

Galilei group
Lorentz group

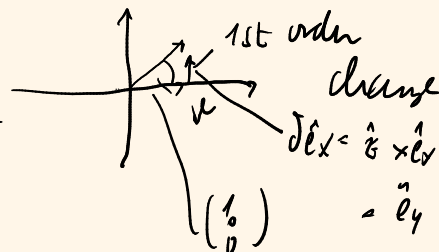
{ rotations
switch between reference frames (inertial observers)

typical parametrisations

rotations $R_n(\theta) \approx 1 + \theta \hat{n} \times \mathbf{L} + \dots = 1 + \theta X_{\hat{n}}$ generates rotations around \hat{n} -axis

2-axis

$$\begin{pmatrix} c & -s & 0 \\ +s & c & 0 \\ 0 & 0 & 1 \end{pmatrix} \approx 1 + \begin{pmatrix} 1 & -\theta & 0 \\ \theta & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \theta$$



n axis vector has 3 components so we have 3 generators

$$X_3 = X_{\hat{z}} = \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

structure constants

$$X_i = X_1 = X_2 = X_3$$

$$X_n = \sum_i n_i X_i$$

$$[X_i, X_j] = ?$$

$$[X_n, X_m] = n \times (m \times \mathbf{L}) - m \times (n \times \mathbf{L})$$

$$b(ca) - c(ba)$$

Jacobi identity
or $ba(-cab)$

$$= m(n \cdot \mathbf{L}) - n(m \cdot \mathbf{L}) - n(m \cdot \mathbf{L}) + m(n \cdot \mathbf{L})$$

$$= \mathbf{L} \times (m \times n) = -\mathbf{L} \times (n \times m)$$

$$= (n \times m) \times \mathbf{L} = X_{n \times m}$$

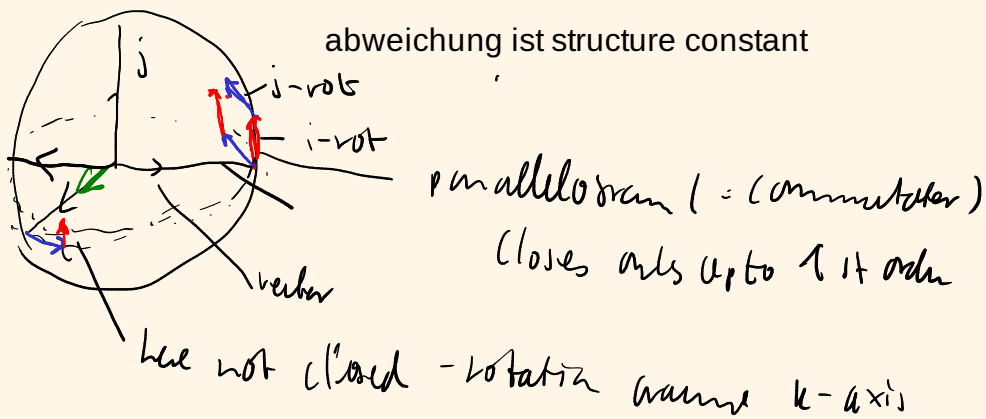
$$\hat{n} = e_i \quad \hat{m} = e_j \quad \hat{n} \times \hat{m} = e_i \times e_j = \epsilon_{ijk} e_k$$

$$\text{hence } [X_i, X_j] = \epsilon_{ijk} X_k$$

$$\text{Structure constants } C_{ij}^k = \epsilon_{ijk}$$

("pure glomerology")

for $SO(3)$ rotations



Hermitian generators

$$J = iX$$

$$[J_k, J_l] = i \epsilon_{klm} J_m$$

same as angular momentum operators

why? angular momentum also elements of a lie algebra (basis of Lie Algebra)

to Lie group of quantum mechanics transformations
which is a representation of rotations

$$\hat{T}_n(\alpha) \approx \hat{T}(\vec{\alpha})$$

if $R(\vec{\alpha})R(\vec{\beta}) = R(\vec{\gamma})$ the T 's must satisfy

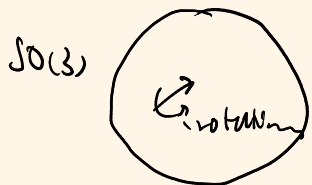
$$T(\vec{\alpha})T(\vec{\beta}) = T(\vec{\gamma})$$

a parametrised form of

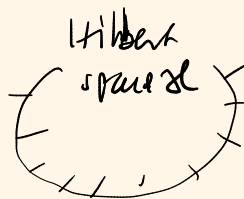
$$T(R)T(R') = T(RR')$$

these Lie algebras have the same structure constants

T "represents" R
space R^3



representation



unitary
transformation group

Wigner theorem:

a group of transformations that leaves scalar products invariant (probabilities are the same in all systems)

$$|\langle T(\psi) | T(\chi) \rangle| = |\langle \psi | \chi \rangle|$$

is a group of unitaries or of anti-unitaries linear transformations.

symmetries in quantum
mechanics with Wigner theorem

in quantum mechanics Hilbert spaces / systems

Quantum mechanics systems are representations ("live in") of the classical symmetry groups

mathematics of Lie algebras - representations are well understood and "just a few"

example rotations $SO(3)$ or $SU(2)$ which unitary representation

for each dimension

$$2j+1 = 1, 2, \dots$$

exactly one irreducible representation

Quantum Zahlen
(Quantum Zahl)

representation

Spin

"Skalar" (invariant) $j=0$

$= 2$ "Spinor"

$j=1/2$

$= 3$ "Vector"

$j=1$

Supergesetz: muss so sein