jj Terms for Non equivalent Electrons in d^xp^ys^z Configurations

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Abstract. jj coupling is predominant in heavier atoms where spin orbit interactions are important than electrostatic interactions. In this manuscript jj coupled terms derived for non equivalent electrons in $d^x p^y s^z$ (x = 1-2, y & z = 0-1) configurations i.e. $d^2 p^1 s^1$, $d^1 p^1 s^1$, $d^1 p^1$ and $d^2 s^1$ configurations, the obtained jj terms are $[(5/2, 5/2, 3/2, 1/2), (5/2, 5/2, 1/2, 1/2), (3/2, 3/2, 1/2), (3/2, 3/2, 1/2), (5/2, 5/2, 3/2, 1/2), (5/2, 5/2, 3/2, 1/2)] for <math>d^2 p^1 s^1$, [(5/2, 3/2, 1/2), (5/2, 1/2), (3/2, 3/2, 1/2)] for $d^1 p^1 s^1$, [(5/2, 3/2), (5/2, 1/2), (3/2, 3/2), (3/2, 1/2)] for $d^1 p^1$ and [(5/2, 5/2, 1/2), (5/2, 3/2, 1/2), (3/2, 3/2, 1/2)] d² s¹ configurations and the ground state terms determined for these configurations are (3/2, 3/2, 1/2), (3/2, 1/2, 1/2), (3/2, 1/2, 1/2) and (3/2, 3/2, 1/2) respectively.

Keywords: Angular momentum, jj coupling, L-S coupling and spin-orbit interaction

INTRODUCTION

LS terms are significant in lower elements which gradually change to jj coupling in going from lighter to heavy atom due to increase nuclear charge (Gauerke & Campbell, 1994). LS terms for equivalent or nonequivalent electrons are derived by different methods i.e. Vector model (Lande, 1921), Quantum mechanical method (Russell & Saunders, 1925), Ford method (Ford, 1972), Hyde method (Hyde, 1975) , Spin factoring method (McDaniel, 1977), Numerical algorithm method (Kiremire, 1987), Slater graphics (Slater, 1960), Partitioning total spin method (Guofan & Ellzey, 1987), Group representation method (Chen, 1989), Group theoretical method (Wybourne, 1966; Judd, 1967) Generating functions derived via group theory method (Curl & Kilpatrik, 1960), Partial term method (Kiremire, 1990), Partitioning technique (Olson, 2011). The microstate building through electronic arrangement method has been used to generating the spectroscopic LS terms for equivalent electrons of f³ and f⁴ configurations (Meena et al., 2011a; 2011b), and for nonequivalent electrons of (n-1) f³nd¹, (n-1) f²nd¹ and d² p¹ s¹ configurations (Meena et al., 2012; Meena et al., 2013).

jj terms can also be determine by using different methods which are described by (Rubio & Perez, 1986), (Tuttle, 1967), (Haigh, 1990), (Gauerke & Campbell, 1994), (Campbell, 1998), (Novak, 1999), (Orofino & Faria, 2010), (Richtmyer et al, 1969) and (Meena et al., 2015). Equivalent electrons have same values of n and l, the electrostatic interaction is expected to be larger than spin-orbit interaction and L-S coupling is favoured and for nonequivalent, j-j coupling is important. In this manuscript the spectroscopic jj coupled terms for non equivalent electrons of $d^x p^y s^z$ configurations (x = 1 - 2, y & z = 0 - 1) were determined and correlated with LS terms (for $d^1 p^1$ and $d^2 s^1$ configuration).

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Methodology

The microstates were built up by arranging electrons with different possible j values for non equivalent electrons of $d^xp^ys^z$ configurations (x= 1-2, y & z = 0-1). Total microstates calculated for d^2 p^1 s^1 , d^1 p^1 s^1 , d^1 p^1 and d^2 s^1 configurations are 540, 120, 60 and 90 respectively. Notations for the jj terms designated by the j's are $[(j_1)^a(j_2)^b(j_3)^c...]$ (Tuttle, 1967 & 1980; Orofino & Faria, 2010) and $[(j_1, j_2)_J]$ (Haigh, 1990). The possible jj terms for non equivalent electrons of $d^xp^ys^z$ configurations (x= 1-2, y & z = 0-1) are $[(5/2, 5/2, 3/2, 1/2), (5/2, 5/2, 1/2, 1/2), (3/2, 3/2, 1/2), (3/2, 3/2, 1/2), (5/2, 3/2, 1/2), (5/2, 3/2, 1/2), (5/2, 3/2, 1/2), (5/2, 3/2, 1/2), (5/2, 3/2, 1/2), (3/2, 3/2, 1/2), (3/2, 3/2, 1/2), (3/2, 3/2, 1/2), (3/2, 3/2, 1/2)] for <math>d^1$ p^1 s^1 , [(5/2, 3/2), [(5/2, 1/2), [(3/2, 3/2), (3/2, 1/2)] for d^1 p^1 and [(5/2, 5/2, 1/2), (5/2, 3/2, 1/2)] for d^2 s^1 configuration.

Microstates for jj Terms for d² p¹ s¹ Configuration

The microstate tables for each term is drawn by arranging four electrons and the M_J values for all microstates are determined. The largest M_J value for each term represents a value of J level for term (Table A1). Number of microstates for a particular term of the of the form $[(l_{\ell-1/2})^i(l_{\ell+1/2})^{n-i}]$ or (j_1, j_2, j_3, j_4) for each sub set of equivalent electrons is given by $(2\ell)!(2\ell+2)!$ $i!(2\ell-i)!(n-i)!(2\ell+2+i-n)!$

J levels for jj terms for d² p¹ s¹ Configuration

J level for jj term are obtained by removing microstates associated with that J level starting from the maximum M_J value in the microstate tables and followed for next levels also. For example, when the 13 microstates associated with maximum M_J =6 for the jj coupled term (5/2, 5/2, 3/2, 1/2) are eliminated from Table A2 results in J=6 level and maximum M_J level remain is 5 that yield another J=5 level for this term when 22 microstat associated with this are eliminated, and further elimination of 27, 28, 20, 9 and 1 microstates associated with M_J 4, 3, 2, 1 and 0, give 4, 3, 2, 1 and 0 J levels for this term. By applying the same procedure to other terms as illustrated in Table A3 for (5/2, 5/2, 1/2, 1/2) term, Table A4 for (3/2, 3/2, 3/2, 1/2) term, Table A5 for (3/2, 3/2, 1/2) term, Table A6 for (5/2, 3/2, 3/2, 1/2) term and Table A7 for (5/2, 3/2, 1/2, 1/2) term.

Number of microstates for jj terms of d² s¹ configuration

$$\frac{1. \text{ Term } (5/2, 5/2, 1/2) \text{ or } [(d_{5/2})^2 (s_{1/2})^1]}{(2x2)!(2x2+2)!} \times \frac{(0x2)!(0x2+2)!}{0!(2x2-0)!(2-0)!(2-0)!(2x2+2+0-2)!} \times \frac{(0x2)!(0x2+2)!}{0!(0x2-0)!(1-0)!(0x2+2+0-1)!} = 30$$
2. Term $(5/2, 3/2, 1/2)$ or $[(d_{5/2})^1 (d_{3/2})^1 (s_{1/2})^1]$

$$\frac{(2x2)!(2x2+2)!}{1!(2x2-1)!(2-1)!(2x2+2+1-2)!} \times \frac{(0x2)!(0x2+2)!}{0!(0x2-0)!(1-0)!(0x2+2+0-1)!} = 48$$
3. Term $(3/2, 3/2, 1/2)$ or $[(d_{3/2})^2 (s_{1/2})^1]$

$$\frac{(2x2)!(2x2+2)!}{2!(2x2-2)!(2-2)!(2x2+2+2-2)!} \times \frac{(0x2)!(0x2+2)!}{0!(0x2-0)!(1-0)!(0x2+2+0-1)!} = 12$$

Running head: jj Terms for Non equivalent Electrons in dxpysz Configurations

Number of microstates for jj terms of d¹ p¹ s¹ configuration

$$\frac{1. \text{ Term } (5/2, 3/2, 1/2) \text{ or } [(d_{5/2})^1 (p_{3/2})^1 (s_{1/2})^1]}{0!(2x2)!(2x2+2)!} x \frac{(1x2)!(1x2+2)!}{0!(1x2-0)!(1-0)!(1x2+2+0-1)!} x \frac{(0x2)!(0x2+2)!}{0!(0x2-0)!(1-0)!(0x2+2+0-1)!} = 48$$

$$2. \text{ Term } (3/2, 3/2, 1/2) \text{ or } [(d_{3/2})^1 (p_{3/2})^1 (s_{1/2})^1]$$

$$\frac{(2x2)!(2x2+2)!}{1!(2x2-1)!(1-1)!(2x2+2+1-1)!} x \frac{(1x2)!(1x2+2)!}{0!(1x2-0)!(1-0)!(1x2+2+0-1)!} x \frac{(0x2)!(0x2+2)!}{0!(0x2-0)!(1-0)!(0x2+2+0-1)!} = 32$$

$$3. \text{ Term } (5/2, 1/2, 1/2) \text{ or } [(d_{5/2})^1 (p_{1/2})^1 (s_{1/2})^1]$$

$$\frac{(2x2)!(2x2+2)!}{0!(2x2-0)!(1-0)!(2x2+2+0-1)!} x \frac{(1x2)!(1x2+2)!}{1!(1x2-1)!(1-1)!(1x2+2+1-1)!} x \frac{(0x2)!(0x2+2)!}{0!(0x2-0)!(1-0)!(0x2+2+0-1)!} = 24$$

$$4. \text{ Term } (3/2, 1/2, 1/2) \text{ or } [(d_{3/2})^1 (p_{1/2})^1 (s_{1/2})^1]$$

$$\frac{(2x2)!(2x2+2)!}{1!(2x2-1)!(1-1)!(2x2+2+1-1)!} x \frac{(1x2)!(1x2+2)!}{0!(0x2-0)!(1-0)!(0x2+2+0-1)!} = 16$$

Number of microstates for jj terms of d¹ p¹ configuration

$$\frac{1. \text{ Term } (5/2, 3/2) \text{ or } [(d_{5/2})^1 (p_{3/2})^1]}{(2x2)!(2x2+2)!} \times \frac{(1x2)!(1x2+2)!}{0!(1x2-0)!(1-0)!(1x2+2+0-1)!} = 24$$

$$2. \text{ Term } (3/2, 3/2) \text{ or } [(d_{3/2})^1 (p_{3/2})^1]$$

$$\frac{(2x2)!(2x2+2)!}{1!(2x2-1)!(1-1)!(2x2+2+1-1)!} \times \frac{(1x2)!(1x2+2)!}{0!(1x2-0)!(1-0)!(1x2+2+0-1)!} = 16$$

$$3. \text{ Term } (5/2, 1/2) \text{ or } [(d_{5/2})^1 (p_{1/2})^1]$$

$$\frac{(2x2)!(2x2+2)!}{0!(2x2-0)!(1-0)!(2x2+2+0-1)!} \times \frac{(1x2)!(1x2+2)!}{1!(1x2-1)!(1-1)!(1x2+2+1-1)!} = 12$$

$$4. \text{ Term } (3/2, 1/2) \text{ or } [(d_{3/2})^1 (p_{1/2})^1]$$

$$\frac{(2x2)!(2x2+2)!}{1!(2x2-1)!(1-1)!(2x2+2+1-1)!} \times \frac{(1x2)!(1x2+2)!}{1!(1x2-1)!(1-1)!(1x2+2+1-1)!} = 8$$

By applying same method as used for $d^2 p^1 s^1$ configuration J levels are determined which are $[(5/2, 3/2, 1/2)_{9/2, 7/2(2), 5/2(2), 3/2(2), 1/2}]$, $[(5/2, 1/2, 1/2)_{7/2, 5/2(2), 3/2}]$, $[(3/2, 3/2, 1/2)_{7/2, 5/2(2), 3/2(2), 1/2}]$ and $[(3/2, 1/2, 1/2)_{5/2, 3/2(2), 1/2}]$ for $d^1 p^1 s^1$ configuration, $[(5/2, 3/2)_{4, 3, 2, 1}]$, $[(5/2, 1/2)_{3, 2}]$, $[(3/2, 3/2)_{3, 2, 1, 0}]$ and $[(3/2, 1/2)_{2, 1}]$ for $d^1 p^1$ configuration and $[(5/2, 5/2, 1/2)_{9/2, 7/2, 5/2, 3/2, 1/2}]$, $[(5/2, 3/2, 1/2)_{9/2, 7/2(2), 5/2(2), 3/2(2), 1/2}]$ and $[(3/2, 3/2, 1/2)_{5/2, 3/2, 1/2}]$ for $d^2 s^1$ configuration.

RESULTS AND DISCUSSION

jj coupled spectroscopic terms obtained for $d^x p^y s^z$ configurations (x= 1-2, y & z = 0-1) are $[\{(5/2, 5/2, 3/2, 1/2)_{6, 5(2), 4(3), 3(4), 2(4), 1(3), 0}\}$, $\{(5/2, 5/2, 1/2, 1/2)_{5, 4(2), 3(2), 2(2), 1(2), 0}\}$, $\{(3/2, 3/2, 3/2, 1/2)_{4, 3(2), 2(3), 1(3), 0}\}$, $\{(3/2, 3/2, 1/2, 1/2)_{3, 2(2), 1(2), 0}\}$, $\{(5/2, 3/2, 3/2, 1/2)_{6, 5(3), 4(5), 3(7), 2(7), 1(5), 0(2)}\}$, $\{(5/2, 3/2, 1/2, 1/2)_{5, 4(3), 3(4), 2(4), 1(3), 0}\}$ for $d^1 p^1 s^1$ configuration, $[\{(5/2, 3/2, 1/2)_{9/2, 7/2(2), 5/2(2), 3/2(2), 1/2}\}$, $\{(5/2, 1/2, 1/2)_{7/2, 5/2(2), 3/2}\}$, $\{(3/2, 3/2, 1/2)_{7/2, 5/2(2), 3/2(2), 1/2}\}$, $\{(3/2, 3/2, 1/2)_{9/2, 7/2, 5/2, 3/2, 1/2}\}$ for $d^2 p^1 s^1$ configuration, $[\{(5/2, 3/2)_{4, 3, 2, 1}\}, \{(5/2, 1/2)_{3, 2}\}, \{(3/2, 3/2)_{3, 2, 1, 0}\}, \{(3/2, 1/2)_{2, 1}\}\}$ for $d^1 p^1$ configuration and $[\{(5/2, 5/2, 1/2)_{9/2, 7/2, 5/2, 3/2, 1/2}\}, \{(5/2, 3/2, 1/2)_{5/2, 3/2, 1/2}\}\}$, $\{(5/2, 3/2, 1/2)_{9/2, 7/2, 5/2, 3/2, 1/2}\}$, $\{(5/2, 3/2, 1/2)_{5/2, 3/2, 1/2}\}\}$ for $d^2 p^1 s^1$ configuration and $[\{(5/2, 5/2, 1/2)_{9/2, 7/2, 5/2, 3/2, 1/2}\}, \{(5/2, 3/2, 1/2)_{5/2, 3/2, 1/2}\}\}$ for $d^2 p^1 s^1$ configuration.

And the ground state jj coupled terms determined for these configurations are $[(3/2, 3/2, 1/2, 1/2)_{3, 2(2), 1(2), 0}]$, $[(3/2, 1/2)_{5/2, 3/2(2), 1/2}]$, $[(3/2, 1/2)_{2, 1}]$ and $[(3/2, 3/2, 1/2)_{5/2, 3/2, 1/2}]$ respectively. In correlation level diagram the L-S and the j-j levels for d^1 p^1 and d^2 s^1 configurations are shown (Figure 1 and Figure 2). Total numbers of final states are same, but their relative energies are different.

CONCLUSSION

Here a simple and systematic method is described to obtain the jj coupled spectroscopic terms for nonequivalent electrons of $d^xp^ys^z$ configurations (x= 1-2, y & z= 0-1). For d^2 p¹ s¹, d^1 p¹ and d^2 s¹ configurations, this procedure will make jj coupled terms more popular in chemistry and also helpful to investigate the atomic and electronic spectra of nonequivalent electron containing atoms or free ions.

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Appendix A

Table A1 $\label{eq:constates} \mbox{Number of Microstates for each jj Coupled Term for d^2 p^1 s^1 Configuration}$

E1 j ₁	E2 j ₂	E3 j ₃	E4 j ₄	jj coupled terms	Microstates	M _J values			
5/2	5/2	3/2	1/2	(5/2, 5/2, 3/2, 1/2)	120	6 to -6			
5/2	5/2	1/2	1/2	(5/2, 5/2, 1/2, 1/2)	60	5 to -5			
3/2	3/2	3/2	1/2	(3/2, 3/2, 3/2, 1/2)	48	4 to -4			
3/2	3/2	1/2	1/2	(3/2, 3/2, 1/2, 1/2)	24	3 to -3			
5/2	3/2	3/2	1/2	(5/2, 3/2, 3/2, 1/2)	192	6 to -6			
5/2	3/2	1/2	1/2	(5/2, 3/2, 1/2, 1/2)	96	5 to -5			
	Total number of microstates for d ² p ¹ s ¹								
	configuration-540								

Table A2 $\label{eq:microstates} \mbox{Microstates and their Removal for J Levels for (5/2, 5/2, 3/2, 1/2) Term for d^2 p^1 s^1 Configuration}$

M _J	No.	MS after	MS after	MS after	MS after	MS after	MS after
	of	removing	removing	removing	removing	removing	removing
	MS	J=6 level	J=5(2)	J=4(3)	J=3(4)	J=2(4)	J=1(3)
			levels	levels	levels	levels	levels
6	1	-	-	-	-	-	-
5	3	2	-	-	-	-	-
4	6	5	3	-	-	-	-
3	10	9	7	4	-	-	-
2	14	13	11	8	4	-	-
1	17	16	14	11	7	3	-
0	18	17	15	12	8	4	1
-1	17	16	14	11	7	3	-
-2	14	13	11	8	4	-	-
-3	10	9	7	4	-	-	-
-4	6	5	3	-	-	-	-
-5	3	2	-	-	-	-	-
-6	1	-	-	-	-	-	-
To	120	107	85	58	30	10	1
tal							

Table A3 $\label{eq:microstates} \mbox{Microstates and their Removal for J Levels for (5/2, 5/2, 1/2, 1/2) Term for d^2 p^1 s^1 Configuration}$

$M_{\rm J}$	No.	MS after	MS after	MS after	MS after	MS after
	of	removing	removing	removing	removing	removing
	MS	J=5 level	J=4(2) levels	J=3(2) levels	J=2(2) levels	J=1(2) levels
5	1	-	-	-	-	-
4	3	2	-	-	-	-
3	5	4	2	-	-	-
2	7	6	4	2	-	-
1	9	8	6	4	2	-
0	10	9	7	5	3	1
-1	9	8	6	4	2	-
-2	7	6	4	2	-	-
-3	5	4	2	-	-	-
-4	3	2	-	-	-	-
-5	1	-	-	-	-	-
To	60	49	31	17	7	1
tal						

Table A4 $\label{eq:microstates} \mbox{Microstates and their Removal for J Levels for (3/2, 3/2, 3/2, 1/2) Term for d^2 p^1 s^1 Configuration}$

Мл	No. of MS	MS after removing J=4	MS after removing J=3(2)	MS after removing J=2(3)	MS after removing J=1(3)
		level	levels	levels	levels
4	1	-	-	-	-
3	3	2	-	-	-
2	6	5	3	-	-
1	9	8	6	3	-
0	10	9	7	4	1
-1	9	8	6	3	-
-2	6	5	3	-	-
-3	3	2	-	-	-
-4	1	-	-	-	-
Tot	48	39	25	10	1
al					

Table A5 $\label{eq:microstates} \mbox{Microstates and their Removal for J Levels for (3/2, 3/2, 1/2, 1/2) Term for d^2 p^1 s^1 Configuration}$

$M_{\rm J}$	No. of	MS after removing	MS after removing	MS after removing
	M. S.	J=3 level	J=2(2) levels	J=1(2) levels
3	1	-	-	-
2	3	2	-	-
1	5	4	2	-
0	6	5	3	1
-1	5	4	2	-
-2	3	2	-	-
-3	1	-	-	-
To	24	17	7	1
tal				

Table A6

Microstates and their Removal for J Levels for (5/2, 3/2, 3/2, 1/2) Term for d^2 p^1 s^1 Configuration

$M_{\rm J}$	No.	MS after	MS after	MS after	MS after	MS after	MS after
	of	removing	removing	removing	removing	removing	removing
	MS	J=6 level	J=5(3)	J=4(5)	J=3(7)	J=2(7)	J=1(5)
			levels	levels	levels	levels	levels
6	1	-	-	-	-	-	-
5	4	3	-	-	-	-	-
4	9	8	5	-	-	-	-
3	16	15	12	7	-	-	-
2	23	22	19	14	7	-	-
1	28	27	24	19	12	5	-
0	30	29	26	21	14	7	2
-1	28	27	24	19	12	5	-
-2	23	22	19	14	7	-	-
-3	16	15	12	7	-	-	-
-4	9	8	5	-	-	-	-
-5	4	3	-	-	-	-	-
-6	1	-	-	-	-	-	-
Tot	192	179	146	101	52	17	2
al							

Table A7 $\label{eq:microstates} \mbox{Microstates and their Removal for J Levels for (5/2, 3/2, 1/2, 1/2) Term for d^2 p^1 s^1 Configuration}$

$M_{\rm J}$	No.	MS after	MS after	MS after	MS after	MS after
	of	removing	removing	removing	removing	removing
	MS	J=5 level	J=4(3)	J=3(4)	J=2(4)	J=1(3)
			levels	levels	levels	levels
5	1	-	-	-	-	-
4	4	3	-	-	-	-
3	8	7	4	-	-	-
2	12	11	8	4	-	-
1	15	14	11	7	3	-
0	16	15	12	8	4	1
-1	15	14	11	7	3	-
-2	12	11	8	4	-	-
-3	8	7	4	-	-	-
-4	4	3	-	-	-	-
-5	1	-	-	-	-	-
Total	96	85	58	30	10	1

Appendix B

Figure B1

Correlation Diagram for LS and jj Coupling Schemes for Levels for d¹ p¹ Configuration



