

## Examples:

\* Pure coherent states:  $\hat{S} = |\gamma \times \beta|$ ,  $\gamma \in \mathbb{C}$

Nomenclature

$$|\mu\rangle \rightarrow |\beta\rangle \quad \langle -\beta | \hat{S} | \beta \rangle = \langle -\beta | \gamma \rangle \times \gamma | \beta \rangle \\ = e^{-|\beta|^2} e^{-|\gamma|^2} e^{-\beta^* \gamma - \gamma^* \beta}$$

$$\Rightarrow P_\gamma(\omega) = e^{|\alpha|^2} e^{-|\gamma|^2} \frac{1}{\pi^2} \int e^{-\beta^* \gamma - \gamma^* \beta + \beta^* \alpha - \beta \alpha^*} d\beta^2$$

$$\boxed{\int_{-\infty}^{\infty} e^{i\omega x} dx = 2\pi \delta(\omega)}$$

$$= e^{|\alpha|^2} e^{-|\gamma|^2} \left( \underbrace{\frac{1}{\pi} \int_{-\infty}^{\infty} d\beta_R}_{= \delta(\alpha_R - \gamma_R)} e^{2i\beta_R(\alpha_R - \gamma_R)} \right) \left( \underbrace{\frac{1}{\pi} \int_{-\infty}^{\infty} d\beta_I}_{= \delta(\alpha_I - \gamma_I)} e^{2i\beta_I(\alpha_I - \gamma_I)} \right)$$

$$= \delta^{(2)}(\alpha - \gamma) \quad [\text{as expected}]$$

\* Pree number state  $\hat{f} = |n \times u\rangle$

$$\langle -\beta | \hat{f} | \beta \rangle = \langle \beta | n \times u | \beta \rangle = e^{-|\beta|^2} \frac{(-\beta^* \beta)^n}{n!}$$

$$e^{-|\beta|^2/2} \beta^n$$

$$\frac{1}{\sqrt{n!}}$$

$$\Rightarrow P_n(x) = \frac{e^{(x)^2}}{n!} \frac{1}{\pi^2} \int (-\beta^* \beta)^n e^{\beta^* x - \beta x^*} d^2 \beta$$

Formal trick

$$= \frac{e^{(x)^2}}{n!} \frac{\partial^{2n}}{\partial x^n \partial x^{*n}} \frac{1}{\pi^2} \int e^{\beta^* x - \beta x^*} d^2 \beta$$

$$\frac{1}{\pi^2} \int e^{2ixn(\beta^* x)} d^2 \beta = \delta^{(2)}(x)$$

$$\Rightarrow \boxed{P_n(x) = \frac{e^{(x)^2}}{n!} \frac{\partial^{2n}}{\partial x^n \partial x^{*n}} \delta^{(2)}(x)}$$

Comments:

- 1) can only be interpreted under integral
- 2) more singular than  $\delta$ -function

## 8.2 Optical Equivalence Theorem

Define : normal ordering superoperator

$$:\hat{O}(\hat{a}, \hat{a}^+): = \hat{O}^N(\hat{a}, \hat{a}^+) \leftarrow$$

$\rightarrow$  often useful!

collect all creation operators to the left and all annihilation operators to the right

$$\text{e.g. } :\hat{n}^2: = :\hat{a}^+ \hat{a} \hat{a}^+ \hat{a}: = \hat{a}^+ \hat{a}^+ \hat{a} \hat{a}$$

Optical Equivalence Theorem

$\rightarrow$  consider a generic normal-ordered function of  $\hat{a}$  and  $\hat{a}^+$

$$\hat{G}^N(\hat{a}, \hat{a}^+) = \sum_n \sum_m C_{nm} \hat{a}^{+n} \hat{a}^m \quad (\text{any expression can be brought into this form})$$

Its expectation value is

$$\langle \hat{G}^N(\hat{a}, \hat{a}^+) \rangle = \text{Tr} \left[ \hat{G}^N(\hat{a}, \hat{a}^+) \hat{\rho} \right]$$

$$\hat{\rho} = \int P(x) |x\rangle \langle x| dx^2$$

$$= \text{Tr} \left[ \underbrace{\int P(x) \sum_{nm} C_{nm} \hat{a}^{+n} \hat{a}^m |x\rangle \langle x| dx^2}_{\text{underbrace}} \right]$$

$$= \int P(x) \sum_{nm} C_{nm} \underbrace{\sum_e \langle e | \hat{a}^{+n} \hat{a}^m |x\rangle \langle x| e \rangle}_{\text{underbrace}} dx^2$$

remove  $\sum_e |ex\rangle \langle ex|$

$$= \int P(\alpha) \sum_{nm} c_{nm} \langle \alpha | \underbrace{\hat{a}^+}_{\langle \alpha | (\alpha^*)^n} \underbrace{\hat{a}^m}_{\alpha^m | \alpha} \rangle d^2 \alpha$$

$$\Rightarrow \langle \hat{G}^{(N)}(\hat{a}, \hat{a}^+) \rangle = \int P(\alpha) \underbrace{\sum_{nm} c_{nm} \alpha^*^n \alpha^m}_{G^{(N)}(\alpha, \alpha^*)} d^2 \alpha$$

→ The expectation value of a normal ordered operator is the P-function weighted average of the function obtained by

$$\hat{a} \rightarrow \alpha ; \quad \hat{a}^+ \rightarrow \alpha^*$$

→ easy way of computing expectation values

$$\text{e.g.: } \langle : \hat{a}^2 : \rangle = \langle (\hat{a})^2 \hat{a}^2 \rangle = \int P(\alpha) |\alpha|^4 d^2 \alpha$$

example:



$$\langle : \hat{a}^2 : \rangle_j = \int P_j(\alpha) |\alpha|^4 d^2 \alpha = |j|^4$$

↑  
coherent state  
↑  
 $\delta^{(2)}(j - \alpha)$

How to make an operator normal ordered?

$$\underbrace{\hat{a}^\dagger \hat{a} \hat{a}^\dagger \hat{a}}_{\hat{a}^\dagger \hat{a} + 1} = \hat{a}^\dagger \hat{a}^\dagger \hat{a} \hat{a} + \hat{a}^\dagger \hat{a}$$

$$\rightarrow \langle : \hat{a}^\dagger \hat{a} \hat{a}^\dagger \hat{a} : \rangle_f$$

$$[\hat{a}, \hat{a}^\dagger] = \hat{a} \hat{a}^\dagger - \hat{a}^\dagger \hat{a} = 1$$

$$= |\psi|^4 + |\psi|^2$$

$$\hat{a} \hat{a}^\dagger = \hat{a}^\dagger \hat{a} + \underbrace{[\hat{a}, \hat{a}^\dagger]}_{} = 1$$

## 5.2 Wigner function

"Phase space" representation of a quantum state

here 1D:  $W(q, p) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} \langle q + \frac{x}{2} | \hat{g} | q - \frac{x}{2} \rangle e^{ipx/\hbar} dx$

Example: pure state  $\hat{g} = |\psi(x)\rangle$

$$W(q, p) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} \underbrace{\langle q + \frac{x}{2} | \psi \rangle}_{\psi(q + \frac{x}{2})} \underbrace{\langle \psi | q - \frac{x}{2} \rangle}_{\psi^*(q - \frac{x}{2})} e^{ipx/\hbar} dx$$

Test 1: Integrate out momentum degree of freedom

$$\int_{-\infty}^{\infty} W(q, p) dp = \int_{-\infty}^{\infty} dx \psi^*(q - \frac{x}{2}) \psi(q + \frac{x}{2}) \underbrace{\frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} e^{ipx/\hbar} dp}_{= \delta(x)} \quad \text{margin: } \begin{array}{c} x_2 \\ \uparrow \\ p(x_1, x_2) \end{array}$$

→ probability density over position  $q$