# Part IB — Electromagnetism Example Sheet 1 $\,$

Supervised by Dr. Warnick (cmw50@cam.ac.uk) Examples worked through by Christopher Turnbull

 $Lent\ 2018$ 

Equation for conservation of charge is

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{J} = 0$$

Have  $\mathbf{J} = C\mathbf{r}e^{-atr^2}$ , so

$$\nabla \cdot \mathbf{J} = C \nabla \cdot (e^{-atr^2} \mathbf{r})$$
$$= C e^{-atr^2} \nabla \cdot \mathbf{r} + C \mathbf{r} \cdot \nabla (e^{-atr^2})$$

Now  $\mathbf{r}_i = x_i$  so  $\nabla \cdot \mathbf{r} = \frac{\partial x_j}{\partial x_j} = 3$ , and

$$\nabla e^{-atr^2} = \frac{\partial e^{-atr^2}}{\partial r} \hat{\mathbf{r}}$$
$$= -2ate^{-atr^2} \mathbf{r}$$

Hence

$$\nabla \cdot \mathbf{J} = 3Ce^{-atr^2} - 2Cr^2ate^{-atr^2}$$

Suppose that  $\rho = (f + tg)e^{-atr^2}$ . Then we have

$$\frac{\partial \rho}{\partial t} = (-ar^2f + g - ar^2tg)e^{-atr^2}$$
$$= (g - ar^2f)e^{-atr^2} - gtar^2e^{-atr^2}$$

Hence we conclude that

$$g - ar^2 f = -3C,$$
  $g = -2C$  
$$\Rightarrow f = \frac{C}{ar^2}$$

Using the continuity equation,

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot \mathbf{J}$$
$$= -\nabla \cdot (-D\nabla \rho)$$
$$= D\nabla^2 \rho$$

showing  $\rho(\mathbf{x},t)$  obeys the heat equation with diffusion constant D. Let  $\rho(\mathbf{r},t)$  be defined as

$$\rho(\mathbf{r},t) = \frac{\rho_0 a^3}{(4D(t-t_0) + a^2)^{3/2}} \exp\left(-\frac{r^2}{4D(t-t_0) + a^2}\right)$$

Taking time derivatives,

$$\frac{\partial \rho}{\partial t} = \left(\frac{-6D\rho_0 a^3}{(4D(t-t_0)+a^2)^{5/2}} + \frac{4Dr^2\rho_0 a^3}{(4D(t-t_0)+a^2)^{7/2}}\right) \exp\left(-\frac{r^2}{4D(t-t_0)+a^2}\right)$$

Now

$$\nabla^2 e^{\lambda r^2} = \frac{1}{r^2} \frac{\mathrm{d}}{\mathrm{d}r} \left( r^2 \frac{\mathrm{d}}{\mathrm{d}r} (e^{\lambda r^2}) \right)$$
$$= \frac{2\lambda}{r^2} \frac{\mathrm{d}}{\mathrm{d}r} \left( r^3 e^{\lambda r^2} \right)$$
$$= \frac{2\lambda}{r^2} \left[ 3r^2 + 2\lambda r^4 \right] e^{\lambda r^2}$$
$$= \lambda (6 + 4r^2) e^{\lambda r^2}$$

Thus with  $\lambda = -\frac{1}{4D(t-t_0)+a^2}$ , we have

$$\begin{split} \nabla^2 \rho &= \frac{-\rho_0 a^3}{(4D(t-t_0)+a^2)^{5/2}} \left(6 - \frac{4r^2}{4D(t-t_0)+a^2}\right) \exp\left(-\frac{r^2}{4D(t-t_0)+a^2}\right) \\ &= \left(-\frac{6\rho_0 a^3}{(4D(t-t_0)+a^2)^{5/2}} + \frac{4r^2\rho_0 a^3}{(4D(t-t_0)+a^2)^{-7/2}}\right) \exp\left(-\frac{r^2}{4D(t-t_0)+a^2}\right) \end{split}$$

Hence we can see that  $\frac{\partial \rho}{\partial t} = D\nabla^2 \rho$ , as required.

Considering the infinite plane z = 0, we see this has uniform charge density  $\rho_0$ . By symmetry, the field points vertically, and the field on the bottom is opposite of that on top, we must have

$$\mathbf{E} = E(z)\hat{\mathbf{z}}$$

with

$$E(z) = -E(-z)$$

Consider a vertical cylinder of height 2h and cross-sectional area A. Now only the end caps contribute.

First,

$$\begin{split} Q &= \int_{V} \rho_{0} e^{-k|z|} \, \mathrm{d}V \\ &= \int_{-h}^{h} \int_{0}^{2\pi} \int_{0}^{R} \rho_{0} e^{-k|z|} \rho \, \mathrm{d}\rho \mathrm{d}\phi \mathrm{d}z \\ &= 2\pi \frac{R^{2}}{2} \rho_{0} \int_{-h}^{h} e^{-k|z|} \, \mathrm{d}z \\ &= A \rho_{0} \int_{0}^{h} 2e^{-kz} \, \mathrm{d}z \\ &= 2A \rho_{0} \left[ -\frac{1}{k} e^{-kz} \right]_{0}^{h} \\ &= 2A \frac{\rho_{0}}{k} (1 - e^{-kh}) \end{split}$$

And

$$\int_{S} \mathbf{E} \cdot d\mathbf{S} = E(h)A - E(-h)A = 2AE(h) = 2A\frac{\rho_0}{k\varepsilon_0} (1 - e^{-kh})$$

Hence

$$E(z) = \frac{\rho_0}{k\varepsilon_0} (1 - e^{-kz})$$

as required.

We have that

$$\rho(r) = \begin{cases} 0 & \text{if } r < a \\ \rho & \text{if } a < r < b \\ 0 & \text{if } r > b \end{cases}$$

First consider r > b. By symmetry, the force is the same in all directions and points outwards radially. So

$$\mathbf{E} = E(r)\hat{\mathbf{r}}$$

Put S to be a sphere of radius r > b. Then the total flux is

$$\int_{S} \mathbf{E} \cdot dS = \int_{S} E(r) \hat{\mathbf{r}} \cdot d\mathbf{S}$$
$$= E(r) \int_{S} \hat{\mathbf{r}} \cdot d\mathbf{S}$$
$$= E(r) \cdot 4\pi r^{2}$$

By Gauss's law, we know this is equal to  $Q/\varepsilon_0$ , and  $Q = \frac{4}{3}\pi(b^3 - a^3)\rho$ . Therefore,

$$E(r) = \frac{(b^3 - a^3)\rho}{3\varepsilon_0 r^2}$$

and

$$\mathbf{E}(r) = \frac{(b^3 - a^3)\rho}{3\varepsilon_0 r^2}\hat{\mathbf{r}}$$

Now suppose we are inside the region, a < r < b. Then

$$\int_{S} \mathbf{E} \cdot dS = E(r) 4\pi r^{2} = \frac{Q}{\varepsilon_{0}} \left( \frac{r^{3} - a^{3}}{b^{3} - a^{3}} \right)$$

So

$$\mathbf{E}(r) = \frac{Q(r^3 - a^3)}{4\pi\varepsilon_0(b^3 - a^3)r^2}$$
$$= \frac{Q(r^3 - a^3)\rho}{3\varepsilon_0 r^2}$$

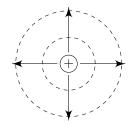
Finally if r < a, Gauss' law tells us that the flux depends only on the total charge contained inside the surface, which in this case is none. So  $\mathbf{E}(r) = 0$ .

Note that the electric field is discontinuous across the surface. We have

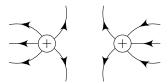
$$E(r \to b+) - E(r \to b-) = \frac{(b-a)(b^2 + 2ab + a^2)\rho}{3\varepsilon_0 b^2}$$
$$= \frac{\sigma}{\varepsilon_0}$$

as expected.

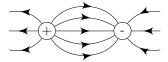
The field lines for a positive charge are:



For two positive charges,



We can also draw field lines for dipoles:



The inverse square law, or Coulomb's Law, states that the electric field generated by a particle with total charge Q (at the origin) is given by

$$\mathbf{E}(r) = \frac{Q}{4\pi\varepsilon_0 r^2}\hat{\mathbf{r}}$$

Now consider an infinite line with uniform charge density per unit length  $\eta$ . We use cylindrical polar coordinates. By symmetry, the field is radial, ie.

$$\mathbf{E}(r) = E(r)\hat{\mathbf{r}}$$

Consider an arbitrary point at  $(r, z_0)$ . We will integrate along the z-axis to find the field at this point.

Here,

$$\begin{split} E(r) &= \int_{-\infty}^{\infty} \frac{Q}{4\pi\varepsilon_0} \frac{1}{r^2 + (z - z_0)^2} \; \mathrm{d}z \\ &= \frac{Q}{4\pi\varepsilon_0} \int_{-\infty}^{\infty} \frac{1}{r^2 + z^2} \; \mathrm{d}z \\ &= \frac{Q}{4\pi\varepsilon_0} \left[ \frac{1}{r} \arctan\left(\frac{z}{r}\right) \right]_{-\infty}^{\infty} \\ &= \frac{Q}{4r\varepsilon_0} \end{split}$$