

Inspiral of binary systems

$$L_{\text{Gw}} = -\frac{dE_{\text{orb}}}{dt} \quad \text{releasing energy due to inspiral and emitting gravitational wave}$$

now $l_0(t)$ kinetic energy

$$E_{\text{orb}} = E_k + U$$

$$E_k = \frac{1}{2} m_1 \omega_1^2 r_1^2 + \frac{1}{2} m_2 \omega_2^2 r_2^2 = \frac{1}{2} \frac{G\mu M}{l_0} \quad \text{reduced mass } \mu = \frac{m_1 m_2}{M}$$

$$r_1 = \frac{m_2 l_0}{M} \quad r_2 = \frac{m_1 l_0}{M} \quad \omega_k = \sqrt{\frac{GM}{l_0^3}}$$

$$E_k = \frac{1}{2} \mu l_0^2 \omega_k^2 \quad U = -\frac{G m_1 m_2}{l_0} = -\frac{G\mu M}{l_0} \quad \text{potential energy}$$

$$\Rightarrow E_{\text{orb}} = -\frac{1}{2} \frac{G\mu M}{l_0}$$

$$\frac{dE_{\text{orb}}}{dt} = +\frac{1}{2} \frac{G\mu M}{l_0^2} \frac{dl_0}{dt} = -E_{\text{orb}} \frac{1}{l_0} \frac{dl_0}{dt}$$

$$\frac{1}{l_0} \frac{dl_0}{dt} = -\frac{\frac{dE_{\text{orb}}}{dt}}{E_{\text{orb}}} = \frac{L_{\text{Gw}}}{E_{\text{orb}}} \quad \text{Luminosity } L_{\text{Gw}} \text{ is derived from the } h \text{ metric}$$

3rd time derivative

$$L_{\text{Gw}} = \frac{G}{5c^5} \sum_{k,lm} \ddot{Q}_{km}^{(l)}(t-\frac{r}{c}) \ddot{Q}^{(l)*}_{km}(t-\frac{r}{c}) = \frac{32\mu^2 G^4 M^3}{5c^5 l_0^5} \quad \text{quadrupole moments}$$

$$\frac{1}{l_0} \frac{dl_0}{dt} = -\frac{32}{5} \frac{G^4 \mu^2 M^3}{c^5 l_0^5} \frac{1}{G\mu M} = -\frac{64}{5} \frac{G^3 \mu M^2}{c^5} \frac{1}{l_0^4}$$

$$\frac{dl_0}{dt} = -\frac{64}{5} \frac{G^3 \mu M^2}{c^5} \frac{1}{l_0^3} \quad \int_{l_0^{\text{init}}}^{l_0(t)} \frac{1}{l_0^3} dl_0 = \int_{t_{\text{Gw}}}^t -\frac{64}{5} \frac{G^3 \mu M^2}{c^5} d\bar{t}$$

$$l_0^4 = (l_0^{\text{init}})^4 - \frac{64}{5} \frac{G^3}{c^3} \mu M^2 t$$

$$l_0(t) = \sqrt[4]{(l_0^{\text{init}})^4 - \frac{256}{5} \frac{G^3}{c^3} \mu M^2 t}$$

$$l_0(t) = l_0^{\text{init}} \left(1 - \frac{t}{t_{\text{coal}}}\right)^{1/4} \quad t_{\text{coal}} = \frac{5}{256} \frac{c^3}{G^3} \frac{(l_0^{\text{init}})^4}{\mu M^2}$$

$$\omega_k(t) = \sqrt{\frac{GM}{l_0^3(t)}} \quad l_0 \text{ insert}$$

$$= \omega_k^{\text{init}} \left(1 - \frac{t}{t_{\text{coal}}}\right)^{-3/8} \quad \omega_k^{\text{init}} = \sqrt{\frac{GM}{(l_0^{\text{init}})^3}}$$

$$f_{\text{GW}}(t) = \frac{2\omega_k(t)}{2\pi} = f_{\text{GW}}^{\text{init}} \left(1 - \frac{t}{t_{\text{coal}}}\right)^{-3/8} \quad f_{\text{GW}}^{\text{init}} = \frac{1}{\pi} \sqrt{\frac{GM}{(l_0^{\text{init}})^3}}$$

$$f_{\text{GW}}(t \rightarrow t_{\text{coal}}) \rightarrow \infty$$

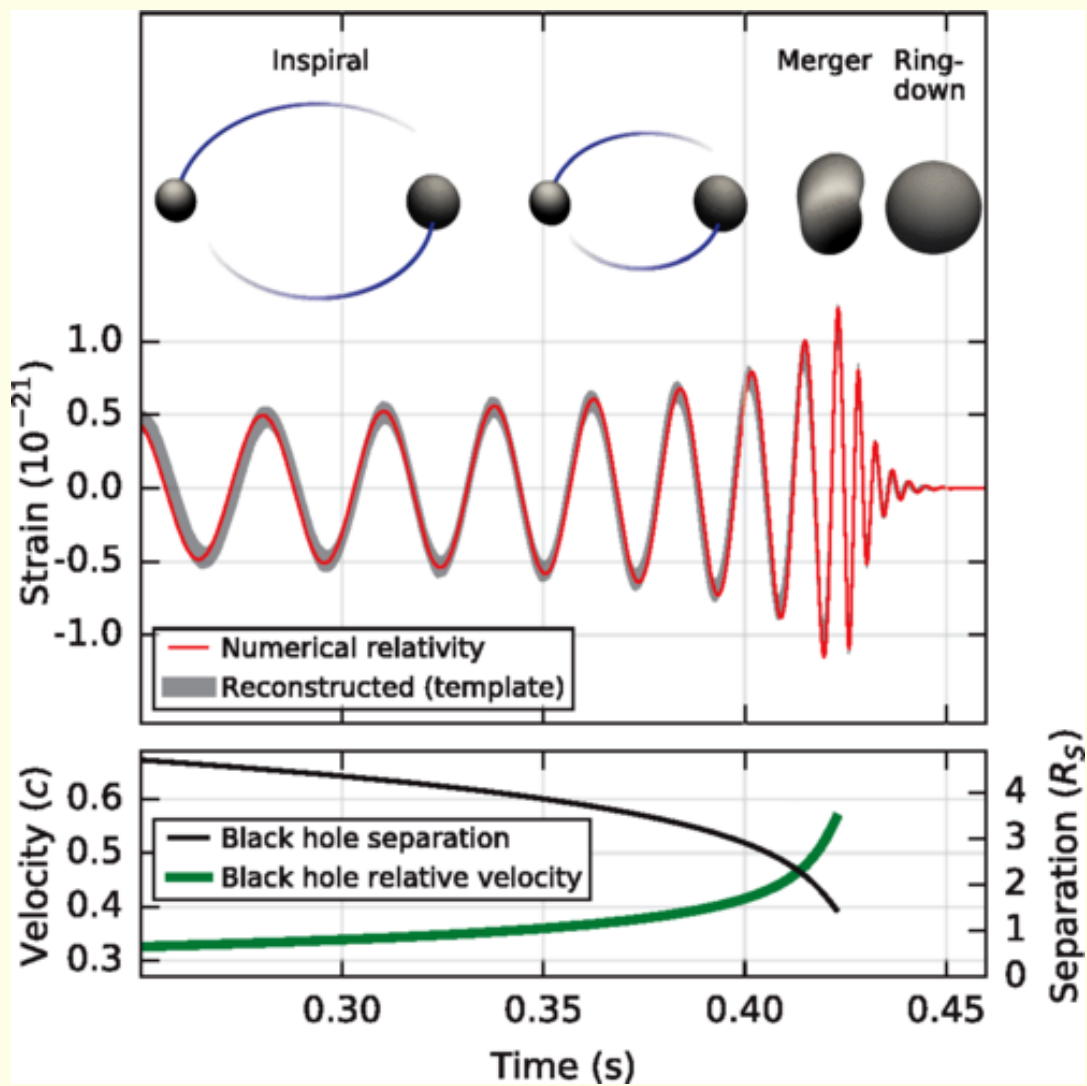
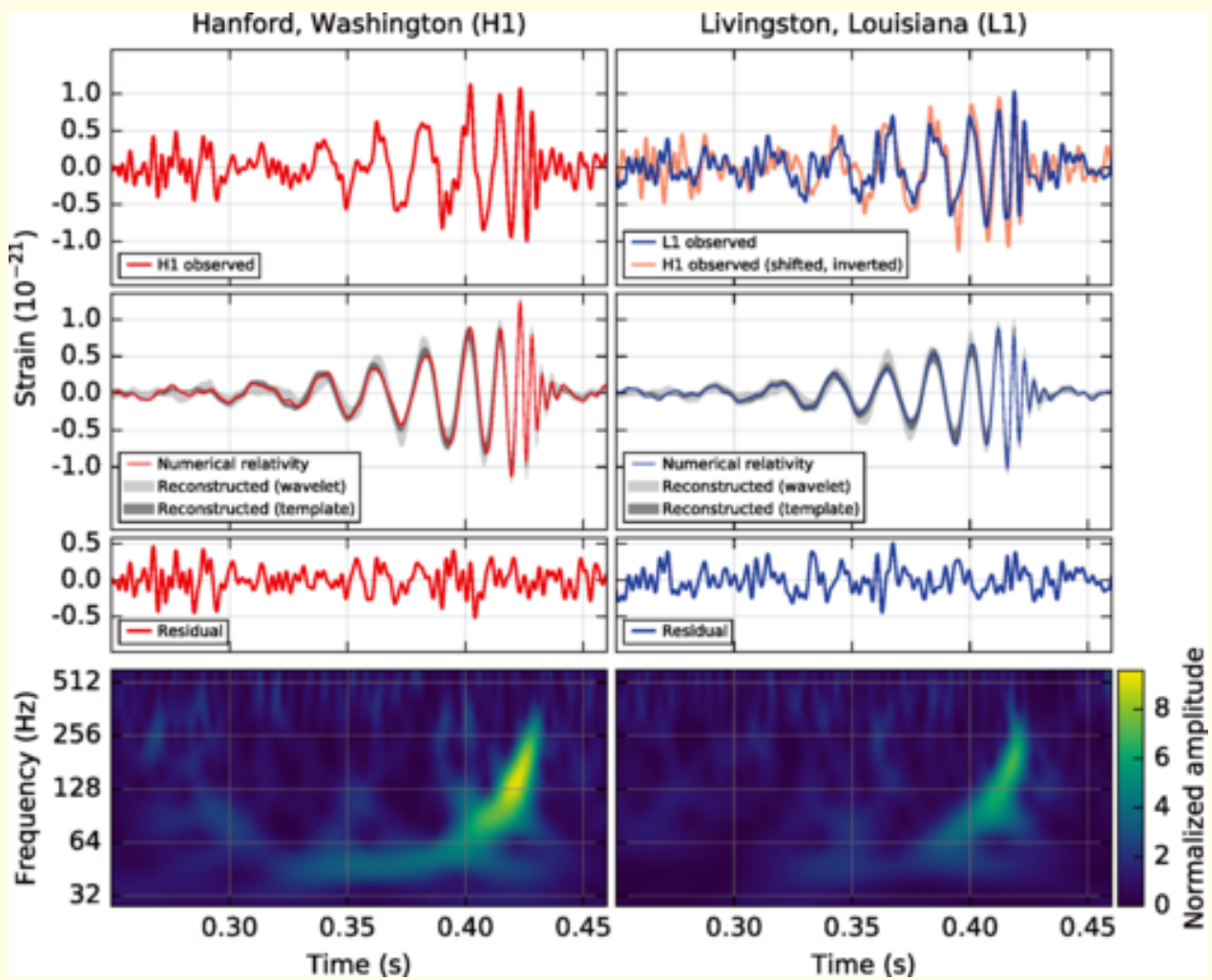
$$h_0(t) = \frac{4\mu M G^2}{c^4 r l_0(t)} = \frac{4\mu M G^2}{c^4 r} \frac{\omega_k(t)^{2/3}}{G^{1/3} M^{1/3}} = \frac{4\mu M G^2}{c^4 r} \frac{f_{\text{GW}}(t)^{2/3}}{G^{1/3} M^{1/3}}$$

$$\text{chirp + mass } \mathcal{M}^{5/3} = \mu M^{2/3}$$

$$h_0(t) = \frac{4\pi^{2/3} G^{5/3}}{c^4 r} f_{\text{GW}}^2(t) \mathcal{M}^{5/3}$$

$$h_0(t \rightarrow t_{\text{coal}}) \rightarrow \infty$$

$$\text{frequency } f_{\text{GW}} \Big|_{\text{max}} \approx 150 \text{ Hz}$$



$$f_{GW}(t) = \frac{1}{\pi} \sqrt{\frac{G\mu}{(t_c - t)^3}} \left(1 - \frac{t}{t_{\text{close}}}\right)^{-3/8} \quad \text{in } \text{in}^{-1}$$

$$f_{GW}^{-8/3}(t) = \frac{(8\pi)^{8/3}}{5} \left(\frac{G\mu}{c^5}\right)^{5/3} (t_c - t)$$

$$\mathcal{M} = \frac{f_{GW}(t)}{(t_c - t)^{3/5}} \frac{5^{3/5}}{(8\pi)^{8/5}} \frac{c^3}{G}$$

example: $f_{GW}(t) = 750 \text{ Hz}$ estimating chirp-mass
 $t = 0.42$ $\mathcal{M} = 30 M_\odot$

$$R = \left(\frac{G\mathcal{M}}{\omega_k^2}\right)^{1/3} = 350 \text{ km} \quad \text{distance}$$

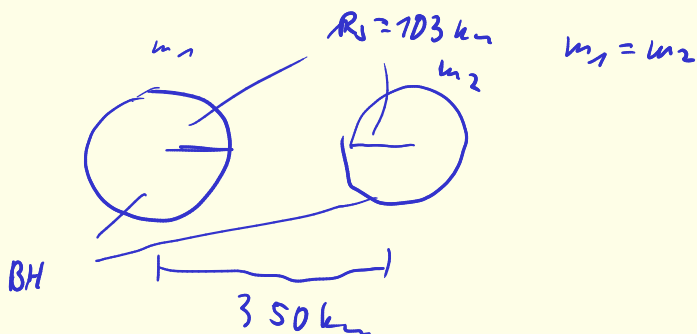
\parallel
 $\frac{2\pi f_{GW}}{2}$

Schwarzschild radius

$$R_S' = \frac{2G}{c^2} m_1 \quad m_1 = m_2$$

$$\mathcal{M} = \mu^{3/5} M^{2/5} = \left(\frac{m_1 m_2}{M}\right)^{3/5} M^{2/5} = 2^{-1/5} m_1 \Rightarrow m_1 = 35 M_\odot$$

$$R_S' \approx 103 \text{ km}$$



It has to be 2 black holes, because if there are two stars, the radii of the stars would be larger and they would have already collided, so we need the black holes

After inspiral of binary systems - remnant black hole

$$e^{i\omega_{\text{aw}} t} = \underbrace{e^{2\pi i f_{\text{ch}} t}}_{\text{wave}} \underbrace{e^{-t/\tau_{\text{damp}}}}_{\text{damping}}$$

$$\omega = \omega_k + i\omega_{\pm} \quad f_{\text{ch}} = \frac{\omega_k}{2\pi} \quad \tau_{\text{damp}} = \frac{1}{\omega_{\pm}}$$

Quasi-normal modes $f_{\text{ch}} \approx 260 \text{ Hz}$ $\tau_{\text{damp}} \approx 4 \text{ ms}$

a test for GR

Using the same assumptions as before the inspiral

$$g_{\mu\nu} = \underbrace{g_{\mu\nu}^0}_{\text{background metric of outside of black hole}} + \underbrace{h_{\mu\nu}}_{\text{perturbation}} \quad \text{non spinning black hole}$$

$G = c = 1$

$$ds^2 = -\left(1 - \frac{2M}{r}\right) dt^2 + \frac{dr^2}{1 - \frac{2M}{r}} + r^2(d\theta^2 + \sin^2\theta d\phi^2)$$

$$\Gamma_{\mu\nu}^{\lambda} = \frac{1}{2} g^{\lambda\epsilon} (h_{\epsilon\mu,\nu} + h_{\epsilon\nu,\mu} - h_{\mu\nu,\epsilon}) + \frac{1}{2} h^{\lambda\epsilon} (g_{\epsilon\mu,\nu}^0 + g_{\epsilon\nu,\mu}^0 - g_{\mu\nu,\epsilon}^0) + O(h)$$

$$\frac{\partial}{\partial x^{\alpha}} \Gamma_{\mu\nu}^{\alpha}(h) - \frac{\partial}{\partial x^{\nu}} \Gamma_{\mu\alpha}^{\alpha}(h)$$

perturbation only scalar field perturbation for the black hole

massless scalar field

$\square \Phi = 0$ for the box derivative we use the Schwarzschild metric

$$\frac{1}{\sqrt{g}} \partial_{\mu} (\sqrt{g} g^{\mu\nu} \partial_{\nu} \Phi) = 0 \quad g = \det g_{\mu\nu}$$

$$\sqrt{g} = r^2 \sin \theta$$

$$\Phi(r, \theta, \phi) = \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{\psi_{lm}(r)}{r} \underbrace{Y_{lm}(\theta, \phi)}_{\substack{\text{spherical harmonics} \\ \text{angular functions}}} e^{-i\omega t}$$

$$Y_{lm} = P_l(\cos \theta) e^{im\phi} = \frac{1}{\sin^2 \theta} \frac{\partial}{\partial \theta} (\sin \theta \partial_{\theta} P_l) - \frac{m^2}{\sin^2 \theta} P_l = -l(l+1) P_l$$

insert Ψ_{lm} in here and the metric

$$\frac{1}{\sqrt{g}} \partial_\mu (\sqrt{g} g^{\mu\nu} \partial_\nu \Psi) = 0$$

$$\left(1 - \frac{2M}{r}\right)^2 \frac{d^2 \Psi}{dr^2} + \frac{2M}{r} \left(1 - \frac{2M}{r}\right) \frac{d\Psi}{dr} + \underbrace{\left(\omega^2 - \left(1 - \frac{2M}{r}\right) \left(\frac{l(l+1)}{r^2} + \frac{2M}{r^3}\right)\right)}_{V_{\text{eff}}(r)} \Psi = 0$$

$$\frac{dr}{dr_*} = 1 - \frac{2M}{r} \quad r_* = r + 2M \log\left(\frac{r}{2M} - 1\right)$$

$$\frac{d\Psi}{dr} = \frac{d\Psi}{dr_*} \frac{dr_*}{dr} = \frac{d\Psi}{dr_*} \frac{1}{\left(1 - \frac{2M}{r}\right)}$$

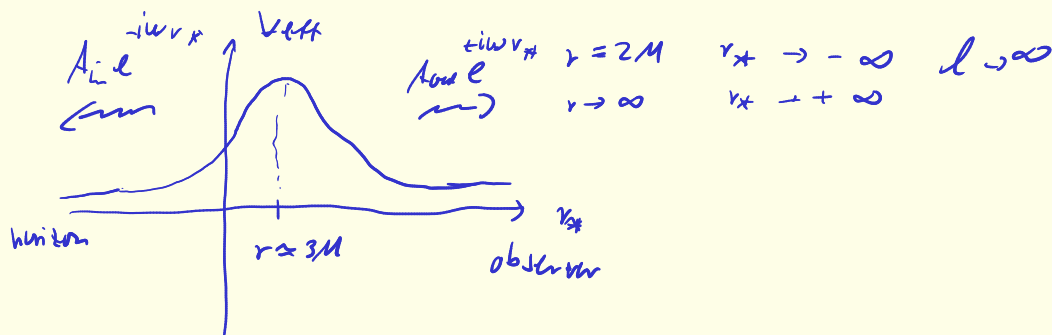
$$\frac{d^2 \Psi}{dr^2} = \frac{d}{dr} \left(\frac{d\Psi}{dr_*} \frac{1}{\left(1 - \frac{2M}{r}\right)} \right) = \frac{d^2 \Psi}{dr_*^2} \frac{dr_*}{dr} \frac{1}{1 - \frac{2M}{r}} + \frac{d\Psi}{dr_*} \frac{d}{dr} \left(\frac{1}{1 - \frac{2M}{r}} \right)$$

rewrite in the new coordinate r_*

$$\frac{d^2 \Psi}{dr_*^2} + (\omega^2 - V_{\text{eff}}(r)) \Psi = 0$$

$$V'_{\text{eff}}(r) = 0 \quad \text{for max and min}$$

$$r \approx 3M \quad \text{-- circular orbit for photons}$$



Considered mode of $l=m=2$

search for eigenvalues of the this equation, for that we need 2 boundary conditions

$$\frac{d^2 \Psi}{dr_*^2} + (\omega^2 - V_{\text{eff}}(r)) \Psi = 0$$

eigenvalues

$$\Psi(r) = A_{\text{in}} e^{-i\omega r_*} + A_{\text{out}} e^{i\omega r_*}$$

$$r \rightarrow 2M \quad \Psi(r) \sim A_{\text{in}} e^{-i\omega r_*}$$

$$r \rightarrow \infty \quad \Psi(r) \sim A_{\text{out}} e^{i\omega r_*}$$

asymptotic region

integrating numerically to get the ω

$$l = 2 \quad f = \frac{\omega_k}{2\pi} \sim 12 \text{ kHz} \quad \frac{M_0}{M} \quad \text{using } 60 M_0$$

$$\tau = \frac{1}{\omega_I} \sim 0.055 \text{ ms} \quad \frac{1}{M_0}$$

k_{eff} for many interactions

$$k_{\text{eff}}(r) = \left(1 - \frac{2M}{r}\right) \left[\frac{l(l+1)}{r^2} + (1 - \delta^2) \frac{2M}{r^3} \right]$$

$\delta = 0$ scalar
 $\delta = -1$ lem
 $\delta = -2$ grav.

for spinning $\chi = 0.7$ $\delta = -2$, $l = m = 2$

$$M_{\omega_k} = 0.5148 \quad M_{\omega_I} = 0.0818$$

$$f = \frac{M_{\omega_k}}{2\pi M} \approx 250 \text{ kHz}$$

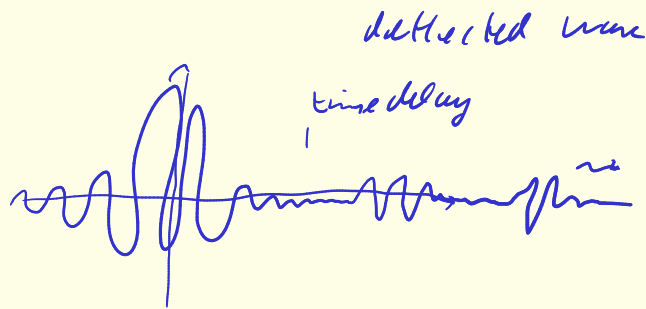
$$\tau = \left(\frac{M_{\omega_k}}{M} \right)^{-1} = 4 \text{ ms}$$

modified solution to have also
contribution from the black hole
not only the outward solution wave

Lecturer:

Research in ring down of the black holes

allowing partial reflection of the wave at the black hole



testing GR at the horizon
with gravitational waves
signals after the ringdown
some kind of echo

maybe quantum reflections

maybe without a horizon