

# Two-dimensional simulation of viscoplastic Drucker–Prager free surface flows, application to granular collapse

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## Drucker–Prager model

- 2-dimensional **viscoplastic** flow with free surface.
- Time dependant domain  $\Omega_t$ .
- Incompressible non-Newtonian dynamics:

$$\begin{aligned} \rho(\partial_t \mathbf{u} + (\mathbf{u} \cdot \nabla_{\vec{X}}) \mathbf{u}) - \operatorname{div}_{\vec{X}} \boldsymbol{\sigma} &= \rho \mathbf{f} & \text{in } (0, T) \times \Omega_t, \\ \operatorname{div}_{\vec{X}} \mathbf{u} &= 0 & \text{in } (0, T) \times \Omega_t, \end{aligned}$$

where  $\mathbf{u}$  is the material velocity,  $\mathbf{f}$  an external force (gravity),  $\rho > 0$  the mass density, and  $\boldsymbol{\sigma}$  the total stress tensor defined by the Drucker-Prager rheology  $\boldsymbol{\sigma} = \boldsymbol{\sigma}' - p \mathbf{Id}$ , with the viscoplastic constitutive equation:

$$\operatorname{trace}(\boldsymbol{\sigma}') = 0, \quad \begin{cases} \boldsymbol{\sigma}' = 2\eta D\mathbf{u} + \kappa \frac{D\mathbf{u}}{\|D\mathbf{u}\|} & \text{if } D\mathbf{u} \neq 0, \\ \|\boldsymbol{\sigma}'\| \leq \kappa & \text{if } D\mathbf{u} = 0. \end{cases}$$

- The plasticity  $\kappa \equiv \kappa(p)$  stands for the yield limit with a **Drucker–Prager** criterion:

$$\kappa(p) = \sqrt{2}\mu[p]_+,$$

with  $\mu$  an internal friction coefficient.

- The force  $\mathbf{f}$  is the gravity force acting on the material. Taking into account the slope  $\theta$ , it is given by  $\mathbf{f} = (g \sin \theta, -g \cos \theta)$  in the coordinates  $(X, Z)$ , with  $g$  the gravity constant.

## Domain geometry and boundary conditions

**(a) periodic flow over an inclined bed:**

- no-slip condition at the bottom  $\Gamma_b$ .
- periodicity condition on the lateral side  $\Gamma_{l,t} \cup \Gamma_{r,t}$ .

**(b) collapse of a granular mass over a rigid or erodible bed:**

- no-penetration condition on  $\Gamma_{b,t} \cup \Gamma_{l,t}$ .
- Coulomb friction condition on  $\Gamma_{b,t} \cup \Gamma_{l,t}$ .

**(a)-(b)**

- no-stress condition at the free surface  $\Gamma_{f,t}$ .
- kinematic condition at the free surface:  $N_t + \vec{N} \cdot \mathbf{u}(t, \vec{X}) = 0$  on  $\Gamma_{f,t}$ .
- initial condition:  $\mathbf{u}(0, \vec{X}) = \mathbf{u}_0(\vec{X})$ .

## Regularization and variational formulation

**Regularized constitutive equation:**

$$\boldsymbol{\sigma}'_{\epsilon} = 2\eta D\mathbf{u} + \kappa(p) \frac{D\mathbf{u}}{\sqrt{\|D\mathbf{u}\|^2 + \epsilon^2}}, \quad 0 < \epsilon \ll 1.$$

**Associated variational formulation:**

$$\begin{aligned} & \int_{\Omega_t} \rho(\partial_t \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u}) \cdot \mathbf{v} + \int_{\Omega_t} 2\eta D\mathbf{u} : D\mathbf{v} + \int_{\Omega_t} \kappa(p) \frac{D\mathbf{u}}{\sqrt{\|D\mathbf{u}\|^2 + \epsilon^2}} : D\mathbf{v} - \int_{\Omega_t} p \operatorname{div} \mathbf{v} \\ & + \int_{\Gamma_{b,t}} \mu_b \frac{\mathbf{u}_T \cdot \mathbf{v}}{\sqrt{|\mathbf{u}_T|^2 + \epsilon_f^2}} [p - \sigma'_N]_+ + \int_{\Gamma_{l,t}} \mu_l \frac{\mathbf{u}_T \cdot \mathbf{v}}{\sqrt{|\mathbf{u}_T|^2 + \epsilon_f^2}} [p - \sigma'_N]_+ = \int_{\Omega_t} \rho \mathbf{f} \cdot \mathbf{v} + \int_{\Gamma_{f,t}} \gamma \mathcal{C} \mathbf{v} \cdot \vec{N}, \\ & \int_{\Omega_t} q \operatorname{div} \mathbf{u} = 0. \end{aligned}$$

**(a):**  $\mu_b = \mu_l = 0$ .

**Displacement of the domain:** ( $\mathbf{w}$  = mesh velocity)

$$-\operatorname{div}(D\mathbf{w}) = 0 \quad \text{in } \Omega_t,$$

boundary conditions in case **(a)**

$$\begin{aligned} (\mathbf{w} - \mathbf{u}) \cdot \vec{N} &= 0 & \text{on } \Gamma_b \cup \Gamma_{f,t}, \\ \mathbf{w} \cdot \vec{N} &= 0 & \text{on } \Gamma_{l,t} \cup \Gamma_{r,t}, \\ (D\mathbf{w} \vec{N})_T &= 0 & \text{on } \Gamma. \end{aligned}$$

boundary conditions in case **(b)**

$$\begin{aligned} (\mathbf{w} - \mathbf{u}) \cdot \vec{N} &= 0 & \text{on } \Gamma, \\ (D\mathbf{w} \vec{N})_T &= 0 & \text{on } \Gamma. \end{aligned}$$

## Discrete algorithm

**Time discretization:** semi-implicit finite difference scheme of first order.

**Space discretization:**  $\mathbb{P}_2/\mathbb{P}_1$  finite element (Taylor-Hood).

**Mesh mapping:**

$$\begin{aligned} \mathcal{A}_{n,n+1} : \Omega_h^n &\rightarrow \Omega_h^{n+1} \\ x &\mapsto y = x + \Delta t_n \mathbf{w}_h^n(x) \end{aligned}$$

**Update algorithm:**

- 1) We suppose that  $\Omega_h^n$ ,  $\mathbf{u}_h^n$ ,  $p_h^n$  are known and we compute  $\mathbf{w}_h^n$ .
- 2) We move the nodes of the mesh according to  $\mathcal{A}_{n,n+1}$ , thus we obtain  $\Omega_h^{n+1}$ .
- 3) Finally, we compute  $(\mathbf{u}_h^{n+1}, p_h^{n+1})$  on  $\Omega_h^{n+1}$ .

## References

- [1] M. FARIN, A. MANGENEY, O. ROCHE, *Fundamental changes of granular flow dynamics, deposition, and erosion processes at high slope angles: Insights from laboratory experiments*, J. Geophys. Res. Earth Surf. 119 (2014).
- [2] I.R. IONESCU, A. MANGENEY, F. BOUCHUT, O. ROCHE, *Viscoplastic modelling of granular column collapse with pressure dependent rheology*, hal-01080456, (2014).
- [3] C. LUSSO, F. BOUCHUT, A. ERN, A. MANGENEY, M. FARIN, O. ROCHE, *Two-dimensional simulation of viscoplastic Drucker–Prager free surface flows by regularization, and application to granular collapse* (2015).

## Uniform flow with plug

**Evaluation of the accuracy of the 2-dimensional regularization method (case (a)):**

- ★ We consider a test case with a 1-dimensional solution  $U(t, Z)$  depending only on the normal variable  $Z$ .
- ★ The height of domain  $h$  is constant, the domain does not depend on time and the pressure is hydrostatic.

**Comparison between the 1-dimensional and the 2-dimensional solution:**

- ★ We evaluate the error between the 1-dimensional profile  $U$  and the longitudinal component of the 2-dimensional velocity  $\mathbf{u} = (U, 0)$ , by extending the 1-dimensional solution on the 2-dimensional mesh.
- ★ We compute the relative error on the velocity in the  $L^2$ -norm.

**Numerical results:**

- Convergence of order 1 in space and time.
- $\epsilon = 10^{-2}$ : regularization error dominates.
- $\epsilon = 10^{-6}$ : regularization error dominated by discretization error.

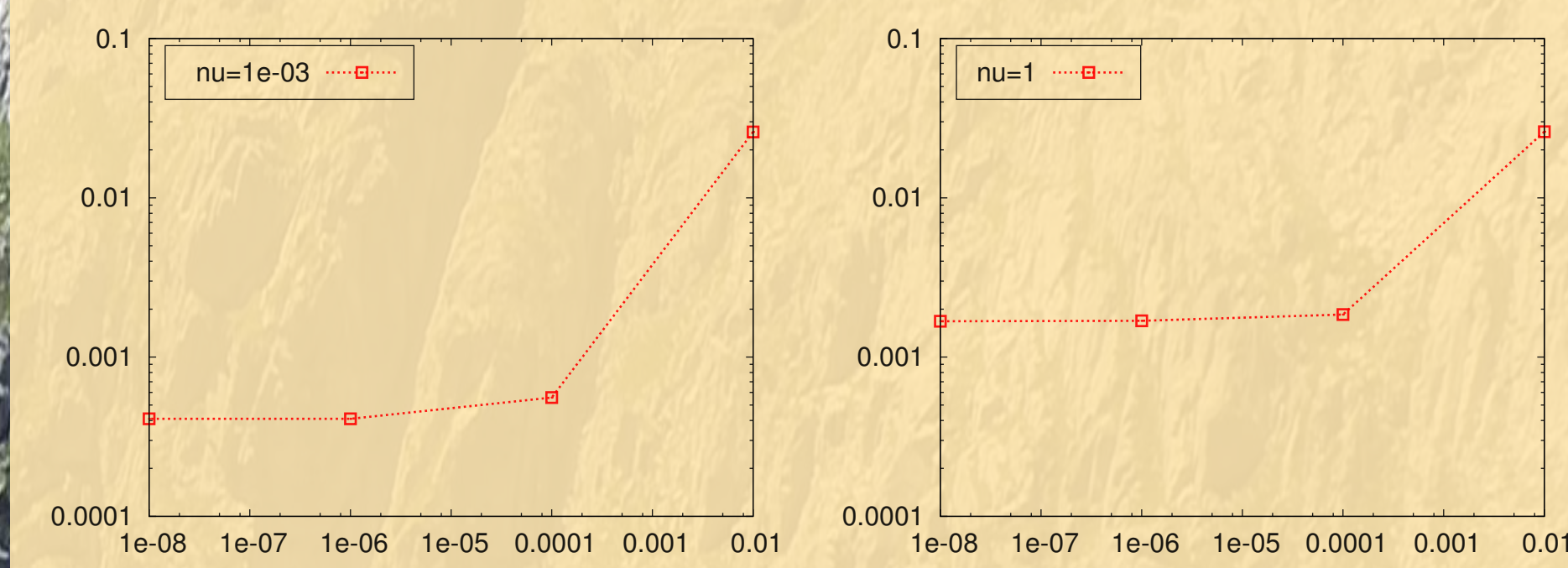


Figure 1: Velocity error with respect to  $\epsilon$ .

$\epsilon$	$10^{-2}$	$10^{-4}$	$10^{-6}$	$10^{-8}$
total iterations	824	2958	5086	5272

$\epsilon = 10^{-8} \Rightarrow$  saturation of iterations cost.

## Granular collapse: comparison with augmented Lagrangian [2]

**Collapse of a trapezoidal mass over a rigid bed (case (b))**

**Static/flowing interface comparison:**

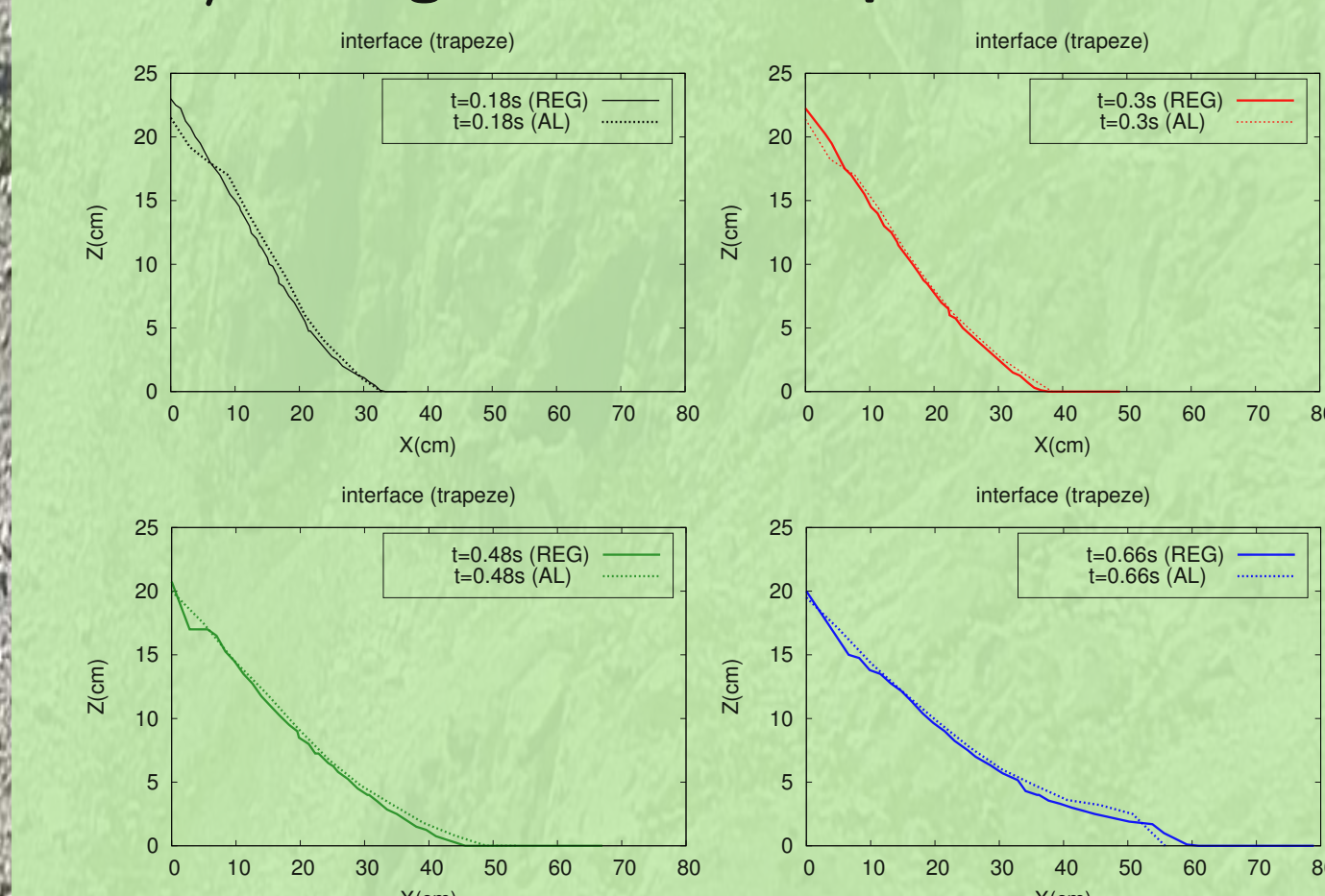


Figure 2: Comparison of the static/flowing interface.

The regularization method leads to:

- very similar profiles (with respect to AL method),
- a slightly lower position of the interface.

**Freesurface comparison:**

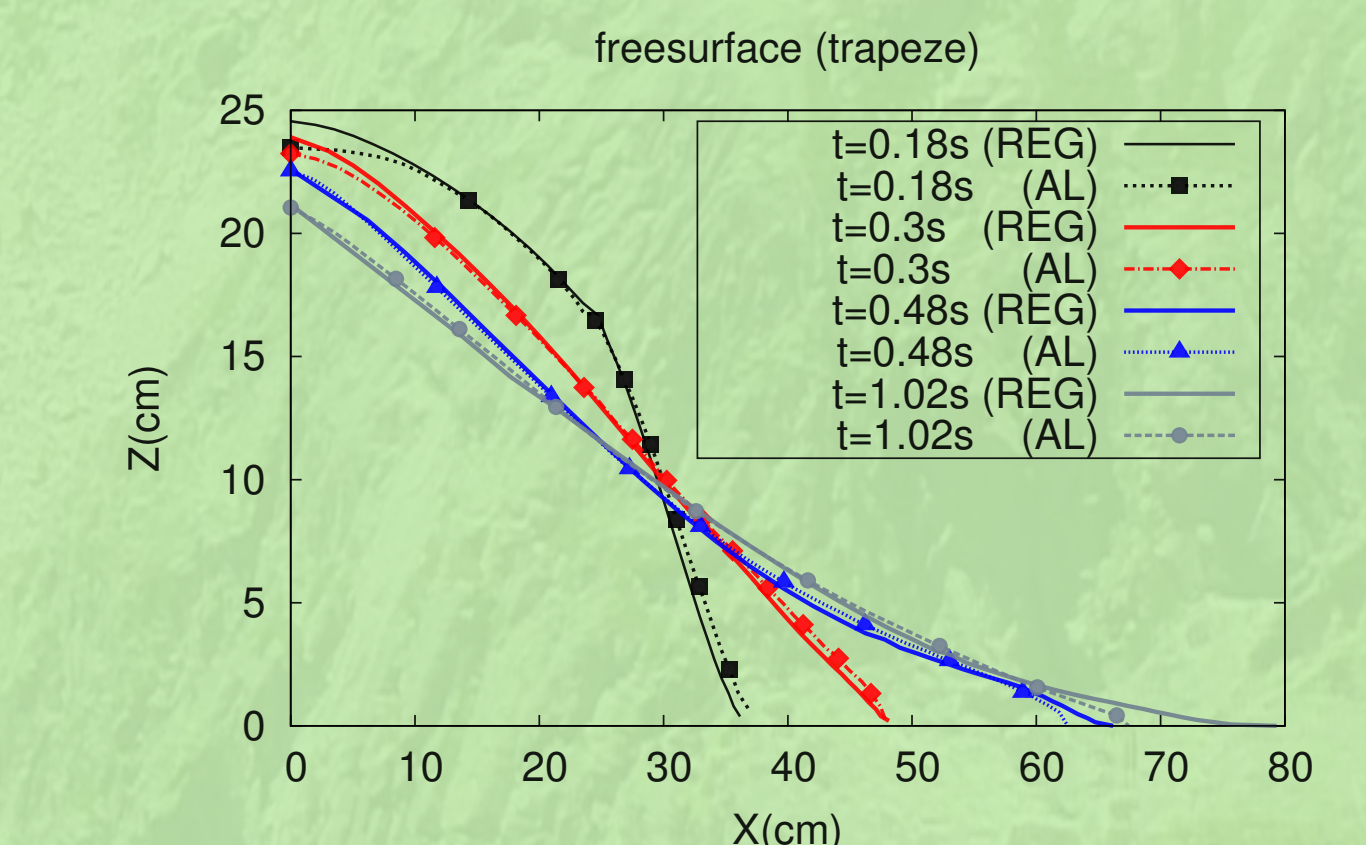


Figure 3: Evolution of the free surface ( $\epsilon = 10^{-6} \text{ s}^{-1}$ ).

The regularization method leads to:

- more decreasing thickness on the left side,
- a slightly faster front propagation (on the right),
- running 7 times faster.

## Granular collapse: comparison with laboratory experiments [1]

**Collapse of a trapezoidal mass over a rigid bed (case (b))**

**Freesurface comparison:**

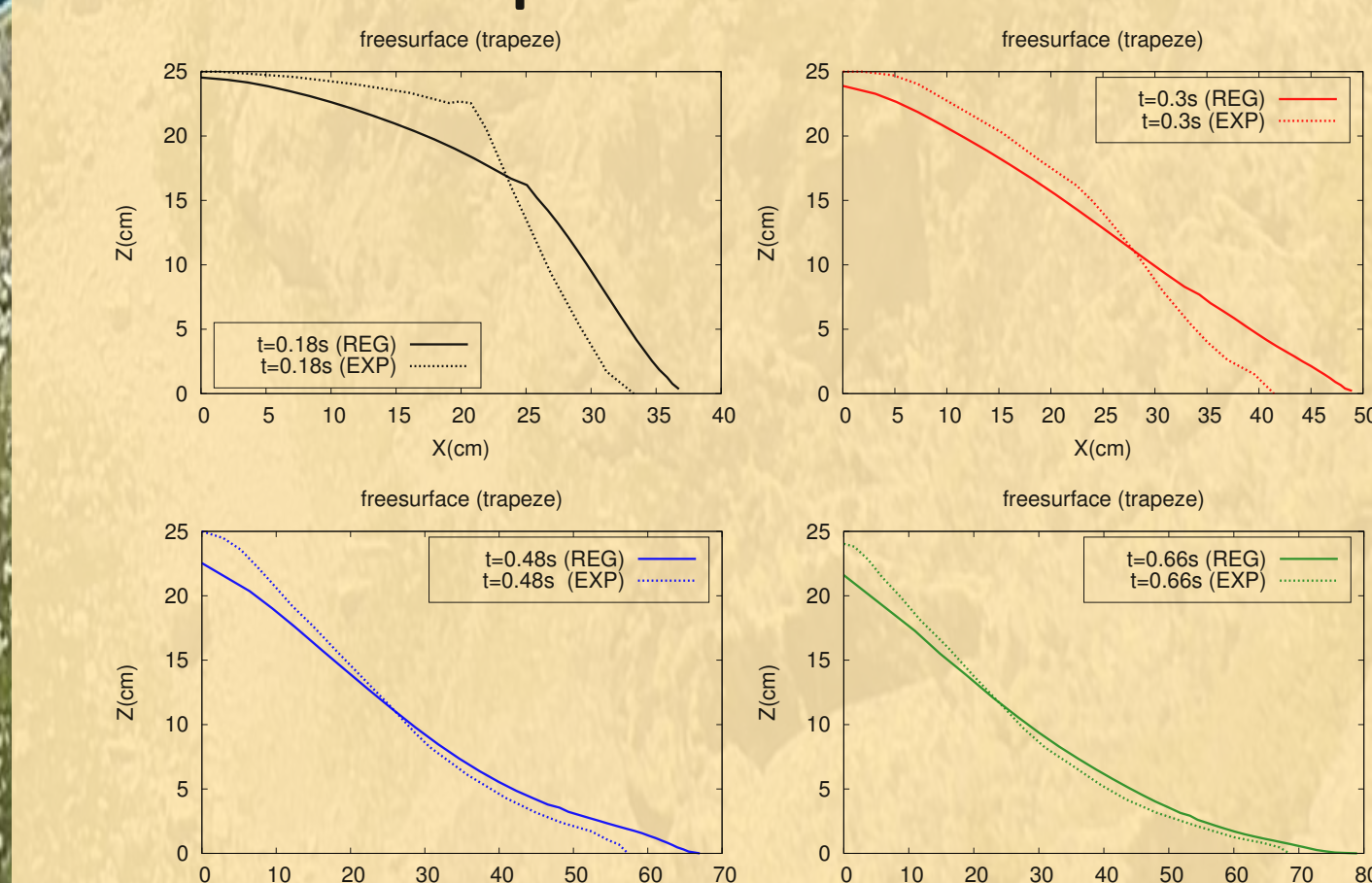


Figure 4: Evolution of the free surface for a trapezoidal mass over a rigid bed.

The regularization method leads to:

- more decreasing thickness on the left side,
- overestimated front position (on the right),
- the shape of the final deposit is very well reproduced.

Numerical simulations  $\Rightarrow$  faster dynamics than in experiments.

**Static/flowing interface comparison:**

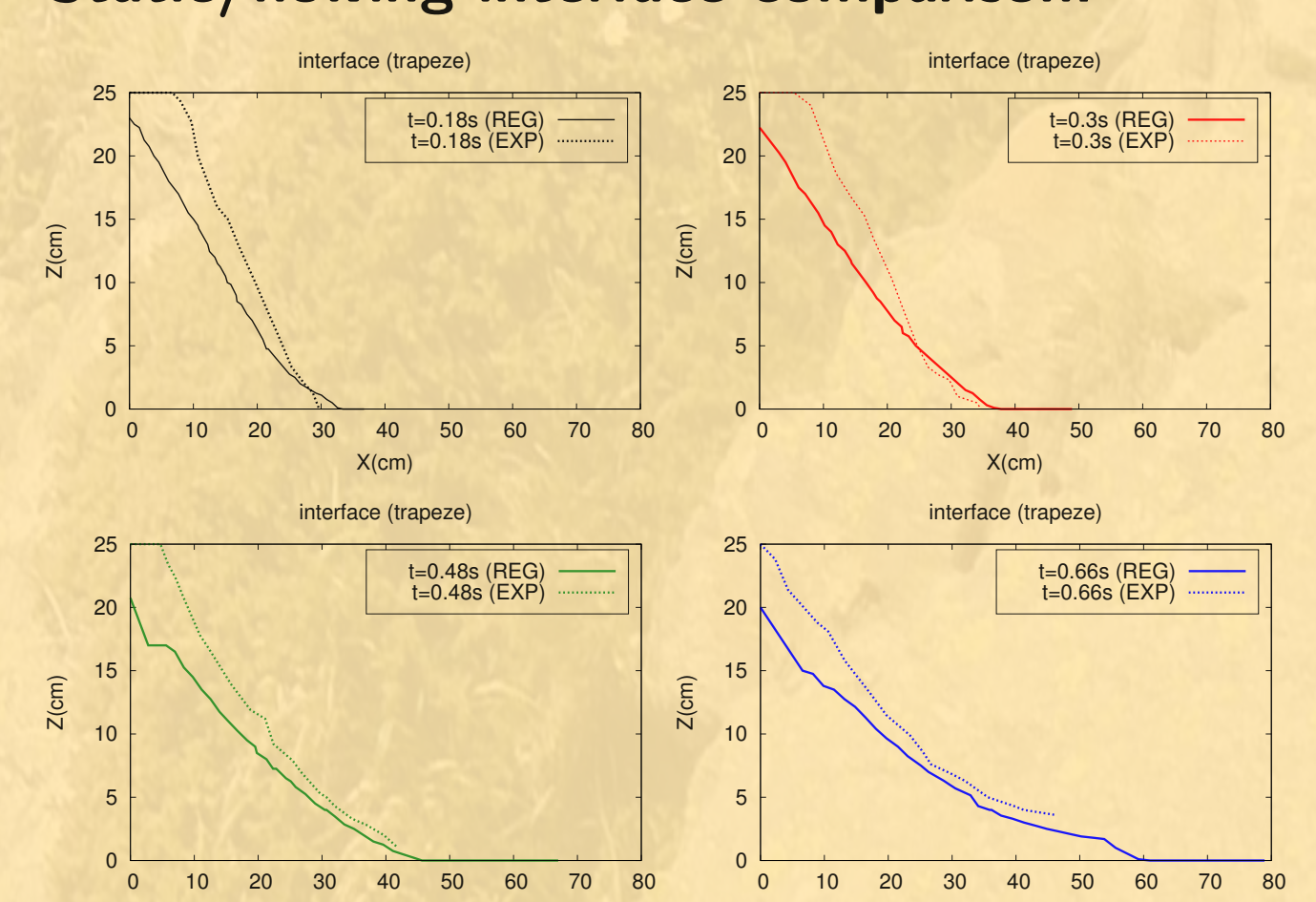


Figure 5: Comparison of the static/flowing interface.

The regularization method leads to:

- pretty well predicted static/flowing interface,
- position underestimated on the top of the left side.

**Collapse of a trapezoidal mass over an erodible bed (case (b),[3])**

**Goal:** to model erosion process in nature.

Simulation of a granular collapse over an erodible bed made of the same material represented by a thin layer of thickness  $h_e = 5 \text{ mm}$  under the trapezoidal column.

**Mean:** local surface tension effects.

The regularization method with surface tension leads to:

- comparable profiles all along the simulation,
- the final shape of the deposit is very well reproduced.

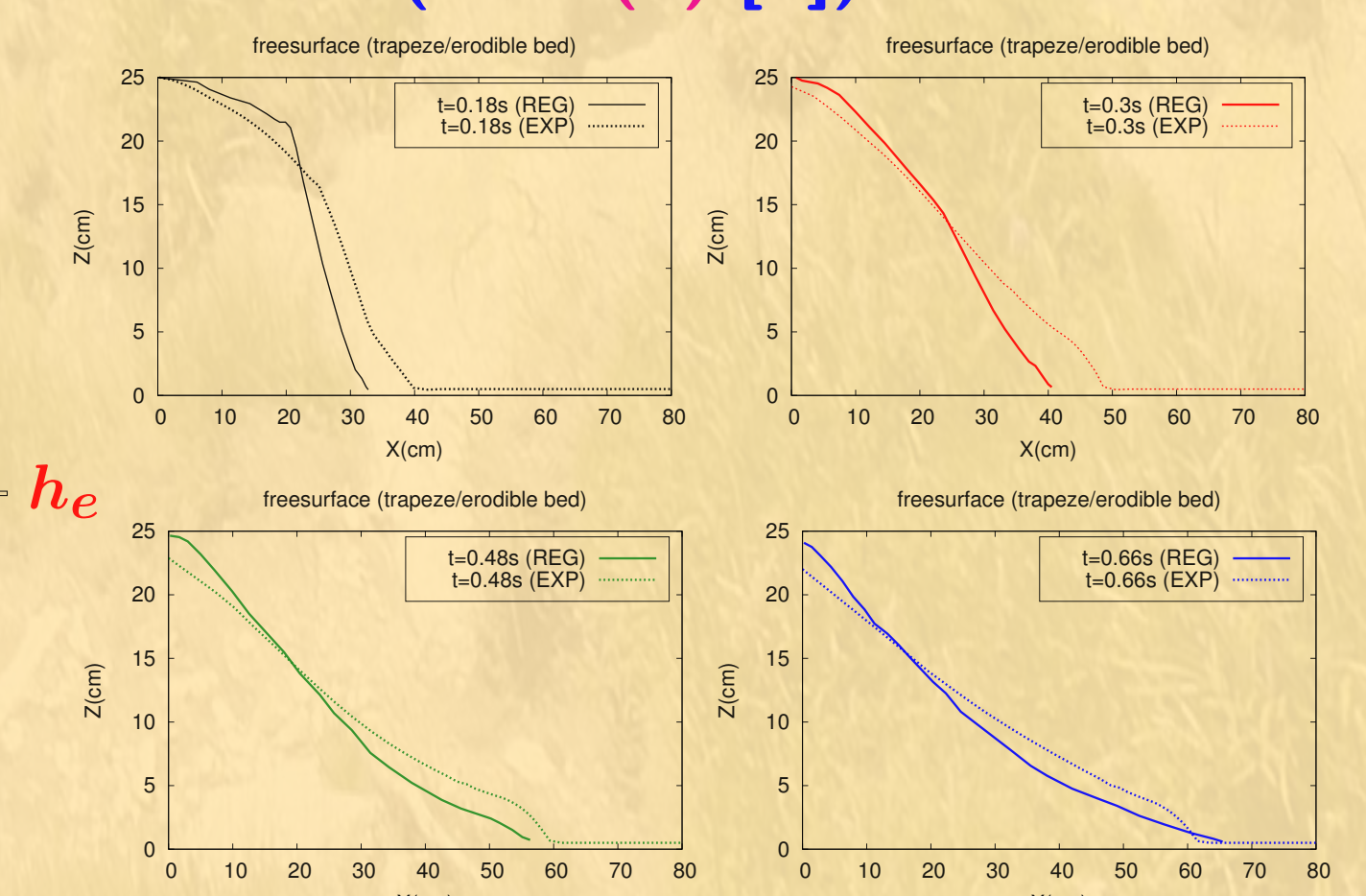


Figure 6: Evolution of the free surface for a trapezoidal mass over an erodible bed of thickness  $h_e = 5 \text{ mm}$ .