# Part IA — Variational Principles Example Sheet 2

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$$F[x + \delta x] - F[x] = \int_{t_1}^{t_2} f(x + \delta x, \dot{x} + \delta \dot{x}, \ddot{x} + \delta \ddot{x}, t) dt - \int_{t_1}^{t_2} f(t, x, \dot{x}, \ddot{x}) dt$$
$$= \int_{t_1}^{t_2} \left\{ \delta x \frac{\partial f}{\partial x} + (\delta \dot{x}) \frac{\partial f}{\partial \dot{x}} + (\delta \ddot{x}) \frac{\partial f}{\partial \ddot{x}} \right\} dt + O(t^2)$$

Discarding the (small )terms of  $O(t^2)$ , we call the first order variation  $\delta F[x]$  and integrating by parts (twice), we have

$$\begin{split} \delta F[x] &= \int_{t_1}^{t_2} \left\{ \delta x \left[ \frac{\partial f}{\partial x} - \frac{\mathrm{d}}{\mathrm{d}t} \left( \frac{\partial f}{\partial \dot{x}} \right) \right] - (\delta \dot{x}) \frac{\mathrm{d}}{\mathrm{d}t} \left( \frac{\partial f}{\partial \ddot{x}} \right) \right\} \, \mathrm{d}t + \left[ \delta x \frac{\partial f}{\partial \dot{x}} + (\delta \dot{x}) \frac{\partial f}{\partial \ddot{x}} \right]_{t_1}^{t_2} \\ &= \int_{t_1}^{t_2} \left\{ \delta x \left[ \frac{\partial f}{\partial x} - \frac{\mathrm{d}}{\mathrm{d}t} \left( \frac{\partial f}{\partial \dot{x}} \right) + \frac{\mathrm{d}^2}{\mathrm{d}t^2} \left( \frac{\partial f}{\partial \ddot{x}} \right) \right] \right\} \, \mathrm{d}t \\ &+ \left[ \delta x \left\{ \frac{\partial f}{\partial \dot{x}} - \frac{\mathrm{d}}{\mathrm{d}t} \left( \frac{\partial f}{\partial \ddot{x}} \right) \right\} + (\delta \dot{x}) \frac{\partial f}{\partial \ddot{x}} \right]_{t_1}^{t_2} \end{split}$$

We have fixed end boundary conditions, so  $\delta x(t_1) = \delta x(t_2) = 0$  and also  $\delta \dot{x}(t_1) = \delta \dot{x}(t_2) = 0$ . Thus the boundary term is zero and we can write  $\delta F[x]$  in the form

$$\delta F[x] = \int_{t_1}^{t_2} \left\{ \delta x(t) \frac{\delta F[x]}{\delta x(t)} \right\} dt$$

where the functional derivative  $\frac{\delta F[x]}{\delta x(t)}$  is defined as

$$\frac{\delta F[x]}{\delta x(t)} := \frac{\partial f}{\partial x} - \frac{\mathrm{d}}{\mathrm{d}t} \left( \frac{\partial f}{\partial \dot{x}} \right) + \frac{\mathrm{d}^2}{\mathrm{d}t^2} \left( \frac{\partial f}{\partial \ddot{x}} \right)$$

The functional F is stationary when its functional derivative is zero ( assuming that this derivative is defined on  $(t_1, t_2)$ ) and the condition for this to be true is the following Euler-Lagrange equation:

$$\frac{\partial f}{\partial x} - \frac{\mathrm{d}}{\mathrm{d}t} \left( \frac{\partial f}{\partial \dot{x}} \right) + \frac{\mathrm{d}^2}{\mathrm{d}t^2} \left( \frac{\partial f}{\partial \ddot{x}} \right) = 0, \qquad t_1 < t < t_2$$

Given the functional

$$L[x] = \int_{1}^{2} t^{4} [\ddot{x}(t)]^{2} dt$$

In this case,  $f=t^4[\ddot{x}(t)]^2$ , so  $\frac{\partial f}{\partial x}=\frac{\partial f}{\partial \dot{x}}=0$  and the EL equation can be immediately twice integrated to give

$$\frac{\partial}{\partial \ddot{x}} \left[ t^4 [\ddot{x}(t)]^2 \right] = At + B$$

for some constants A and B. ie.

$$2t^4(\ddot{x})(t) = At + B$$

$$\Rightarrow \ddot{x}(t) = At^{-3} + Bt^{-4} \qquad t \neq 0$$

$$\Rightarrow \dot{x}(t) = A't^{-2} + B't^{-3} + C$$

Using b.c.s we have

$$-2 = A' + B' + C$$
 
$$-\frac{1}{4} = \frac{1}{4}A' + \frac{1}{8}B' + C \Rightarrow -2 = 2A' + B' + 8C$$

Subtracting immediately yields A' = 7C:

$$-2 = 8C + B'$$

$$-2 = 22C + B'$$

Therefore C' = 0, B' = -2, and  $\dot{x}(n) = -2t^{-3}$ . Integrating, we get

$$x(t) = t^{-2} + D$$

b.c.s  $\Rightarrow D=0$ , so  $x(t)=\frac{1}{t^2}$ . This function is a global minimum? Hard to see.

Our aim is to maximise A[x,y] subject to the constraint P[x,y]=L, where L is the fixed length and  $P[x,y]=\int_0^{2\pi}\sqrt{(x')^2+(y')^2}\,\mathrm{d}\theta$ 

Using a Lagrange multiplier  $\lambda$  to impose this, we seek to maximize the functional

$$\phi_{\lambda}[x,y] = A[x,y] - \lambda(P[y] - L)$$

$$= \int_{0}^{2\pi} \underbrace{\frac{1}{2}(xy' - yx') - \lambda\sqrt{(x')^{2} + (y')^{2}}}_{=f_{\lambda}(\mathbf{x},\mathbf{x}')} d\theta + \lambda L$$

where  $\mathbf{x} = (x, y)$ . The boundary conditions fix x, x', y, y' at the endpoints, so the functional is stationary for solutions of the E-L equation.  $f_{\lambda}(\mathbf{x}, \mathbf{x}')$  has no explicit  $\theta$  dependence, so considering the y E-L equation we have that:

$$f_{\lambda}(y, y') - y' \frac{\partial f_{\lambda}}{\partial y'} = \text{constant}$$

$$\begin{split} &\Rightarrow f - y' \left( \frac{1}{2} x - \frac{\lambda y'}{\sqrt{(x')^2 + (y')^2}} \right) = \text{constant} \\ &\Rightarrow -\frac{1}{2} y x' - \lambda \frac{((x')^2 + (y')^2)}{\sqrt{(x')^2 + (y')^2}} + \lambda \frac{(y')^2}{\sqrt{(x')^2 + (y')^2}} = \text{constant} \\ &\Rightarrow -\frac{1}{2} y x' - \lambda \frac{(x')^2}{\sqrt{(x')^2 + (y')^2}} = C \end{split}$$

for some constant C. Similarly considering  $f_{\lambda}(x,x')$  we have

$$-\frac{1}{2}xy' - \lambda \frac{(y')^2}{\sqrt{(x')^2 + (y')^2}} = D$$

for some constant D. Squaring and adding,

$$(C + \frac{1}{2}yx')^2 + (D + \frac{1}{2}xy')^2 = \lambda^2$$

Feel like the x' and y' shouldn't be here, and it would be a circle?

Using Lagrange multiplier  $\lambda$ , wish to minimize

$$\Phi_{\lambda}[\psi] = I[\psi] - \lambda \left[ \int_{-\infty}^{\infty} \psi^2 \, dx = 1 \right]$$
$$= \int_{-\infty}^{\infty} \underbrace{(\psi')^2 + (x^2 - \lambda)\psi^2}_{f_{\lambda}(\psi,\psi';x)} \, dx + \lambda$$

Normalisation condition, can assume  $\phi = 0$  at endpoints, so the functional is stationary for solutions of the E-L equation. Euler-Lagrange equations imply:

$$2(x^{2} - \lambda)\psi - \frac{\mathrm{d}}{\mathrm{d}x} [2\psi'] = 0$$
$$\Rightarrow \psi'' + (x^{2} - \lambda)\psi = 0$$

Note that

$$I[\psi] = \int_{-\infty}^{\infty} (\psi' + x^2 \psi^2)^2 - 2x\psi\psi' \, \mathrm{d}x$$

$$= \int_{-\infty}^{\infty} (\psi' + x^2 \psi^2)^2 \, \mathrm{d}x - \int_{-\infty}^{\infty} x \frac{\mathrm{d}}{\mathrm{d}x} [\psi^2] \, \mathrm{d}x$$

$$= \int_{-\infty}^{\infty} (\psi' + x^2 \psi^2)^2 \, \mathrm{d}x - \left[x\psi^2\right]_{-\infty}^{\infty} + \int_{-\infty}^{\infty} \psi^2 \, \mathrm{d}x \quad \text{by parts}$$

$$= \int_{-\infty}^{\infty} (\psi' + x^2 \psi^2)^2 \, \mathrm{d}x + 1$$

As  $(psi' + x^2\psi^2)$  is real valued, its square gives a positive function, thus the integral is positive and  $I[\psi] \geq 1$ . Equality holds for

$$(\psi' + x\psi) = 0$$

$$\Rightarrow \psi' + x\psi = 0$$

$$\Rightarrow \frac{d}{dx} \left( \psi e^{x^2/2} \right) = 0$$

$$\Rightarrow \psi = Ce^{-x^2/2}$$

for some constant C (which we recognise as the Gaussian Wave Function). The normalisation condition implies that  $C = \left(\frac{1}{\pi}\right)^{1/4}$ . Showing it satisfies E-L:  $\psi'' = -Ce^{-x^2/2} + x^2Ce^{-x^2/2}$ 

Using Lagrange multiplier  $\lambda$  with constraint  $|\mathbf{x}| = 1$  wish to minimize

$$\begin{split} \Phi_{\lambda}[\mathbf{x}] &= I[\mathbf{x}] - \lambda(|\mathbf{x}| - 1) \\ &= \int_{t_1}^{t_2} \underbrace{|\dot{\mathbf{x}}|^2 - \lambda|\mathbf{x}|}_{f_{\lambda}(\mathbf{x}, \dot{\mathbf{x}}; t)} \, \mathrm{d}t + \lambda \end{split}$$

The E-L equation for  $x_i$  component is  $\frac{\partial f}{\partial x_i} - \frac{\mathrm{d}}{\mathrm{d}t} \left( \frac{\partial f}{\partial \dot{x}_i} \right) = 0$ , giving

$$-\frac{\lambda x_i}{|\mathbf{x}|} - \frac{\mathrm{d}}{\mathrm{d}t}(2\dot{x}_i) = 0$$

$$2\ddot{x}_i + \lambda x_i = 0$$
 as  $|\mathbf{x}| = 1$ 

$$\Rightarrow \ddot{\mathbf{x}} + \frac{\lambda}{2}\mathbf{x} = 0$$

Also, there is no t dependence in  $f_{\lambda}$ , so E-L equations imply that

$$f - \dot{x}_i \frac{\partial f}{\partial \dot{x}_i} = \text{constant}$$

$$f - 2\dot{x}_i^2 = \text{constant}$$

Have the Lagrangian

$$L = \underbrace{\frac{1}{2}ma^2\dot{\theta}^2 + \frac{1}{2}ma^2\dot{\theta}^2}_{T} + \underbrace{mga\cos\theta}_{-V}$$

Note that  $\frac{\partial L}{\partial \phi} = 0$  so we have the first integral

const. 
$$=\frac{\partial L}{\partial \dot{\phi}} = ma^2 \sin^2 \theta$$

Note too that  $\frac{\partial L}{\partial t} = 0$  (no t dependence) so we have another first integral

const. = 
$$L - \dot{\phi} \frac{\partial L}{\partial \dot{\phi}} - \dot{\theta} \frac{\partial L}{\partial \dot{\theta}}$$
  
=  $-T - V$ 

from which we deduce that

$$\frac{1}{2}ma^2\dot{\theta}^2 + \frac{1}{2}ma^2\dot{\theta}^2 - mga\cos\theta = E$$

for the total energy.

The Hamiltonian is defined as the Legendre transform of the Lagrangian with respect to velocity  $\mathbf{v} = \dot{\mathbf{x}}$ :

$$H(\mathbf{x}, \mathbf{p}; t) = [\mathbf{p} \cdot \mathbf{v} - L(\mathbf{x}, \mathbf{v})]_{\mathbf{v} = \mathbf{v}(\mathbf{p})}$$

where  $\mathbf{v}(\mathbf{p})$  is the solution to  $\frac{\partial L}{\partial \mathbf{v}} = \mathbf{p}$ . The momentum  $p_{\theta}$  is given by

$$p_{\theta} = \frac{\partial L}{\partial \dot{\theta}}$$
$$= ma^2 \dot{\theta}$$

Similarly,

$$p_{\phi} = \frac{\partial L}{\partial \dot{\phi}}$$
$$= ma^{2}(\sin^{2}\theta)\dot{\phi}$$

Thus  $\mathbf{v}=(\dot{\theta},\dot{\phi})=\left(\frac{p_{\theta}}{ma^2},\frac{p_{\phi}}{ma^2\sin^2\theta}\right)$  and

$$H = \frac{p_{\theta}^{2}}{ma^{2}} + \frac{p_{\phi}^{2}}{ma^{2}\sin^{2}\theta} - \left(\frac{1}{2}ma^{2}\left(\frac{p_{\theta}}{ma^{2}}\right)^{2} + \frac{1}{2}ma^{2}\left(\frac{p_{\phi}}{ma^{2}\sin^{2}\theta}\right)^{2} + mga\cos\theta\right)$$

$$= \frac{1}{2}\frac{p_{\theta}^{2}}{ma^{2}} + \frac{1}{2}\frac{p_{\phi}^{2}}{ma^{2}\sin^{2}\theta} - mga\cos\theta$$

Hamilton's equations are given by

$$\dot{\mathbf{x}} = \frac{\partial H}{\partial \mathbf{p}}, \qquad \dot{\mathbf{p}} = -\frac{\partial H}{\partial \mathbf{x}}$$

ie.

$$\dot{\theta} = \frac{p_{\theta}}{ma^2}$$

$$\dot{\phi} = \frac{p_{\phi}}{ma^2 \sin^2 \theta}$$

$$\dot{p_{\theta}} = \frac{-p_{\phi}^2 \sin(2\theta)}{2ma^2 \sin^4 \theta} - mga \cos \theta$$
$$\dot{p_{\phi}} = 0$$

(i) Consider the variation directly,  $u_t \to u_t + (\delta u)_t$  and  $u_x \to u_x + (\delta u)_x$ . Then

$$\begin{split} \delta I[u] &= I[u + \delta u] - I[u] \\ &= \int \left[ \frac{1}{2} \left( u_t + (\delta u)_t \right)^2 - F(u_x + (\delta u)_x) \right] \, \mathrm{d}x \, \mathrm{d}t - \int \left[ \frac{1}{2} u_t^2 - F(u_x) \, \mathrm{d}x \, \mathrm{d}t \right] \\ &= \int u_t (\delta u)_t - (\delta u)_x \frac{\partial F(u_x)}{\partial x} \, \mathrm{d}x \, \mathrm{d}t + O(t^2) \\ &= \int u_{tt} (\delta u) - (\delta u) \frac{\partial F(u_x)}{\partial x^2} \, \mathrm{d}x \, \mathrm{d}t + \text{boundary terms} \quad \text{by parts} \end{split}$$

ignoring boundary terms

So the Euler-Lagrange equation is

$$\frac{\partial^2 u}{\partial t^2} - \frac{\partial F(u_x)}{\partial x^2} = 0$$

(ii) Discarding second order terms throughout:

$$\begin{split} \delta I[u] &= \int \left[ (u_x + (\delta u)_x)^2 + (u_y + (\delta u)_y)^2 + e^{2(u + \delta u)} \right] \, \mathrm{d}x \, \mathrm{d}y - \int u_x^2 + u_y^2 + e^{2u} \, \mathrm{d}x \, \mathrm{d}y \\ &= \int 2u_x (\delta u)_x + 2u_y (\delta u)_y + \left( e^{2(u + \delta u)} - e^{2u} \right) \, \mathrm{d}x \, \mathrm{d}y + O(u^2) \\ &= \int 2u_x (\delta u)_x + 2u_y (\delta u)_y + (2\delta u) \, \mathrm{d}x \, \mathrm{d}y \qquad \text{using } e^x \approx 1 + x \\ &= 2 \int (\delta u) \{ u_{xx} + u_{yy} + 1 \} \, \mathrm{d}x \, \mathrm{d}y \end{split}$$

So E-L equation is

$$u_{xx} + u_{yy} = -1$$

The E-L equations are of the form  $\frac{\partial L}{\partial x_i} - \frac{\mathrm{d}}{\mathrm{d}t} \left( \frac{\partial L}{\partial \dot{x}_i} \right) = 0$ 

We have

$$\frac{\partial L}{\partial x_i} = -q\nabla_i \phi + q \frac{\partial}{\partial x_i} (\mathbf{v} \cdot \mathbf{A})$$

and

$$\frac{\partial L}{\partial \dot{x}_i} = m\gamma v_i + qA_i$$

E-L equations imply

$$-q\nabla_{i}\phi + q\frac{\partial}{\partial x_{i}}(\mathbf{v}\cdot\mathbf{A}) - \frac{\mathrm{d}}{\mathrm{d}t}(m\gamma v_{i} + qA_{i}) = 0$$
  
$$\Rightarrow \frac{\mathrm{d}}{\mathrm{d}t}(m\gamma v_{i}) = -q\nabla_{i}\phi + q\frac{\partial}{\partial x_{i}}(\mathbf{v}\cdot\mathbf{A}) - q\frac{\mathrm{d}}{\mathrm{d}t}(A_{i})$$

Using the chain rule:

$$\frac{\mathrm{d}A_i}{\mathrm{d}t} = \frac{\partial A_i}{\partial t} \frac{\mathrm{d}t}{\mathrm{d}t} + \frac{\partial A_i}{\partial x_j} \frac{\mathrm{d}x_j}{\mathrm{d}t}$$
$$= \frac{\partial A_i}{\partial t} + v_j \frac{\partial A_i}{\partial x_j}$$

and

$$\frac{\partial}{\partial x}(\mathbf{v} \cdot \mathbf{A}) = v_j \frac{\partial A_j}{\partial x_i}$$

$$\Rightarrow \frac{\mathrm{d}}{\mathrm{d}t} (m\gamma v_i) = -q\nabla_i \phi + qv_j \frac{\partial A_j}{\partial x_i} - q \left[ \frac{\partial A_i}{\partial t} + v_j \frac{\partial A_i}{\partial x_j} \right]$$
$$= q \left( \underbrace{\nabla_i \phi - \frac{\partial A_i}{\partial t}}_{E_i} + \left[ v_j \frac{\partial A_j}{\partial x_i} - v_j \frac{\partial A_i}{\partial x_j} \right] \right)$$