Part IB — Complex Methods Example Sheet 3 $\,$

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Know that

$$\tilde{f}(\omega) = \int_{-\infty}^{\infty} f(t)e^{-i\omega t} dt, \quad f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{f}(\omega)e^{i\omega t} d\omega$$

First, let $f(t) = e^{-a|t|}$. Then

$$\begin{split} \tilde{f}(\omega) &= \int_{-\infty}^{\infty} e^{-a|t|} e^{-i\omega t} \, \mathrm{d}t \\ &= \int_{-\infty}^{0} e^{at} e^{-i\omega t} \, \mathrm{d}t + \int_{0}^{\infty} e^{-at} e^{-i\omega t} \, \mathrm{d}t \\ &= \frac{1}{a - i\omega} \left[e^{at} e^{-i\omega t} \right]_{-\infty}^{0} + \frac{1}{-a - i\omega} \left[e^{-at} e^{-i\omega t} \right]_{0}^{\infty} \\ &= \frac{1}{a - i\omega} + \frac{1}{a + i\omega} \\ &= \frac{2a}{a^2 + \omega^2} \end{split}$$

So using the inverse Fourier transform relation we have

$$e^{-a|t|} = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{2a}{a^2 + \omega^2} e^{i\omega t} d\omega$$

$$\iff \frac{\pi}{a} e^{-a|t|} = \int_{-\infty}^{\infty} \frac{1}{a^2 + \omega^2} e^{i\omega t} d\omega$$

as required.

Next, let $f(t) = e^{-at} \sin bt H(t)$. Then

$$\begin{split} \tilde{f}(\omega) &= \int_{-\infty}^{\infty} e^{-at} \sin bt e^{-i\omega t} H(t) \, \mathrm{d}t \\ &= \int_{0}^{\infty} e^{-at} e^{-i\omega t} \sin bt \, \mathrm{d}t \\ &= \int_{0}^{\infty} e^{-at} e^{-i\omega t} \frac{i}{2} \left(e^{-ibt} - e^{ibt} \right) \, \mathrm{d}t \\ &= \frac{i}{2} \int_{0}^{\infty} e^{-at} e^{-i(\omega + b)t} - e^{-at} e^{-i(\omega - b)t} \, \mathrm{d}t \\ &= \frac{i}{2} \left(-\frac{1}{-a - i(\omega + b)} - \frac{-1}{-a - i(\omega - b)} \right) \\ &= \frac{i}{2} \left(\frac{1}{a + i\omega + ib} - \frac{1}{a + i\omega - ib} \right) \\ &= \frac{i(-ib)}{(a + i\omega)^2 + b^2} \\ &= \frac{b}{(a + i\omega)^2 + b^2} \end{split}$$

So using the inverse Fourier transform relation we have

$$e^{-at}\sin bt H(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{b}{(a+i\omega)^2 + b^2} e^{i\omega t} d\omega$$

$$\iff 2\pi e^{-at} \sin bt H(t) = \int_{-\infty}^{\infty} \frac{b}{(a+i\omega)^2 + b^2} e^{i\omega t} d\omega$$

as required.

Not sure about the a < 0, a = 0 parts.

Given

$$f(x) = \begin{cases} 1 & \text{if } |x| < \frac{1}{2}a \\ 0 & \text{otherwise} \end{cases}$$

We compute the Fourier transform as

$$\tilde{f}(k) = \int_{-\infty}^{\infty} f(x)e^{-ikx} dx$$

$$= \int_{-a/2}^{a/2} e^{-ikx} dx$$

$$= \left[\frac{1}{-ik}e^{-ikx}\right]_{-a/2}^{a/2}$$

$$= \frac{i}{k} \left(e^{-ika/2} - e^{ika/2}\right)$$

$$= \frac{2}{k} \sin\left(\frac{ak}{2}\right)$$

as required. Similarly, given

$$g(x) = \begin{cases} a - |x| & \text{if } |x| < a \\ 0 & \text{otherwise} \end{cases}$$

We compute the Fourier transform as

$$\begin{split} \tilde{g}(k) &= \int_{-\infty}^{\infty} g(x)e^{-ikx} \, \mathrm{d}x \\ &= \int_{-a}^{a} (a - |x|)e^{-ikx} \, \mathrm{d}x \\ &= \underbrace{\int_{-a}^{a} ae^{-ikx} \, \mathrm{d}x}_{(1)} - \underbrace{\underbrace{\int_{-a}^{0} -xe^{-ikx} \, \mathrm{d}x}_{(2)} + \underbrace{\int_{0}^{a} xe^{-ikx} \, \mathrm{d}x}_{(3)}}_{(3)} \Big) \\ &= \frac{2a}{k} \sin ak + \frac{a}{ik} (e^{ika} + e^{-ika}) - \frac{1}{(ik)^2} \left[1 - e^{ika} \right] + \frac{1}{(ik)^2} \left[e^{-ika} - 1 \right] \end{split}$$

Now

$$(1) = a \left[\frac{1}{-ik} e^{-ikx} \right]_{-a}^{a}$$
$$= a \frac{i}{k} \left(e^{-ika} - e^{ika} \right)$$
$$= \frac{2a}{k} \sin ak,$$

$$(2) = \left[-x \left(-\frac{1}{ik} e^{-ikx} \right) - \int -\frac{1}{-ik} e^{-ikx} \right]_{-a}^{0}$$

$$= \left[-\frac{a}{ik} e^{ika} + \frac{1}{(ik)^2} \left[e^{-ikx} \right]_{-a}^{0} \right]$$

$$= -\frac{a}{ik} e^{ika} - \frac{1}{k^2} \left(1 - e^{ika} \right)$$

$$(3) = \left[x \left(-\frac{1}{ik} e^{-ikx} \right) - \int -\frac{1}{ik} e^{-ikx} \right]_0^a$$

$$= \left[-\frac{a}{ik} e^{-ika} - \frac{1}{(ik)^2} \left[e^{-ikx} \right]_0^a \right]$$

$$= -\frac{a}{ik} e^{-ika} + \frac{1}{k^2} \left(e^{-ika} - 1 \right)$$

The convolution of $g(x) = e^{-|x|}$ with itself is given by

$$g * g(x) = \int_{-\infty}^{\infty} e^{-|x-y|} e^{-|y|} dy$$

If x > 0 then we can split up the integral as

$$g * g(x) = \int_{-\infty}^{0} e^{x-y} e^{y} dy + \int_{0}^{x} e^{x-y} e^{-y} dy + \int_{x}^{\infty} e^{-(x-y)} e^{-y} dy$$
=

Similarly if x < 0 then

If x > 0 then we can split up the integral as

$$g * g(x) = \int_{-\infty}^{x} e^{x-y} e^{y} dy + \int_{x}^{0} e^{-(x-y)} e^{y} dy + \int_{0}^{\infty} e^{-(x-y)} e^{-y} dy$$

Next, the convolution theorem for Fourier transforms states that $\mathcal{F}[g*g(x)] = \mathcal{F}[g]\mathcal{F}[g]$. Applied here, we have

$$\int_{-\infty}^{\infty} (1+|x|)e^{-|x|}e^{-ikx} dx = \left[\int_{-\infty}^{\infty} e^{-|x|}e^{ikx} dx\right]^2$$

Starting with

$$\mathcal{L}(1) = \frac{1}{s}$$

(i) By shifting,

$$\mathcal{L}(e^{-2t}) = \frac{1}{s+2}$$

(ii) Starting with shifting,

$$\mathcal{L}(e^{-3t}) = \frac{1}{s+3}$$

Then

$$\mathcal{L}(t^3 e^{-3t}) = -\frac{d^3}{ds^3} \mathcal{L}(e^{-3t})$$
$$= -\frac{d^3}{ds^3} \frac{1}{s+3}$$
$$= \frac{6}{(s+3)^4}$$

(iii) Write $\sin 4t = \frac{1}{2i}(e^{4it} - e^{-4it})$. Then

$$\mathcal{L}(e^{3t}\sin 4t) = \frac{1}{2i}\mathcal{L}(e^{(3+4i)t}) - \frac{1}{2i}\mathcal{L}(e^{(3-4i)t})$$
$$= \frac{1}{2i}\left[\frac{1}{s - (3+4i)} - \frac{1}{s - (3-4i)}\right]$$
$$= \frac{4}{(s-3)^2 + 16}$$

(iv) Writing $\cosh 4t = \frac{1}{2}(e^{4t} + e^{-4t})$, we have

$$\mathcal{L}(e^{-4t}\cosh 4t) = \frac{1}{2}\mathcal{L}(1) + \frac{1}{2}\mathcal{L}(e^{-8t})$$
$$= \frac{1}{2}\left[\frac{1}{s} + \frac{1}{s+8}\right]$$
$$= \frac{s+4}{s(s+8)}$$

(v) First by shifting,

$$\mathcal{L}(e^{-(t+1)}) = e^{-1}\mathcal{L}(e^{-t}) = \frac{e^{-1}}{s+1}$$

Then by translation,

$$\mathcal{L}(e^{-t}H(t-1)) = \frac{e^{-(s+1)}}{s+1}$$

By partial fractions we have

$$\hat{f}(s) = \frac{s+3}{(s-2)(s^2+1)} = \frac{1}{s-2} - \frac{s+1}{s^2+1}$$

Have

$$\mathcal{L}(e^{2t}) = \frac{1}{s-2}, \quad \mathcal{L}(\cos t) = \frac{s}{s^2+1}, \quad \mathcal{L}(sint) = \frac{1}{s^2+1}$$

Hence by linearity the inverse Laplace transform of $\hat{f}(s)$ is given by

$$f(t) = e^{2t} - \cos t - \sin t$$

Alternatively we can use the Bromwich inversion formula,

$$f(t) = \sum_{k=1}^{n} \operatorname{res}_{p=p_k}(\hat{f}(p)e^{pt}),$$

as $\hat{f}(s)$ has only a finite number of singularities, namely s=2,i,-i, and $\hat{f}(s)\to 0$ as $|s|\to \infty$. First,

$$\operatorname{res}_{s=2} \left(\frac{s+3}{(s-2)(s^2+1)} e^{st} \right) = \lim_{s \to 2} \left(\frac{s+3}{(s^2+1)} e^{st} \right)$$
$$= e^{2t}$$

Next,

$$\operatorname{res}_{s=i} \left(\frac{s+3}{(s^2+1)} e^{st} \right) = \lim_{s \to i} \left(\frac{s+3}{(s-2)(s+i)} e^{st} \right) \\
= \frac{i+3}{2i(i-2)} e^{it} \\
= -\frac{i+3}{4i+2)} e^{it} \\
= \frac{-10+10i}{20} e^{it} = \frac{-1+i}{2} (\cos t + i \sin t)$$

Similarly,

$$\operatorname{res}_{s=-i} \left(\frac{s+3}{(s^2+1)} e^{st} \right) = \lim_{s \to -i} \left(\frac{s+3}{(s-2)(s-i)} e^{st} \right) \\
= \frac{-i+3}{2i(i+2)} e^{-it} \\
= \frac{-10-10i}{20} e^{-it} = \frac{-1-i}{2} (\cos t - i \sin t)$$

Hence adding the residues we achieve

$$f(t) = e^{2t} - \cos t - \sin t$$

in agreement with what we had above.

Consider the differential equation

$$\frac{\mathrm{d}^3 y}{\mathrm{d}t^3} - 3\frac{\mathrm{d}^2 y}{\mathrm{d}t^2} + 3\frac{\mathrm{d}y}{\mathrm{d}t} - y = t^2 e^t$$

$$y(0) = 1, \dot{y}(0) = 0, \ddot{y}(0) = -2$$

Taking the Laplace transform of this equation, where $\mathcal{L}(y) = \hat{y}$, we have

$$\mathcal{L}(\dot{y}) = p\hat{y} + y(0)$$
$$= p\hat{y} + 1$$

Similarly

$$\mathcal{L}(\ddot{y}) = p\mathcal{L}(\dot{y}) + \dot{y}(0)$$
$$= p^2 \hat{y} + py(0) + \dot{y}(0)$$
$$= p^2 \hat{y} + p$$

and,

$$\mathcal{L}\left(\frac{\mathrm{d}^3 y}{\mathrm{d}t^3}\right) = p\mathcal{L}(\ddot{y}) + \ddot{y}(0)$$
$$= p^3 \hat{y} + p^2 y(0) + p\dot{y}(0) + \ddot{y}(0)$$
$$= p^3 \hat{y} + p^2 + -2$$

The term on the RHS gives

$$\mathcal{L}(t^2 e^t) = \frac{\mathrm{d}^2}{\mathrm{d}p^2} \mathcal{L}(e^t)$$
$$= \frac{\mathrm{d}^2}{\mathrm{d}p^2} \left(\frac{1}{p-1}\right)$$
$$= \frac{2}{(p-1)^3}$$

Thus substituting in gives:

$$(p^3 - 3p^2 + 3p - 1)\hat{y} + p^2 - 3p + 1 = \frac{2}{(p-1)^3}$$

Hence

$$\hat{y} = \frac{2}{(p-1)^6} - \frac{1}{(p-1)} + \frac{p}{(p-1)^3}$$

We know $\mathcal{L}(e^t)=1/(p-1)$. For the other two terms, we will use the Bromwich inversion formula from the previous question, noting both tend to zero as $|p| \to \infty$.

$$f(t) = \sum_{k=1}^{n} \underset{p=p_k}{\text{res}} (\hat{f}(p)e^{pt}),$$

For the first term, the singularity p = 1 has a pole of order 6. Hence the residue here is given by

$$\lim_{p \to 1} \frac{1}{5!} \frac{d^5}{dp^5} (p-1)^5 \left(\frac{2}{(p-1)^6} e^{pt} \right) = \lim_{p \to 1} \frac{1}{5!} \frac{d^5}{dp^5} \left(-2e^{pt} (1-p)^{-1} \right)$$

$$= \lim_{p \to 1} \frac{1}{5!} \frac{d^5}{dp^5} \left[-2e^{pt} \left(1 + p + p^2 + p^3 + \cdots \right) \right]$$

$$= \cdots$$

The last term has a singularity at p=3 with a pole of order 3 there; the residue is given by

$$\begin{split} \lim_{p \to 1} \frac{1}{2!} \frac{\mathrm{d}^2}{\mathrm{d}p^2} (p-1)^2 \left(\frac{p}{(p-1)^3} e^{pt} \right) &= \lim_{p \to 1} \frac{1}{2!} \frac{\mathrm{d}^2}{\mathrm{d}p^2} \left(\frac{p}{1-p} e^{pt} \right) \\ &= \lim_{p \to 1} \frac{1}{2!} \frac{\mathrm{d}}{\mathrm{d}p} \left(\frac{1}{(1-p)^2} e^{pt} + \frac{p^2}{(1-p)} e^{pt} \right) \\ &= \lim_{p \to 1} \frac{1}{2!} \left(\frac{2}{(1-p)^3} e^{pt} + \frac{p}{(1-p)^2} e^{pt} + \frac{1-p^2}{(1-p)^2} e^{pt} + \frac{p^3}{(1-p)} e^{pt} \right) \end{split}$$

Not sure how to take limit.

The differential equation we are investigating is

$$\ddot{y} - 2\dot{y} - 2y = \delta(t) - \delta(t - t_0)$$

Note that

$$\mathcal{L}(\delta(t)) = \int_0^\infty \delta(t)e^{\mathrm{pt}} dt = 1, \quad \mathcal{L}(\delta(t - t_0)) = e^{-pt_0}$$

Taking the Laplace transform of the equation gives

$$p^2\hat{y} + 2(p\hat{y}) + 2\hat{y} = 1 + e^{-pt_0}$$

which gives

$$\hat{y} = \frac{1 + e^{-pt_0}}{(p^2 + 2p + 2)}$$

Taking the Laplace transform of the equation gives

$$\hat{f} + 4\frac{1}{p}\hat{f} = \frac{1}{p^2}$$

Thus

$$\hat{f} = \frac{1}{p^2 + 4p} = \frac{1}{(p+2)^2 - 4}$$

 $(whenever\ I\ can\ be\ bothered)\ http://www.robots.ox.ac.uk/\ jmb/lectures/pdelecture4.pdf$