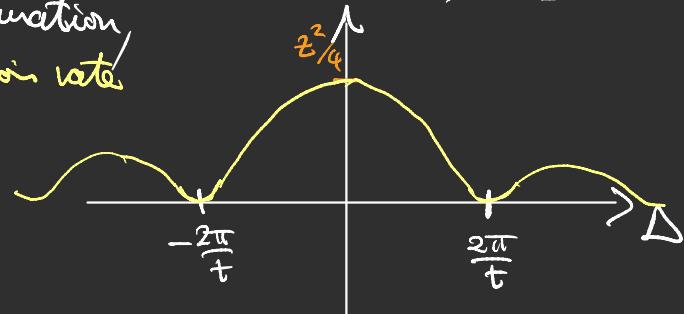


$$P_{i \rightarrow m}(t) = \frac{|\vec{d}_{mi} \cdot \vec{E}_0|^2}{t^2} \frac{\sin^2(t\Delta/2)}{\Delta^2}$$

For $t \gg \frac{2\pi}{|\Delta|}$:

$$\sin^2(t\Delta/2)/\Delta^2 \approx \frac{\pi}{2} t \delta(\Delta) \quad \text{under this approximation}$$

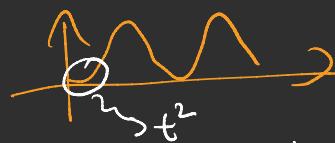
we obtain a transition rate



$$W_{i \rightarrow m}(t) = \frac{P_{i \rightarrow m}(t)}{t} = \frac{\pi}{2} \frac{|\vec{d}_{mi} \cdot \vec{E}_0|^2}{t^2} \delta(\omega - \omega_{mi})$$

→ typically applies to
discrete → continuous
transition

(Fermi's Golden Rule)

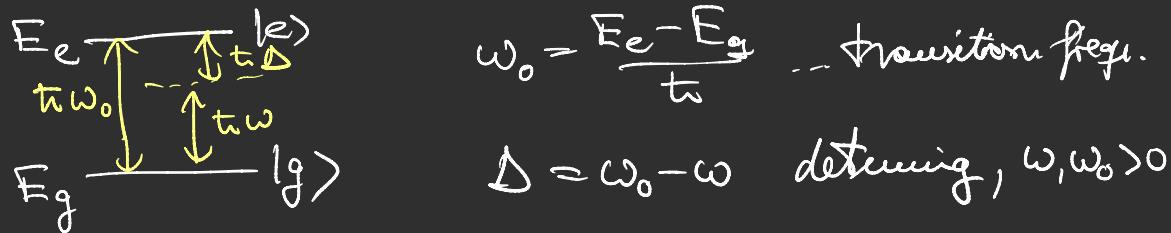


→ Comment: energy is conserved for $t \gg \frac{\hbar}{|\Delta E|}$

→ Comment 2: applies to time-independent perturbations as well: $\delta(E_f - E_i)$

6.3 Rabi model

→ one of the rare examples of an exactly solvable model



$$\hat{H}'(t) = -\vec{d} \cdot \vec{E}_0 \cos(\omega t) = \hat{V}_0 \cos(\omega t)$$

Express general atomic state as a superposition of $|g\rangle, |e\rangle$

$$|\Psi(t)\rangle = c_g(t) e^{-i\omega_g t} |g\rangle + c_e(t) e^{-i\omega_e t} |e\rangle$$

$$c_g(0) = 1, c_e(0) = 0$$

Insert into TDSE: $\partial_t |\Psi(t)\rangle = [\hat{H}_0 + \hat{H}'(t)] |\Psi(t)\rangle$

and project onto state $\langle g|$ and $\langle e|$

→ We assume atomic states to be of definite parity:

$$\langle g | \hat{d} | g \rangle = \langle e | \hat{d} | e \rangle = 0$$

$$\text{and define } \mathcal{V} = \langle e | \hat{V}_0 | g \rangle = -\vec{d}_{eg} \cdot \vec{E}_0$$

→ We can assume \mathcal{V} as real: $\mathcal{V} = \mathcal{V}^*$

→ Equations of motion:

$$\dot{c}_g = -\frac{i}{\hbar} \mathcal{J} \cos(\omega t) e^{-i\omega_0 t} c_e$$

$$\dot{c}_e = -\frac{i}{\hbar} \mathcal{J} \cos(\omega t) e^{i\omega_0 t} c_g$$

$\nwarrow \nearrow$
 $e^{i(\omega_0 \pm \omega)t}$

RWA:

$$\dot{c}_g = -\frac{i}{2\hbar} \mathcal{J} e^{i(\omega - \omega_0)t} c_e$$

$$\dot{c}_e = -\frac{i}{2\hbar} \mathcal{J} e^{-i(\omega - \omega_0)t} c_g$$

→ form an equivalent 2nd-order ODE:

$$\ddot{c}_e e^{i(\omega - \omega_0)t} + i(\omega - \omega_0) \dot{c}_e e^{i(\omega - \omega_0)t} = -\frac{i}{2\hbar} \dot{c}_g$$
$$= -\frac{\mathcal{J}^2}{(2\hbar)^2} e^{i(\omega - \omega_0)t} c_e$$

$$\ddot{c}_e + i(\omega - \omega_0) \dot{c}_e + \frac{1}{4} \frac{\mathcal{J}^2}{\hbar^2} c_e = 0$$

→ Try solutions of type $c_e(t) = A_+ e^{i\lambda_+ t} + A_- e^{i\lambda_- t}$
from initial conditions

$$-\lambda^2 - (\omega - \omega_0)\lambda + \frac{1}{4} \frac{\sigma^2}{t_h^2} = 0$$

$$\Leftrightarrow \lambda = \frac{(\omega - \omega_0) \pm \sqrt{(\omega - \omega_0)^2 + \frac{\sigma^2}{t_h^2}}}{-2\Delta}$$

$$\lambda_{\pm} = \frac{1}{2} (\Delta \pm \sqrt{\Delta^2 + \frac{\sigma^2}{t_h^2}})$$

We define : $\Omega_R = \sqrt{\Delta^2 + \frac{\sigma^2}{t_h^2}}$ generalized Rabi frequency

→ with initial conditions $\Delta = \sqrt{3} \frac{\sigma}{t_h}$

$$c_e(t) = i \frac{\sigma}{\Omega_R t_h} e^{i t \Delta / 2} \sin\left(\frac{\Omega_R t}{2}\right)$$

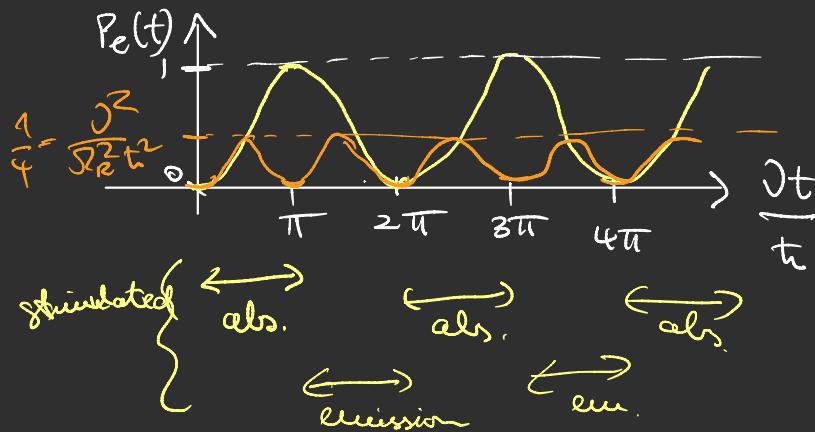
The probability that the atom is in the excited state is:

$$P_e(t) = |c_e(t)|^2 = \frac{\sigma^2}{\Omega_R^2 t_h^2} \sin^2\left(\frac{\Omega_R t}{2}\right) \xrightarrow[\Delta=0]{\text{resonance}} \sin^2\left(\frac{\sigma t}{2 t_h}\right)$$

↑ transfer is reduced by detuning

Rabi oscillations

$$\Delta = 0$$



$$\Delta = \sqrt{3} \frac{\sigma}{\pi}$$

$$\Omega_R \rightarrow 2 \Omega_R$$

compared to $\Delta = 0$

Rabi frequency is not applied frequency, but rather a measure of the coupling strength of the EM field with the "atom" (2-level system)

Typical values of Ω_R

$$1 \text{ eV} \sim 2.4 \cdot 10^{14} \text{ Hz}$$

* In InGaAs QD (quantum dots)

\sim PRL 87,

$$\text{In } \Omega_R \sim 100 \mu\text{eV} \text{ at } 3.8K, \text{ cw } 246401 \text{ (2001)}$$

$$\hookrightarrow \frac{\Omega_R}{2\pi} \sim 246 \text{ Hz (microwave)} \quad I \approx 28 \frac{K_w}{\text{cm}^2}$$

$$\frac{1 \text{ eV}}{20.2 \mu\text{eV}} = \frac{|x^* \rangle}{|x\rangle}$$

* transmon qubit \sim PRX 10, 031032 (2020)

transition $\sim 7.86 \text{ Hz}$, typical Rabi freq.

$$\frac{\Omega_R}{2\pi} \sim 1 \text{ Hz}$$