Benchmarking Atomistic Simulations against the ThermoML Data Archive: **Neat Liquid Densities and Static Dielectric Constants**

Kyle A. Beauchamp⁺,^{1,*} Julie M. Behr⁺,^{2,†} Patrick B. Grinaway,^{3,‡} Arien S. Rustenburg,^{3,§} Kenneth Kroenlein,^{4,¶} and John D. Chodera^{1,**}

¹Computational Biology Program, Memorial Sloan Kettering Cancer Center, New York, NY ²Tri-Institutional Program in Computational Biology and Medicine, Weill Cornell Medical College, New York, NY 3 Graduate Program in Physiology, Biophysics, and Systems Biology, Weill Cornell Medical College, New York, NY ⁴Themodynamics Research Center, NIST, Boulder, CO (Dated: February 22, 2015)

Useful atomistic simulations in the condensed phase require accurate depictions of solvent. While experimental measurements of fundamental physical properties offer a straightforward approach for evaluating forcefield quality, the bulk of this information has been tied up in formats that are not machine-readable. These formats require substantial human effort to compile benchmark datasets which are prone to accumulation of human errors, hindering the development of reproducible benchmarks of forcefield accuracy. Here, we examine the feasibility of benchmarking atomistic forcefields against the NIST ThermoML data archive of physicochemical measurements, which aggregates thousands of experimental measurements in a portable, machine-readable, self-annotating format. As a proof of concept, we present a detailed benchmark of the generalized Amber small molecule forcefield (GAFF) using the AM1-BCC charge model against measurements (specifically liquid densities and static dielectric constants at ambient pressure) automatically extracted from the archive, and discuss the extent of available data for neat liquids. The results of this benchmark highlights a general problem with fixed-charge forcefields in the representation of liquids of low dielectric.

Keywords: molecular mechanics forcefields; forcefield parameterization; forcefield accuracy; forcefield validation; mass density; static dielectric constant

I. INTRODUCTION

Recent advances in hardware and software for molecular dynamics simulation now permits routine access to atomistic simulations at the 100 ns timescale and beyond [CITE e.g. http://pubs.acs.org/doi/abs/ 10.1021/ct400314y]. Leveraging these advances in combination with consumer GPU clusters, distributed 40 crosecond and milliseconds within reach. These dramatic advances in sampling, however, have revealed forcefields as a critical barrier for truly predictive simulation.

of numerous benchmarks [1] and enhancements [2–4], with key outcomes including the ability to fold fastfolding proteins, improved fidelity of water thermodynamic properties, and improved prediction of NMR observables. Although small molecule forcefields have also been the subject of benchmarks [5] and improve-29 ments [6], such work has typically focused on small perturbations to specific functional groups. For example, a recent study found that modified hydroxyl nonbonded parameters led to improved prediction of static dielec-

33 tric constants and hydration free energies [6]. There are 34 also outstanding questions of generalizability of param-35 eters. Will changes to a specific chemical moiety be compatible with seemingly unrelated improvements? Ad-37 dressing these questions requires agreement on shared 38 benchmarks that can be easily replicated with proposed 39 forcefield enhancements.

A key barrier in forcefield development is that many 17 computing, or custom hardware has brought the mi- 41 experimental datasets are heterogeneous, paywalled, 42 and unavailable in machine-readable formats (although 43 notable counterexamples exist, e.g. RCSB [7], Free-44 Solv [8] and BMRB [9]). While this inconvenience 45 is relatively minor for benchmarking a single target Protein and water forcefields have been the subject 46 (e.g. water), it becomes prohibitive for studies span-47 ning chemical space. To ameliorate problems of data 48 archival, the NIST Thermodynamics Research Center 49 has developed a IUPAC standard XML-based format— 50 ThermoML [10]—for storing physicochemical measure-51 ments, uncertainties, and metadata. Experimental re-52 searchers publishing measurements in several journals 53 (J. Chem. Eng. Data, J. Chem. Therm., Fluid Phase 54 Equil., Therm. Acta, and Int. J. Therm.) are now guided 55 through a data archival process that involves sanity 56 checks and archival at the TRC (http://trc.nist. 57 gov/ThermoML.html).

> Here we examine the ThermoML archive as a poten-59 tial source for neat liquid density and static dielectric 60 constant measurements, with the goal of developing a 61 standard benchmark for validating these properties in 62 fixed-charge forcefields of drug-like molecules. These 63 two observables provide sensitive tests of forcefield ac-64 curacy that are nonetheless straightforward to calculate.

6

10

^{*} kyle.beauchamp@choderalab.org

[†] julie.behr@choderalab.org

[‡] patrick.grinaway@choderalab.org

[§] bas.rustenburg@choderalab.org

[¶] kenneth.kroenlein@nist.gov

^{**} Corresponding author; john.chodera@choderalab.org

65 Using these data, we evaluate the generalized Amber small molecule forcefield (GAFF) [11] with the AM1-BCC charge model [12, 13] and identify systematic biases that might be improved upon.

II. RESULTS

A. Neat Liquid Measurements in the ThermoML Data 70 Archive 71

69

81

82

85

86

89

90

91

92

93

95

We performed a number of sequential queries to sum-72 73 marize the ThermoML content relevant for benchmark-74 ing organic molecule forcefields. Our aim is to explore 75 neat liquid data with functional groups relevant to druglike molecules. We therefore applied the following ordered sequence filters, starting with all data containing density or static dielectric constants:

- The measured solution contains only a single com-79 ponent (e.g. no binary mixtures)
 - 2. The molecule contains only the druglike elements (H, N, C, O, S, P, F, Cl, Br)
- 3. The molecule has fewer than or equal to 10 heavy 83 atoms 84
 - 4. The measurement was performed under ambient temperature [K] $(270 \le T \le 330)$
- The measurement was performed under ambient 87 pressure [kPA] $(100 \le P \le 102)$ 88
 - 6. Measured densities below 300 kg m^{-3} were discarded; this criterion eliminated all non-liquid data in the collection.
 - 7. The temperature and pressure were rounded to 111 nearby values (as described below), averaging all 112 measurements within each group of like conditions.
- 8. Only conditions (molecule, temperature, pressure) 96 for which both density and dielectric constants were available were retained.

99 The temperature and pressure rounding step was motivated by common data reporting variations; for example, an experiment performed at water's freezing point at ambient pressure might be entered as either 101.325 103 kPA or 100 kPA, with a temperature of either 273 K 104 or 273.15 K. Therefore all pressures within the range 121 105 [kPA] $(100 \le P \le 102)$ were rounded to exactly one at-106 mosphere. Temperatures were rounded to one decimal place. The application of these filters (Table I) leaves 245 123 108 conditions for which both density and dielectric data are 109 available. The functional groups present are summa- 125 mark of electrostatic models. We therefore compare sim-110 rized in Table II.

Filter	Mass Density	Static Dielectric
1. Single Component	130074	1649
2. Druglike Elements	120410	1649
3. Heavy Atoms	67897	1567
4. Temperature	36827	962
5. Pressure	13598	461
6. Liquid state	13573	461
7. Aggregate T, P	3573	432
8. Density+Dielectric	245	245

TABLE I. Number of ThermoML measurements matching sequentially applied filters.

Functional Group	Counts
1,2-aminoalcohol	4
1,2-diol	3
alkene	3
aromatic compound	1
carbonic acid diester	2
carboxylic acid ester	4
dialkyl ether	7
heterocyclic compound	3
ketone	2
lactone	1
primary alcohol	19
primary aliphatic amine (alkylamine)	2
primary amine	2
secondary alcohol	4
secondary aliphatic amine (dialkylamine)	2
secondary aliphatic/aromatic amine (alkylarylamine)	1
secondary amine	3
sulfone	1
sulfoxide	1
tertiary aliphatic amine (trialkylamine)	3
tertiary amine	3

TABLE II. Functional group counts present in the dataset. The number of unique compounds is 44. Functional group classification was performed using checkmol version 0.5 [14].

B. Benchmarking GAFF against the ThermoML Data **Archive: Mass Density**

Mass density has been widely used for parameter-114 izing and testing forcefields, particularly the Lennard Jones parameters [15, 16]. We therefore used the present ThermoML compilation as a benchmark of the GAFF / 117 AM1-BCC forcefield (Fig. 1). Overall, the densities show 118 reasonable accuracy (RMS percent error: 3 % \pm 0.1%), consistent with previous studies [5] reporting agreement of 4 % on a different benchmark set.

C. Benchmarking GAFF / AM1-BCC against the ThermoML Data Archive: Static Dielectric

As a measure of the electronic medium, the static dielectric constant of neat liquids provides a critical bench-126 ulations against the measurements in our ThermoML

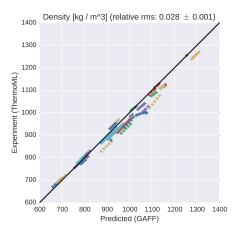


FIG. 1. Comparison of liquid densities between experiment and simulation. Liquid density measurements extracted from ThermoML are compared against densities predicted using the GAFF / AM1-BCC small molecule fixed-charge forcefield. Color groupings represent identical chemical species. Simulation error bars represent one standard error of the mean, with the number of effective (uncorrelated) samples estimated using pymbar. Experimental error bars indicate the standard deviation between independently reported measurements, when available, or author-reported standard deviations in ThermoML entries; for some measurements, neither uncertainty estimate is available. See **Section B** for further discussion of error.

127 compilation. Overall, we find the dielectric constants to be qualitatively reasonable, but with clear deviations from experiment. In particular, GAFF / AM1-BCC systematically underestimates the dielectric constants for nonpolar organics, with the predictions of $\epsilon \approx 1.0 \pm 0.05$ being substantially smaller than the measured $\epsilon \approx 2$. Because this deviation likely stems from the lack of electronic polarization, we added a simple empirical correction for polarization [17] that is based on counting the 136 elements in a molecule:

$$\frac{\alpha}{\dot{A}} = 1.53n_C + 0.17n_H + 0.57n_O + 1.05n_N + 2.99n_S + 2.48n_P + 0.22n_F + 2.16n_{Cl} + 3.29n_{Br} + 5.45n_I + 0.32$$
(1)

From the polarizability, one can correct the static dielectric using the following equation (from ref. [18]):

$$\epsilon_{corrected} = \epsilon_{MD} + 4\pi N \frac{\alpha}{\langle V \rangle}$$

A similar polarization correction was used in the devel- 163 opment of the TIP4P-EW water model [18]; however, the need is much greater for the nonpolar organics, as 165 unanticipated consequences. the missing polarizability is the dominant contribution 166

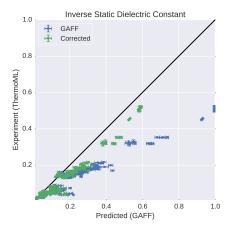


FIG. 2. Measured (ThermoML) versus predicted (GAFF / AM1-BCC) inverse static dielectrics (a). Simulation error bars represent one standard error of the mean estimated via block averaging with block sizes of 200 ps [20]. Experimental error bars indicate the larger of standard deviation between independently reported measurements and the authors reported standard deviations; for some measurements, neither uncertainty estimate is available. See section B for further discussion of error. The inverse dielectric $\frac{1}{\epsilon}$ is plotted instead of ϵ because $\frac{1}{\epsilon}$ is directly proportional to energy in continuum dielectric models: e.g. $U(r) = \frac{1}{4\pi\epsilon} \frac{q_1 q_2}{r} \propto \frac{1}{\epsilon}$.

and saw similarly improved agreement with experiment for both the GAFF and OPLS forcefields (Fig. 7).

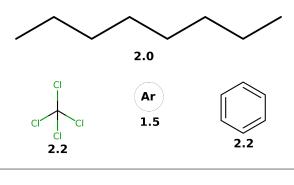
III. DISCUSSION

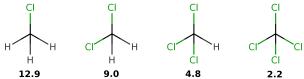
Fitting Forcefields to Dielectric Constants

Recent forcefield development has seen a resurgence 151 of papers fitting dielectric constants as primary data 152 [6, 21]. However, a number of authors have pointed out potential challenges in constructing self-consistent fixed-charge forcefields [22, 23].

Interestingly, a recent work by Dill [22] pointed out that, for CCl₄, reasonable choices of point charges are incapable of recapitulating the observed dielectric of $\epsilon = 2.2$, instead producing dielectric constants in the range of $1.0 \le \epsilon \le 1.05$. This behavior is quite general: a fixed charge monopole force field predicts $\epsilon \approx 1$ for several nonpolar or symmetric molecules, but the measured dielectric constants are instead $\epsilon \approx 2$ (Fig. 3). While this behavior is well-known and results from missing physics of polarizability, we suspect it may have several

Suppose, for example, that one attempts to fit forceto the static dielectric constant. In the case of water, 167 field parameters to match the static dielectric constants the Sales polarizability model predicts a dielectric cor- 168 of CCl₄, CHCl₃, CH₂Cl₂, and CH₃Cl. In moving from rection of 0.52, while 0.79 was used for the TIP4P-EW 169 the tetrahedrally-symmetric CCl₄ to CHCl₃, it suddenly model. For comparison, we also applied the same em- 170 becomes possible to achieve the observed dielectric con-145 pirical correction to the VirtualChemistry dataset [5, 19] 171 stant of 4.8 by an appropriate choice of point charges.





nonpolar or symmetric molecules [?]; fixed-charge forcefields give $\epsilon \approx 1$ for each species. (b). A congeneric series of chloromethanes show dielectric constants between 2 and 13.

172 However, the model for CHCl₃ uses fixed point charges to account for both the permanent dipole moment and the electronic polarizability, whereas the CCl₄ model 175 contains no treatment of polarizability. We hypothesize that this inconsistency in parameterization may lead to 177 strange mismatches, where symmetric molecules (e.g. benzene, CCl₄) have qualitatively different properties than closely related asymmetric molecules (e.g. toluene,

As a possible real-world example, we imagine that 181 182 the missing polarizability could be important in accurate transfer free energies involving low-dielectric solvents. Using the Onsager model for the transfer free energy of a dipole (Eqn. 2) gives an error of $\Delta\Delta G =$ 186 $\Delta G(\epsilon=2.2) - \Delta G(\epsilon=1)$ of -2 kcal / mol for the transfer of water ($a=1.93~\text{Å}~\mu=2.2\text{D}$) into a low dielectric 188 medium such as tetrachloromethane or benzene.

$$\Delta G = -\frac{\mu^2}{a^3} \frac{\epsilon - 1}{2\epsilon + 1} \tag{2}$$

Similarly, we calculated the mean polarization error 235 the 643 molecules, suggesting that the missing polariz- 242 (http://media.iupac.org/namespaces/ dicted solvation properties of druglike molecules.

201 electric constants are potentially usable as primary data 247 data and apply the data filters described above.

202 for forcefield parameterization efforts. Although this will require the use of forcefields with explicit polarizability, the inconsistency of fixed-charge models in low-dielectric media is sufficiently alarming to motivate further study of polarizable forcefields. In particular, continuum methods [24–26], induced dipole methods [27, 28], and drude methods [29, 30] have been matur-209 ing rapidly. Finding the optimal balance of accuracy 210 and performance remains an open question; however, 211 the use of experimentally-parameterized direct polar-212 ization methods [31] may provide polarizability physics 213 at a cost not much greater than fixed charge forcefields.

ThermoML as a Data Source

The present work has focused on the neat liquid den-FIG. 3. (a). Measured static dielectric constants of various 216 sity and dielectric measurements present in the Ther-217 moML Data Archive [10, 32, 33] as a target for molecular dynamics forcefield validation. While densities and 219 dielectric constants have been widely used in forcefield 220 work, several aspects of ThermoML make it a unique 221 resource for the forcefield community. First, the ag-222 gregation, support, and dissemination of ThermoML is 223 supported by NIST, whose mission makes these tasks a ²²⁴ long-term priority. Second, ThermoML is actively grow-225 ing, through partnerships with several journals—new 226 experimental measurements published in these journals 227 are critically examined by the TRC and included in the 228 archive. Finally, the files in the ThermoML Data Archive 229 are machine readable via a formal XML schema, allow-230 ing facile access to thousands of measurements. In the 231 future, we hope to examine additional measurement 232 classes, including both mixture and two-phase data.

IV. METHODS

ThermoML Processing

ThermoML XML files were obtained from the 190 for solvation free energies of druglike molecules in cy- 236 the NIST TRC. To explore their content, we creclohexane. For each molecule in the FreeSolv database 297 ated a python (version 2.7.9) tool (ThermoPyl: [8], we estimated a as the half the maximum interatomic 238 https://github.com/choderalab/ThermoPyL) distance and calculated $\mu = \sum_i q_i r_i$ using the provided 239 that munges the XML content into a spreadsheetmol2 coordinates and AM1-BCC charges. This calcula- 240 like format accessible via the Pandas (version tion predicts a mean error of -0.9 ± 0.07 kcal / mol for $_{241}$ 0.15.2) library. First, we obtained the XML schema abilty physics contributes substantially to errors in pre- 243 ThermoML.xsd) defining the layout of the 244 data. This schema was converted into a Python object Given their ease of measurement and direct connec- 245 via PyXB 1.2.4 (http://pyxb.sourceforge.net/). 200 tion to long-range electrostatic interactions, static di- 246 Finally, this schema and Pandas was used to extract the

B. Simulation

248

261

280

281

282

283

284

286

287

289

Boxes of 1000 molecules were constructed using 249 ²⁵⁰ PackMol [34]. AM1-BCC [12, 13] charges were gen-251 erated using OpenEye Toolkit 2014-6-6 [35], using the oequacpac.OEAssignPartialCharges module with the OECharges_AM1BCCSym. The selected conformer was then processed using antechamber in AmberTools 14 [36]. The resulting AMBER files were con- 302 conda. 260

Molecular dynamics simulations were performed us- 308 be found to serve as well. 262 ing OpenMM 6.2 [37] using a Langevin integrator (with 263 collision rate 1 ps^{-1}) and a 1 fs timestep; interestingly, ²⁶⁴ we found that a 2 fs timestep led to insufficient accuracy 265 in equilibrium densities (Table III). [JDC: Cite Langevin 266 integrator used in OpenMM.] Pressure coupling at 1 at-267 mosphere was achieved with a Monte Carlo barostat uti-268 lizing molecular scaling and automated step size adjust-269 ment during equilibration, applied every 25 steps. Particle mesh Ewald [39] was used with a long-range cutoff 271 of 0.95 nm and an long-range isotropic dispersion cor-²⁷² rection. [JDC: Can we report the automatically-selected 273 PME parameters?] Simulations were continued until density standard errors were less than 2×10^{-4} g / mL, 275 as estimated using the equilibration detection module in 276 pymbar 2.1 [40]. Trajectory analysis was performed us-277 ing OpenMM [37] and MDTraj [38]. Density data was 278 output every 250 fs, while trajectory data was stored ev-279 ery 10 ps.

V. CONCLUSIONS

- ThermoML is a potentially useful resource for the forcefield community
- We have curated a subset of the ThermoML Data Archive for neat liquids with druglike atoms, with thousands of densities and hundreds of dielectrics
- Empirical polarization models correct a systematic bias in comparing fixed-charge forcefields to static dielectric constants

VI. ACKNOWLEDGEMENTS

We thank Vijay S. Pande (Stanford University), Lee-291 Ping Wang (Stanford University), Peter Eastman (Stan-292 ford University), Robert McGibbon (Stanford University), Jason Swails (Rutgers University), David L. Mob-294 ley (University of California, Irvine), Christopher 295 I. Bayly (OpenEye Software), Michael R. Shirts (Uni-296 versity of Virginia), and members of Chodera lab for

297 helpful discussions. Support for JMB was provided 298 by the Tri-Institutional Training Program in Computa-299 tional Biology and Medicine (via NIH training grant 300 1T32GM083937).

VII. DISCLAIMERS

This contribution of the National Institute of Stanverted to OpenMM [37] ffxml forcefield XML files. Sim- 303 dards and Technology is not subject to copyright in the ulation code used libraries gaff2xml 0.6, TrustButVer- 304 United States. Products or companies named here are ify 0.1, OpenMM 6.2 [37], and MDTraj 1.2 [38]. [TODO: 305 cited only in the interest of complete technical descrip-Provide a script to install all of these versions via 306 tion, and neither constitute nor imply endorsement by 307 NIST or by the U.S. government. Other products may

TABLE III. To probe the systematic error from finite time-step integration, we examined the timestep dependence of butyl acrylate density. The number of effective samples was estimated using pymbar's statistical inefficiency routine [40]. To approximate the timestep bias, we compare the density expectation ($\langle \rho \rangle$) to values calculated with a 0.5fs timestep. We find a 2fs timestep leads to systematic biases in the density on the order of 0.2%, while 1fs reduces the systematic bias to less than 0.1%—we therefore selected a 1fs timestep for the present work, where we aimed to achieve three digits of accuracy in density predictions.

Appendix A: Supplementary Information

309

315

316

317

318

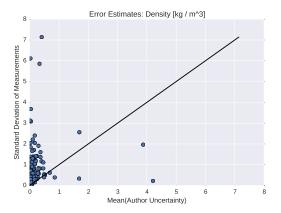
319

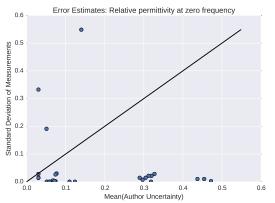
320

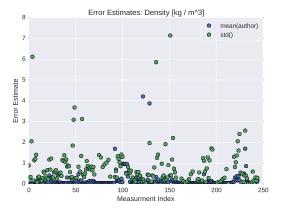
All information below this point will eventually be pulled into a separate SI. This will happen closer to sub-mission, as the formatting may be journal-specific. The references may be split in two as well, depending on journal.

- Table: Timestep-dependence of density
- Figure: Error analysis for ThermoML dataset
- Table (CSV File): ThermoML Dataset used in present analysis.

Appendix B: Assessment of experimental error in ThermoML measurements







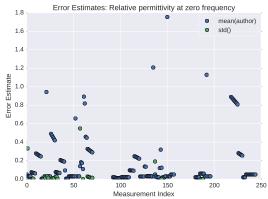


FIG. 4. Assessment of experimental error in ThermoML data. To assess the experimental error in our benchmark set, we consider two orthogonal error measurements. The first is the mean of the uncertainties reported by the measurement authors. The second is the standard deviation of measurements in ThermoML. We see that author-reported uncertainties appear to be overly optimistic for densities (a, c), but author-reported uncertainties of dielectrics (b, d) appear consistent with the standard deviations. A simple psychological expla-

nation might be that because density measurements are



FIG. 5. Comparison of simulated and experimental densities for all compounds. Measured (blue) and simulated (green) densities are shown in units of kg/m^3 .

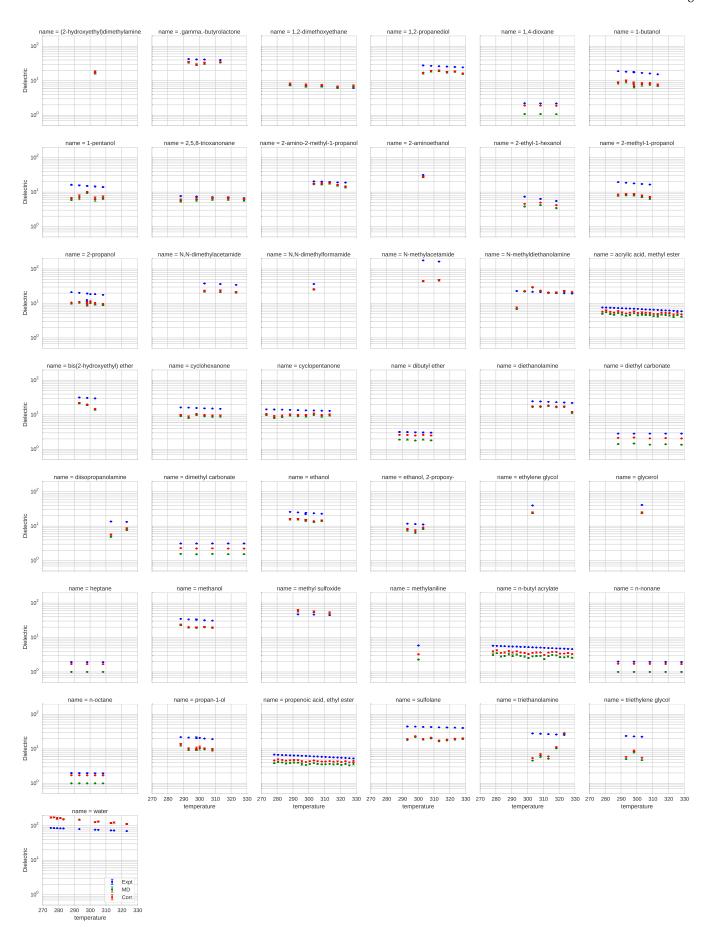
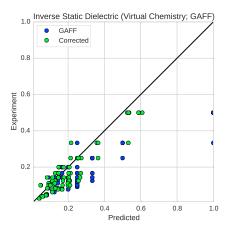


FIG. 6. Comparison of simulated and experimental static dielectric constants for all compounds. Measured (blue), simulated (green), and polarizability-corrected simulated (red) static dielectric constants are shown for all compounds. Note that dielectric constants, rather than inverse dielectric constants, are plotted here.



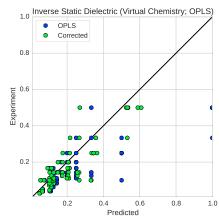


FIG. 7. Comparison of measured and simulated dielectric constants from the virtualchemistry dataset with and without polarizability correction. Measured (blue), MD (green), and MD + polarizability-corrected (red) dielectrics for the Virtual Chemistry dataset [5, 19].

[1] K. Lindorff-Larsen, P. Maragakis, S. Piana, M. Eastwood, 375 R. Dror, and D. Shaw, PloS one 7, e32131 (2012).

321

322

323

324

325

326 327

328

329

330

331

332

333

334

335

336

337

338

339

340

341

- D.-W. Li and R. Bruschweiler, J. Chem. Theory Comput. 377 7, 1773 (2011).
- [3] R. B. Best, X. Zhu, J. Shim, P. E. Lopes, J. Mittal, M. Feig, and A. D. MacKerell, J. Chem. Theory Comput. (2012).
- [4] K. Lindorff-Larsen, S. Piana, K. Palmo, P. Maragakis, 381 J. Klepeis, R. Dror, and D. Shaw, Proteins: Struct., Funct., 382 Bioinf. 78, 1950 (2010).
- C. Caleman, P. J. van Maaren, M. Hong, J. S. Hub, L. T. 384 [27] Costa, and D. van der Spoel, Journal of chemical theory 385 and computation 8, 61 (2011).
- of Physical Chemistry B (2014).
- [7] H. M. Berman, J. Westbrook, Z. Feng, G. Gilliland, T. N. 389 [29] Bhat, H. Weissig, I. N. Shindyalov, and P. E. Bourne, Nu- 390 cleic Acids Res. 28, 235 (2000).
- [8] D. L. Mobley, Experimental and calculated small molecule 392 hydration free energies, Retrieved from: http://www. 393 Department of Pharmaceutical Sciences, UCI.
- [9] E. Ulrich, H. Akutsu, J. Doreleijers, Y. Harano, Y. Ioanni- 396 342 dis, J. Lin, M. Livny, S. Mading, D. Maziuk, and Z. Miller, 397 [32] 343 Nucleic Acids Res. 36, D402 (2008). 344
- M. Frenkel, R. D. Chiroco, V. Diky, Q. Dong, K. N. 399 345 Marsh, J. H. Dymond, W. A. Wakeham, S. E. Stein, 400 346 E. Königsberger, and A. R. Goodwin, Pure and applied 401 347 chemistry 78, 541 (2006). 348
- J. Wang, R. M. Wolf, J. W. Caldwell, P. A. Kollman, and 403 349 D. A. Case, J. Comput. Chem. 25, 1157 (2004). 350
- 351 12] A. Jakalian, B. L. Bush, D. B. Jack, and C. I. Bayly, J. Com- 405 put. Chem. 21, 132 (2000). 352
- 13] A. Jakalian, D. B. Jack, and C. I. Bayly, J. Comput. Chem. 407 353 **23**, 1623 (2002). 354
- [14] N. Haider, Molecules 15, 5079 (2010). 355
- [15] W. L. Jorgensen, J. Chandrasekhar, J. D. Madura, R. W. 410 356 Impey, and M. L. Klein, The Journal of chemical physics 411 357 79, 926 (1983). 358
- [16] W. L. Jorgensen, J. D. Madura, and C. J. Swenson, Journal 413 359 of the American Chemical Society 106, 6638 (1984). 360
- R. Bosque and J. Sales, Journal of chemical information 415 361 and computer sciences 42, 1154 (2002). 362
- H. Horn, W. Swope, J. Pitera, J. Madura, T. Dick, G. Hura, 417 363 and T. Head-Gordon, J. Chem. Phys. 120, 9665 (2004). 364
- D. van der Spoel, P. J. van Maaren, and C. Caleman, Bioin- 419 365 formatics 28, 752 (2012). 366
- H. Flyvbjerg and H. G. Petersen, J. Chem. Phys. 91, 461 421 367 (1989).368
- L.-P. Wang, T. J. Martínez, and V. S. Pande, The Journal of 423 369 Physical Chemistry Letters (2014). 370
- C. J. Fennell, L. Li, and K. A. Dill, The Journal of Physical 425 371 Chemistry B 116, 6936 (2012). 372
- I. V. Leontyev and A. A. Stuchebrukhov, The Journal of 427 373 chemical physics 141, 014103 (2014). 374

- [24] J.-F. Truchon, A. Nicholl's, J. A. Grant, R. I. Iftimie, B. Roux, and C. I. Bayly, Journal of computational chemistry 31, 811 (2010).
- J.-F. Truchon, A. Nicholls, B. Roux, R. I. Iftimie, and C. I. Bayly, Journal of chemical theory and computation 5, 1785 (2009).
- J.-F. Truchon, A. Nicholls, R. I. Iftimie, B. Roux, and C. I. Bayly, Journal of chemical theory and computation 4, 1480 (2008).
- J. Ponder, C. Wu, P. Ren, V. Pande, J. Chodera, M. Schnieders, I. Haque, D. Mobley, D. Lambrecht, R. DiStasio Jr, et al., J. Phys. Chem. B 114, 2549 (2010).
- C. J. Fennell, K. L. Wymer, and D. L. Mobley, The Journal 387 [28] P. Ren and J. W. Ponder, The Journal of Physical Chemistry B 108, 13427 (2004).
 - G. Lamoureux and B. Roux, The Journal of Chemical Physics 119, 3025 (2003).
 - 391 [30] V. M. Anisimov, G. Lamoureux, I. V. Vorobyov, N. Huang, B. Roux, and A. D. MacKerell, Journal of Chemical Theory and Computation 1, 153 (2005).
- escholarship.org/uc/item/6sd403pz, uC Irvine: 394 [31] L.-P. Wang, T. L. Head-Gordon, J. W. Ponder, P. Ren, J. D. 395 Chodera, P. K. Eastman, T. J. Martínez, and V. S. Pande, J. Phys. Chem. B 117, 9956 (2013).
 - M. Frenkel, R. D. Chirico, V. V. Diky, Q. Dong, S. Frenkel, P. R. Franchois, D. L. Embry, T. L. Teague, K. N. Marsh, and R. C. Wilhoit, Journal of Chemical & Engineering Data 48, 2 (2003).
 - [33] R. D. Chirico, M. Frenkel, V. V. Diky, K. N. Marsh, and 402 R. C. Wilhoit, Journal of Chemical & Engineering Data 48, 1344 (2003).
 - [34] L. Martínez, R. Andrade, E. G. Birgin, and J. M. Martínez, Journal of computational chemistry **30**, 2157 (2009).
 - [35] Openeye toolkits 2014, URL http://www.eyesopen. 406
 - [36] D. Case, V. Babin, J. Berryman, R. Betz, Q. Cai, D. Cerutti, 408 T. Cheatham III, T. Darden, R. Duke, H. Gohlke, et al., 409 University of California, San Francisco (2014).
 - P. Eastman, M. S. Friedrichs, J. D. Chodera, R. J. Radmer, C. M. Bruns, J. P. Ku, K. A. Beauchamp, T. J. Lane, L.-P. 412 Wang, D. Shukla, et al., J. Chem. Theory Comput. 9, 461 (2012).414
 - [38] R. T. McGibbon, K. A. Beauchamp, C. R. Schwantes, L.-P. Wang, C. X. Hernández, M. P. Harrigan, T. J. Lane, J. M. Swails, and V. S. Pande, bioRxiv p. 008896 (2014).
 - T. Darden, D. York, and L. Pedersen, J. Chem. Phys. 98, 418 10089 (1993).
 - M. R. Shirts and J. D. Chodera, J. Chem. Phys. 129, 124105 420 (2008).
 - Mettler toledo density meters, [Online; accessed 15-Jan-422 URL http://us.mt.com/us/en/home/ products/Laboratory_Analytics_Browse/ Density_Family_Browse_main/DE_Benchtop. tabs.models-and-specs.html.
 - R. D. Chirico, M. Frenkel, J. W. Magee, V. Diky, C. D. 428 Muzny, A. F. Kazakov, K. Kroenlein, I. Abdulagatov, G. R. Hardin, and W. E. Acree Jr, Journal of Chemical & Engi-429 neering Data 58, 2699 (2013). 430