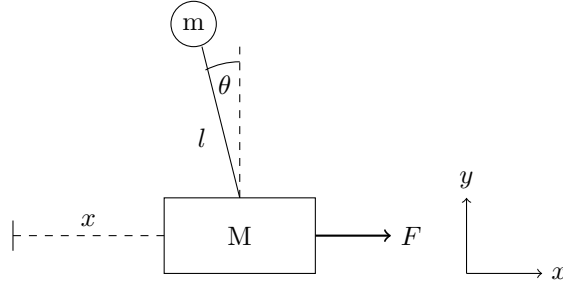


# Contents

## 1 Equations of Motion

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- A cart of mass  $M$  is constrained to move along the  $x$  axis with its distance from an arbitrary point on the  $x$  axis denoted  $x$ . A driving force of magnitude  $F$  is applied to the cart in the  $x$  direction. A simple pendulum consisting of a mass  $m$  and a massless rod of length  $l$  is connected to the cart with its angle from the positive  $y$  axis denoted  $\theta$ .
- The kinetic energy of the cart is

$$T_{\text{cart}} = \frac{1}{2}M\dot{x}^2.$$

- The  $x$  and  $y$  coordinates of the pendulum are

$$\begin{aligned} X &= x - l \sin \theta \\ Y &= l \cos \theta, \end{aligned}$$

thus its  $x$  and  $y$  velocities are

$$\begin{aligned} \dot{X} &= \dot{x} - l\dot{\theta} \cos \theta \\ \dot{Y} &= -l\dot{\theta} \sin \theta \end{aligned}$$

and its kinetic energy is

$$\begin{aligned} T_{\text{pendulum}} &= \frac{1}{2}m\dot{v}^2 \\ &= \frac{1}{2}(\dot{X}^2 + \dot{Y}^2) \\ &= \frac{1}{2}m[(\dot{x} - l\dot{\theta} \cos \theta)^2 + (-l\dot{\theta} \sin \theta)^2] \\ &= \frac{1}{2}m(\dot{x}^2 - 2l\dot{x}\dot{\theta} \cos \theta + l^2\dot{\theta}^2). \end{aligned}$$

- The total kinetic energy of the system is

$$\begin{aligned} T &= T_{\text{cart}} + T_{\text{pendulum}} \\ &= \frac{1}{2}(m + M)\dot{x}^2 + \frac{1}{2}m(l^2\dot{\theta}^2 - 2l\dot{x}\dot{\theta}\cos\theta). \end{aligned}$$

- The potential energy of the system is equal to the gravitational potential energy of the pendulum. If its potential energy is 0 when  $\theta = \frac{\pi}{2}$  then

$$U = mgl \cos \theta.$$

- The Lagrangian of the system is

$$\begin{aligned} \mathcal{L} &= T - U \\ &= \frac{1}{2}(m + M)\dot{x}^2 + \frac{1}{2}m(l^2\dot{\theta}^2 - 2l\dot{x}\dot{\theta}\cos\theta) - mgl \cos \theta. \end{aligned}$$

- By d'Alembert's principle the generalized forces associated with the  $\theta$  and  $x$  coordinates are 0 and  $F$ , respectively.
- The Euler-Lagrange equation for the  $\theta$  coordinate is

$$\begin{aligned} \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\theta}} - \frac{\partial \mathcal{L}}{\partial \theta} &= 0 \\ \frac{d}{dt}(ml^2\dot{\theta} - ml\dot{x}\cos\theta) - ml\dot{x}\sin\theta - mgl\sin\theta &= 0 \\ l\ddot{\theta} - \ddot{x}\cos\theta - g\sin\theta &= 0. \end{aligned}$$

As the goal is to keep  $\theta$  close to 0 this can be linearized by using the approximations  $\cos\theta \approx 1$  and  $\sin\theta \approx \theta$  giving

$$l\ddot{\theta} - \ddot{x} - g\theta = 0.$$

- The Euler-Lagrange equation for the  $x$  coordinate is

$$\begin{aligned} \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{x}} - \frac{\partial \mathcal{L}}{\partial x} &= F \\ \frac{d}{dt}[(m + M)\dot{x} - ml\dot{\theta}\cos\theta] &= F \\ (m + M)\ddot{x} - ml\ddot{\theta}\cos\theta + ml\dot{\theta}^2\sin\theta &= F. \end{aligned}$$

Using the same approximations as above gives

$$(m + M)\ddot{x} - ml\ddot{\theta} + ml\dot{\theta}^2\theta = F.$$