac.org/ Given their ease of measurement and direct connection to mamespaces/ThermoML/ThermoML.xsd) defining the laylong-range electrostatic interactions, static dielectric con- 301 out of the data. This schema was converted into a Python stants have high potential utility as primary data for force- 302 object via PyXB 1.2.4 (http://pyxb.sourceforge.net/). field parameterization efforts. Although this will require the 303 Finally, this schema and Pandas was used to extract the use of forcefields with explicit treatment of atomic polar- 304 data and apply the successive data filters described in

Simulation

Using an automated tool, boxes of 1000 molecules 308 were constructed using PackMol [39] [JDC: Which AM1-BCC [17, 18] charges were genervide polarizability physics at a cost not much greater than 310 ated using OpenEye Toolkit 2014-6-6 [40], using the 311 oequacpac.OEAssignPartialCharges module with the OECharges_AM1BCCSym option, which utilizes a conformational expansion procedure prior to charge fitting to 314 minimize artifacts from intramolecular contacts. 315 selected conformer was then processed using antechamber The present work has focused on the neat liquid den- 316 in AmberTools 14 [41]. The resulting AMBER files were sity and dielectric measurements present in the ThermoML 317 converted to OpenMM [42] ffxml forcefield XML files. Simu-Archive [15, 37, 38] as a target for molecular dynamics force- 318 lation code used libraries gaff2xml 0.6, TrustButVerify 0.1, field validation. While liquid mass densities and static di- 319 OpenMM 6.2 [42], and MDTraj 1.2 [43]. [TODO: Provide a

Molecular dynamics simulations were performed with resource for the forcefield community. First, the aggre- 322 OpenMM 6.2 [42] using a Langevin integrator (with collision gation, support, and dissemination of ThermoML datasets $_{323}$ rate 1 ps $^{-1}$) and a 1 fs timestep, as we found that timesteps of through the ThermoML Archive is supported by NIST, whose 324 2 fs timestep or greater led to a significant timestep depenmission makes these tasks a long-term priority. Second, 325 dence in computed equilibrium densities (Table III). [JDC: the ThermoML Archive is actively growing, through part- 326 Cite Langevin integrator used in OpenMM.] Pressure connerships with several journals, and new experimental mea- 327 trol to 1 atm was achieved with a Monte Carlo barostat utisurements published in these journals are critically exam- 328 lizing molecular scaling and automated step size adjustined by the TRC and included in the archive. [JDC: Is the 329 ment during equilibration, with volume moves attempted was used with a long-range cutoff of 0.95 nm and a long-range isotropic dispersion correction. [JDC: Can we report the automatically-selected PME parameters to aid reproducibility in other codes?] Simulations were continued until density standard errors were less than 2×10^{-4} g / mL, as estimated using the equilibration detection module in pymbar 2.1 [45]. Trajectory analysis was performed using OpenMM [42] and MDTraj [43]. [JDC: Which versions?] Instantaneous densities were stored every 250 fs, while trajectory snapshots were stored every 10 ps. [JDC: Did we plan to make this data available somewhere, or is it sufficient to put out the scripts?]

V. CONCLUSIONS

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- ThermoML is a potentially useful resource for the forcefield community
- We have curated a subset of the ThermoML Data Archive for neat liquids with druglike atoms, with thousands of densities and hundreds of dielectrics
- Empirical polarization models correct a systematic bias in comparing fixed-charge forcefields to static dielectric constants

VI. ACKNOWLEDGEMENTS

We thank Vijay S. Pande (Stanford University), Lee-Ping Wang (Stanford University), Peter Eastman (Stanford University), Robert McGibbon (Stanford University), Jason Swails (Rutgers University), David L. Mobley (University of California, Irvine), Christopher I. Bayly (OpenEye Software), Michael R. Shirts (University of Virginia), and members of Chodera lab for helpful discussions. Support for JMB was provided by the Tri-Institutional Training Program in Computational Biology and Medicine (via NIH training grant 1T32GM083937). [JDC: Need support acknowledgments for Patrick and Bas.]

VII. DISCLAIMERS

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Appendix A: Supplementary Information

All information below this point will eventually be pulled into a separate SI. This will happen closer to submission, as the formatting may be journal-specific. The references may be split in two as well, depending on journal. [JDC: It may be fine to leave this as an Appendix.]

• Table: Timestep-dependence of density

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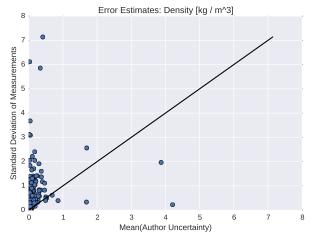
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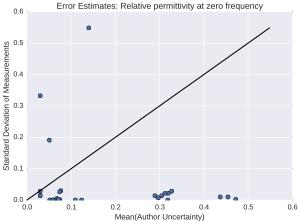
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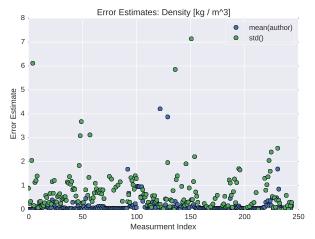
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- Figure: Error analysis for ThermoML dataset
- Table (CSV File): ThermoML Dataset used in present analysis.

382 Appendix B: Assessment of experimental error in ThermoML measurements







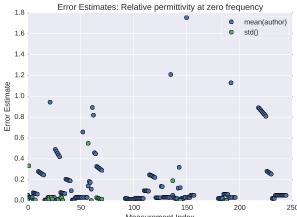


FIG. 4. Assessment of experimental error in ThermoML data. To assess the experimental error in our ThermoML extract, we considered two different approaches. In the first approach, we computed the mean of the uncertainties reported by the measurement

Δt	$\langle \rho \rangle ({\rm g/cm^3})$	n	neff	$stddev(\rho)$	stderr	abs error (g/cm 3)	rel error (%)
0.5	0.903701	145510	20358.0	0.007362	0.000052	0.000000	0.0000
1.0	0.903114	159515	21988.5	0.007415	0.000050	-0.000588	-0.0650
2.0	0.901811	108346	15964.1	0.007494	0.000059	-0.001891	-0.2092

TABLE III. Timestep dependence in computed equilibrium density of butyl acrylate. To probe the systematic error from finite timestep integration, we examined the timestep dependence of butyl acrylate density. The number of effective samples was estimated using pymbar's statistical inefficiency routine [45]. To approximate the timestep bias, we compare the density expectation $(\langle \rho \rangle)$ to values calculated with a 0.5 fs timestep. We find a 2 fs timestep leads to systematic biases in the density on the order of 0.2%, while 1fs reduces the systematic bias to less than 0.1%—we therefore selected a 1 fs timestep for the present work, where we aimed to achieve three digits of accuracy in density predictions. [JDC: I've reformatted this table a bit, paying more attention to sig figs. I think this might actually be better presented as a figure showing the timestep dependence, perhaps for 4 or 5 timesteps from 0.5 to 2.5 fs, rather than just 3.]



FIG. 5. Comparison of simulated and experimental densities for all compounds. Measured (blue) and simulated (green) densities are shown in units of kg/m³.

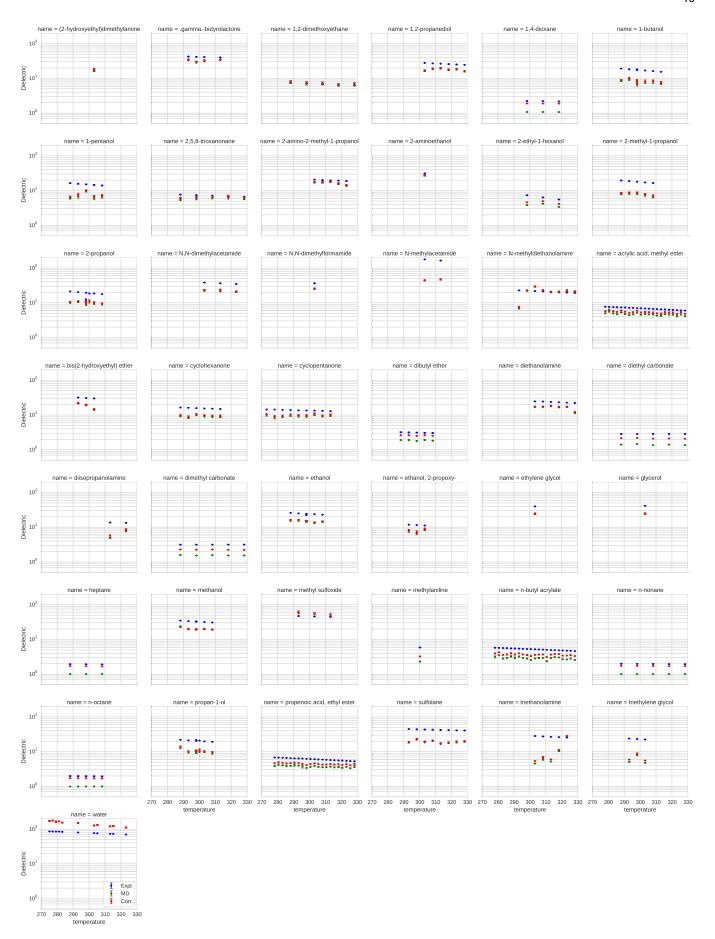
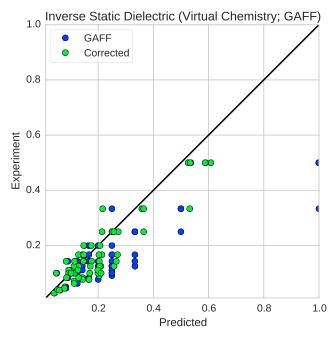


FIG. 6. Comparison of simulated and experimental static dielectric constants for all compounds. Measured (blue), simulated (green), and polarizability-corrected simulated (red) static dielectric constants are shown for all compounds. Note that dielectric constants, rather than inverse dielectric constants, are plotted here. [JDC: Let's plot these as in Fig. 1 and Fig. 2, maybe only four plots across so they are larger and more legible. We can also shorten "name = compound" to just "compound".]



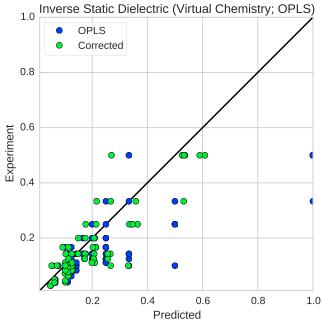


FIG. 7. Atomic polarizability corrections for the Virtual Chemistry dataset. A comparison of measured (*blue*) and simulated dielectric constants from the Virtual Chemistry dataset [10, 23]. with (*red*) and without (*green*) atomic polarizability corrections are shown. [JDC: What from the Virtual Chemistry dataset did we end up using? The computed values? The parameterized files? Since the story is about the ThermoML Archive, do we really want this here, or do we just want to do a comparison between our computed results for some molecules and the Virtual Chemistry results as a validation or sanity check on our computational pipeline?] [JDC: Plot needs to be adjusted to be more legible if we decide to keep it.]

Walker, Journal of Chemical Theory and Computation 9, 3878 (2013).

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- K. Lindorff-Larsen, P. Maragakis, S. Piana, M. Eastwood, R. Dror, and D. Shaw, PloS one 7, e32131 (2012).
- [3] D.-W. Li and R. Bruschweiler, J. Chem. Theory Comput. 7, 1773 389 390
 - [4] R. B. Best, X. Zhu, J. Shim, P. E. Lopes, J. Mittal, M. Feig, and A. D. MacKerell, J. Chem. Theory Comput. (2012).
- K. Lindorff-Larsen, S. Piana, K. Palmo, P. Maragakis, J. Klepeis, 393 R. Dror, and D. Shaw, Proteins: Struct., Funct., Bioinf. 78, 1950 395
 - [6] K. Lindorff-Larsen, S. Piana, R. Dror, and D. Shaw, Science 334, 517 (2011).
- [7] D. Ensign, P. Kasson, and V. Pande, J. Mol. Biol. 374, 806 398
- [8] V. Voelz, G. Bowman, K. Beauchamp, and V. Pande, J. Am. 459 400 Chem. Soc. 132, 1526 (2010). 401
- [9] H. Horn, W. Swope, J. Pitera, J. Madura, T. Dick, G. Hura, and 402 T. Head-Gordon, J. Chem. Phys. 120, 9665 (2004). 403
- [10] C. Caleman, P. J. van Maaren, M. Hong, J. S. Hub, L. T. Costa, 463 404 and D. van der Spoel, Journal of chemical theory and compu- 464 405 tation 8, 61 (2011). 406
- [11] C. J. Fennell, K. L. Wymer, and D. L. Mobley, The Journal of 466 407 Physical Chemistry B (2014). 408
- [12] H. M. Berman, J. Westbrook, Z. Feng, G. Gilliland, T. N. Bhat, 468 409 H. Weissig, I. N. Shindyalov, and P. E. Bourne, Nucleic Acids 469 410 Res. 28, 235 (2000). 411
- [13] D. L. Mobley, Experimental and calculated small 471 412 molecule hydration free energies, Retrieved from: 472 413 http://www.escholarship.org/uc/item/6sd403pz, uC Irvine: Department of Pharmaceutical Sciences, UCI. 415
- [14] E. Ulrich, H. Akutsu, J. Doreleijers, Y. Harano, Y. Ioannidis, 475 J. Lin, M. Livny, S. Mading, D. Maziuk, and Z. Miller, Nucleic 476 417 Acids Res. 36, D402 (2008). 418
- [15] M. Frenkel, R. D. Chiroco, V. Diky, Q. Dong, K. N. Marsh, J. H. 478 419 Dymond, W. A. Wakeham, S. E. Stein, E. Königsberger, and A. R. 479 420 Goodwin, Pure and applied chemistry 78, 541 (2006). 421
- J. Wang, R. M. Wolf, J. W. Caldwell, P. A. Kollman, and D. A. 422 Case, J. Comput. Chem. 25, 1157 (2004). 423
- A. Jakalian, B. L. Bush, D. B. Jack, and C. I. Bayly, J. Comput. 424 Chem. 21, 132 (2000). 425
- A. Jakalian, D. B. Jack, and C. I. Bayly, J. Comput. Chem. 23, [18] 426 1623 (2002). 427
- [19] N. Haider, Molecules 15, 5079 (2010). 428
- W. L. Jorgensen, J. Chandrasekhar, J. D. Madura, R. W. Im-[20] 429 pey, and M. L. Klein, The Journal of chemical physics 79, 926 430 431
- [21] W. L. Jorgensen, J. D. Madura, and C. J. Swenson, Journal of 432 the American Chemical Society 106, 6638 (1984). 433
- R. Bosque and J. Sales, Journal of chemical information and 434 computer sciences 42, 1154 (2002). 435
- D. van der Spoel, P. J. van Maaren, and C. Caleman, Bioinfor-436 matics 28, 752 (2012). 437
- [24] H. Flyvbjerg and H. G. Petersen, J. Chem. Phys. **91**, 461 (1989). 438
- L.-P. Wang, T. J. Martínez, and V. S. Pande, The Journal of Phys-439 ical Chemistry Letters (2014). 440
- C. J. Fennell, L. Li, and K. A. Dill, The Journal of Physical Chem-441 istry B 116, 6936 (2012).

- [1] R. Salomon-Ferrer, A. W. GolĹtz, D. Poole, S. Le Grand, and R. C. 443 [27] I. V. Leontyev and A. A. Stuchebrukhov, The Journal of chemical physics 141, 014103 (2014).
 - [28] D. L. Mobley, Experimental and calculated small molecule hy-445 dration free energies, Retrieved from: https://github.com/ choderalab/FreeSolv, uC Irvine: Department of Pharmaceutical Sciences, UCI.
 - 449 [29] J.-F. Truchon, A. Nicholl's, J. A. Grant, R. I. Iftimie, B. Roux, and C. I. Bayly, Journal of computational chemistry 31, 811 (2010).
 - [30] J.-F. Truchon, A. Nicholls, B. Roux, R. I. Iftimie, and C. I. Bayly, 451 Journal of chemical theory and computation 5, 1785 (2009). 452
 - 453 [31] J.-F. Truchon, A. Nicholls, R. I. Iftimie, B. Roux, and C. I. Bayly, Journal of chemical theory and computation 4, 1480 (2008).
 - J. Ponder, C. Wu, P. Ren, V. Pande, J. Chodera, M. Schnieders, [32] 455 I. Haque, D. Mobley, D. Lambrecht, R. DiStasio Jr, et al., J. 456 Phys. Chem. B 114, 2549 (2010). 457
 - 458 [33] P. Ren and J. W. Ponder, The Journal of Physical Chemistry B **108**, 13427 (2004).
 - 460 [34] G. Lamoureux and B. Roux, The Journal of Chemical Physics **119**, 3025 (2003). 461
 - 462 [35] V. M. Anisimov, G. Lamoureux, I. V. Vorobyov, N. Huang, B. Roux, and A. D. MacKerell, Journal of Chemical Theory and Computation 1, 153 (2005).
 - L.-P. Wang, T. L. Head-Gordon, J. W. Ponder, P. Ren, J. D. 465 Chodera, P. K. Eastman, T. J. Martínez, and V. S. Pande, J. Phys. Chem. B 117, 9956 (2013).
 - M. Frenkel, R. D. Chirico, V. V. Diky, Q. Dong, S. Frenkel, P. R. Franchois, D. L. Embry, T. L. Teague, K. N. Marsh, and R. C. Wilhoit, Journal of Chemical & Engineering Data 48, 2 (2003).
 - R. D. Chirico, M. Frenkel, V. V. Diky, K. N. Marsh, and R. C. Wilhoit, Journal of Chemical & Engineering Data 48, 1344 (2003).
 - L. Martínez, R. Andrade, E. G. Birgin, and J. M. Martínez, Journal of computational chemistry **30**, 2157 (2009).
 - [40] Openeye toolkits 2014, URL http://www.eyesopen.com.
 - [41] D. Case, V. Babin, J. Berryman, R. Betz, Q. Cai, D. Cerutti, T. Cheatham III, T. Darden, R. Duke, H. Gohlke, et al., Univer-477 sity of California, San Francisco (2014).
 - [42] P. Eastman, M. S. Friedrichs, J. D. Chodera, R. J. Radmer, C. M. Bruns, J. P. Ku, K. A. Beauchamp, T. J. Lane, L.-P. Wang, 480 D. Shukla, et al., J. Chem. Theory Comput. 9, 461 (2012).
 - [43] R. T. McGibbon, K. A. Beauchamp, C. R. Schwantes, L.-P. Wang, 482 C. X. Hernández, M. P. Harrigan, T. J. Lane, J. M. Swails, and V. S. Pande, bioRxiv p. 008896 (2014).
 - T. Darden, D. York, and L. Pedersen, J. Chem. Phys. 98, 10089 485 (1993).486
 - [45] M. R. Shirts and J. D. Chodera, J. Chem. Phys. 129, 124105 487 (2008).488
 - [46] Mettler toledo density meters, [Online; accessed 15-Jan-489 2015], URL http://us.mt.com/us/en/home/products/ 490 491 Laboratory_Analytics_Browse/Density_Family_ Browse_main/DE_Benchtop.tabs.models-and-specs. 492 493
 - [47] R. D. Chirico, M. Frenkel, J. W. Magee, V. Diky, C. D. Muzny, A. F. Kazakov, K. Kroenlein, I. Abdulagatov, G. R. Hardin, and W. E. Acree Jr, Journal of Chemical & Engineering Data 58, 2699 (2013).

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