# Development of a Computer Program for Electronic Structure Calculation using Lagrange Basis Functions

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#### Outline

- ▶ Introduction
- ► Kohn-Sham equations
- ► Lagrange basis functions
- ► Implementation
- ► Numerical results

#### Introduction

- Electronic structure calculations play important role for investigation of materials properties.
- Much of electronic structure calculations are based on Kohn-Sham density functional theory.
- ► Several computer packages can already carry out electronic structure calculations for systems containing of few atoms (characteristic length in angstrom) to ten thousands atoms (10-10³ nm)



#### **Problems**

- ▶ These computer programs can "simplify" electronic structure calculations, espescially for non-specialists. However, they are not suitable for development of new methodologies. It is quite difficult to extend the program if we want to do some customized calculations.
- ▶ Much details are hidden beneath the programs. They usually contain thousands of lines codes or more which can be very intimidating for beginner developers to work with.

#### This research

- ▶ We will to write our own electronic structure calculations program.
- ▶ Downside: It might takes several years to reach the same level of maturity with state-of-the-art well-established electronic structure programs.
- ▶ In this presentation, I will describe some of our preliminary works to implement electronic calculation based on Kohn-Sham equations and Lagrange basis functions and results of the calculations for several simple systems.

## Kohn-Sham total energy functional

Total energy of a system of interacting electrons according to Kohn-Sham can be written as:

$$\begin{split} E_{\mathrm{tot}}\left[\{\psi_{i_{st}}(\mathbf{r})\},\rho(\mathbf{r})\right] &= E_{\mathrm{kin}} + E_{\mathrm{ion}} + E_{\mathrm{Ha}} + E_{\mathrm{xc}} \\ \rho(\mathbf{r}) &= \sum_{i_{st}=1}^{N_{st}} f_{i_{st}} \psi_{i_{st}}^{*}(\mathbf{r}) \psi_{i_{st}}(\mathbf{r}) \\ E_{\mathrm{kin}} &= -\frac{1}{2} \sum_{i_{st}} \int f_{i_{st}} \psi_{i_{st}}^{*}(\mathbf{r}) \nabla^{2} \psi_{i_{st}}(\mathbf{r}) \, \mathrm{d}\mathbf{r} \\ E_{\mathrm{ion}} &= \int V_{\mathrm{ion}}(\mathbf{r}) \, \rho(\mathbf{r}) \, \mathrm{d}\mathbf{r} \\ E_{\mathrm{Ha}} &= \frac{1}{2} \int \frac{\rho(\mathbf{r}) \rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} \mathrm{d}\mathbf{r} \mathrm{d}\mathbf{r}' \\ E_{\mathrm{xc}} &= \int \epsilon_{\mathrm{xc}} \left[\rho(\mathbf{r})\right] \rho(\mathbf{r}) \, \mathrm{d}\mathbf{r} \end{split}$$

### Kohn-Sham equations

Ground state energy can be found by minimizing the Kohn-Sham energy functional or by solving the Kohn-Sham equation:

$$\left[ -\frac{1}{2} \nabla^2 + V_{KS}(\mathbf{r}) \right] \psi_{i_{st}}(\mathbf{r}) = \epsilon_{i_{st}} \psi_{i_{st}}(\mathbf{r})$$
 (1)

with the following potential terms

$$egin{align} V_{
m KS}({f r}) &= V_{
m ion}({f r}) + V_{
m Ha}({f r}) + V_{
m xc}({f r}) \ V_{
m Ha}({f r}) &= \int rac{
ho({f r}')}{{f r}-{f r}'} \, {
m d}{f r}' \ 
onumber \ 
abla^2 V_{
m Ha}({f r}) &= -4\pi
ho({f r}) \ 
onumber \$$

### Kohn-Sham equations

- ▶ In the the implementation to solve Kohn-Sham equations, single-electron wave function must be represented as an expansion of basis functions or discretized using some discretization schemes such as finite-difference or finite-elements
- Several existing computer programs for DFT calculations, such as QUANTUM ESPRESSO and GAUSSIAN09, expand single-electron wave function using plane wave and Gaussian functions, respectively.
- ▶ Others, such as Octopus and GPAW use discretization based on finite-difference scheme.
- ▶ In this work, we will use expansion based on Lagrange basis functions, which is relatively new and currently there are no widely-available programs which implement this.

### Lagrange basis functions

For a given interval [0, L], with L > 0, the grid points  $x_{\alpha}$  appropriate for periodic Lagrange function are given by:

$$x_{\alpha} = \frac{L}{2} \frac{2\alpha - 1}{N} \tag{2}$$

with  $\alpha=1,\ldots,N$ . Number of points N should be an odd number. The periodic cardinal functions  $\phi_{\alpha}(x)$ , defined at grid point i are given by:

$$\phi_{\alpha}(x) = \frac{1}{\sqrt{NL}} \sum_{n=1}^{N} \cos\left(\frac{\pi}{L} (2n - N - 1)(x - x_{\alpha})\right). \tag{3}$$

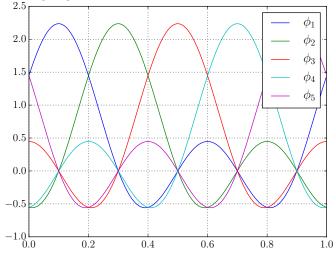
The expansion of periodic function in terms of Lagrange functions:

$$f(x) = \sum_{\alpha=1}^{N} c_{\alpha} \phi_{\alpha}(x) \tag{4}$$

with expansion coefficients  $c_{\alpha} = \sqrt{L/N} f(x_{\alpha})$ .

### Lagrange basis functions

Plot of for Lagrange basis functions for  $N=5,\ L=1.0$ 



# Using Lagrange basis function for Schrodinger equation

Given the 1D Schrodinger equation:

$$\left[-\frac{1}{2}\frac{\mathrm{d}^2}{\mathrm{d}x^2} + V(x)\right]\psi(x) = \epsilon\psi(x)$$

Lagrange basis function to expand the one-electron wavefunction:

$$\psi(x) = \sum_{\alpha}^{N} c_{\alpha} \phi_{\alpha}(x)$$

Matrix equation:

$$(T + V)C = \epsilon C$$

Analytic expression for matrix  ${\bf T}$  can be obtained from reference  $^1$ . Matrix  ${\bf V}$  is diagonal. Eigenvalue  $\epsilon$  can be found using standard eigenvalue solver.

¹See for example: *J. Phys. Chem. A* **110**, 5549-5560, (2006) → ⟨ ₹ ⟩ ⟨

#### Extension to 3D

Expansion of function in 3D using Lagrange basis function:

$$\psi(\mathbf{r} = (x, y, z)) = \sum_{\alpha} \sum_{\beta} \sum_{\gamma} C_{\alpha\beta\gamma} \phi_{\alpha}(x) \phi_{\beta}(y) \phi_{\gamma}(z)$$

Similar matrix equation can be obtained in 3D case. Potential matrix is still diagonal and kinetic matrix now is expressed as:

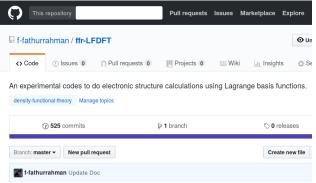
$$\mathbf{T}_{\alpha\beta\gamma}^{\alpha'\beta'\gamma'} = \mathbf{T}_{\alpha\alpha'}\delta_{\beta\beta'}\delta_{\gamma\gamma'} + \mathbf{T}_{\beta\beta'}\delta_{\alpha\alpha'}\delta_{\gamma\gamma'} + \mathbf{T}_{\gamma\gamma'}\delta_{\alpha\alpha'}\delta_{\beta\beta'}$$

where  $\mathbf{T}_{ii'}$ ,  $i=\alpha,\beta,\gamma$  are kinetic matrix for 1D case. Kinetic matrix is now sparse.

Similar matrix representation can be found in for finite-difference methods  $^{2}.$ 

### Our implementation

Public repository: https://github.com/f-fathurrahman/ffr-LFDFT.



#### Numerical results

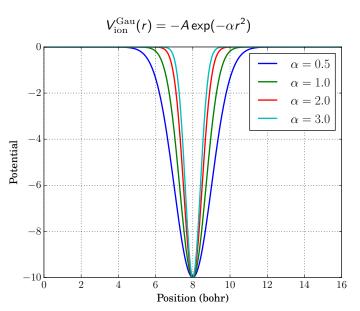
Total energy calculation on the following simple potentials:

- Gaussian potentials: convergence of total energy with respect to grid spacing
- atomic pseudotentials: hydrogen and lithium pseudopotentials

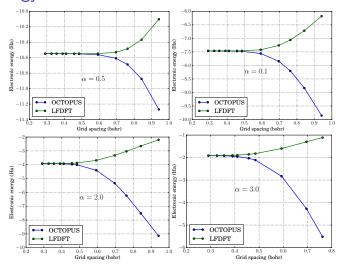
All calculations are done in  $16\times16\times16$  bohr periodic box. Center of the potential is set to the center of the box.

We will also validate the results of our program against result from well-established program  ${\rm OCTOPUS}$  which uses finite difference method.

### Gaussian potentials

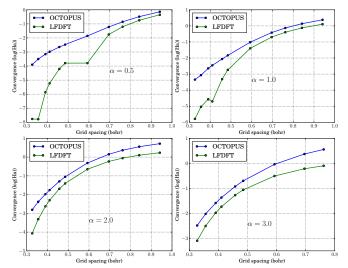


#### Total energy



Both our program (LFDFT) and OCTOPUS converge to the same total energy value

### Total energy convergence



LFDFT converges faster than Octopus (finite difference) For smoother potential (smaller  $\alpha$ ) both methods converge faster (typicall for real-space methods)

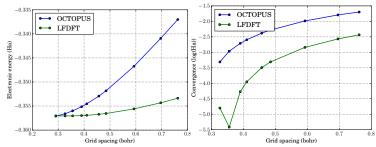


## Hydrogen atom

$$V_{\text{loc}}^{\text{H}}(\mathbf{r}) = -\frac{Z_{\text{ion}}}{r} \operatorname{erf}\left(\frac{r}{\sqrt{2}r_{\text{loc}}}\right) + \exp\left[-\frac{1}{2}\left(\frac{r}{r_{\text{loc}}}\right)^{2}\right] \times \left[C_{1} + C_{2}\left(\frac{r}{r_{\text{loc}}}\right)^{2} + \right]$$
(5)

with the following parameters:  $Z_{\rm ion}=1$ ,  $r_{\rm loc}=0.2$ ,  $C_1=-4.1802372$ , and  $C_2=0.725075$ .  $^3$ 





LFDFT converges faster than OCTOPUS (finite difference)

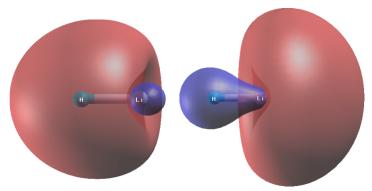
#### Lithium atom

$$\begin{split} V_{\rm loc}(\mathbf{r}) &= -\frac{Z_{\rm ion}}{r} \mathrm{erf}\left(\frac{r}{\sqrt{2}r_{\rm loc}}\right) + \\ \exp\left[-\frac{1}{2}\left(\frac{r}{r_{\rm loc}}\right)^2\right] \times \left[C_1 + C_2\left(\frac{r}{r_{\rm loc}}\right)^2 + C_3\left(\frac{r}{r_{\rm loc}}\right)^4 + C_4\left(\frac{r}{r_{\rm loc}}\right)^6\right] \end{split}$$

with the following parameters:  $Z_{\rm ion}=$  3,  $r_{\rm loc}=$  0.4,  $C_1=-$ 14.034868,  $C_2=$  9.553476,  $C_3=-$ 1.7664885 and  $C_4=$  0.084370.  $^4$ 



#### LiH HOMO and LUMO



HOMO (left image) and LUMO (right image) of LiH

#### Remarks and future works

#### Remarks

- ▶ Use of Lagrange basis functions to Kohn-Sham equation gives matrix-representation which is very similar to the one obtained by using finite-difference methods
- Compared to finite difference method, use of Lagrange basis functions gives better total energy convergence with respect to grid size (number of basis functions).

#### Future works will focus on:

- ► Larger systems
- Parallization