Benchmarking Atomistic Simulations against the ThermoML Data Archive: **Neat Liquid Densities and Static Dielectric Constants**

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Useful atomistic simulations in the condensed phase require accurate depictions of solvent. While experimental measurements of fundamental physical properties offer a straightforward approach for evaluating forcefield quality, the bulk of this information has been tied up in formats that are not machine-readable. These formats require substantial human effort to compile benchmark datasets which are prone to accumulation of human errors, hindering the development of reproducible benchmarks of forcefield accuracy. Here, we examine the feasibility of benchmarking atomistic forcefields against the NIST ThermoML data archive of physicochemical measurements, which aggregates thousands of experimental measurements in a portable, machine-readable, self-annotating format. As a proof of concept, we present a detailed benchmark of the generalized Amber small molecule forcefield (GAFF) using the AM1-BCC charge model against measurements (specifically liquid densities and static dielectric constants at ambient pressure) automatically extracted from the archive, and discuss the extent of available data for neat liquids. The results of this benchmark highlights a general problem with fixed-charge forcefields in the representation of liquids of low dielectric.

Keywords: molecular mechanics forcefields; forcefield parameterization; forcefield accuracy; forcefield validation; mass density; static dielectric constant

I. INTRODUCTION

Recent advances in hardware and software for molec-13 ular dynamics simulation now permits routine access to atomistic simulations at the 100 ns timescale and beyond [1]. Leveraging these advances in combination with consumer GPU clusters, distributed computing, or custom hardware has brought microsecond and millisecond simulation timescales within reach of many laboratories. These ¹⁹ dramatic advances in sampling, however, have revealed deficiencies in forcefields as a critical barrier to enabling truly predictive simulations of physical properties of biomolecular systems.

Protein and water forcefields have been the subject of numerous benchmarks [2] and enhancements [3-5], with key outcomes including the ability to fold fast-folding proteins [6-8], improved fidelity of water thermodynamic properties [9], and improved prediction of NMR observables. Although small molecule forcefields have also been the subject of benchmarks [10] and improvements [11], such 30 work has typically focused on small perturbations to spe-31 cific functional groups. For example, a recent study found

32 that modified hydroxyl nonbonded parameters led to im-33 proved prediction of static dielectric constants and hydra-34 tion free energies [11]. There are also outstanding ques-35 tions of generalizability of these targeted perturbations; it is uncertain whether changes to the parameters for a specific chemical moiety will be compatible with seemingly un-38 related improvements to other groups. Addressing these 39 questions requires establishing a community agreement on 40 shared benchmarks that can be easily replicated among 41 laboratories to test proposed forcefield enhancements and 42 expanded as the body of experimental data grows.

A key barrier to establishing reproducible and extensible 44 forcefield accuracy benchmarks is that many experimen-45 tal datasets are heterogeneous, paywalled, and unavail-46 able in machine-readable formats (although notable coun-47 terexamples exist, e.g. the RCSB [12], FreeSolv [13], and 48 the BMRB [14]). While this inconvenience is relatively mi-49 nor for benchmarking forcefield accuracy for a single tar-50 get (e.g. water), it becomes prohibitive for studies span-51 ning the relevant chemical space. To ameliorate prob-52 lems of data archival, the NIST Thermodynamics Research 53 Center (TRC) has developed a IUPAC standard XML-based 54 format—ThermoML [15]—for storing physicochemical mea-55 surements, uncertainties, and metadata. Experimental 56 researchers publishing measurements in several journals 57 (J. Chem. Eng. Data, J. Chem. Therm., Fluid Phase Equil., 58 Therm. Acta, and Int. J. Therm.) are guided through a data 59 archival process that involves sanity checks, conversion to a 50 standard machine-readable format, and archival at the TRC 61 (http://trc.nist.gov/ThermoML.html).

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Here, we examine the ThermoML archive as a potential 112 Cite Langevin integrator used in OpenMM.] Pressure conlation codes—neat liquid density and static dielectric condard benchmark for validating these properties in fixed-70 charge forcefields of drug-like molecules and biopolymer 120 ducibility in other codes?] Simulations were continued un-71 residue analogues. These two properties provide sensitive tests of forcefield accuracy that are nonetheless straightforward to calculate. Using these data, we evaluate the gener-₇₅ AM1-BCC charge model [17, 18] and identify systematic bi- 125 sities were stored every 250 fs, while trajectory snapshots ₇₆ ases to aid further forcefield refinement.

METHODS

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ThermoML Processing

A tarball archive of the ThermoML Archive was obtained from the the NIST TRC on 13 Sep 2014. explore the content of this archive, we created a Python (version 2.7.9) tool (ThermoPyL: https://github.com/choderalab/ThermoPyL) formats the XML content into a spreadsheet-like format accessible via the Pandas (version 0.15.2) library. First, we obtained the XML schema (http://media.iupac.org/ namespaces/ThermoML/ThermoML.xsd) defining the layout of the data. This schema was converted into a Python object via PyXB 1.2.4 (http://pyxb.sourceforge.net/). 90 Finally, this schema and Pandas was used to extract the 91 data and apply the successive data filters described in 92 Section III A.

Simulation

Using an automated tool, boxes of 1000 molecules were constructed using PackMol version 14-225 [19]. AM1-BCC [17, 18] charges were generated using OpenEye Toolkit 2014-6-6 [20], using the oequacpac. OEAssignPartialCharges module with the OECharges_AM1BCCSym option, which utilizes a conformational expansion procedure prior to charge fitting to minimize artifacts from intramolecular contacts. The selected conformer was then processed using antechamber in AmberTools 14 [21]. The resulting AMBER files 150 were converted to OpenMM [22] ffxml forcefield XML files. Simulation code used libraries gaff2xml 0.6, TrustButVerify 0.1, OpenMM 6.2 [22], and MDTraj 1.3 [23]. [TODO: Provide a script to install all of these versions via conda.]

Molecular dynamics simulations were performed with OpenMM 6.2 [22] using a Langevin integrator (with collision $_{109}$ rate 1 ps $^{-1}$) and a 1 fs timestep, as we found that timesteps of $_{155}$ ¹¹⁰ 2 fs timestep or greater led to a significant timestep depenun dence in computed equilibrium densities (Table 4). [JDC: 157

source for providing the foundation for a reproducible, ex- 113 trol to 1 atm was achieved with a Monte Carlo barostat utitensible accuracy benchmark of biomolecular forcefields. 114 lizing molecular scaling and automated step size adjust-In particular, we concentrate on two important physical 115 ment during equilibration, with volume moves attempted property measurements easily computable in many simu- 116 every 25 steps. The particle mesh Ewald (PME) method [24] was used with a long-range cutoff of 0.95 nm and a longstant measurements—with the goal of developing a stan- 118 range isotropic dispersion correction. [JDC: Can we report 119 the automatically-selected PME parameters to aid reprotil density standard errors were less than 2×10^{-4} g / mL, as estimated using the equilibration detection module in pymbar 2.1 [25]. Trajectory analysis was performed using alized Amber small molecule forcefield (GAFF) [16] with the 124 OpenMM 6.2 [22] and MDTraj 1.3 [23]. Instantaneous denwere stored every 10 ps. [JDC: Did we plan to make this data available somewhere, or is it sufficient to put out the scripts?] Static cielectric constants were calculated using the dipole fluctuation approach [9], with box dipoles calcu-130 lated from trajectory files using MDTraj 1.3 [23].

$$\epsilon = 1 + \frac{4\pi}{3} \frac{\langle \mu \cdot \mu \rangle - \langle \mu \rangle \cdot \langle \mu \rangle}{\langle V \rangle k_B T} \tag{1}$$

III. RESULTS

Extracting neat liquid measurements from the NIST TRC ThermoML Archive

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We retrieved a copy of the ThermoML Archive from 135 the NIST TRC (http://trc.nist.gov/ThermoML.html ac-136 cessed 13 Sep 2014) and performed a number of sequen-137 tial filtering steps to produce an extract of the ThermoML 138 Archive relevant for benchmarking forcefields describing 139 small organic molecules. As our aim is to explore neat liquid data with functional groups relevant to biopolymers and drug-like molecules, we applied the following ordered filters, starting with all data containing density or static dielec-143 tric constants:

- 1. The measured solution contains only a single component (e.g. no binary mixtures)
- 2. The molecule contains only the druglike elements (defined here as H, N, C, O, S, P, F, Cl, Br)
- 3. The molecule has \leq 10 heavy atoms
- 4. The measurement was performed in a biophysically relevant temperature range [K] $(270 \le T \le 330)$
- 5. The measurement was performed at ambient pressure [kPA] $(100 \le P \le 102)$
- 6. Measured densities below 300 kg m^{-3} were discarded to eliminate gas-phase measurements
- 7. The temperature and pressure were rounded to nearby values (as described below), averaging all measurements within each group of like conditions

	Number of measurements remaining	
Filter step	Mass density	Static dielectric
1. single component	130074	1649
2. only druglike elements	120410	1649
3. ≤10 heavy atoms	67897	1567
4. $(270 \le T \le 330)$ [K]	36827	962
5. ambient pressure	13598	461
6. liquid state	13573	461
7. aggregate T, P	3573	432
8. density and dielectric	245	245

TABLE I. Successive filtration of the ThermoML Archive. A set of successive filters were applied to all measurements in the ThermoML Archive (accessed 13 Sep 2014) that contained either mass density or static dielectric constant measurements. Each column reports the number of measurements remaining after successive application of the corresponding filtration step.

 Only conditions (molecule, temperature, pressure) for which both density and dielectric constants were available were retained

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The temperature and pressure rounding step was motivated by common data reporting variations; for example, an experiment performed at water's freezing point at ambient pressure might be entered as either 101.325 kPA or 100 kPA, with a temperature of either 273 K or 273.15 K. Therefore all pressures within the range [kPA] $(100 \le P \le 102)$ were rounded to exactly one atmosphere. Temperatures were rounded to one decimal place. The application of these filters (Table I) leaves 245 conditions—where a condition here indicates a (molecule, temperature, pressure) tuple—for which both density and dielectric data are available. The functional groups present in the resulting dataset are summarized in Table II; see Methods for further description of the software pipeline used.

B. Benchmarking GAFF/AM1-BCC against the ThermoML Archive

1. Mass density

Mass density has been widely used for parameterizing and testing forcefields, particularly the Lennard-Jones parameters representing dispersive and repulsive interactions [27, 28]. We therefore used the present ThermoML extract as a benchmark of the GAFF/AM1-BCC forcefield (Fig. 1). We estimate the mass density via $\rho = \langle \frac{M}{V} \rangle$.

Overall, the densities show reasonable accuracy, with a root-mean square (RMS) relative error over all measurements of $2.8\pm0.1\%$ (with one standard error of the mean determined by bootstrapping over all measurements), especially encouraging given that this forcefield was not designed with the intention of modeling bulk liquid properties of organic molecules [16] This is reasonably consistent with previous studies reporting agreement of 4% on a different benchmark set [10].

[JDC: Discuss outliers here. There must be more things

Functional Group	Occurrences
1,2-aminoalcohol	4
1,2-diol	3
alkene	3
aromatic compound	1
carbonic acid diester	2
carboxylic acid ester	4
dialkyl ether	7
heterocyclic compound	3
ketone	2
lactone	1
primary alcohol	19
primary aliphatic amine (alkylamine)	2
primary amine	2
secondary alcohol	4
secondary aliphatic amine (dialkylamine)	2
secondary aliphatic/aromatic amine (alkylarylamine)	1
secondary amine	3
sulfone	1
sulfoxide	1
tertiary aliphatic amine (trialkylamine)	3
tertiary amine	3

TABLE II. Functional groups present in filtered dataset. The filtered ThermoML dataset contained 245 distinct (molecule, temperature, pressure) conditions, spanning 44 unique compounds. The functional groups represented in these compounds (as identified by the program checkmol v0.5 [26]) is summarized here.

we can say about densities. Some of the densities are quite good, while others seem poor, with systematic bias toward higher densities than experiment. We can also point out that densities at different temperatures for a given molecule seem to be biased in a consistent way.]

2. Static dielectric constant

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As a measure of the dielectric response, the static dielectric constant of neat liquids provides a critical benchmark of the accuracy electrostatic treatment in forcefield models. We therefore compare simulations against the measurements in our ThermoML extract. Overall, we find the dielectric constants to be qualitatively reasonable, but with clear deviations from experiment. In particular, GAFF/AM1-BCC 207 systematically underestimates the dielectric constants for 208 nonpolar organics, with the predictions of $\epsilon \approx 1.0 \pm 0.05$ being substantially smaller than the measured $\epsilon \approx 2$. Be-210 cause this deviation likely stems from the lack of an explicit treatment of electronic polarization, we used a sim-212 ple empirical polarization model that computes the molecular electronic polarizability α a sum of elemental atomic polarizability contributions [29]. From the computed molecular electronic polarizability lpha, an additive correction to the 216 simulation-derived static dielectric constant accounting for the missing electronic polarizability can be computed [9]

$$\Delta \epsilon = 4\pi N \frac{\alpha}{\langle V \rangle} \tag{2}$$

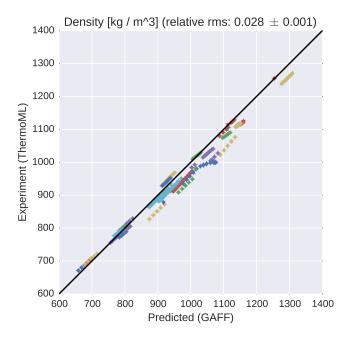


FIG. 1. Comparison of liquid densities between experiment and simulation. Liquid density measurements extracted from ThermoML are compared against densities predicted using the GAFF/AM1-BCC small molecule fixed-charge forcefield. Color groupings represent identical chemical species. Simulation error bars represent one standard error of the mean, with the number of effective (uncorrelated) samples estimated using pymbar. Experimental error bars indicate the standard deviation between independently reported measurements, when available, or authorreported standard deviations in ThermoML entries; for some measurements, neither uncertainty estimate is available. See SI Fig. 5 for further discussion of error.

While a similar polarization correction was used in the development of the TIP4P-Ew water model, where it had a minor effect [9], missing polarizability is a dominant contribution to the static dielectric constant of nonpolar organic molecules; in the case of water, the empirical atomic polarizability model predicts a dielectric correction of 0.52, while 0.79 was used for the TIP4P-Ew model. Considering all predictions in the present work leads to polarizability corrections to the static dielectric of 0.74 ± 0.08 .

DISCUSSION

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Fitting Forcefields to Dielectric Constants

Recent forcefield development has seen a resurgence 252 of papers fitting dielectric constants during forcefield parameterization [11, 32]. However, a number of authors have pointed out potential challenges in constructing self- 255 consistent fixed-charge forcefields [33, 34].

served that, for CCl₄, reasonable choices of point charges ²⁵⁸ lated asymmetric molecules (e.g. toluene and CHCl₃). are incapable of recapitulating the observed dielectric of 259

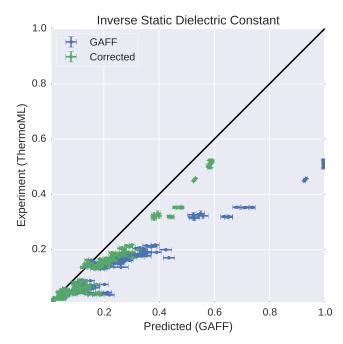
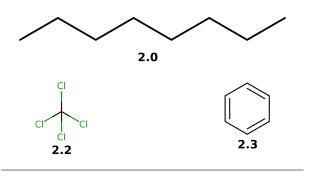


FIG. 2. Measured (ThermoML) versus predicted (GAFF/AM1-BCC) inverse static dielectrics (a). Simulation error bars represent one standard error of the mean estimated via circular block averaging [30] with block sizes automatically selected to maximize the error [31]. Experimental error bars indicate the larger of standard deviation between independently reported measurements and the authors reported standard deviations; for some measurements, neither uncertainty estimate is available. See SI Fig. 5 for further discussion of error. The inverse dielectric constant ϵ^{-1} is plotted instead of ϵ because ϵ^{-1} is directly proportional to the Coulomb interaction energy between point charges embedded in a dielectric material [e.g. $U(r) \propto q_1 q_2 / r \propto \epsilon^{-1}$].

of $1.0 \le \epsilon \le 1.05$. This behavior is quite general: fixed point charge forcefields will predict $\epsilon \approx 1$ for many nonpolar or symmetric molecules, but the measured dielectric ₂₄₁ constants are instead $\epsilon \approx 2$ (Fig. 3). While this behavior 242 is well-known and results from missing physics of polariz-243 ability, we suspect it may have several unanticipated consequences, which we discuss below.

Suppose, for example, that one attempts to fit forcefield parameters to match the static dielectric constants of CCl₄, CHCl₃, CH₂Cl₂, and CH₃Cl. In moving from the tetrahedrally-symmetric CCl₄ to the asymmetric CHCl₃, it suddenly becomes possible to achieve the observed dielectric constant of 4.8 by an appropriate choice of point charges. However, the model for CHCl3 uses fixed point charges to account for both the permanent dipole moment 253 and the electronic polarizability, whereas the CCl₄ model contains no treatment of polarizability. We hypothesize that this inconsistency in parameterization may lead to strange 256 mismatches, where symmetric molecules (e.g. benzene and Interestingly, recent work by Dill and coworkers [33] ob- 257 CCl₄) have qualitatively different properties than closely re-

How important is this effect? As a possible real-world $\epsilon = 2.2$, instead producing dielectric constants in the range 260 example, we imagine that the missing atomic polarizabil-



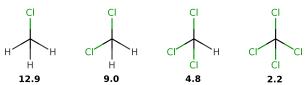


FIG. 3. Typical experimental static dielectric constants of some **nonpolar compounds.** (a). Measured static dielectric constants of various nonpolar or symmetric molecules [35, 36]. Fixed-charge forcefields give $\epsilon \approx 1$ for each species; for example, we calculated $\epsilon = 1.003 \pm 0.0002$ for octane. (b). A congeneric series of chlorosubstituted methanes have static dielectric constants between 2 and 13. Reported dielectric constants are at near-ambient temperatures.

261 ity could be important in accurate transfer free energies 262 involving low-dielectric solvents. The Onsager model for the transfer free energy of a dipole (Eq. 3) gives an error of $\Delta\Delta G = \Delta G(\epsilon=2.2) - \Delta G(\epsilon=1)$ of -2 kcal / mol for the transfer of water (a=1.93 Å, $\mu=2.2$ D) into a low dielectric 266 medium such as tetrachloromethane or benzene.

$$\Delta G = -\frac{\mu^2}{a^3} \frac{\epsilon - 1}{2\epsilon + 1} \tag{3}$$

Similarly, we calculated the mean polarization error for solvation free energies (gas to solvent transfer free energies) of druglike molecules in cyclohexane. For each molecule in 322 the latest (Feb. 20) FreeSolv database [13, 37], we took the cavity radius a to be the half the maximum interatomic distance and calculated $\mu=\sum_i q_i r_i$ using the provided mol² coordinates and AM1-BCC charges. This calculation predicts a mean error of -0.91 ± 0.07 kcal / mol for the 643 molecules 325 (where the standard error is computed from bootstrapping 326 over measurements), suggesting that the missing atomic 327 polarizabilty unrepresentable by fixed point charge forcefields could contribute substantially to errors in predicted 328 transfer and solvation properties of druglike molecules. We 329 also conjecture that

Given their ease of measurement and direct connection to long-range electrostatic interactions, static dielectric constants have high potential utility as primary data for force- 331 field parameterization efforts. Although this will require the use of forcefields with explicit treatment of atomic polar-286 izability, the inconsistency of fixed-charge models in low- 333 Ping Wang (Stanford University), Peter Eastman (Stanford 287 dielectric media is sufficiently alarming to motivate fur- 334 University), Robert McGibbon (Stanford University), Jason

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288 ther study of polarizable forcefields. In particular, continuum methods [38-40], point dipole methods [41, 42], and Drude methods [43, 44] have been maturing rapidly. Finding the optimal balance of accuracy and performance remains an open question; however, the use of experimentallyparameterized direct polarization methods [45] may provide polarizability physics at a cost not much greater than fixed charge forcefields.

ThermoML as a data source

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The present work has focused on the neat liquid density and dielectric measurements present in the ThermoML Archive [15, 46, 47] as a target for molecular dynamics forcefield validation. While liquid mass densities and static dielectric constants have already been widely used in forcefield work, several aspects of ThermoML make it a unique resource for the forcefield community. First, the aggregation, support, and dissemination of ThermoML datasets 305 through the ThermoML Archive is supported by NIST, whose 306 mission makes these tasks a long-term priority. Second, 307 the ThermoML Archive is actively growing, through partnerships with several journals, and new experimental mea-309 surements published in these journals are critically examined by the TRC and included in the archive. [JDC: Is the 311 number of journal here also expanding?] Finally, the files in the ThermoML Archive are portable and machine readable via a formal XML schema, allowing facile access to hundreds of thousands of measurements. Numerous additional physical properties contained in ThermoML—including ac-316 tivity coefficients, diffusion constants, boiling point temper-317 atures, critical pressures and densities, coefficients of expansion, speed of sound measurements, viscosities, excess molar enthalpies, heat capacities, and volumes—for neat (3) 320 phases and mixtures represent a rich dataset of high utility 321 for forcefield validation and parameterization.

CONCLUSIONS

- ThermoML is a potentially useful resource for the forcefield community
- We have curated a subset of the ThermoML Data Archive for neat liquids with druglike atoms, with thousands of densities and hundreds of dielectrics
- Empirical polarization models correct a systematic bias in comparing fixed-charge forcefields to static dielectric constants

VI. ACKNOWLEDGEMENTS

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VII. DISCLAIMERS

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Appendix A: Supplementary Information

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All information below this point will eventually be pulled into a separate SI. This will happen closer to submission, as the formatting may be journal-specific. The references may be split in two as well, depending on journal. [JDC: It may be fine to leave this as an Appendix.]

- Figure: Timestep-dependence of density
- Figure: Error analysis for ThermoML dataset
- Table (CSV File): ThermoML Dataset used in present analysis.

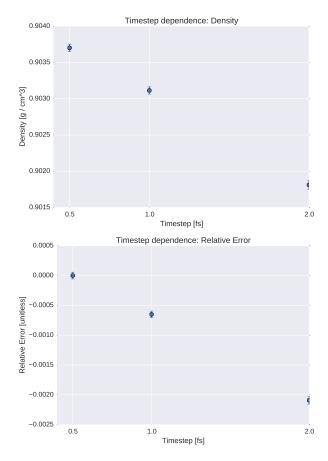


FIG. 4. To probe the systematic error from finite time-step integration, we examined the timestep dependence of butyl acrylate density. (a). The density is shown for several choices of timestep. (b). The relative error, as compared to the 0.5 fs value, is shown for several choices of timestep. Error bars represent stand errors of the mean, with the number of effective samples estimated using pymbar's statistical inefficiency routine [25]. We find a 2 fs timestep leads to systematic biases in the density on the order of 0.2%, while 1 fs reduces the systematic bias to less than 0.1%—we therefore selected a 1 fs timestep for the present work, where we aimed to achieve three digits of accuracy in density predictions.

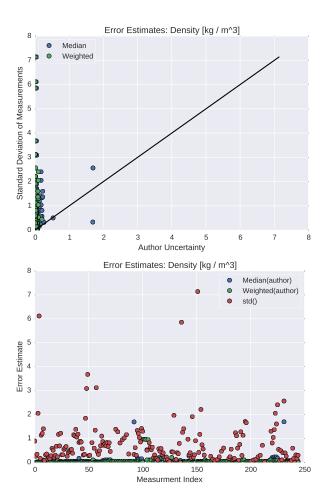
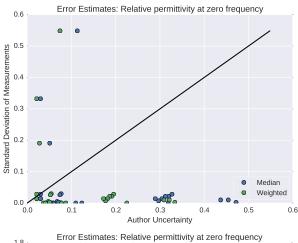


FIG. 5. Assessment of experimental error: Density To assess the experimental error in our ThermoML extract, we compared three difference estimates of uncertainty. In the first approach (Weighted), we computed the standard deviation of the optimally weighted average of the measurements, using the uncertainties reported by authors ($\sigma_{Weighted} = [\sum_k \sigma_k^{-2}]^{-0.5}$). This uncertainties tainty estimator places the highest weights on measurements with small uncertainties and is therefore easily dominated by small outliers and uncertainty under-reporting. In the second approach (Median), we estimated the median of the uncertainties reported by authors; this statistic should be robust to small and large outliers of author-reported uncertainties. In the third approach (Std), we calculated at the standard deviation of independent measurements reported in the ThermoML extract, completely avoiding the author-reported uncertainties. Plot (a) compares the three uncertainty estimates. We see that author-reported uncertainties appear to be substantially smaller than the scatter between the observed measurements. A simple psychological explanation might be that because density measurements are more routine, the authors simply report the accuracy limit of their hardware (e.g. 0.0001 g/mL for a Mettler Toledo DM40 [48]). However, this hardware limit is not achieved due to inconsistencies in sample preparation; see Appendix in Ref. [49]. Panel (b) shows the same information as (a) but as a function of the measurement index, rather than as a scatter plot—because not all measurements have authorsupplied uncertainties, panel (c) contains slightly more data points than (a, b).



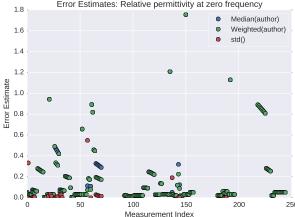


FIG. 6. Assessment of experimental error: Static Dielectric Constant To assess the experimental error in our ThermoML extract, we compared three difference estimates of uncertainty. In the first approach (Weighted), we computed the standard deviation of the optimally weighted average of the measurements, using the uncertainties reported by authors ($\sigma_{Weighted} = [\sum_k \sigma_k^{-2}]^{-0.5}$). This uncertainty estimator places the highest weights on measurements with small uncertainties and is therefore easily dominated by small outliers and uncertainty under-reporting. In the second approach (Median), we estimated the median of the uncertainties reported by authors; this statistic should be robust to small and large outliers of author-reported uncertainties. In the third approach (Std), we calculated at the standard deviation of independent measurements reported in the ThermoML extract, completely avoiding the author-reported uncertainties. Plot (a) compares the three uncertainty estimates. Unlike the case of densities, author-reported uncertainties appear to be somewhat larger than the scatter between the observed measurements. Panel (b) shows the same information as (a) but as a function of the measurement index, rather than as a scatter plot—because not all measurements have author-supplied uncertainties, panel (c) contains slightly more data points than (a, b).



FIG. 7. Comparison of simulated and experimental densities for all compounds. Measured (blue) and simulated (green) densities are shown in units of kg/m³.

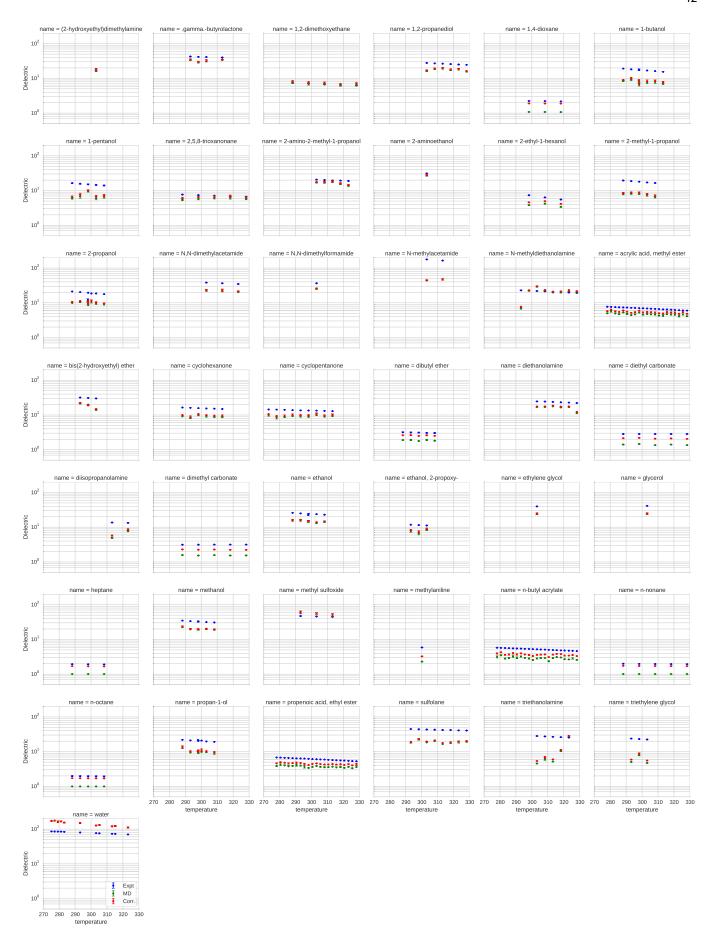


FIG. 8. Comparison of simulated and experimental static dielectric constants for all compounds. Measured (blue), simulated (green), and polarizability-corrected simulated (red) static dielectric constants are shown for all compounds. Note that dielectric constants, rather than inverse dielectric constants, are plotted here. [JDC: Let's plot these as in Fig. 1 and Fig. 2, maybe only four plots across so they are larger and more legible. We can also shorten "name = compound" to just "compound".] KAB: We should discuss this more before I rebuild this figure several times.

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