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Investigations on two-way coupling effects of particle-laden decaying isotropic turbulent flows

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I Nomenclature

Greek Symbols

γ	Isentropic exponent
η	Kolmogorov length scale
λ	Taylor microscale
μ	Dynamic viscosity
ν	Kinematic viscosity
ρ	Fluid density
ρ_p	Particle density
$\bar{\tau}$	Stress tensor
τ_η	Kolmogorov time scale
τ_L	Eddy turnover time
τ_p	Particle relaxation time
τ_{ps}	Stokes relaxation time
ψ	Coupling rate

Operators

∇	Nabla operator
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Roman Symbols

c_p	Specific isobaric heat capacity
c_v	Specific isochoric heat capacity
d_p	Particle diameter
f_D	Drag correction
E	Specific inner energy
e	Specific internal energy

\boldsymbol{F}	Force per unitary volume of the particle acting on the fluid
\boldsymbol{f}_a	Added mass force
\boldsymbol{f}_d	Drag force
\boldsymbol{f}_h	History force
\boldsymbol{f}_l	Lift force
\boldsymbol{f}^n	Sum of forces acting on the fluid
$\bar{\boldsymbol{H}}$	Flux tensor
$\bar{\boldsymbol{H}}^i$	Inviscid part of the flux tensor
$\bar{\boldsymbol{H}}^v$	Viscous part of the flux tensor
$\bar{\boldsymbol{I}}$	Identity tensor
k_t	Thermal conductivity
L	Integral length scale
m_c	Number of clustered particles
N_c	Number of computational particles
N_p	Number of physical particles
p	Pressure
Pr	Prandtl number
\boldsymbol{Q}	Vector of conservative Eulerian variables
\boldsymbol{q}	Heat conduction vector
R	Specific gas constant
Re	Reynolds number
Re_p	Particle Reynolds number
r_p	Particle radius
\boldsymbol{S}	Rate-of-strain tensor
S	Sutherland temperature

T	Temperature
t^*	Time normalized by the initial eddy turnover time
\boldsymbol{u}	Fluid velocity
U	Characteristic fluid velocity
V_p	Particle volume
\boldsymbol{v}_p	Particle velocity vector
\boldsymbol{x}_p	Particle position vector

2 Introduction

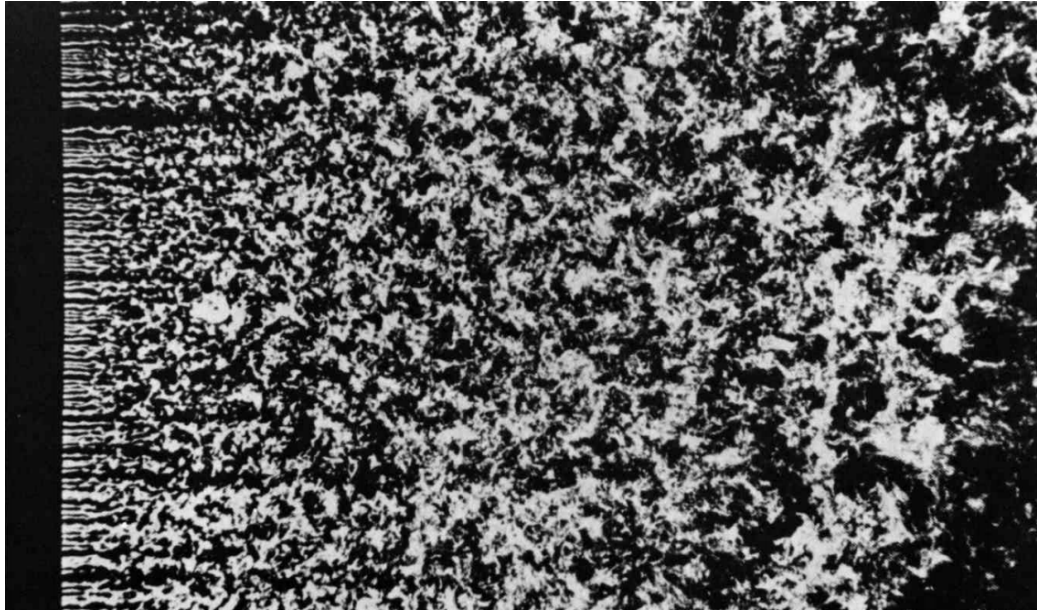


Figure 1: Photograph courtesy of Hassan Nagib and Thomas Corke: Formation of an nearly isotropic turbulent flow field behind a grid. Source: [?]

iiiiii HEAD This work deals with the effects of particles on the flow conditions of isotropic turbulent flows. There are many examples of these particle-laden turbulent flows in nature, e.g. in volcanic eruptions and in the "white water" of breaking waves.

Multiphase flows also occur in many technical applications. Spray atomization in fuel injectors, cyclonic particle separation in oil refineries and sediment accumulation in pipelines are just three of many cases, where it is of huge interest to predict the influence of particles on the turbulence of the flow. Of major industrial interest are the fuel dispersion and combustion in an petrol-powered engine. To get a homogeneous and fast combustion, both should take place at turbulent flow conditions. One could say it is an important question, in which way and intensity the particles influence the turbulence. ===== Particle laden turbulent flows are ubiquitous in nature, e.g. in volcanic eruptions and in the "white water" of breaking waves. Spray atomization in fuel injectors, cyclonic particle separation in oil refineries and sediment accumulation in pipelines are examples of technical applications, where it is of huge interest to predict the impact of particles on turbulent flows. Turbulence augmentation or attenuation by particles is therefore a decisive factor. 4113d64895656e6ba27b195eb1861d9a73879796

A study on the impact of particles on the flow conditions of isotropic turbulent

flows is presented at hand. In this work the influence of the particles on the conditions is presented. Analyzing the energy transfer between particles and fluid will be focused.

iiiiii HEAD In this work the influence of the particles on the streaming conditions is presented by using DNS and LES simulations. The focus lies on analyzing the energy transfer between particles and fluid as well as the intensity of the decrease of turbulent kinetic energy in the fluid. At the beginning some mathematical background information about modeling turbulent single phase and particle-laden flows are given. Additionally the Navier-Stokes equations and the different scales of turbulent motion are presented as an approach to modelling turbulence. For particle-laden flows the equations of motion are expanded by a new term that describes the forces acting between fluid and particles. Also the coupling rate that describes the exchanged energy and the computational basics of methods used in this work are introduced. According to the computational basics the simulation methods DNS and LES, how they work and in which case which one is used is explained. Then the used discretization method to integrate the particle tracking equations, which is a predictor-corrector-scheme, is described. Finally the computational method 'particle clustering' is introduced. This approach aims to decrease the computational effort of particle-laden simulations. It works the following way: the particles are coupled to clusters and each cluster is considered as one particle. In the following the deviations caused by using this method are determined and the applicability is evaluated. Therefore the influence of the number of particles per cluster on the accuracy of the plots of the turbulent kinetic energy and the exchanged energy between particles and fluid per time is identified. =====

First, some mathematical background information about modeling single phase and particle laden flows are given. Additionally, the scales of turbulent motion are described. According to the computational basics the simulation methods DNS and LES, how they work and in which case which one is used is explained. Then the used discretization method to integrate the particle tracking equations is described. Finally the computational method 'particle clustering' is introduced. In the following the deviations caused by using this method are determined and the applicability is evaluated. ~~~~~ 4113d64895656e6ba27b195eb1861d9a73879796 Also the number of particles that is needed to get consistent average values for the flow characteristics is investigated.

-Due to the small particle concentration, a popular method to describe these flows is the point-particle approach, which means that every particle is treated as a mathematical point source of mass, momentum and energy.

3 Mathematical models

3.1 Single phase flow

In this section, the mathematical models are introduced, which describe turbulent flows. For a more detailed description, see e.g. [16].

3.1.1 Equations governing the fluid phase

The conservation of mass, momentum and energy for a control volume reads

$$\int_V \frac{\partial \mathbf{Q}}{\partial t} dV + \int_{\partial V} \bar{\mathbf{H}} \cdot \mathbf{n} dA = \mathbf{0}. \quad (3.1)$$

with time t and the flux tensor $\bar{\mathbf{H}}$. The vector \mathbf{Q} contains the variables fluid density ρ , fluid velocity \mathbf{u} and specific inner energy E :

$$\mathbf{Q} = \begin{pmatrix} \rho \\ \rho \mathbf{u} \\ \rho E \end{pmatrix}. \quad (3.2)$$

$\bar{\mathbf{H}}$ is the flux tensor which contains the inviscid and viscous flux, i.e.:

$$\bar{\mathbf{H}} = \bar{\mathbf{H}}^i + \bar{\mathbf{H}}^v = \begin{pmatrix} \rho \mathbf{u} \\ \rho \mathbf{u} \mathbf{u} + p \\ \mathbf{u}(\rho E + p) \end{pmatrix} - \frac{1}{Re} \begin{pmatrix} 0 \\ \bar{\boldsymbol{\tau}} \\ \bar{\boldsymbol{\tau}} \mathbf{u} + \mathbf{q} \end{pmatrix} \quad (3.3)$$

$\bar{\mathbf{H}}^i$ contains the variables that are independent of the fluid's viscosity and $\bar{\mathbf{H}}^v$ represents the effects of viscosity $\bar{\boldsymbol{\tau}}$ and heat conduction \mathbf{q} . The Reynolds number $Re = \frac{\rho v d}{\mu}$ is defined to be the ratio of inertia forces to viscous forces. This is also due to the fact that two familiar objects with Re being identical behave similar in flows. To solve the Navier-Stokes equations, more information regarding some variables is required. For calculating the specific inner energy E and the heat conduction \mathbf{q} , the following equations are used

$$E = e + \frac{1}{2} |\mathbf{u}|^2, \quad (3.4)$$

$$\mathbf{q} = -\frac{\mu}{Pr(\gamma - 1)} \nabla T, \quad (3.5)$$

with

$$\gamma = \frac{c_p}{c_v} \quad (3.6)$$

and the Prandtl number $Pr = \frac{\mu_\infty c_p}{k_t}$ using the specific heat capacities of the fluid c_v and c_p . Assuming that the fluid is a newtonian fluid, the linear correlation between stress and the rate of strain results in:

$$\bar{\tau} = 2\mu\bar{\mathbf{S}} - \frac{2}{3}\mu(\nabla \cdot \mathbf{u})\bar{\mathbf{I}}, \quad (3.7)$$

in which $\bar{\mathbf{S}} = \frac{(\nabla \mathbf{u})(\nabla \mathbf{u})^T}{2}$ denotes the rate-of-strain tensor. Additionally, the viscosity μ can be approximated by Sutherland's law, which is based on the ideal gas-theory

$$\mu(T) = \mu_\infty \left(\frac{T}{t_\infty} \right)^{3/2} \frac{T_\infty + S}{T + S}, \quad (3.8)$$

where S is the Sutherland temperature. To achieve closure the caloric state equation $e = c_v T$ and the state equation for an ideal gas $p = \rho R T$ are used. The specific gas constant is determined by $R = c_p - c_v$.

3.2 Scales of turbulent flows

Turbulent flows contain eddies, which are induced by reverse currents in the flow, of all sizes and shapes. Large-scale eddies bring energy to the flow which is then passed down to smaller-scales, and finally dissipated into heat by viscous effects. This behavior is called the 'energy cascade' and was first described by Richardson in the year 1920 [17]. The theory then was further developed by Kolmogorov [12]. The main difference is that Kolmogorov added hypotheses like the statistical isotropy of small-scale turbulent motions and the decrease of the velocity and timescales with the decrease of the length scale along the energy cascade.

The first set of scales describe the large eddies. At these scales the energy is brought into the flow, creating the 'energy-containing range'. The corresponding timescale, which is most times called 'eddy turnover time', is defined as

$$\tau_L = \frac{L}{U}, \quad (3.9)$$

with the length scale L and the characteristic fluid velocity U .

The smallest scales in a turbulent flow are the Kolmogorov length (η) and time (τ_η) scale. At these scales, the effects of viscosity take place and the energy dissipates into heat. With the estimate $\epsilon \approx \frac{U^3}{L}$ they can be written as

$$\eta = \left(\frac{\nu^3 L}{U^3} \right)^{1/4}, \quad (3.10)$$

$$\tau_\eta = \left(\frac{\nu L}{U^3} \right). \quad (3.11)$$

Both these scales are coupled by the Reynolds numbers

$$\frac{L}{\eta} = Re^{3/4}, \quad (3.12)$$

$$\frac{\tau_L}{\tau_\eta} = Re_L^{1/2}. \quad (3.13)$$

It can be observed from these equations that the spacing between the scales increases for higher Reynolds numbers.

A scale between these two is the Taylor microscale, often referred to as 'turbulence length scale'. It is often used to describe the intermediate range between integral and Kolmogorov scales. This scale's definition is:

$$\lambda = \sqrt{15 \frac{\nu}{\epsilon}} |\mathbf{u}'| \quad (3.14)$$

with $|\mathbf{u}'|$ denoting the absolute value of the velocity's fluctuation. This scale can be used to compute the Taylor-scale Reynolds number Re_λ :

$$Re_\lambda = \frac{|\mathbf{u}'| \lambda}{\nu}. \quad (3.15)$$

3.3 Particle dynamics

This work deals with particle-laden fluids, therefore the interaction between particles and the carrier fluid needs to be considered in the equation of motion. The phenomenon described in this study is called two-way-coupling. This chapter is based on [1].

In this thesis, small and heavy rigid particles with a spherical shape are regarded. Their diameter d_p is even smaller than the Kolmogorov scale η , but also large enough to neglect the Brownian motion.

Due to the small particle concentration, a popular method to describe these flows is the point-particle approach, which means that every particle is treated as a mathematical point source of mass, momentum and energy. In this case the focus lies on the momentum exchange, effects like particle collision are neglected.

The assumptions that the fluid is incompressible and that mass exchange over the particle surface as well as gravity are negligible are used. After including all the simplifications mentioned before, the Navier-Stokes equation for the carrier fluid "Gleichung xy" - described in chapter 1 - is extended by the force \mathbf{F} that acts on the ambient cells with different intensity regarding on the distance between these cells and the particle:

$$\mathbf{F} = \mathbf{F}_{pp} \cdot \frac{e^{-(d_i^2/(\sigma\Delta^2))}}{\sum_i e^{-(d_i^2/(\sigma\Delta^2))}}, \quad (3.16)$$

with the distance between particle position and midpoint of the cell d_i , the length of the cells Δ and a smoothing parameter σ that controls the distribution of \mathbf{F} on the adjacent cells. \mathbf{F}_{pp} is the sum of pressure and shear forces acting on the particles and is described explicit in the Maxey Riley equation that can be achieved from [14]. For the equation of motion for the particles follows

$$V_p \rho_p \frac{d\mathbf{v}_p}{dt} = \mathbf{F}_{pp}. \quad (3.17)$$

\mathbf{F}_{pp} could be divided into several forces. Reduced to the governing forces acting on the particles the equation of motion for the particles becomes

$$m_p \frac{d\mathbf{v}_p}{dt} = \rho V_p \frac{D\mathbf{u}}{Dt} - \frac{3}{4} \rho V_p \frac{C_d}{d_p} |\mathbf{v}_p - \mathbf{u}| (\mathbf{v}_p - \mathbf{u}) + \frac{1}{2} \rho V_p \left(\frac{D\mathbf{u}}{Dt} - \frac{d\mathbf{v}_p}{dt} \right) + \frac{3}{2} d_p^2 \rho \sqrt{\pi \nu} \int_{t_0}^t \frac{dt'}{(t-t')^{1/2}} \left(\frac{D\mathbf{u}}{Dt'} - \frac{d\mathbf{v}_p}{dt'} \right). \quad (3.18)$$

\mathbf{u} is the velocity of the uniform stream, which is distant enough from the particle that it is not disturbed by it. The individual forces on the right side of the equa-

tion are stated in the following:

- $\rho V_p \frac{D\mathbf{u}}{Dt}$:
- $-\frac{3}{4}\rho v_p \frac{C_d}{d_p} |\mathbf{v}_p - \mathbf{u}| (\mathbf{v}_p - \mathbf{u})$:
hydrodynamical drag force that is parallel to the undisturbed streamlines, which depends on an empirical drag coefficient C_d defined by Schiller and Naumann as $C_d = \frac{24}{Re_p} (1 + 0.15 Re_p^{0.687})$.
- $\frac{1}{2}\rho v_p \left(\frac{D\mathbf{u}}{Dt} - \frac{d\mathbf{v}_p}{dt} \right)$:
added mass force, that represents the influence of the fluid's inertia that has an impact on the particle, if it has a different acceleration than the mean flow.
- $\frac{3}{2}d_p^2\rho\sqrt{\pi\nu} \int_{t_0}^t \frac{dt'}{(t-t')^{1/2}} \left(\frac{D\mathbf{u}}{Dt'} - \frac{d\mathbf{v}_p}{dt'} \right)$:
The history force takes diffusion and convection into account that results from the vortices behind the particles

The kinematic equation of a particle is

$$\frac{d\mathbf{x}_p^n}{dt} = \mathbf{v}_p^n, \quad (3.19)$$

with \mathbf{v}_p^n being the velocity of the n-th particle.

As presented in -quelle- due to the density relation $\rho_p/\rho \gg 1$ all forces except the drag force are negligible. Then the equation becomes

$$m_p \frac{d\mathbf{v}_p}{dt} = \frac{3}{4}\rho v_p \frac{C_d}{d_p} |\mathbf{v}_p - \mathbf{u}| (\mathbf{v}_p - \mathbf{u}) \quad (3.20)$$

In this case of gas-solid suspensions several simplifications could be made, because of the density ratio $\rho_p/\rho \gg 1$.

In the following the fact that the added mass force, the lift force and the history force are negligible small compared to the drag force is shown. Vincenzo Armenio and Virgilio Fiorotto [2] show that the added mass force is much smaller than the drag force for all density ratios and that ratios among the forces and the Stokes drag decrease with rising particle density. The relation of the history force and the added mass force with the assumption of comparable accelerations, with the relaxation time of the order of τ results in

Hence it is of the order of the ratio of the diffusion length to the particle diameter, which might be of the order of one, if $\tau \propto \frac{d_p^2}{\nu}$.

It shows that the particle motion is only depending on the the hydrodynamical drag force which results from the actual difference between the fluid velocity $\mathbf{u}(\mathbf{x}_p)$ at particle position \mathbf{x}_p and the particle velocity \mathbf{v}_p . Due to these assumptions (3.18) becomes

$$\rho_p \frac{d\mathbf{v}_p}{dt} = \mathbf{f}_d. \quad (3.21)$$

This equation can be rewritten with help of the relaxation time τ_p , which physically represents the time scale over which the drag force decreases the particle's relative velocity to zero. At one point in the fluid with the proportionality

$$\frac{\rho_p v_p}{\tau_p^n} \propto f_d \quad (3.22)$$

it follows the equation for the relaxation time

$$\tau_p^n = \frac{4}{3} \frac{\rho_p}{\rho} \frac{d_p}{C_d^n} \frac{1}{v_r^n}, \quad (3.23)$$

where v_r^n is the relative velocity. With equation (??) and the Reynolds number in relation to the particles $Re_p^n = \frac{v_r^n d_p}{\nu}$, with the particle diameter d_p it follows:

$$\tau_p^n = \frac{\rho_p}{\rho} \frac{d_p^2}{18\nu} \frac{1}{1 + 0.15 Re_p^{0.687}}. \quad (3.24)$$

Then the equation of motion of the particles becomes

$$\rho_p \frac{d\mathbf{v}_p^n}{dt} = -\rho_p \frac{\mathbf{v}_p^n - \mathbf{u}^n}{\tau_p^n}. \quad (3.25)$$

Using the Stokes response time $\tau_{ps} = \frac{d_p^2}{18\nu} \frac{\rho_p}{\rho}$, that is the particle response time at resting fluid conditions, there finally exists another form of this equation

$$\rho_p \frac{d\mathbf{v}_p^n}{dt} = \rho_p - \rho_p \frac{\mathbf{v}_p^n - \mathbf{u}^n}{\tau_{ps}} f_d, \quad (3.26)$$

with $f_d = 1 + 0.15 Re_p^{0.687}$, which is a correction factor describing the flow conditions and not the drag force mentioned before.

The coupling rate ψ , that describes the transferred energy per time between particles and fluid, is defined as

$$\psi = \mathbf{F} \cdot \mathbf{u}. \quad (3.27)$$

Hence it results from the force per unitary volume exerted between the fluid and the particles and the main velocity of the fluid. It is an useful variable to describe how much the main fluid and the particles influence each other.

4 Numerical methods

Two numerical methods are discussed and their main differences pointed out in the following chapter, the DNS (direct numerical simulation) and the LES (large-eddy simulation). The basis of both are the Navier-Stokes equations as described above.

4.1 Discretization of the particle dynamics

To integrate the Lagrangian particle tracking equations, discussed above, a predictor-corrector scheme based on the trapezoidal rule for numerical integration

$$f(t + \delta t) \approx f(t) + \frac{\delta t}{2} \left[\frac{\partial f(t)}{\partial t} + \frac{\partial f(t + \delta t)}{\partial t} \right] \quad (4.1)$$

is used.

The first step is the prediction of the new particle position $\mathbf{x}_{p,n+1}$ using a Taylor expansion for a small time step δt

$$\mathbf{x}_{p,n+1} = \mathbf{x}_{p,n} + \delta t \mathbf{v}_{p,n} + \frac{1}{2} \delta t^2 \mathbf{a}_{p,n}, \quad (4.2)$$

with $\mathbf{a}_{p,n}$ particle acceleration.

To avoid filtering effects, we will set the fluid velocity $\mathbf{u}(\mathbf{x}_{p,n+1})$ at the particle position $\mathbf{x}_{p,n+1}$ equal to the nearest cell fluid velocity.

The updated velocity and acceleration are calculated as

$$\mathbf{v}_{p,n+1} = \frac{\mathbf{v}_{p,n} + \frac{1}{2} \delta t \left(\mathbf{a}_{p,n} + \frac{f_D}{\tau_p} \mathbf{u}(\mathbf{x}_{p,n+1}) \right)}{1 + \frac{1}{2} \frac{f_D}{\tau_p} \delta t}, \quad (4.3)$$

$$\mathbf{a}_{p,n+1} = \frac{\frac{f_D}{\tau_p} (\mathbf{u}(\mathbf{x}_{p,n+1}) - \mathbf{v}_{p,n}) - \frac{1}{2} \delta t \mathbf{a}_{p,n}}{1 + \frac{1}{2} \frac{f_D}{\tau_p} \delta t}. \quad (4.4)$$

The updated particle position must be corrected by an additional term according to the trapezoidal rule

$$\mathbf{x}_{p,n+1} = \mathbf{x}_{p,n} + \frac{1}{2} \delta t (\mathbf{v}_{p,n+1} + \mathbf{v}_{p,n}) + \frac{1}{12} \delta t^2 (\mathbf{a}_{p,n+1} - \mathbf{a}_{p,n}). \quad (4.5)$$

4.2 Direct numerical simulation

With DNS, the Navier-Stokes equations are solved completely. This provides a very accurate result, as all scales of motion are being resolved. Still it requires an immense level of computational resources which increases rapidly with the Reynolds number: $N^3 \propto 4.4 Re_L^{9/4} \approx 0.06 Re_\lambda^{9/2}$ with N size of the simulation. These computational resources were not available until the 1970s. With the LES,

as described below, the computational effort is 99.98 % less compared to DNS, which indeed is the fraction of the dissipative scale. This leaves 0.02 % of the flow, which is correlative with the fraction of the energy-containing larger-scale [16].

4.3 Large-eddy simulation

Due to the fact that DNS is effortful, LES was created to save time and resources. The energy containing larger-scale motion is completely resolved and the small effects of the smaller-scale motion are just modeled. Otherwise in DNS resolving the small dissipative scale would require most of the computational resources.

Simulating only the larger-scale motions is called filtering, which means that the smaller-scale motions, also known as fluctuation, are filtered out. $\int_{-\infty}^{\infty} \frac{\partial \mathbf{Q}}{\partial t} dV + \int_{\partial V} \bar{\mathbf{H}} \cdot \mathbf{n} dA = \mathbf{0}$ For further information on filter functions, the works of Pope [16] should be considered. To model the filtered smaller-scale motions usually a subgrid-scale (SGS) model is used. According to Hickel (2007) the interference between explicit SGS and the truncation error can be exploited, i.e. the truncation error can serve as model of the effects of unresolved scales, which is therefore an implicit SGS model. Thus we call it implicit LES (ILES) [18].

4.4 Applied simulation

The simulations were carried out using ZFS, the simulation tool developed and implemented at the Institute of Aerodynamics at RWTH Aachen University [4] [5]. The tool is capable of simulating finite-volume flows of compressible fluids. In the simulations, which results the reader can see at hand, the Mach-number Ma was set to 0.1 to simulate an almost incompressible fluid.

5 Results

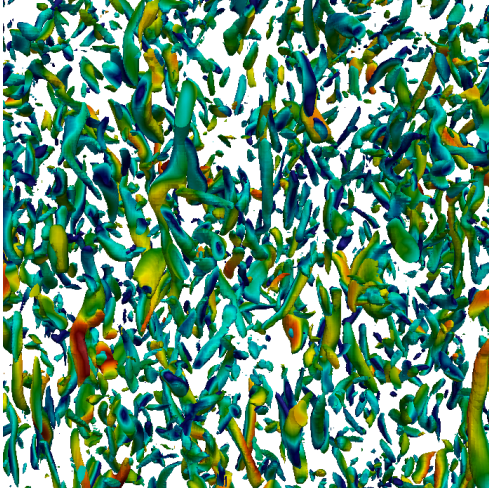


Figure 2: ???

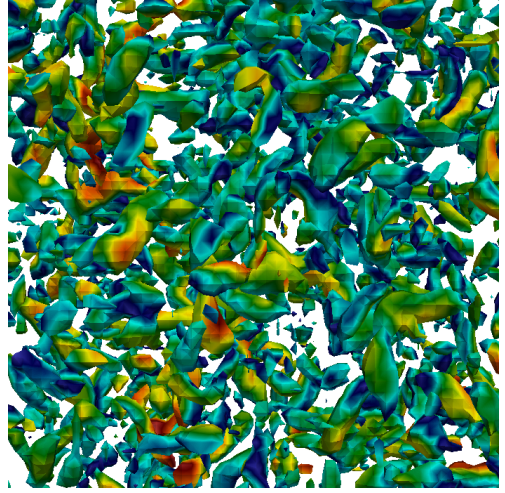


Figure 3: ??? Appendix A

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5.1 Computational point particles

The high number of point particles requires even more computational resources for the particle-laden simulations than for the single-phase simulations. The main idea to reduce this requirement is to create clusters of point particles. For this purpose the ratio λ_c of physical point particles N_p to computational point particles N_c is introduced ($\lambda_c = \frac{N_p}{N_c}$). To compensate this lack of particles, the coupling force is multiplied by λ_c , due to the λ_c -fold mass of the (cluster-)particles.

6 Results

In this section the properties and results of the simulations that have been carried out will be presented. Special emphasis will be put on the turbulent kinetic energy budgets and their use to interpret the findings.

6.1 Simulation properties & assumptions

In this section the results of the performed simulation will be pointed out with special emphasis on turbulent kinetic energy budgets and their influence on particle-laden turbulent flows.

All cases were simulated on a cubic domain and a cartesian grid using multiple grid refinement levels, 64^3 , 96^3 and 128^3 representing the LES hence the need to model smaller scales with SGS-models is present. Due to higher grid refinement in the 256^3 -case, these smaller scales can be directly simulated and therefore this simulation is a DNS.

The particle-free case was initialized using a seed-based random generator. At $t^* \approx 0.27$ a restart file is written out to initialize the subsequent simulation of the particle-laden isotropic turbulence. This file is then used to set up a second simulation including particle-load. Both these cases, DNSs of a single-phase and a particle-laden flow will be used as reference for analyzing other results. Single-phase simulations will in the following be referred to as sP, and particle-laden ones as PP. The PP-simulation was set up to match the volume fraction $\phi_v = 10^{-3}$ and mass fraction of $\phi_m = 1$. Both were constant in all simulations regarding particle clustering. The ratio of the densities for this set of simulations was set to $\frac{\rho_p}{\rho} = 1000$, the particle diameter was set to match $d_p \approx 0.6\eta$. The injection of a million spherical point-particles was carried out interpolating the velocities of the surrounding fluid. At the timestep of injection the Stokes response time τ_{ps} was 0.03497, the Prandtl number was 0.72 and Re_λ was 57.9757.

6.2 The turbulent kinetic energy budget

The change in kinetic energy of the fluid is in particle-laden decaying isotropic turbulence determined by the coupling rate Ψ , which describes the energy transfer between both fluid and particle phase, and the integral dissipation rate ϵ :

$$\frac{\partial E_k}{\partial t} = \Psi(t) - \epsilon(t). \quad (6.1)$$

As mentioned before, the flow field is considered nearly incompressible, therefore the equation for the viscous dissipation rate can be approximated by

$$\epsilon(t) \approx 2\mu \bar{\mathbf{S}} : \bar{\mathbf{S}}, \quad (6.2)$$

where $:$ denotes the inner tensor product. This rate can then be integrated over the full fluid domain Υ_f , leading to

$$\varepsilon(t) = \int_{\Upsilon_f} \epsilon(t) dV. \quad (6.3)$$

As the dissipation rate is always of positive value, it acts a sink for the fluid's turbulent kinetic energy. In difference to that, the coupling rate can serve either as source or sink, depending on the acceleration of the particles [8]. The coupling rate for point-particles is defined as

$$\Psi(t) = \sum_{p=1}^{N_p} \psi_p = - \sum_{p=1}^{N_p} \mathbf{F}_p \bullet \mathbf{v}_p, \quad (6.4)$$

describing the transfer of kinetic energy resulting from surface forces at each particle.

$$\frac{dE_{kB}}{dt} = -\Psi(t) \quad (6.5)$$

therefore is the rate of change for global kinetic energy of the particles. Additionally the particles change the fluid's rate of dissipation ε due to their volume forces. Schneiders et al. [13] therefore proposed a separation of dissipation rates for suspended heavy particles:

$$\frac{\partial E_k}{\partial t} = \Psi(t) - \bar{\varepsilon}(t) - \varepsilon'(t). \quad (6.6)$$

In this equation, $\bar{\varepsilon}(t)$ displays the dissipation rate of the flow field and $\varepsilon'(t)$ is the additional dissipation rate introduced by the injected point-particles. For $\rho_p \gg \rho$ and the point-particle approach the equation

$$\varepsilon' = \sum_{p=1}^{N_p} \mathbf{F}_p \bullet (\mathbf{u}_p - \mathbf{v}_p) \quad (6.7)$$

can be formulated using the velocity of the fluid \mathbf{u}_p at the particle position.

6.3 Simulation results

The first set of simulation was set up to investigate how different numbers of injected particles have an influence on the difference of turbulent kinetic energy of

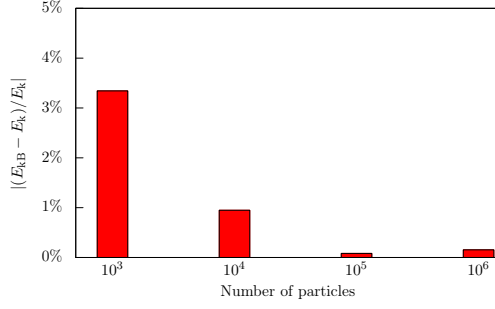


Figure 4: Initializing different numbers of point-particles on a 128^3 -grid: An unsatisfactory inaccuracy can be observed for small numbers of particles

both particle and fluid. As mathematical description of turbulent flows is based on statistics, small numbers of particles can lead to questionable results. An example for the deviation in kinetic energy for different numbers of particles can be found in Fig. 4. The normalized difference in kinetic energy E_{kB} of the particles and the flow itself E_k shows in this single simulation a correlation between particle number and accuracy in the simulation. Although this was just a single initialization of particles in a flow, it can be stated that simulations using only 10^3 or even up to 10^4 particles are not accurate enough for technical or scientific use of data. Simulations in other grid sizes show similar results.

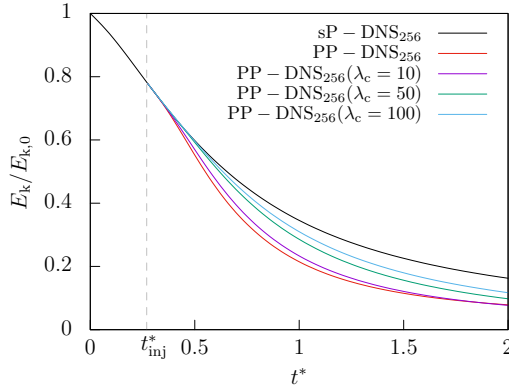


Figure 5: Kinetic energy E_k normalized by its initial value $E_{k,0}$ over time normalized by initial eddy turnover time. Shortly after the injection the PP-cases with separate from the sP-flow. The higher-clustered cases show less distinctions from the sP-case.

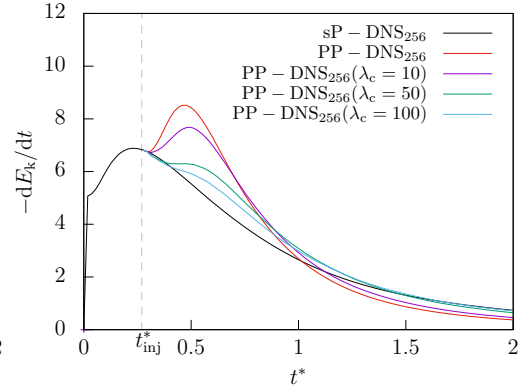


Figure 6: Change in kinetic energy over time normalized by initial eddy turnover time. After $t^* \approx 1$ the unclustered PP-case shows a lower decay rate in kinetic energy than the sP-case.

The second set of simulations conducted in this work is set up to investigate the influence of particle clustering on the accuracy of the simulations with the goal to minimize computational effort, while still achieving high quality results. For this purpose, the variable λ_c , which was established before, was implemented in

the code. The simulations were then set up with the overall same number of a million particles, altering just λ_c . It can be seen in Fig. 5 and 6 that the decay in kinetic energy from the injection point depends highly on the number of clustered particles. It can be stated that the flow's statistics for high λ_c converge approach the sP-case.

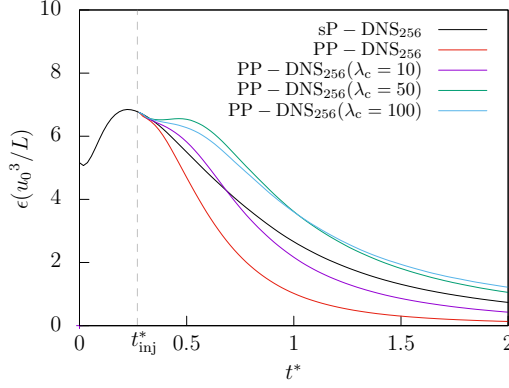


Figure 7: Normalized dissipation rate ϵ over time normalized by initial eddy turnover time. The unclustered PP-case shows a lower dissipation rate than the sP-case. Highly clustered PP-cases show a higher dissipation rate.

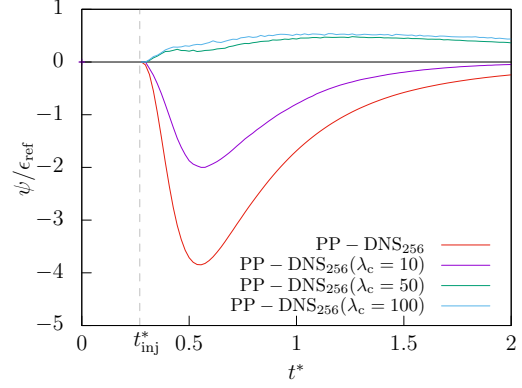


Figure 8: Normalized coupling rate Ψ over time normalized by initial eddy turnover time. The PP-case without clustering show the lowest coupling rate. Both highly-clustered PP-cases show physically questionable behavior.

Fig. 7 and 8 can be interpreted using the turbulent kinetic energy budget introduced earlier. The PP-case without clustering shows a lower dissipation rate than the sP-case, both highly-clustered cases show a higher rate over the monitored period of time. The coupling rate is negative all this time for the PP-cases with λ_c and positive for the higher-clustered cases. Both these effects form the change in kinetic energy in particle-laden isotropic flow found in Fig. 6.

Being very similar in the time shortly after the injection, the flow statistics diverge for different λ_c rapidly. The variables of these simulations one turnover time after the injection can be found in table 1. 9. Plot fuer kinetische Energie der Partikel und Vergleichsplot DNS i-1 LES (Coupling Rate)

12. Computational savings

λ_c	$\frac{m_c}{m_{V,\text{cell}}}$	$\epsilon \frac{u_0^3}{L}$	$\frac{\lambda}{L}$	$\frac{\eta}{L}$	Re_λ
1	16.78	0.97	0.039	0.0032	38.73
10	167.78	2.08	0.028	0.0026	28.66
50	838.86	3.51	0.024	0.0023	27.31
100	1677.72	3.51	0.025	0.0023	29.42

Table 1: Variables of the first set of simulations one turnover time after injection for the 256^3 -case

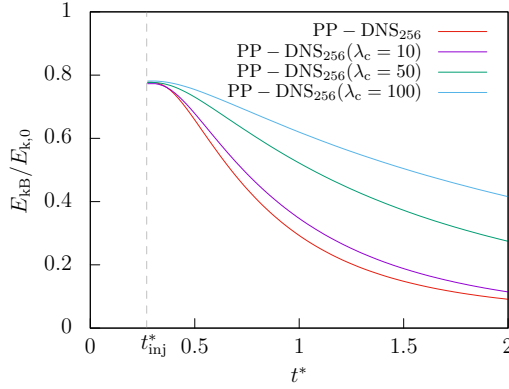


Figure 9: Kinetic energy of the particles E_{kB} normalized by initial kinetic energy. The PP-case without clustering shows the biggest decay in kinetic energy.

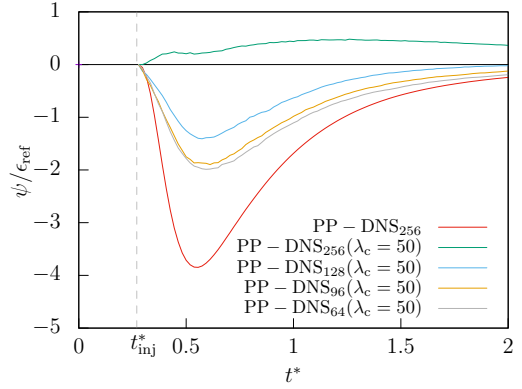


Figure 10: Coupling Rate Ψ with constant λ_c for various resolutions. A relation of resolution and accuracy can be observed.

7 Conclusion and outlook

In this work, two sets of simulations were carried out to evaluate a method for lowering computational effort.

The results presented in this work show, that for a sufficiently exact simulation, particle clustering has to be treated with caution. Depending on the application maybe small numbers of clustered particles could be used, but the savings in computing time would not compensate the loss in accuracy. This is the case in particular high numbers of clustered particles at which the results get highly inaccurate. Maybe investigations on smaller numbers of clustered particles could follow, the space between 2 and 10 could be closed by simulations in the future to determine the real border for inaccurate results.

Further investigation is needed regarding the second set of simulation, in which the goal was to find out which number of particles is necessary to get accurate results for the turbulent kinetic energy of the particles. The deviation of the av-

eraged kinetic energy of the particles and the fluid should be zero. This difference converges to zero with high numbers of particles. The distribution of the experiment's outcomes should match the well-known normal distribution, which leads to analyzing the standard deviation. Concluding, further simulations should be carried out until a sufficient standard deviation can be computed, from which an assumption about the accuracy of the initialization could be made.

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This work was supervised by Konstantin Fröhlich, we would like to express our gratitude. Thank you for the chance of learning about turbulent flows and simulations, the advice and the deep insights in scientific work. We also appreciate the chance of writing this work at the Institute of Aerodynamics of the RWTH Aachen University.

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9 Appendix A

Creating of pictures showing tubular structures

The pictures used in the Introduction were generated using ParaView, an open-source-software developed by a joint-venture of Kitware and the Los Alamos National Laboratory. More information about the software can be found at www.paraview.org. To show the tubular structures in a turbulent flow, two filters were used: One was the AIALambda2Criterion1-Filter and the other one was the ISOVolume1-Filter. These filters were then set to visualize the velocity of the flow colored by magnitude. To diversify the different velocity-magnitudes, a rainbow-colorscheme was used.