

Multiplicative Hitchin Systems and Supersymmetric Gauge Theory

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Abstract

1 Introduction

(Chris: There are a lot of notes in red: that just indicates that there are a lot of arguments that are only partially written. They're notes for myself.)

2 Multiplicative Higgs Bundles and q -Connections

We'll begin with an abstract definition of the moduli spaces we'll investigate in this paper using the language of derived algebraic geometry.

Let G be a reductive complex algebraic group, let C be a smooth complex algebraic curve and fix a finite set $D = \{z_1, \dots, z_k\}$ of closed points in C . We write $\mathrm{Bun}_G(C)$ for the moduli stack of G -bundles on C , which we view as a mapping stack $\underline{\mathrm{Map}}(C, BG)$ into the classifying stack of G .

Definition 2.1. The moduli stack of *multiplicative G -Higgs fields* on C with singularities at D is the fiber product

$$\mathrm{mHiggs}_G(C, D) = \mathrm{Bun}_G(C) \times_{\mathrm{Bun}_G(C \setminus D)} \underline{\mathrm{Map}}(C \setminus D, G/G)$$

where G/G is the adjoint quotient stack.

Remark 2.2. A closed point of $\mathrm{mHiggs}_G(C, D)$ consists of a principal G -bundle P on C along with an automorphism of the restriction $P|_{C \setminus D}$, i.e. a section of $\mathrm{ad}(P)$ away from D .

The adjoint quotient stack can also be described as the derived loop space $\underline{\mathrm{Map}}(S_B^1, BG)$ of the classifying stack, where S_B^1 is the “Betti stack” of S^1 , i.e. the constant derived stack at the simplicial set S^1 . We can therefore view $\underline{\mathrm{Map}}(C \setminus D, G/G)$ instead as the mapping stack $\underline{\mathrm{Map}}((C \setminus D) \times S_B^1, BG)$, and the moduli stack of multiplicative Higgs bundles instead as

$$\mathrm{mHiggs}_G(C, D) = \underline{\mathrm{Map}}((C \times S_B^1) \setminus (D \times \{0\}), BG).$$

The source of this mapping stack can be q -deformed. Indeed, let q denote an automorphism of the curve C . Write $C \times_q S_B^1$ for the *mapping torus* of q , i.e the derived fiber product

$$C \times_q S_B^1 = C \times_{C \times C} C$$

where the two maps $C \rightarrow C \times C$ are given by the diagonal and the q -twisted diagonal $(1, q)$ respectively.

Definition 2.3. The moduli stack of q *difference connections* for the group G on C with singularities at D is the mapping space

$$q\text{-Conn}_G(C, D) = \underline{\text{Map}}((C \times_q S_B^1) \setminus (D \times \{0\}), BG).$$

In particular when $q = 1$ this recovers the moduli stack of multiplicative Higgs bundles.

Remark 2.4. A closed point of $q\text{-Conn}_G(C, D)$ consists of a principal G -bundle P on C along with a q *difference connection*: an isomorphism of G -bundles $P|_{C \setminus D} \rightarrow q^*P|_{C \setminus D}$ away from the divisor D .

2.1 Residue Conditions

These moduli stacks are typically of infinite type. In order to obtain finite type stacks, and later in order to define symplectic rather than only Poisson structure, we can fix the behaviour of our multiplicative Higgs fields and q difference connections near the punctures $D \subseteq C$.

We'll write \mathbb{D} to denote the *formal disk* $\text{Spec } \mathbb{C}[[z]]$. Likewise we'll write \mathbb{D}^\times for the *formal punctured disk* $\text{Spec } \mathbb{C}((z))$. We'll then write \mathbb{B} for the derived pushout $\mathbb{D} \sqcup_{\mathbb{D}^\times} \mathbb{D}$. Let $LG = \underline{\text{Map}}(\mathbb{D}^\times, G)$ and let $L^+G = \underline{\text{Map}}(\mathbb{D}, G)$.

There is a canonical inclusion $\mathbb{B}^{\sqcup k} \rightarrow (C \times_q S_B^1) \setminus (D \times \{0\})$, the inclusion of the formal punctured neighbourhood of $D \times \{0\}$. This induces a restriction map on mapping spaces

$$\text{res}_D: q\text{-Conn}_G(C, D) \rightarrow \text{Bun}_G(\mathbb{B})^k.$$

The following is well-known (see e.g. the expository article [Zhu17]).

Lemma 2.5. The set of closed points of $\text{Bun}_G(\mathbb{B})$ is in canonical bijection with the set of dominant coweights of G .

Definition 2.6. Choose a map from D to the set of dominant coweights and denote it by $\omega^\vee: z_i \mapsto \omega_{z_i}^\vee$. Write Λ_i for the isotropy group of the point $\omega_{z_i}^\vee$ in $\text{Bun}_G(\mathbb{B})$. The moduli stack of q difference connections on C with singularities at D and fixed residues given by ω^\vee is defined to be the fiber product

$$q\text{-Conn}_G(C, D, \omega^\vee) = q\text{-Conn}_G(C, D) \times_{\text{Bun}_G(\mathbb{B})^k} (B\Lambda_1 \times \cdots \times B\Lambda_k).$$

Examples 2.7. The most important examples for our purposes are given by the following rational/trigonometric/elliptic trichotomy.

- **Rational:** We can enhance the definition of our moduli space by including a framing at a point $c \in C$ not contained in D . We always assume that such framed points are fixed by the automorphism q .

Definition 2.8. The moduli space of q -difference connections on C with a framing at c is defined to be the relative mapping space

$$q\text{-Conn}_G^{\text{fr}}(C) = \underline{\text{Map}}(C, G/G; f)$$

where $f: \{c\} \rightarrow G/G$ is a choice of adjoint orbit. We can define the framed mapping space with singularities and fixed residues in exactly the same way as above.

In this paper we'll be most interested in the following example. Let $C = \mathbb{CP}^1$ with framing point $c = \infty$, and consider automorphisms of the form $z \mapsto z + \varepsilon$ for $\varepsilon \in \mathbb{C}$. Choose a finite subset $D \subseteq \mathbb{A}^1$ and label the points $z_i \in D$ by dominant coweights $\omega_{z_i}^\vee$. We can then study the moduli space $\varepsilon\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D, \omega^\vee)$. The main object of study in this paper will be the holomorphic symplectic structure on this moduli space.

Note that the motivation for this definition comes in part from Spaide's formalism [Spa16] of AKSZ symplectic structures on relative mapping spaces – in this formalism \mathbb{CP}^1 with a single framing point is relatively 1-oriented, so mapping spaces out of it with 1-shifted symplectic targets have AKSZ 0-shifted symplectic structures.

- **Trigonometric:** Alternatively, we can enhance our definition by including a reduction of structure group at a point $c \in C$ not contained in D , again fixed by the automorphism q .

Definition 2.9. The moduli space of q -difference connections on C with an H -reduction at c for a subgroup $H \subseteq G$ is defined to be the fiber product

$$q\text{-Conn}_G^{H,c}(C) = \underline{\text{Map}}(C, G/G) \times_{G/G} H/H$$

associated to the evaluation at c map $\underline{\text{Map}}(C, G/G) \rightarrow G/G$. We can define the moduli space with H -reduction with singularities and fixed residues in the same way as above.

Again let $C = \mathbb{CP}^1$. Fix a pair of opposite Borel subgroups B_+ and $B_- \subseteq G$ with unipotent radicals N_\pm and consider the moduli space of q -connections with B_+ -reduction at 0 and N_- -reduction at ∞ . We'll now take q to be an automorphism of the form $z \mapsto qz$ for $q \in \mathbb{C}^\times$. We'll defer in depth analysis of this example to future work.

- **Elliptic:** Finally, let $C = E$ be a smooth curve of genus one. In this case we won't fix any additional boundary data, but just consider the moduli space $q\text{-Conn}_G(E, D, \omega^\vee)$. In the case $q = 1$ this space – or rather its polystable locus – was studied by Hurtubise and Markman [HM02], who proved that it can be given the structure of an algebraic integrable system with symplectic structure related to the elliptic R-matrix of Etingof and Varchenko [EV98].

(Vasily: What shall we say about higher genus? The definition of mHiggs as above makes sense. Will the resulting space be Poisson? I guess not because to define the Poisson structure we need to fix a section of canonical bundle on C , that is a one-form “dz”. However, in higher genus such 1-form necessarily has zeroes. May be here is a good place to comment or restrict attention to the trichotomy with an explanation why higher genus is on different grounds)

Remark 2.10. In the elliptic case it's natural to ask to what extent Hurtubise and Markman's integrable system structure can be extended from the variety of polystable multiplicative Higgs bundles to the full moduli stack. If D is empty then it's easy to see that we have a symplectic structure given by the AKSZ construction of Pantev-Toën-Vaquié-Vezzosi [PTVV13]. Indeed, E is compact 1-oriented and the quotient stack G/G is 1-shifted symplectic, so the mapping stack $\underline{\text{Map}}(E, G/G) = \text{mHiggs}_G(E)$ is equipped with a 0-shifted symplectic structure by [PTVV13, Theorem 2.5]. The role of the Hitchin fibration is played by the Chevalley map $\chi: G/G \rightarrow T/W$, and therefore

$$\underline{\text{Map}}(E, G/G) \rightarrow \underline{\text{Map}}(E, T/W).$$

The fibers of this map over regular points in T/W are given by moduli stacks of the form $\text{Bun}_T(\tilde{E})^W$ where \tilde{E} is a W -fold cover of E (the cameral cover). (Chris: is the base just T/W ? Come back to this, at least by dimension counting.) (Vasily: If there are no singularities, I think that's correct, but perhaps we can extend the comment above for the case of non-empty D ?)

2.2 Stability Conditions

For comparison to results in the literature it is important that we briefly discuss the role of stability conditions for difference connections. In our main example of interest – the rational case – these conditions won't play a role, but they do appear in the comparison results between q -connections and monopoles in the literature for more general curves. For definitions for general G we refer to [Smi15], although see also [AB01] on polystable G -bundles. In what follows we fix a choice of $0 < t_0 < 2\pi R$.

Definition 2.11. Let (P, ϕ) be a q -connection on a curve C , and let χ be a character of G . The χ -degree of (P, ϕ) is defined to be

$$\deg_\chi(P, \phi) = \deg(P \times_\chi \mathbb{C}) - \frac{t_0}{2\pi R} \sum_{i=1}^k (\chi \circ \omega_{z_i}^\vee)_*(0).$$

(Vasily: Sorry, what that last zero “...(0)” denotes above?)

A q -connection (P, ϕ) on C is *stable* if for every maximal parabolic subgroup $H \subseteq G$ with Levi decomposition $H = LN$ and every reduction of structure group (P_H, ϕ) to H , we have

$$\deg_\chi(P_H, \phi) < 0$$

for the character $\chi = \det(\text{Ad}_L^n)$ defined to be the determinant of the adjoint representation of L on \mathfrak{n} .

The q -connection (P, ϕ) is *polystable* if there exists a (not necessarily maximal) parabolic subgroup H with Levi factor L and a reduction of structure group (P_L, ϕ) to L so that (P_L, ϕ) is a stable q -connection and so that the associated H -bundle is admissible, meaning that for every character χ of H which is trivial on $Z(G)$ the associated line bundle $P_H \times_\chi \mathbb{C}$ has degree zero.

Below we'll write $q\text{-Conn}_G^{\text{ps}}(C, D, \omega^\vee)$ for the moduli space of polystable q -connections. This moduli space is a smooth algebraic variety of finite type. (Chris: Reference for this other than for q trivial?)

When $C = \mathbb{CP}^1$ every principal G -bundle on C admits an essentially unique (up to the action of the Weyl group) holomorphic reduction of structure group to a maximal torus [Gro57]. Since q -connections for an abelian group are automatically stable, polystability on the sphere is equivalent to admissibility of the torus reduction. As a consequence, for our main example of interest – the rational case – the moduli space of polystable q -connections is equivalent to the moduli space of all q -connections of admissible degree. For instance for $G = \text{SL}_n$ the moduli space of polystable q -connection is equivalent to the moduli space of q -connections on the trivial bundle.

3 Periodic Monopoles

Moduli spaces of q -connections on a Riemann surface C are closely related to moduli spaces of periodic monopoles, i.e. monopoles on 3-manifolds that fiber over the circle (more specifically, with fiber C and monodromy determined by q). Let $G_{\mathbb{R}}$ be a compact Lie group whose complexification is G . The discussion in this section will mostly follow that of [CH10, Smi15].

Definition 3.1. A *monopole* on the Riemannian 3-manifold $M = C \times_q S_R^1$ is a smooth principal $G_{\mathbb{R}}$ -bundle \mathbf{P} equipped with a connection A and a section Φ of the associated bundle $\mathfrak{g}_{\mathbf{P}}$ satisfying the Bogomolny equation

$$*F_A = d_A \Phi.$$

Remark 3.2. We should emphasise the difference between the Riemannian 3-manifold $M = C \times_q S_R^1$ appearing in this section and the derived stack $C \times_q S_B^1$ (the mapping torus) appearing in the previous section. These should be thought of as smooth and algebraic realizations of the same object (justified by the comparison Theorem 3.6) but they are a priori defined in different mathematical contexts.

(Vasily: We should specify the Riemannian metric on M in order to define the $*$ operation on forms. I guess our convention for the metric on M which leads to the symplectomorphism with mHiggs is

$$ds^2 = v \otimes \bar{v} + dt \otimes dt \tag{1}$$

where $v = "dz"$ is a chosen section of the canonical bundle on C ; I suggest to give some fixed coordinate-free name to this holomorphic volume form on C , like v or whatever, and may be even consider that that this volume form is fixed as an attribute of C through out the paper))

We can rephrase the data of a monopole on M as follows. Let $C_0 = C \times \{0\}$ be the fiber over 0 in S^1 , viewed as a Riemann surface. Let P be the restriction of the complexified bundle $\mathbf{P}_{\mathbb{C}}$ to C_0 . Consider first the restriction of the complexification of A to a connection A_0 on P over C_0 . The $(0, 1)$ part of A_0 automatically defines a holomorphic

structure on P . We can introduce an additional piece of structure on this holomorphic G -bundle. In order to do so we can decompose the Bogomolny equation into one real and one complex equation as follows.

$$\begin{aligned} F_{A_0} - \nabla_t \Phi &= 0 \\ [\bar{\partial}_{A_0}, \nabla_t - i\Phi dt] &= 0 \end{aligned}$$

where ∇_t is the component of the covariant derivative d_A normal to C_0 . From now on we'll denote $\mathcal{A} = \nabla_t - i\Phi dt$.

Let's now introduce singularities into the story. We'll keep the description brief, referring the reader to [CH10, Smi15] for details.

Definition 3.3. Let $D \subseteq M$ be a finite subset. Let ω^\vee be a choice of coweight for G . A monopole on $M \setminus D$ has *Dirac singularity* at $z \in D$ with charge ω^\vee if locally on a neighbourhood of z in M it is obtained by pulling back under ω^\vee the standard Dirac monopole solution to the Bogomolny equation where Φ is spherically symmetric with a simple pole at z (see e.g. [CH10, Section 2.2]). (Vasily: ... and the restriction of a connection A to a two-sphere S^2 enclosing the singularity defines a $U(1)$ bundle on this S^2 of degree 1 so that

$$\frac{1}{2\pi} \int_{S^2} F = 1 \quad (2)$$

)

We can also introduce a framing (or a reduction of structure group as in the trigonometric example, though we won't consider the latter in this paper). As usual let $c \in C$ be a point fixed by the automorphism q .

Definition 3.4. A monopole on M with *framing* at the point $c \in C$ is a monopole (P, A, Φ) on M (possibly with Dirac singularities at D) along with a trivialization of the restriction of P to the circle $\{c\} \times S_R^1$, with the condition that the holonomy of \mathcal{A} around this circle is 1. (Vasily: Perhaps, instead of “1” we should pick framing data like in 2.8 which was $f \in G/G$: a choice of coadjoint orbit)

The moduli theory of monopoles on 3-manifolds of this form has been studied by Foscolo [Fos16]. (Vasily: need to add refs to Kapustin Cherkis and others before Foscolo) In the cases of interest to us it can be obtained as a hyperkähler quotient. Let's focus initially on the rational case, so monopoles on $M = \mathbb{CP}^1 \times_\varepsilon S_R^1$ with Dirac singularities at $D \times \{t_0\}$ and a framing at ∞ . Consider the infinite-dimensional vector space \mathcal{V} consisting of pairs (A, Φ) where A is a connection on a fixed principal $G_{\mathbb{R}}$ -bundle P on M , Φ is a section of \mathfrak{g}_P , and (A, Φ) have a Dirac singularity with charge $\omega_{z_i}^\vee$ at each (z_i, t_0) in $D \times \{t_0\}$. Let \mathcal{G} be the group of gauge transformations of the bundle P .

The hyperkähler moment map is given by the Bogomolny functional, namely

$$\begin{aligned} \mu: \mathcal{V} &\rightarrow \Omega^1(M; (\mathfrak{g}_{\mathbb{R}})_P) \\ (A, \Phi) &\mapsto *F_A - d_A \Phi. \end{aligned}$$

Definition 3.5. Let D be a finite subset $\{(z_1, t_1), \dots, (z_k, t_k)\}$ of points in $M = \mathbb{CP}^1 \times_\varepsilon S_R^1$, and let ω_i^\vee be a choice of coweight for each point in D . The moduli space $\text{Mon}_G(M, D, \omega^\vee)$ is the hyperkähler quotient

$$\text{Mon}_G(M, D, \omega^\vee) = \mu^{-1}(0)/\mathcal{G}.$$

Now let's address the relationship between periodic monopoles and q -connections. Suppose from now on that q is in the identity component of the group of automorphisms of C (fixing the framing point c if present).

Theorem 3.6. There is an analytic isomorphism between the moduli space of polystable monopoles on $C \times_q S^1$ with Dirac singularities at $D \times \{t_0\}$ (and a possible framing on $\{c\} \times S^1$) and the moduli space of q -connections on C with singularities at D and framing at $\{c\}$. More precisely there is an analytic isomorphism

$$H: \text{Mon}_G^{(\text{fr})}(C \times_q S^1, D \times \{t_0\}, \omega^\vee) \rightarrow q\text{-Conn}^{\text{ps}, (\text{fr})}(C, D, \omega^\vee)$$

given by the holonomy map around S^1 , i.e. sending a monopole (P, \mathcal{A}) to the holomorphic bundle $P = (P_C)|_{C_0}$ with q -connection $\phi = \text{Hol}_{S^1}(\mathcal{A}): P \rightarrow q^*(P)$. (Vasily: Do we assume that $C_0 = C \times t_0$ or not necessarily the same t_0 where are the singularities?)

Proof. (Chris: check) This follows by the same argument as that given by Charbonneau-Hurtubise [CH10] and Smith [Smi15]. More explicitly, first let's think about injectivity, so let $(\mathbf{P}, \mathcal{A})$ and $(\mathbf{P}', \mathcal{A}')$ be a pair of periodic monopoles on $C \times_q S^1$ with images (P, ϕ) and (P', ϕ') respectively, and choose a bundle isomorphism $\tau: P \rightarrow P'$ intertwining the q -connections ϕ and ϕ' . One observes first that \mathbf{P} and \mathbf{P}' are also isomorphic G -bundles since, by intertwining with the q -connections, we have an isomorphism $\mathbf{P}|_{C \times \{t\}} \rightarrow \mathbf{P}'|_{C \times \{t\}}$ for every $t \in S^1$. That the monopole structures also match up follows by the same argument as in [CH10, Proposition 4.7].

For surjectivity, again we'll match the argument in the case where $q = \text{id}$. We begin by extending a holomorphic G -bundle P on C_0 with q -connection ϕ to a G -bundle on $M \setminus (D \times \{t_0\}) = (C \times_q S^1_R) \setminus (D \times \{t_0\})$. Let $\gamma: [-2\pi R, 2\pi R] \rightarrow G$ (Vasily: Sorry, what is the target space of map γ ? Perhaps not G but the group of automorphisms of the curve C ? Also, now we are assuming that there exists group homomorphism from real numbers to the automorphism group of C . It is true for the rational, and trigonometric case. However, what if automorphism group of the curve C includes some discrete factor, and q is that discrete automorphism. In this case we will not be able to smoothly interpolate. Perhaps we need to restrict explicitly to a subgroup of automorphism group of C which is covered by the group of real numbers or circle?)

be a geodesic with $\gamma(-2\pi R) = q^{-1}$, $\gamma(0) = 1$ and $\gamma(2\pi R) = q$. Let \widetilde{M} be the 3-manifold

$$\widetilde{M} = ((-2\pi R, 2\pi R) \times C) \setminus \bigcup_{j=1}^k (A_j^+ \cup A_j^-)$$

where A_j^+ is the arc $\{(t + t_0, \gamma(t)(z_j)) : t \in (0, 2\pi R - t_0)\}$ and A_j^- is the arc $\{(t + t_0 - 4\pi R, \gamma(t)(z_j)) : t \in [2\pi R - t_0, 2\pi R)\}$.

Let $\pi: \widetilde{M} \rightarrow C$ be the projection sending (t, z) to $\gamma(t)(z)$. The bundle P pulls back to a bundle $\pi^*(P)$ on \widetilde{M} . We obtain a bundle on $M \setminus (D \times t_0)$ by applying the identification $(t, z) \sim (t - 2\pi R, q(z))$. This bundle extends to an S^1 -invariant holomorphic G -bundle on $M \times S^1$. The remainder of the proof – verifying the existence of the monopole structure associated to an appropriate choice of hermitian structure – consists of local analysis which is independent of the value of the parameter q .

It remains to remark on the compatibility of framing data on the two sides. A trivialization of the bundle \mathbf{P} along the circle $\{c\} \times S^1$ yields a trivialization of the fiber of the bundle P at c . The condition that the holonomy around the circle at c is 1 fixes the value of the q -connection at c to be 1.

(Vasily: Again, it seems that it is natural to fix arbitrary framing called f in 2.8, not necessarily 1?) □

Remark 3.7. Mochizuki [Moc17] proved a stronger result in the rational case for the group $G = \text{GL}_n$. He allows not just a framing at infinity in \mathbb{CP}^1 but also a singularity encoded in terms of a B -reduction of the bundle.

3.1 Deformation Theory

(Vasily: Do we assume in this section the case of rational C ? I think that some of our statements hold for generic C and some for rational case)

In the next section we'll compare symplectic forms on these moduli spaces. In order to do so it will be important to understand the tangent spaces at a point of the source and target. There's a natural description of these tangent spaces in terms of the hypercohomology of certain cochain complexes.

Let's first describe the tangent complex to the moduli space of monopoles. For simplicity we'll restrict attention to the rational case where we can use the description as a hyperkähler reduction. For more general deformation theory calculations we refer to Foscolo [Fos16]. Recall that in this case we can write

$$\text{Mon}_G^{(\text{fr})}(C \times_q S^1, D \times \{t_0\}, \omega^\vee) \cong \mu^{-1}(0)/\mathcal{G}$$

where \mathcal{G} is the group of gauge transformations of \mathbf{P} and $\mu: \mathcal{V} \rightarrow \Omega^1(M; (\mathfrak{g}_R)_{\mathbf{P}})$ is the Bogomolny functional $(A, \Phi) \mapsto *F_A - d_A \Phi$. The tangent complex to this hyperkähler quotient at a point $(\mathbf{P}, \mathcal{A})$ can be written as $\Omega^0(M \setminus D; (\mathfrak{g}_R)_{\mathbf{P}})[1] \rightarrow \mathbb{T}_{\mu^{-1}(0)}$ where $\mathbb{T}_{\mu^{-1}(0)}$ is the tangent complex to the zero locus of the moment map, concentrated in non-negative degrees. Roughly speaking $\mathbb{T}_{\mu^{-1}(0)} = \mathbb{T}_{\mathcal{V}} \xrightarrow{d\mu} \Omega^1(M; (\mathfrak{g}_R)_{\mathbf{P}})[-1]$.

More explicitly, following [Fos16] let

$$\mathcal{F}_{\mathbf{P}, \mathcal{A}}^{\text{mon}} = \left(\Omega^0(M \setminus D; (\mathfrak{g}_R)_{\mathbf{P}}) \xrightarrow{d_1} \Omega^1(M \setminus D; (\mathfrak{g}_R)_{\mathbf{P}}) \oplus \Omega^0(M \setminus D; (\mathfrak{g}_R)_{\mathbf{P}}) \xrightarrow{d_2} \Omega^1(M \setminus D; (\mathfrak{g}_R)_{\mathbf{P}}) \right) \otimes_{\mathbb{R}} \mathbb{C}$$

placed in degrees $-1, 0$ and 1 where $d_1(g) = -(d_A(g), [\Phi, g])$ and $d_2(a, \psi) = *d_A(a) - d_A(\psi) + [\Phi, a]$. Write d_{mon} for the total differential. (Vasily: Somewhere here we should clarify that the hyperKähler moment map takes value in the space imaginary quaternions which is isomorphic as real space to $\mathbb{R}^3_{\text{quat}}$. After we pick a complex structure on the twistor sphere, there is natural split $\mathbb{R}^3_{\text{quat}} = \mathbb{C} \oplus \mathbb{R}$. If we care only about complex structure and holomorphic symplectic form on the tangent space to the moduli space of monopoles with respect to a certain fixed point on a twistor sphere, we trade the real component of the equation in the split $\mathbb{R}^3_{\text{quat}} = \mathbb{C} \oplus \mathbb{R}$ in the third term with a complexification of the gauge group transformation in the first term in which we would replace \mathfrak{g}_R by \mathfrak{g}_C)

Remark 3.8. If we restrict $\mathcal{F}_{\mathbf{P}, \mathcal{A}}^{\text{mon}}$ to a slice $C_t = C \times \{t\}$ in the t -direction we can identify it with a complex of the form

$$\Omega^\bullet(C_t; \mathfrak{g}_P)[1] \xrightarrow{\mathcal{A}} i\Omega^\bullet(C_t; \mathfrak{g}_P)dt$$

with differential d_A on each of the two factors along with the differential \mathcal{A} (Vasily: sorry, what is differential \mathcal{A} ?) mixing the two factors (Chris: careful here, need to justify how we get this). These two summands each split up into the sum of a Dolbeault complex on C with its dual. That is there's a natural subcomplex of the form

$$\Omega^{0,\bullet}(C_t; \mathfrak{g}_P)[1] \xrightarrow{\mathcal{A}} i\Omega^{0,\bullet}(C_t; \mathfrak{g}_P)dt$$

where the internal differentials on the two factors are now given by $\bar{\partial}_{A_0}$. (Chris: so why does it make sense to restrict to this subcomplex?) (Vasily: Sorry, I am afraid I did not understand the notations of the remark and hence the statement. Anyways, I was expecting, in the line of our discussions, that the relevant differential, from which we would find the tangent bundle to Mon and then determine symplectic form, is given by

$$D = dtD_t + d\bar{z}D_{\bar{z}}, \quad D_t = \partial_t + A_t + i\Phi, \quad D_{\bar{z}} = \partial_{\bar{z}} + A_{\bar{z}} \quad (3)$$

that is a mixture of Dol differential along C and de Rham differential along S^1 . And we should have been referring here to the choice of the Riemmanian metric on M (1) to make the claims and the twistor sphere point corresponding to the split of $T_M \simeq \mathbb{R}^3_{\text{quat}} \simeq \mathbb{C} \oplus \mathbb{R} = T_C \oplus T_{S^1}$)

Remark 3.9. If one introduces a framing at a point $c \in C$ then we must correspondingly twist the complex \mathcal{F}^{mon} above by the line bundle $\mathcal{O}(-c)$ on C . So in that case we define (Chris: check)

$$\mathcal{F}_{(\mathbf{P}, \mathcal{A})}^{\text{mon, fr}} = \mathcal{F}_{\mathbf{P}, \mathcal{A}}^{\text{mon}} \otimes (\mathbb{C}_{S^1} \boxtimes \mathcal{O}(-c)).$$

The following is proved in [Fos16].

Proposition 3.10. The tangent complex of $\text{Mon}_G(S^1 \times C, D \times \{t_0\}, \omega^\vee)$ at the point $(\mathbf{P}, \mathcal{A})$ is quasi-isomorphic to the hypercohomology $\mathbb{H}^\bullet(C \times S^1; \mathcal{F}'_{(\mathbf{P}, \mathcal{A})})$ of a subsheaf $\mathcal{F}' \subseteq \mathcal{F}^{\text{mon}}$ where growth conditions are imposed on the degree 0 part of \mathcal{F}^{mon} near the singularities.

In order to compare this tangent complex with the tangent complex for q -connections we'll take advantage of the following Poincaré/Dolbeault lemma.

Lemma 3.11. The space of sections of the sheaf $\mathcal{F}'_{(\mathbf{P}, \mathcal{A})}$ on a contractible set $U_i \subseteq M$ containing exactly one $z_i \in D$ has vanishing cohomology in degree 0.

Proof. The cochain complex $\mathcal{F}'_{(\mathbf{P}, \mathcal{A})}(U_i)$ is the tangent complex to the (derived) moduli space of monopoles on U_i with Dirac singularity at z_i with prescribed charge $\omega_{z_i}^\vee$. There's a unique such monopole by analytic continuation, which means $H^0(\mathcal{F}'_{(\mathbf{P}, \mathcal{A})}(U_i)) = 0$. \square

(Vasily: I am not sure if I agree with the statement of the lemma (or may be I misunderstand), and the proof, namely the wording "there is a unique such monopole by analytic continuation". I am confused by comparison with the fact that the moduli space of monopoles on a contractable space \mathbb{R}^3 with a prescribed singularity ω^\vee at a single point (z, t) is a non-trivial space of dimension $2(\rho, \omega^\vee)$ studied by BFN and Kamnitzer et al in the papers about slices in affine Grassmanian. If we replace \mathbb{R}^3 by a ball U with a reasonable boundary conditions like Dirichlet or Neumann, we would still get a non-trivial moduli space of monopoles of dimension $2(\rho, \omega^\vee)$)

Now let's consider the tangent complex to the moduli space of q -connections. For the arguments in this article we'll only need to carefully consider the case $q = \text{id}$ of multiplicative Higgs bundles, but we'll include some remarks regarding the more general case. In this case the calculation was performed by Bottacin [Bot95], see also [HM02, Section 4]. Fix a multiplicative Higgs bundle (P, ϕ) on C . We consider the sheaf of cochain complexes on C

$$\mathcal{F}_{(P, \phi)} = (\mathfrak{g}_P[1] \xrightarrow{\text{Ad}_\phi} \mathfrak{g}_P(D))$$

in degrees -1 and 0 with differential given by the adjoint action of ϕ . More precisely let L_ϕ and R_ϕ be the bundle maps $\mathfrak{g}_P \rightarrow \mathfrak{g}_P$ obtained as the derivative of left- and right-multiplication. Then $\text{Ad}_\phi = L_\phi - R_\phi$. We can alternatively phrase this, as in [HM02, Section 4], as follows. Define $\text{ad}(\phi)$ to be the vector bundle

$$\text{ad}(\phi) = (\mathfrak{g}_P \oplus \mathfrak{g}_P) / \{(X, -\phi X \phi^{-1}) : X \in \mathfrak{g}_P\}.$$

Then we can write \mathcal{F} as the sheaf of complexes

$$\mathcal{F}_{(P, \phi)} = (\mathfrak{g}_P[1] \xrightarrow{\text{Ad}_\phi} \text{ad}(\phi))$$

where now Ad_ϕ is just the map $X \mapsto [(X, -X)]$.

Remark 3.12. If one introduces a framing at a point $c \in C$ then we must correspondingly twist the complex \mathcal{F} above by the line bundle $\mathcal{O}(-c)$ on C (Vasily: perhaps we could add a comment like that or similar to explain: "...since we fix the values of all fields at the framing point c , their deformations should vanish at c , and this is accounted by the tensoring the deformation complex with $\mathcal{O}(-c)$ ")

So in that case we define

$$\mathcal{F}_{(P, \phi)}^{\text{fr}} = (\mathfrak{g}_P[1] \xrightarrow{\text{Ad}_\phi} \mathfrak{g}_P(D)) \otimes \mathcal{O}(-c).$$

Remark 3.13. For more general q we should modify this description by replacing ϕ by a q -connection. Note that one can still define the (q -twisted) adjoint action $X \mapsto \phi X \phi^{-1}$ using a q -connection, and so we can still define the complex

$$\mathcal{F}_{(P, \phi)} = (\mathfrak{g}_P[1] \xrightarrow{\text{Ad}_\phi} \text{ad}(\phi))$$

just as in the untwisted case (Chris: check). (Vasily: I wonder what would replace the rational r -matrix that defines the Poisson structure on the Lie group of rational Higgs fields given the explicit formula for the symplectic structure on the moduli space of q -connections (6) for non-trivial q ? And if we quantize the space of functions on q -deformed big mHiggs what do we get instead of completed Yangian when q is trivial, say for the case of a single singularity at 0)

Proposition 3.14. The tangent complex of $\text{mHiggs}_G(C, D)$ at the point (P, ϕ) is quasi-isomorphic to the hypercohomology $\mathbb{H}^\bullet(C; \mathcal{F}_{(P, \phi)})$ of the sheaf \mathcal{F} .

Proof. We'll explain the proof in different language to that used in [Bot95]. Recall that for a mapping stack $\text{Map}(X, Y)$ where X and Y are derived Artin stacks the tangent complex at a point f can be computed as the cohomology $\mathbb{T}_{\text{Map}(X, Y), f} \cong \mathbb{H}^*(X, f^* \mathbb{T}_Y)$. Let $X = C \setminus D$ and let $Y = G/G$, so $\mathbb{T}_{(P, \phi), \text{Map}(C \setminus D, G/G)} \cong \mathbb{H}^*(C \setminus D; \mathfrak{g}_P[1]|_{C \setminus D} \rightarrow \mathfrak{g}_P|_{C \setminus D})$ where the map $\mathfrak{g}_P|_{C \setminus D} \rightarrow \mathfrak{g}_P|_{C \setminus D}$ is given by the adjoint action of ϕ . In order

to compute the tangent complex to $\mathrm{mHiggs}_G(C, D)$ we must impose the conditions that the deformation of the G -bundle P extends across the divisor D and that the deformation of the Higgs field has a simple pole at D . In other words we find

$$\mathbb{T}_{\mathrm{mHiggs}_G(C, D), (P, \phi)} = \mathbb{H}^\bullet(C; \mathfrak{g}_P[1] \rightarrow \mathfrak{g}_P(D)) \subseteq \mathbb{H}^*(C \setminus D; \mathfrak{g}_P[1]|_{C \setminus D} \rightarrow \mathfrak{g}_P|_{C \setminus D}).$$

(Chris: This last step isn't rigorous, needs improving.) □

Remark 3.15. The embedding of the moduli space $\mathrm{mHiggs}_G(C, D, \omega^\vee)$ into $\mathrm{mHiggs}_G(C, D)$ induces an isomorphism of tangent complexes (in degree 0). In other words, all infinitesimal deformations of $(P, \phi) \in \mathrm{mHiggs}_G(C, D, \omega^\vee)$ preserve the residue condition at D . (Chris: check and justify. To insert here: the fact that the tangent space is concentrated in degree 0.)

In order to calculate hypercohomology we'll use a Čech resolution. We fix the following open cover of the manifold $M = C \times_q S^1$. Let \mathbb{D}_i be a contractible analytic open neighbourhood in C of each $z_i \in D$. Let $U_i = \mathbb{D}_i \times (0, 2\pi)$ be the corresponding contractible open subset of M . Likewise let \mathbb{D}_∞ be a contractible analytic open neighbourhood of $c \in C$ and let $U_\infty = \mathbb{D}_\infty \times (0, 2\pi)$ be a contractible open subset of M . Choose the \mathbb{D}_i so that the U_i and U_∞ are pairwise disjoint. Finally let $U_0 = M \setminus (D \times \{0\} \cup \{(c, 0)\})$.

Let's explain why these covers suffice to compute hypercohomology. On the multiplicative Higgs side we need only note that the \mathbb{D}_i are contractible and the remaining subset $C \setminus (D \cup \{c\})$ is an affine algebraic curve, which means that for any quasi-coherent sheaf of cochain complexes the higher cohomology groups vanish. Likewise for the intersections – the punctured discs \mathbb{D}_i^\times – which are analytic open sets of an affine curve.

On the monopole side we need to verify that our sheaf $\mathcal{F}'_{(P, \mathcal{A})}$ has no higher cohomology after restricting to open sets U_i, U_0 and their intersections. This is clear for the contractible U_i . For U_0 and the intersections the restriction of the sheaf $\mathcal{F}'_{(P, \mathcal{A})}$ is coherent after restriction to $(C \setminus (D \cup \{c\})) \times \{t_0\}$ and locally constant in the t direction. The sheaf $\mathcal{F}'_{(P, \mathcal{A})}|_{U_0}$ has no higher cohomology after restricting to $(C \setminus (D \cup \{c\})) \times \{t_0\}$, in other words any higher global sections must be zero on this slice. But then local constancy implies that they're zero on all of U_0 . (Chris: this local constancy business seems fishy to me. For coherence I want to restrict to the $(0, 1)$ -part after restricting to a slice as in Remark 3.8, then use GAGA.)

On the multiplicative Higgs side we specify a representative 0-cocycle $(\alpha_\infty, \{\alpha_i\}, \alpha_0, \beta_\infty, \{\beta_i\})$ for the Čech cohomology group with respect to our chosen cover. Explicitly a 0-cochain is given by the following data:

$$\begin{aligned} \alpha_\infty &\in \mathrm{ad}(g)(-1)(\mathbb{D}_\infty) \\ \alpha_i &\in \mathrm{ad}(g)(1)(\mathbb{D}_i) \text{ for } i = 1, \dots, k \\ \alpha_0 &\in \mathrm{ad}(g)(C \setminus (D \cup \{\infty\})) \\ \beta_\infty &\in \mathfrak{g}_P(\mathbb{D}_\infty^\times) \\ \beta_i &\in \mathfrak{g}_P(\mathbb{D}_i^\times) \text{ for } i = 1, \dots, k. \end{aligned}$$

(Vasily: Please, could you clarify the notations (1) and (-1) in the above lines?)

(Vasily: Did we switch notation here from ϕ to g compared to the previous page?) Being a 0-cocycle means that $(\alpha_\infty - \alpha_0)|_{U_\infty^\times} = \mathrm{Ad}_g(\beta_\infty)$ and $(\alpha_i - \alpha_0)|_{U_i^\times} = \mathrm{Ad}_g(\beta_i)$ for each i . We consider 0-cocycles modulo 0-coboundaries of the form $(\mathrm{Ad}_g(f_\infty), \{\mathrm{Ad}_g(f_i)\}, \mathrm{Ad}_g(f_0), f_\infty - f_0, \{f_i - f_0\})$.

Now Ad_g is an isomorphism on \mathbb{D}_0 for the sections $\tilde{\alpha}$ of \mathfrak{g}_P that occur: those with no poles or zeroes in \mathbb{D}_0 . That means that we can add a coboundary to force $\tilde{\alpha} = 0$. We can likewise use the freedom in f_i to find a representative cocycle where the β_i are all zero. We now have a unique representative cocycle of the form $(\alpha_\infty, \{\alpha_i\}, 0, 0, 0)$.

Likewise on the monopole side we can describe a 0-cocycle with similar data. Explicitly a 0-cochain in our Čech

complex is given by $(\alpha_\infty, \{\alpha_i\}, \alpha_0, \beta_\infty, \{\beta_i\})$ where now

$$\begin{aligned}\alpha_\infty &\in \mathcal{F}_0^{\text{mon}}(U_\infty) \\ \alpha_i &\in \mathcal{F}_0^{\text{mon}}(U_i) \text{ for } i = 1, \dots, k \\ \alpha_0 &\in \mathcal{F}_0^{\text{mon}}(U_0) \\ \beta_\infty &\in \Omega_{\mathbb{C}}^0(U_\infty^\times; \mathfrak{g}_{P_{\mathbb{C}}}) \\ \beta_i &\in \Omega_{\mathbb{C}}^0(U_i^\times; \mathfrak{g}_{P_{\mathbb{C}}}) \text{ for } i = 1, \dots, k.\end{aligned}$$

Here we write $\mathcal{F}_0^{\text{mon}}$ to indicate the degree 0 term in the cochain complex. There's a similar condition for being a 0-cocycle involving the differential d_{mon} . Again we can add a 0-coboundary to find a representative cocycle of the form $(\alpha_\infty, \{\alpha_i\}, 0, 0, 0)$.

It will also be important to have an explicit description of the derivative of the holonomy map H as a map between tangent spaces. We can describe this map using our Čech resolutions on each contractible open set U_i individually.

(Vasily: Does the proof above applies to the case of generic q or for the trivial q ?)

Proposition 3.16. The derivative $dH: \mathbb{H}^\bullet(U_i; \mathcal{F}'_{(P, \mathcal{A})}) \rightarrow \mathbb{H}^\bullet(\mathbb{D}_i; \mathcal{F}_{(P, \phi)})$ is given on an open patch U_i by the formula

$$dH(\alpha_i) = dH(d_{\text{mon}} b_i) = b_i(2\pi)H(\mathcal{A}) - H(\mathcal{A})b_i(0)$$

where $i = 1, \dots, k$ or ∞ . More precisely by $b_i(2\pi)$ and $b_i(0)$ we mean the limit of $b_i(t)$ as $t \rightarrow 2\pi$ or 0 respectively.

Proof. Note that the right-hand side is the derivative at \mathcal{A} of the map $B_i \mapsto B_i(2\pi)H(\mathcal{A})B_i(0)^{-1}$ where $B_i \in \Omega^0((0, 2\pi) \times \mathbb{D}_i) \setminus \{(t_0, z_i)\}; \mathfrak{g}_P$. This is the definition of the action of the group of gauge transformations on the holonomy $H(\mathcal{A})$ from $t = 0$ to 2π . (Chris: say more?) \square

4 Symplectic Structures

From now on we'll exclusively study the rational situation where $C = \mathbb{CP}^1$ and we fix a framing point $c = \infty$.

4.1 The Example of GL_2

(Chris: Ultimately we might move this, but we should probably include a discussion of what happens for GL_2 before or in parallel to discussing the general story.)

(Vasily: Chris, is the idea to put here explicit expressions for Higgs field in terms of some chart Darboux coordinates (p, q) that we have discussed in the past drafts draft?. If all singularities are miniscule, then for GL_2 the leaves could be described quite explicitly as products of regular semi-simple co-adjoint orbits in \mathfrak{gl}_2 . Rouven is finishing a draft with explicit parametrization by Darboux coordinates of a certain family of symplectic leaves for GL_n . Perhaps here we can put a reference on it.)

4.2 Symplectic Structures for General G

We begin the more anstract general analysis by briefly discussing the holomorphic symplectic structure on the moduli space of periodic monopoles on \mathbb{CP}^1 following the analysis of Cherkis and Kapustin [CK98, CK02]. This structure arises from the description we gave as a hyperkähler quotient. To describe it specifically, let $\delta^{(1)}\mathcal{A}$ and $\delta^{(2)}\mathcal{A}$ be two tangent vectors at (P, \mathcal{A}) to the moduli space of monopoles. Choose representatives for these two tangent vectors of the form α_i and α'_i respectively in the Čech resolution we described above. Then we can write

the holomorphic symplectic form coming from the hyperkähler reduction in terms of the symplectic pairing on the infinite-dimensional vector space \mathcal{V} , which is given by the Killing form on \mathfrak{g} along with the wedge pairing of differential forms. So summing over the local patches in our Čech resolution we can write it as.

$$\omega(\delta^{(1)}\mathcal{A}, \delta^{(2)}\mathcal{A}) = \sum_{i=1}^k \int_{U_i} \kappa(\alpha_i \wedge \alpha'_i) dz.$$

(Chris: todo: fix this expression)

(Vasily: we should start from the expression like

$$\omega(\delta^{(1)}\mathcal{A}, \delta^{(2)}\mathcal{A}) = \int_M \kappa(\delta^{(1)}\mathcal{A} \wedge \delta^{(2)}\mathcal{A}) \wedge v \quad (4)$$

where $v = "dz"$ is the CY 1-form and where we should have had defined before (did we ?)

$$\mathcal{A} = A_{\bar{z}} d\bar{z} + A_z dz + (A_t + i\Phi) dt \quad (5)$$

)

Let's now consider the moduli space of framed q -connections on \mathbb{CP}^1 . This moduli space has an algebraic symplectic structure with a natural abstract origin, via *derived Poisson geometry*. We refer to [CPT⁺17] for the theory of derived Poisson structures and to [MS16, MS17, Spa16] for that of derived coisotropic structures.

Lemma 4.1. The restriction map $r: \text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D) \rightarrow BLG^k$ has a 2-shifted Lagrangian structure with respect to the residue pairing on BLG .

Remark 4.2. More precisely, BLG is a Tate stack in the sense of Hennion [Hen15]. Although we won't check that the residue pairing symplectic structure ω is compatible with the Tate structure, this problem doesn't arise when defining a Lagrangian structure as a potential for $r^*\omega$ in $\Omega^{2, \text{cl}}(\text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D), 2)$ since the stack in question is of ind-finite type.

Proof. We can describe a potential for the pullback $r^*\omega$ as follows. The tangent complex to $\text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D)$ at P is quasi-isomorphic to $\Omega^{0, \bullet}(C \setminus D; \mathfrak{g}_P)[1]$. There is a pairing of degree 1 on this complex defined by

$$\nu((f, \alpha), (g, \beta)) = \int_{C \setminus D} \kappa(f, \alpha) + \kappa(\beta, g) dz$$

where (f, α) and (g, β) are general non-homogenous elements of $\Omega^{0, \bullet}(C \setminus D; \mathfrak{g}_P)[1]$. Applying the internal differential on $\bigwedge^2(\mathbb{T}_{\text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D), P})$ to this pairing induces the degree 2 pairing

$$\begin{aligned} d\nu(f, g) &= \int_{C \setminus D} \kappa(f, \bar{\partial}g) + \kappa(g, \bar{\partial}f) dz \\ &= \int_{C \setminus D} \bar{\partial}\kappa(f, g) dz \\ &= 2\pi i \sum_{i=1}^k \oint_{S_i^1} \kappa(f, g) dz \end{aligned}$$

which is the pullback $r^*\omega$ of the residue pairing on $(BLG)^k$. (Chris: Rewrite more rigorously.) \square

Construction 4.3. Recall that we can identify the moduli space of singular q -connections on a curve as a fiber product: $q\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D) \cong \text{Bun}_G^{\text{fr}}(C) \times_{\text{Bun}_G(\mathbb{CP}^1 \setminus D)^2} \text{Bun}_G(\mathbb{CP}^1 \setminus D)$ where the map $g: \text{Bun}_G^{\text{fr}}(\mathbb{CP}^1) \rightarrow \text{Bun}_G(\mathbb{CP}^1 \setminus D)$ is given by $P \mapsto (P|_{\mathbb{CP}^1 \setminus D}, q^*P|_{\mathbb{CP}^1 \setminus D})$. Consider the following commutative cube:

$$\begin{array}{ccccc}
& q\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D) & \xrightarrow{f_1} & \text{Bun}_G^{\text{fr}}(\mathbb{CP}^1) & \\
& \swarrow f_2 & \downarrow & \swarrow g & \downarrow \\
\text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D) & \xrightarrow{\quad} & \text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D)^2 & & \\
& \downarrow \text{res} & \downarrow r & & \\
& \text{Bun}_G(\mathbb{B})^k & \xrightarrow{\quad} & BL^+G^{2k} & \\
& \swarrow & \downarrow & \swarrow & \\
BLG^k & \xrightarrow{\quad} & BLG^{2k} & &
\end{array}$$

Here the top and bottom faces are homotopy Cartesian. The front and the right faces are not homotopy Cartesian but there are canonical 1-shifted Lagrangian structures on the induced maps from the source of the face into the homotopy fiber product, i.e.

$$\begin{aligned}
& \text{Bun}_G^{\text{fr}}(\mathbb{CP}^1) \rightarrow \text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D)^2 \times_{BLG^{2k}} BL^+G^{2k} \cong \text{Bun}_G^{\text{fr}}(\mathbb{CP}^1)^2 \\
& \text{and } \text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D) \rightarrow \text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D)^2 \times_{BLG^{2k}} BLG^k.
\end{aligned}$$

Call these Lagrangian structures μ_1 and μ_2 respectively. Denote the 2-shifted Lagrangian structure on the map r from Lemma 4.1 by ν . The 1-shifted Lagrangian structure that we want on the map res is then given by

$$\lambda = f_1^*(\mu_1) + f_2^*(\mu_2) + \eta$$

where η is a potential for the pullback $(f_1 \circ g)^*\nu$.

This 1-shifted Lagrangian structure λ defines a quasi-isomorphism $\mathbb{T}_{q\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D)} \rightarrow \mathbb{L}_{\text{res}}$ where I'm writing \mathbb{L}_{res} for the dual of $\text{res}^*(\mathbb{T}_{q\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D)}) \rightarrow \mathbb{T}_{\text{Bun}_G(\mathbb{B})^k}[-1]$: the relative tangent complex to the map res . Write λ^{-1} for a quasi-inverse. This defines a Poisson bivector on $q\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D)$ where the associated map from the cotangent complex to the tangent complex is given by

$$\pi = \lambda^{-1} \circ \iota: \mathbb{L}_{q\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D)} \rightarrow \mathbb{L}_{\text{res}} \rightarrow \mathbb{T}_{q\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D)}$$

where ι is the canonical inclusion of the cotangent complex into the relative cotangent complex.

We obtain the 0-shifted symplectic structure on a Poisson leaf as the Lagrangian intersection

$$\begin{array}{ccc}
q\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D, \omega^\vee) & \longrightarrow & q\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D) \\
\downarrow & & \downarrow \\
\coprod B\Gamma_i & \longrightarrow & \text{Bun}_G(\mathbb{B})^k
\end{array}$$

where $B\Gamma_i$ is the classifying space of the group of symmetries of the G -bundle on \mathbb{B} associated to the coweight ω_i^\vee . If we restrict the Poisson bivector to the leaf $q\text{-Conn}_G^{\text{fr}}(\mathbb{CP}^1, D, \omega^\vee)$ it becomes non-degenerate and so defines a 0-shifted symplectic structure.

In order to work with this symplectic structure we'll need to give a local description as a pairing on the tangent complex. Using the description of the tangent complex in terms of hypercohomology from Section 3.1 it will suffice to describe the pairing on each summand in the Čech resolution.

Lemma 4.4. As a pairing on the tangent complex at a point (P, ϕ) the symplectic form on the moduli space of q -connections pairs α_i with β_i for $i = 1, \dots, k$. It is given on this coordinate patch by the formula

$$\omega(\alpha_i, \beta'_i) = \langle \rho_\phi^*(\alpha_i + \tilde{\alpha})|_{U_i^\times}, \rho_\phi^*(\text{Ad}_\phi^*)^{-1}(\beta'_i) \rangle$$

where we use the Killing form to identify β'_i with a \mathfrak{g}^* -valued form, and where ρ_ϕ^* is the pullback along the right multiplication by ϕ .

Proof. The formula given above is saying that the Poisson bivector – viewed as a map $\pi: \mathbb{L}_{(P,\phi)} \rightarrow \mathbb{T}_{(P,\phi)}$ – acts in degree -1 by identifying \mathfrak{g} with \mathfrak{g}^* then applying the map Ad_ϕ^* . Recall that to describe the Poisson bivector we must locally on U_i describe the sum

$$\lambda = f_1^*(\mu_1) + f_2^*(\mu_2) + \eta$$

where η is a potential for the pullback $(f_1 \circ g)^*\nu$. The Lagrangian structure μ_1 on the diagonal is 0, and the Lagrangian structure μ_2 vanishes upon restriction to the open set U_i since the map $\text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D) \rightarrow \text{Bun}_G^{\text{fr}}(\mathbb{CP}^1 \setminus D)^2 \times_{BLG^{2k}} BLG^k$ becomes an isomorphism after restricting to a formal neighbourhood of z_i ([Chris: justify the difference between a formal neighbourhood and a contractible analytic neighbourhood here](#)). We must therefore find a local potential for $(f_1 \circ g)^*\nu$ where ν is the structure from Lemma 4.1. On the cotangent complex, the pullback g^* acts by the dual to the q -twisted diagonal map $\mathfrak{g}[[z]] \rightarrow \mathfrak{g}((z))^2$. The pullback $(f_1 \circ g)^*$ acts by the projection map $\mathfrak{g}^*((z))^2 \rightarrow \text{ad}(\phi)^*|_{\mathbb{D}^\times}$ onto formal local section of the bundle $\text{ad}(\phi)^*$.

We obtain a potential for the image of the residue pairing ν under this pullback as follows. ([Chris: todo: complete. We end up with the residue pairing on \$\text{ad}\(\phi\)^*\$, and we'll get a potential by applying \$\text{Ad}_\phi^*\$ to one of the two factors. Can write it down as in the lemma above.](#)) \square

Proposition 4.5. The symplectic form on the moduli space of q -connections can be written in the form

$$\omega(\delta\phi, \delta\phi') = - \sum_{i=1}^d \langle b_i^L \phi^{-1}, b_i'^R \phi^{-1} \rangle - (b \leftrightarrow b') \quad (6)$$

([Vasily: I think some repairs are required here if we are to compare to the expression \(8\) below: 1\) we should expect sandwiches between \$g\$ and \$g^{-1}\$ and not \$\phi^{-1}\$ and \$\phi^{-1}\$ 2\) the summation should also include a contribution from \$\infty\$ 3\) Pick either \$\langle \rangle\$ or \$\kappa\$ notation for the Killing form](#)) where $\delta\phi$ is represented on U_i by a pair $(b^L, b^R) \in \text{Ad}_\phi$. Here and throughout the notation “ $-(b \leftrightarrow b')$ ” denotes antisymmetrization.

Proof. Specify a representative cocycle $(\alpha_\infty, \{\alpha_i\}, \alpha_0, \beta_\infty, \{\beta_i\})$ for the first Čech cohomology group. By addition of a coboundary we can assume that $\alpha_0 = 0$, and to force α_∞ to land in $z\mathfrak{g}[[z]]$ and each α_i to land in $z^{-1}\mathfrak{g}[[z]]$. We've now fixed a representative cocycle – there's no further gauge freedom. Ignoring the factor at infinity – since after we add these coboundaries the residue pairing there vanishes – the pairing we end up with looks like

$$\begin{aligned} \omega(\delta\phi, \delta\phi') &= \sum_{i=1}^d \langle \rho_\phi^*(\alpha_i + \tilde{\alpha})|_{U_i^\times}, \rho_\phi^*(\text{Ad}_\phi^*)^{-1}(\beta'_i) \rangle - ((\alpha, \beta) \leftrightarrow (\alpha', \beta')) \\ &= \sum_{i=1}^d \langle \rho_\phi^*(\alpha_i + \tilde{\alpha})|_{U_i^\times}, \rho_\phi^*(\text{Ad}_\phi^*)^{-1}(\text{Ad}_\phi^{-1}(\alpha'_i - \tilde{\alpha}')) \rangle - (\alpha \leftrightarrow \alpha') \\ &= \sum_{i=1}^d \langle \rho_\phi^*(\alpha_i)|_{U_i^\times}, \rho_\phi^*(\text{Ad}_\phi^*)^{-1}(\text{Ad}_\phi^{-1}(\alpha'_i)|_{U_i^\times}) \rangle - (\alpha \leftrightarrow \alpha'). \end{aligned}$$

We can compute the composite operator $(\text{Ad}_\phi^*)^{-1}\text{Ad}_\phi^{-1}$ on $\text{ad}(P)(U_i^\times)$. It's given by the inverse of the operator $\text{Ad}_\phi\text{Ad}_\phi^*$ which sends a pair $(b^L, b^R) \in \text{ad}(P)(U_i^\times)$ to $(b^L - b^R, b^R - b^L) \sim (b^L + \phi^{-1}b^L\phi, b^R - \phi b^R\phi^{-1})$. Denote this by $((1 + A_\phi)b^L, (1 - A_{\phi^{-1}})b^R)$. We can describe the inverse using the expansion $(1 + A_\phi)^{-1} = 1 - A_\phi + A_\phi^2 + \dots$. After applying our gauge transformation above the remaining degree of freedom in α_i is its z^{-1} term. Each time we apply A it raises the order in z of this term by one, so only the linear summand A_ϕ of $-(1 + A_\phi)^{-1}$ contributes to the residue pairing. In the pairing we need to use the invariant pairing on $\mathfrak{g}((z)) \oplus \mathfrak{g}((z))$ that vanishes on the subalgebra spanned by $(X, A_\phi(X))$, so we take the difference of the residue / Killing form pairings on the two summands.

$$\begin{aligned} \omega(\delta\phi, \delta\phi') &= \sum_{i=1}^d -\langle \rho_\phi^*A_\phi b_i^L, \rho_\phi^*b_i'^R \rangle - \langle \rho_\phi^*b_i^L, \rho_\phi^*A_{\phi^{-1}}b_i'^R \rangle - (b \leftrightarrow b') \\ &= - \sum_{i=1}^d \langle b_i^L \phi^{-1}, b_i'^R \phi^{-1} \rangle - (b \leftrightarrow b') \end{aligned}$$

using cyclic invariant to identify the two terms. \square

We can now establish our main result.

Theorem 4.6. The symplectic structure on $\text{Mon}_G^{\text{fr}}(\mathbb{CP}^1 \times S^1, D \times \{0\}, \omega^\vee)$ and the pullback of the symplectic structure on $\text{mHiggs}_G^{\text{ps,fr}}(\mathbb{CP}^1, D, \omega^\vee)$ under H coincide.

Proof. This is straightforward now, by combining the two local descriptions of the symplectic structure on monopoles and multiplicative Higgs bundles on a neighbourhood of a puncture with the description in Proposition 3.16. So, let's begin by taking the symplectic form ω_{mHiggs} on the moduli space $\text{mHiggs}_G^{\text{ps,fr}}(\mathbb{CP}^1, D, \omega^\vee)$ and evaluating it at the image under dH of two elements $(\alpha_i, \alpha'_i) = (d_{\text{mon}}(b_i), d_{\text{mon}}(b'_i))$. According to Proposition 3.16 we have

$$\begin{aligned} \omega_{\text{mHiggs}}(dH(\alpha_i), dH(\alpha'_i)) &= \omega_{\text{mHiggs}}(dH(d_{\text{mon}}(b_i)), dH(d_{\text{mon}}(b'_i))) \\ &= \omega_{\text{mHiggs}}(b_i(2\pi)H(\mathcal{A}) - H(\mathcal{A})b_i(0), b'_i(2\pi)H(\mathcal{A}) - H(\mathcal{A})b'_i(0)). \end{aligned}$$

We can write this in terms of a pairing on the bundle $\text{ad}(\phi)$. Represent the class $b_i(2\pi)H(\mathcal{A}) - H(\mathcal{A})b_i(0)$ by the pair $(b_i(2\pi)H(\mathcal{A}), -H(\mathcal{A})b_i(0))$. Then applying the description of ω_{mHiggs} provided by Proposition 4.5 we have

$$\begin{aligned} \omega_{\text{mHiggs}}(dH(\alpha_i), dH(\alpha'_i)) &= \omega_{\text{mHiggs}}((b_i(2\pi)H(\mathcal{A}), -H(\mathcal{A})b_i(0)), (b'_i(2\pi)H(\mathcal{A}), -H(\mathcal{A})b'_i(0))) \\ &= \langle b_i(2\pi), H(\mathcal{A})b'_i(0)H(\mathcal{A})^{-1} \rangle - (b \leftrightarrow b') \end{aligned}$$

where the remaining terms are killed by the antisymmetrization.

On the other hand, we can evaluate the pairing $\omega_{\text{Mon}}(\alpha_i, \alpha'_i) = \omega_{\text{Mon}}(d_{\text{mon}}(b_i), d_{\text{mon}}(b'_i))$ via integration by parts. We'll compute the symplectic pairing for the full Čech complex, but it will split into a sum over open sets U_i . The result is that (Vasily: the expression (4) expands as) (Vasily: 1. let us unify the notations for the Higgs field in the equations above and the previous sections: somewhere holonomy or the Higgs field is called $H(\mathcal{A})$, somewhere ϕ , somewhere g . I would favour the notation g . 2. Let us introduce somewhere before the notation $b^L = b(0), b^R = b(2\pi)$ to make the equations more readable, and break the multiline computation into separate lines and insert between the lines the explanations of what we are doing)

$$\begin{aligned} \omega_{\text{Mon}}(\{\alpha_i\}, \{\alpha'_i\}) &= \sum_{i=0, \dots, k, \infty} \int_{U_i} \kappa(d_{\text{mon}}(b_i) \wedge d_{\text{mon}}(b'_i)) dz - (b \leftrightarrow b') \\ &= \sum_{i=1, \dots, k, \infty} \int_{\partial \mathbb{D}_i \times [0, 2\pi]} \kappa(b_i - b_0, d_{\text{mon}}(b'_i)) dz - (b \leftrightarrow b') \\ &= \sum_{i=1, \dots, k, \infty} \oint_{\partial \mathbb{D}_i} \kappa(b_i(2\pi) - b_0(2\pi), b'_i(2\pi)) - \kappa(b_i(0) - b_0(0), b'_i(0)) - (b \leftrightarrow b') \\ &= \sum_{i=1, \dots, k, \infty} \oint_{\partial \mathbb{D}_i} -\kappa(b_0(2\pi), b'_i(2\pi)) + \kappa(b_0(0), b'_i(0)) - (b \leftrightarrow b'). \end{aligned}$$

On the second line we used the fact that $d_{\text{mon}}b_i = d_{\text{mon}}b_0$ on $U_i \cap U_0$ and that $d_{\text{mon}}b_i(t) = 0$ when $t = 0$ or 2π . Pick out the summand corresponding to U_i . Choose our potentials so that on the boundary $\partial \mathbb{D}_i$ we have $b_0(0) = 0$ and $b_i(2\pi) = 0$ (Chris: possible why?), (Vasily: Instead, I would suggest that here we choose the potentials such that $b_0(0) = 0$ and $b'_0(0) = 0$ everywhere on U_0 , and in particular on the boundary $\partial \mathbb{D}_i$. There is no obstruction to this choice because for any variation of the holonomy $\delta g = b^R g - g b^L$ with the choice $b^L = 0$ we set $b^R = \delta g g^{-1}$ which is a holomorphic section on U_0 for the reason that U_0 does not contain any singularities and because of the Bogomolny equation. Given this choice the last line becomes

$$\sum_{i=1, \dots, k, \infty} \oint_{\partial \mathbb{D}_i} -\kappa(b_0^R, b_i^{R'}) - (b \leftrightarrow b') \quad (7)$$

Next, from $\delta g = b_i^R g - g b_i^L = b_0^R g$ on $\partial \mathbb{D}_i = U_0 \cap U_i$ we find $b_0^R = b_i^R - g b_i^L g^{-1}$ and plug in the above line. Because of the antisymmetrization $b \leftrightarrow b'$ the term containing $b_i^R b_i^{R'}$ vanishes and we are left with

$$\sum_{i=1, \dots, k, \infty} \oint_{\partial \mathbb{D}_i} \kappa(g b_i^L g^{-1}, b_i^{R'}) - (b \leftrightarrow b') \quad (8)$$

)

which means that on $U_0 \cap U_i$ we can identify $b_0(2\pi) = H(\mathcal{A}) b_i(0) H(\mathcal{A})^{-1}$. Then

$$\omega_{\text{Mon}}(\{\alpha_i\}, \{\alpha'_i\}) = \oint_{\partial \mathbb{D}_i} \kappa(b_i(2\pi), H(\mathcal{A}) b'_i(0) H(\mathcal{A})^{-1}) - (b \leftrightarrow b')$$

agreeing with the expression coming from mHiggs. (Vasily: I think the minus sign is missed if compared to (8)) \square

Remark 4.7. (Chris: connection to Shapiro's symplectic leaves in [Sha16]. The claim is that taking the Taylor expansion at infinity defines a Poisson map to $\tilde{G} \subseteq G_1[[z^{-1}]]$ sending symplectic leaves to (discrete) unions of symplectic leaves. One can do the classification more completely in the type A case.)

5 Hyperkähler Structures

(Vasily: As of July 19, 2018 I have not yet passed this and following sections)

The results of the previous section imply that the symplectic structure on $\text{mHiggs}_G^{\text{ps}, \text{fr}}(\mathbb{CP}^1, D, \omega^\vee)$ extends canonically to a hyperkähler structure. In this section we'll compare the twistor rotation with the deformation to the moduli space of q -connections. (Chris: todo: expand on this)

Proposition 5.1. (Chris: Monopoles in the large R (Gaiotto) limit – twistor rotation and ε -deformation coincide.)

Proof.

\square

Corollary 5.2. If we equip $\text{mHiggs}_G^{\text{ps}, \text{fr}}(\mathbb{CP}^1, D, \omega^\vee)$ with the complex structure at the twistor parameter $\eta \in \mathbb{C}$, the resulting complex manifold is isomorphic to $\varepsilon\text{-Conn}_G^{\text{ps}, \text{fr}}(\mathbb{CP}^1, D, \omega^\vee)$.

(Chris: comment on the holomorphic symplectic structure vs the one discussed above.)

(Chris: comment on the example of the Nahm transform.)

6 Quantization and Yangians

(Chris: We're still discussing how this story goes.)

7 Duality

(Chris: For a first draft I'm inserting some notes I prepared for a talk. Todo: talk about motivation and the content we discussed regarding non-simply-laced groups)

Pseudo-Conjecture 7.1 (Multiplicative Geometric Langlands). Let G be a Langlands self-dual group. There is an equivalence of categories

$$\text{A-Branes}_{q^{-1}}(\text{mHiggs}_G(C, D, \omega^\vee)) \cong \text{B-Branes}(q\text{-Conn}_G(C, D, \omega^\vee))$$

where the category on the right-hand side depends on the value q .

What does this mean, and are there situations in which we can make it precise? We'll discuss three examples where we can say something more concrete. In each case, by "B-branes" we'll just mean the category $\text{Coh}(q\text{-Conn}_G(C, D, \omega^\vee))$ of coherent sheaves. By "A-branes" we'll mean some version of q^\vee -difference connections on the stack $\text{Bun}_G(C)$.

Remark 7.2. This equivalence is supposed to interchange objects corresponding to branes of opers on the two sides, and introduce an analogue of the Feigin-Frenkel isomorphism between deformed W-algebras (see [FRSTS98, STSS98]). This isomorphism only holds for self-dual groups, which motivates the restriction to the self-dual case here.

7.1 The Abelian Case

Suppose $G = \text{GL}(1)$ (more generally we could consider a higher rank abelian gauge group). In general for an abelian group the moduli spaces we've defined are trivial – for instance the rational and trigonometric spaces are always discrete. However there is one interesting non-trivial example: the elliptic case. For simplicity let's consider the abelian situation with $D = \emptyset$: the case with no punctures.

Definition 7.3. A q -difference module on a variety X with automorphism q is a module for the sheaf $\Delta_{q,X}$ of non-commutative rings generated by \mathcal{O}_X and an invertible generator Φ with the relation $\Phi \cdot f = q^*(f) \cdot \Phi$. Write $\text{Diff}_q(X)$ for the category of q -difference modules on X .

In the abelian case the space $q\text{-Conn}_{\text{GL}(1)}(E)$ is actually a stack, but one can split off the stacky part to define difference modules on it. Indeed, for any q one can write

$$\text{Bun}_{\text{GL}(1)}(E) \cong \text{BGL}(1) \times \mathbb{Z} \times E^\vee$$

and so

$$q\text{-Conn}_{\text{GL}(1)}(E) \cong \text{BGL}(1) \times \mathbb{Z} \times (E^\vee \times_q \mathbb{C}^\times)$$

which means one can define difference modules on these stacks associated to an automorphism of E^\vee or $E^\vee \times_q \mathbb{C}^\times$ respectively.

Conjecture 7.4. There is an equivalence of categories for any $q \in \mathbb{CP}^1$

$$\text{Diff}_q(\text{Bun}_{\text{GL}(1)}(E)) \cong \text{Coh}(q^{-1}\text{-Conn}_{\text{GL}(1)}(E)).$$

In this abelian case we can go even farther and make a more sensitive 2-parameter version of the conjecture.

Conjecture 7.5. There is an equivalence of categories for any $q_1, q_2 \in \mathbb{CP}^1$

$$\text{Diff}_{q_1}(q_2\text{-Conn}_{\text{GL}(1)}(E)) \cong \text{Diff}_{q_2^{-1}}(q_1^{-1}\text{-Conn}_{\text{GL}(1)}(E))$$

where q_1 is the automorphism of $E^\vee \times_{q_2} \mathbb{C}^\times$ acting fiberwise over each point of \mathbb{C}^\times .

This conjecture should be provable using the same techniques as the ordinary geometric Langlands correspondence in the abelian case, i.e. by a (quantum) twisted Fourier-Mukai transform (as constructed by Polishchuk and Rothstein [PR01]).

7.2 The Classical Case

Now, let's consider the limit $q \rightarrow 0$. This will give a conjectural statement involving coherent sheaves on both sides analogous to the classical limit of the geometric Langlands conjecture as conjectured by Donagi and Pantev [DP10]. The existence of an equivalence isn't so interesting in the self-dual case (where both sides are the same), but the classical multiplicative Langlands functor should be an *interesting* non-trivial equivalence. For example we can make the following conjecture

Conjecture 7.6. Let G be a Langlands self-dual group and let E be an elliptic curve. There is an automorphism of categories (for the rational, trigonometric and elliptic moduli spaces)

$$F: \text{Coh}(\text{mHiggs}_G(E)) \cong \text{Coh}(\text{mHiggs}_G(E))$$

so that the following square commutes:

$$\begin{array}{ccc} \text{Coh}(\text{mHiggs}_T(E)) & \xrightarrow{\text{FM}} & \text{Coh}(\text{mHiggs}_T(E)) \\ p_* q^! \downarrow & & \downarrow p_* q^! \\ \text{Coh}(\text{mHiggs}_G(E)) & \xrightarrow{F} & \text{Coh}(\text{mHiggs}_G(E)). \end{array}$$

Here we're using the natural morphisms $p: \text{mHiggs}_B(E) \rightarrow \text{mHiggs}_G(E)$ and $q: \text{mHiggs}_B(E) \rightarrow \text{mHiggs}_T(E)$, and FM is the Fourier-Mukai transform.

It ought to be possible to state something a bit stronger that includes singularities in this auto-duality.

7.3 The Rational Type A Case

There's one more example where we can say something precise, and even draw a connection to the ordinary geometric Langlands correspondence. We already mentioned the Nahm transform in the previous section: in the case where $G = \text{GL}(n)$ and C is \mathbb{C} (where as usual we fix framing data at infinity) the Nahm transform identifies multiplicative Higgs bundles of degree k with *ordinary* Higgs bundles on \mathbb{CP}^1 for the group $\text{GL}(k)$ with $n + 2$ tame singularities (with appropriate fixed locations and residues).

Claim. Under the Nahm transform, Pseudo-Conjecture 7.1 in the rational case for the group $\text{GL}(n)$ becomes the ordinary geometric Langlands conjecture on \mathbb{CP}^1 with tame ramification.

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