# Part IB — Methods Example Sheet 1

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$$\frac{f(x_{+}+f(x_{-}))}{2} = \frac{1}{2}a_{0} + \sum_{n=1}^{\infty} \left[ a_{n} \cos\left(\frac{n\pi x}{L}\right) + b_{n} \sin\left(\frac{n\pi x}{L}\right) \right]$$

For  $f(x)=(x-1)^2$  on the interval  $-1\leq x\leq 1,$  f(x) is an even function, thus  $b_n=0$ . We have L=1, and

$$\frac{1}{2}a_0 = \frac{1}{2L} \int_{-L}^{L} f(x) dx$$

$$= \frac{1}{2} \int_{-1}^{1} x^4 - 2x^2 + 1 dx$$

$$= \int_{0}^{1} x^4 - 2x^2 + 1 dx$$

$$= \frac{8}{15}$$

and

$$a_n = \frac{1}{L} \int_{-1}^{1} f(x) \cos\left(\frac{n\pi x}{L}\right) dx$$
$$= \int_{-1}^{1} x^4 \cos n\pi x dx - 2 \int_{-1}^{1} x^2 \cos n\pi x dx + \int_{-1}^{1} \cos n\pi x dx$$

Evaluating each integral separately, we have:

(i) 
$$\int_{-1}^{1} \cos n\pi x \, dx = \left[ \frac{\sin n\pi x}{n\pi} \right]_{-1}^{-1} = 0$$

as  $\sin n\pi x = 0 \ \forall \ n$ 

(ii) By parts,

$$\int_{-1}^{1} x^{2} \cos n\pi x \, dx = \left[ \frac{x^{2} \sin n\pi x}{n\pi} - \frac{1}{n\pi} \int 2x \sin n\pi x \, dx \right]_{-1}^{1}$$
$$= -\frac{2}{n\pi} \int_{-1}^{1} x \sin n\pi x \, dx$$

and

$$\int_{-1}^{1} x \sin n\pi x \, dx = \left[ \frac{-x \cos n\pi x}{n\pi} + \frac{1}{n\pi} \int \cos n\pi x \, dx \right]_{-1}^{1}$$
$$= \frac{-2 \cos n\pi x}{(n\pi)^{2}}$$

Thus the second integral contributes to give

$$-\frac{8cosn\pi x}{(n\pi)^2}$$

(iii)

$$\int_{-1}^{1} x^{4} \cos n\pi x \, dx = \left[ \frac{x^{4} \sin n\pi x}{n\pi} - \frac{1}{n\pi} \int 4x^{3} \sin n\pi x \, dx \right]_{-1}^{1}$$
$$= -\frac{4}{n\pi} \int_{-1}^{1} x^{3} \sin n\pi x \, dx$$

and

$$\int_{-1}^{1} x^{3} \sin n\pi x \, dx = \left[ \frac{-x^{3} \cos n\pi x}{n\pi} + \frac{1}{n\pi} \int 3x^{2} \cos n\pi x \right]_{-1}^{1}$$
$$= \frac{-2 \cos n\pi}{n\pi} + \frac{3}{n\pi} \int_{-1}^{1} x^{2} \cos n\pi x \, dx$$

Whence

$$\int_{-1}^{1} x^{4} \cos n\pi x \, dx = \frac{8 \cos n\pi}{n\pi} - \frac{12}{(n\pi)^{2}} \int_{-1}^{1} x^{2} \cos n\pi x \, dx$$
$$= \frac{8 \cos n\pi}{n\pi} - \frac{48 \cos n\pi}{(n\pi)^{4}}$$

using (ii).

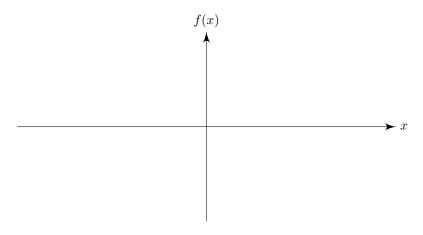
Finally,

$$a_n = -\frac{48\cos n\pi}{(n\pi)^4}$$
$$= \frac{48(-1)^{n+1}}{(n\pi)^4}$$

as  $\cos n\pi x = (-1)^n$ 

Hence the Fourier Series is given by

$$f(x) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} a_n \cos n\pi x$$
$$= \frac{8}{15} + \frac{48}{\pi^4} \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n^4} \cos n\pi x$$



f(x) satisfies the Dirichlet conditions. The 1<sup>st</sup> derivative is the lowest derivative which is discontinuous (at the endpoints, as f(x) even fn  $\Rightarrow f'(x)$  odd), so Fourier coefficients are  $\mathcal{O}(\frac{1}{n^2})$  as  $n \to \infty$ 

Extending on range  $(-\pi, \pi)$  so  $L = \pi$  and

(a)  $\frac{f(x_{+} + f(x_{-}))}{2} = \sum_{n=0}^{\infty} b_{n} \sin nx$ 

where

$$b_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx$$
$$= \frac{2}{\pi} \int_0^{\pi} x^2 \sin nx dx$$

Integrating by parts,

$$\int_0^\pi x^2 \sin nx \, dx = \left[ \frac{-x^2 \cos nx}{n} + \frac{1}{n} \int 2x \cos nx \, dx \right]_0^\pi$$
$$= \frac{-\pi^2 \cos n\pi}{n} + \frac{2}{n} \int_0^\pi x \cos nx \, dx$$

and once again,

$$\int_0^{\pi} x \cos nx \, dx = \left[ \frac{x \sin nx}{n} - \frac{1}{n} \int \sin nx \, dx \right]_0^{\pi}$$
$$= -\frac{1}{n} \int_0^{\pi} \sin nx \, dx$$
$$= -\frac{1}{n} \left[ -\frac{\cos nx}{n} \right]_0^{\pi}$$
$$= \frac{1}{n^2} (\cos n\pi - 1)$$

Back substituting in,

$$b_n = \frac{2}{\pi} \left( \frac{-\pi^2 \cos n\pi}{n} + \frac{2}{n^3} (\cos n\pi - 1) \right)$$
$$= \frac{2}{\pi n^3} \left( -2 + (2 - (\pi n)^2) \cos n\pi \right)$$

Hence Fourier sine series given by:

$$f(x)_s = \sum_{n=1}^{\infty} \frac{2}{\pi n^3} \left( -2 + (2 - (\pi n)^2)(-1)^n \right) \sin nx$$
$$= \sum_{n=1}^{\infty} \left\{ \frac{2\pi (-1)^{n+1}}{n} + \frac{4[(-1)^n - 1]}{n^3 \pi} \right\} \sin nx$$

(b) Similarly,

$$\frac{f(x_{+} + f(x_{-}))}{2} = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos nx$$

where

$$\frac{a_0}{2} = \frac{1}{L} \int_0^L f(x) dx$$
$$= \frac{1}{\pi} \int_0^{\pi} x^2 dx$$
$$= \frac{\pi^2}{3}$$

and

$$a_n = \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx$$
$$= \frac{2}{\pi} \int_0^{\pi} x^2 \cos nx dx$$

Integrating by parts,

$$\int_0^\pi x^2 \cos nx \, dx = \left[ \frac{x^2 \sin nx}{n} - \frac{1}{n} \int 2x \sin nx \, dx \right]_0^\pi$$
$$= \frac{-2}{n} \int_0^\pi x \sin nx \, dx$$

and once again,

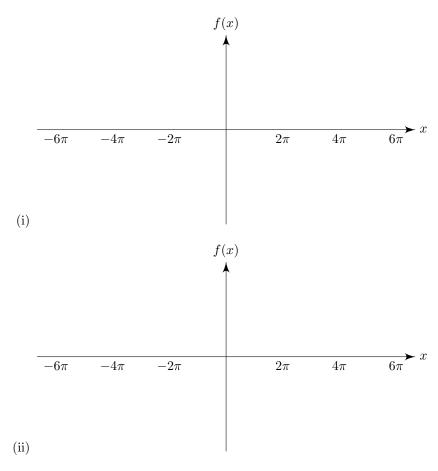
$$\int_0^{\pi} x \sin nx \, dx = \left[ \frac{-x \cos nx}{n} + \frac{1}{n} \int \cos nx \, dx \right]_0^{\pi}$$
$$= \frac{-\pi \cos n\pi}{n} + \frac{1}{n} \left[ \frac{\sin nx}{n\pi} \right]_0^{\pi}$$
$$= \frac{-\pi \cos n\pi}{n}$$

Thus

$$a_n = \frac{4}{n^2} \cos n\pi$$

and the Fourier cosine series is given by

$$f(x) = \frac{\pi^2}{3} + \sum_{n=0}^{\infty} \frac{4}{n^2} (-1)^n \cos nx$$



Fourier series for g(x)=2x (odd function) in the range  $(-\pi,\pi)$  given by

$$\frac{f(x_+ + f(x_-))}{2} = \sum_{n=1}^{\infty} b_n \sin nx$$

where

$$b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin\left(\frac{n\pi x}{L}\right) dx$$
$$= \frac{2}{\pi} \int_{-\pi}^{\pi} x \sin nx dx$$

Integrating by parts,

$$\int_{-\pi}^{\pi} x \sin nx \, dx = \left[ \frac{-x \cos nx}{n} + \frac{1}{n} \int \cos nx \, dx \right]_{-\pi}^{\pi}$$
$$= \frac{-2\pi \cos n\pi}{n} + \frac{1}{n} \left[ \frac{\sin nx}{n\pi} \right]_{-\pi}^{\pi}$$
$$= \frac{2\pi (-1)^{n+1}}{n}$$

Whence

$$g(x) = \sum_{n=1}^{\infty} \frac{4\pi^2 (-1)^{n+1}}{n^2} \sin nx$$

Fourier series for h(x)=2|x| (even function) in the range  $(-\pi,\pi)$  given by

$$\frac{f(x_{+} + f(x_{-}))}{2} = \frac{1}{2}a_{0} + \sum_{n=1}^{\infty} a_{n} \cos nx$$

where

$$\frac{1}{2}a_0 = \frac{1}{2L} \int_{-L}^{L} f(x) dx$$
$$= \frac{1}{2\pi} \int_{-\pi}^{\pi} 2|x| dx$$
$$= \frac{2}{\pi} \int_{0}^{\pi} x dx$$
$$= \pi$$

and

$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{n\pi x}{L}\right) dx$$
$$= \frac{2}{\pi} \int_{-\pi}^{\pi} |x| \cos nx dx$$

Integrating by parts,

$$\int_{-\pi}^{\pi} |x| \cos nx \, dx = 2 \int_{0}^{\pi} x \cos nx \, dx$$

$$= 2 \left[ \frac{x \sin nx}{n} - \frac{1}{n} \int \sin nx \, dx \right]_{0}^{\pi}$$

$$= -\frac{2}{n} \int_{0}^{\pi} \sin nx \, dx$$

$$= -\frac{2}{n} \left[ -\frac{\cos nx}{n} \right]_{0}^{\pi}$$

$$= \frac{2}{n^{2}} (\cos n\pi - 1)$$

Whence

$$h(x) = \pi + \sum_{n=1}^{\infty} \frac{4[(-1)^n - 1]}{n^2 \pi} \cos nx$$

Note that differentiating the Fourier sine series for  $x^2$  gives

$$\frac{\mathrm{d}}{\mathrm{d}x}[f_s(x)] = \sum_{n=1}^{\infty} \left\{ 2\pi (-1)^{n+1} + \frac{4[(-1)^n - 1]}{n^2 \pi} \right\} \cos nx$$

These don't quite match up: what is  $\sum_{n=1}^{\infty} 2\pi (-1)^{n+1} \cos nx$  the Fourier series for?

Note that the cos coefficients  $a_n = O(1)$  as  $n \to \infty$ , so this function is terrible. Using the direchlet conditions,  $a_n = O(\frac{1}{n})$ , so f is discontinuous. This motivates us to check the Fourier series for the direct delta function  $\delta(x)$ ,

with period  $(-\pi,\pi)$ 

$$\delta(x) \sim \frac{a_0}{2} + \sum_{n=1}^{\infty} \left\{ a_n \cos nx + b_n \sin nx \right\}$$

We find that

$$a_n = \frac{1}{\pi} \int_{-pi}^{\pi} \delta(x) \cos nx \, dx$$
$$= \frac{\cos 0}{\pi}$$
$$= \frac{1}{\pi}$$

and similarly

$$b_n = \frac{\sin 0}{\pi} = 0$$

ie.

$$\delta(x) \sim \sum_{n=1}^{\infty} \frac{1}{\pi} \cos nx + \frac{1}{2\pi}$$

This isn't quite what we wanted, but making a small adjustment:

$$\delta(x-\pi): a_n = \frac{\cos n\pi}{\pi} = \frac{(-1)^n}{\pi} \quad b_n = 0$$

Finally, we conclude that

$$f_s'(x) = 2|x| - 2\pi^2 \delta(x - \pi)$$

I guess the morale is, don't differentiate term by term if f is discontinuous...

 $f(x) = e^x$  on  $(-\pi, \pi)$  has Fourier series given by

$$\frac{f(x_{+} + f(x_{-}))}{2} = \frac{1}{2}a_{0} + \sum_{n=1}^{\infty} \left[a_{n} \cos nx + b_{n} \sin nx\right]$$

where

$$\frac{1}{2}a_0 = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^x dx$$
$$= \frac{1}{2\pi} \left( e^{\pi} - e^{-\pi} \right)$$
$$= \frac{1}{\pi} \sinh \pi$$

and

$$a_n = \frac{1}{\pi} \underbrace{\int_{-\pi}^{\pi} e^x \cos nx \, dx}_{I_a}$$

$$I_a = \left[ e^x \cos nx + \int e^x n \sin nx \, dx \right]_{-\pi}^{\pi}$$

$$= \left( e^{\pi} - e^{-\pi} \right) \cos n\pi + n \int_{-\pi}^{\pi} e^x \sin nx \, dx$$

$$= 2 \sinh \pi (-1)^n + n \left[ e^x \sin nx - \int e^x n \cos x \, dx \right]_{-\pi}^{\pi}$$

$$= 2 \sinh \pi (-1)^n + -n^2 \int_{-\pi}^{\pi} e^x \cos nx \, dx$$

$$= 2 \sinh \pi (-1)^n + -n^2 I_a$$

Hence

$$a_n = \frac{1}{\pi} I_a$$
  $I_a = \frac{2}{1+n^2} \sinh \pi (-1)^n$ 

Also,

$$b_n = \frac{1}{\pi} \underbrace{\int_{-\pi}^{\pi} e^x \sin nx \, \mathrm{d}x}_{I_b}$$

$$I_b = \left[ e^x \sin nx - \int e^x n \cos nx \, dx \right]_{-\pi}^{\pi}$$
$$= -nI_a$$

$$b_n = -\frac{n}{\pi}I_a$$

Combining these results, the Fourier series for  $e^x$  is given by

$$f(x) = \frac{1}{\pi} \sinh \pi + \sum_{n=1}^{\infty} \left[ \left( \frac{1}{\pi} \cos nx - \frac{n}{\pi} \sin nx \right) I_a \right]$$
$$= \frac{1}{\pi} \sinh \pi + \frac{2}{\pi} \sinh \pi \sum_{n=1}^{\infty} \left[ (\cos nx - n \sin nx) \frac{(-1)^n}{1 + n^2} \right]$$

Setting  $x = \pi$  yields

$$e^{\pi} = \frac{1}{\pi} \sinh \pi + \frac{2}{\pi} \sinh \pi \sum_{n=1}^{\infty} \frac{1}{1+n^2}$$

Thus

$$\sum_{n=1}^{\infty} \frac{1}{1+n^2} = \frac{\pi e^{\pi} - \sinh \pi}{2 \sinh \pi}$$

Setting  $x = -\pi$  similarly yields

$$\sum_{n=1}^{\infty} \frac{1}{1+n^2} = \frac{\pi e^{-\pi} - \sinh \pi}{2 \sinh \pi}$$

Adding and dividing by two,

$$\sum_{n=1}^{\infty} \frac{1}{1+n^2} = \frac{\pi(e^{\pi} + e^{-\pi}) - 2\sinh \pi}{4\sinh \pi}$$
$$= \frac{2\pi \cosh \pi - 2\sinh \pi}{4\sinh \pi}$$
$$= \frac{1}{2}(\pi \coth \pi - 1)$$

(i) Reposing the Fourier Series of f(t) using complex variables,

$$\begin{split} f(t) &= \frac{a_0}{2} + \sum_{n=1}^{\infty} \left[ \frac{a_n}{2} \left( e^{\frac{in\pi t}{L}} + e^{\frac{-in\pi t}{L}} \right) + \frac{b_n}{2i} \left( e^{\frac{in\pi t}{L}} - e^{\frac{-in\pi t}{L}} \right) \right] \\ &= \sum_{n=-\infty}^{\infty} c_n e^{\frac{in\pi t}{L}}, \\ c_n &= \frac{a_n - ib_n}{2} \ n > 0; \\ c_{-n} &= \frac{a_n + ib_n}{2} \ n > 0; \\ c_0 &= \frac{a_0}{2} \end{split}$$

Using the orthogonality of complex exponentials and the properties of complex Fourier coefficients, we deduce that

$$\int_{-L}^{L} [f(t)]^{2} dt = \sum_{n=-\infty}^{\infty} \sum_{m=-\infty}^{\infty} c_{n} c_{m} \int_{-T}^{T} \exp\left[\frac{i\pi t(n+m)}{L}\right] dt$$

$$= \sum_{n=-\infty}^{\infty} \sum_{m=-\infty}^{\infty} c_{n} c_{m} 2T \delta_{n[-m]}$$

$$= 2T \sum_{n=-\infty}^{\infty} c_{n} c_{-n}$$

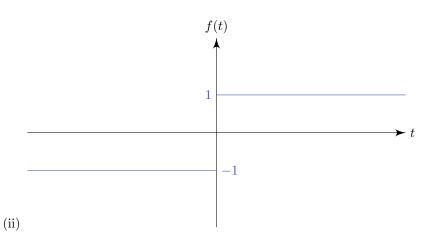
$$= 2T \sum_{n=-\infty}^{\infty} c_{n} c_{n}^{*}$$

$$= 2T \sum_{n=-\infty}^{\infty} |c_{n}|^{2}$$

This can be then re-expressed in terms of the  $a_n$  and  $b_n$  as

$$\int_{-L}^{L} \left[ f(t) \right]^{2} dt = L \left[ \frac{a_{0}^{2}}{2} + \sum_{n=1}^{\infty} (a_{n}^{2} + b_{n}^{2}) \right]$$

as required.



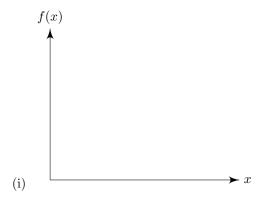
The unit amplitude square wave has Fourier series (odd function)

$$f(t) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi t}{T}\right)$$

Frequencies less than  $\frac{9}{2}\pi T^{-1}$  correspond to terms in the Fourier series with  $\frac{n\pi}{T}<\frac{9}{2}\pi T^{-1}$ , ie. n=1,2,3,4.

Also,

$$b_n = \frac{1}{T} \int_{-T}^{T}$$



f(x) on  $(0,2\pi)$  has Fourier series given by

$$\frac{f(x_{+} + f(x_{-}))}{2} = \frac{1}{2}a_{0} + \sum_{n=1}^{\infty} \left[a_{n} \cos nx + b_{n} \sin nx\right]$$

where

$$\frac{1}{2}a_0 = \frac{1}{2\pi} \int_0^{2\pi} f(x) \, dx$$
$$= \frac{1}{2\pi} \int_{\pi}^{2\pi} 1 \, dx$$
$$= \frac{1}{2}$$

and

$$a_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \cos nx \, dx$$
$$= \frac{1}{\pi} \int_{\pi}^{2\pi} \cos nx \, dx$$
$$= \frac{1}{\pi} \left[ \frac{1}{n} \sin nx \right]_{\pi}^{2\pi}$$
$$= 0$$

and

$$b_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \sin nx \, dx$$

$$= \frac{1}{\pi} \int_{\pi}^{2\pi} \sin nx \, dx$$

$$= \frac{1}{\pi} \left[ -\frac{1}{n} \cos nx \right]_{\pi}^{2\pi}$$

$$= -\frac{1}{\pi n} \left[ \cos nx \right]_{\pi}^{2\pi}$$

$$= 0 \text{ if } n \text{ even or } -\frac{2}{n\pi} \text{ if } n \text{ odd}$$

Hence

$$f(x) = \frac{1}{2} - \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{\sin(2n-1)x}{2n-1}$$

(ii) Taking the hint, differentiating term by term gives

$$\frac{\mathrm{d}}{\mathrm{d}x}[S_n(x)] = \frac{2}{\pi} \sum_{n=1}^{N} \cos(2n-1)x$$

Now

$$\sum_{n=1}^{N} \cos(2n-1)x = \text{Re}\left[\sum_{n=1}^{N} e^{(2n-1)i}\right]$$
$$= \text{Re}\left[\frac{e^{(2N+1)i} - e^{i}}{e^{2i} - 1}\right]$$

Assuming the solution takes the form  $y \propto e^{\sigma x}$ ,, we have

$$y(x) = A\cos(\mu x) + B\sin(\mu x)$$

where A and B are constants, and  $\mu^2 = \lambda$ . Applying the boundary conditions, y(0) = 0 implies that A = 0. The other boundary condition implies

$$B\sin\mu + B\mu\cos\mu = 0$$

$$\mu = -\tan \mu$$

This eigenvalue equation has an infinite number of solutions,  $\mu_n$  (and hence there an infinite number of positive eigenvalues  $\lambda_n = \mu_n^2$ ).

As  $n \to \infty$ ,  $\mu_n \to \infty$ , so  $\mu$  is close to an odd multiple of  $\frac{\pi}{2}$ , ie.  $\mu_n \approx (2n+1)\pi/2$ , and hence  $\lambda_n \approx (2n+1)^2\pi^2/4$ 

(i)  $p(x) = \exp\left(\int^x \frac{-2u}{1-u^2} du\right) = (1-x^2)$ , thus integrating factor is  $-\frac{1}{1-x^2} \left((1-x^2)\right) = -1$ . We can then rewrite the equation as

$$-(1-x^2)y'' + 2xy' - n(n+1)y = 0$$

and

$$-\frac{\mathrm{d}}{\mathrm{d}x}\left((1-x^2)\frac{\mathrm{d}y}{\mathrm{d}x}\right) - n(n+1)y = 0$$

(ii)

$$p(x) = \exp\left(\int_{-\infty}^{x} \frac{(1+a+b)u - c}{u(u-1)} du\right)$$

$$= \exp\left(\int_{-\infty}^{x} \frac{c}{u} + \frac{1+a+b-c}{u-1} du\right)$$

$$= \exp\left(c\log x + (1+a+b-c)\log(x-1)\right)$$

$$= x^{c} + (x-1)^{1+a+b-c}$$

Thus the required integrating factor is

$$-\frac{x^{c} + (x-1)^{1+a+b-c}}{x(x-1)}$$

The equation becomes

$$-(x^{c}+(x-1)^{1+a+b-c})y'' + -[(1+a+b)x-c]\frac{x^{c}+(x-1)^{1+a+b-c}}{x(x-1)}y' - \frac{x^{c}+(x-1)^{1+a+b-c}}{x(x-1)}aby = 0$$

which, in Sturm-Liouville form, is

$$-\frac{\mathrm{d}}{\mathrm{d}x}\left[(x^c + (x-1)^{1+a+b-c})\frac{\mathrm{d}y}{\mathrm{d}x}\right] - \frac{x^c + (x-1)^{1+a+b-c}}{x(x-1)}aby = 0$$

(iii) Self-adjoint form, integrating factor  $-e^{4x}$ ,

$$-\frac{\mathrm{d}}{\mathrm{d}x}\left(e^{4x}\frac{\mathrm{d}y}{\mathrm{d}x}\right) - 4e^{4x}y = \lambda e^{4x}y$$

weight function is hence  $e^{4x}$ .

Easier to consider original equation; assuming the solution takes the form  $y \propto e^{\sigma x}$ ,  $\sigma$  satisfies the auxillary equation

$$\sigma^2 + 4\sigma + 4 + \lambda = 0 \Rightarrow \sigma = -2 \pm i\sqrt{\lambda},$$

$$y(x) = Ae^{-2x}\cos(\mu x) + Be^{-2x}\sin(\mu x)$$

where A and B are constants, and  $\mu^2 = \lambda$ . Applying the boundary conditions, y(0) = 0 implies that A = 0. The other boundary condition implies

$$Be^{-2}\sin\mu = 0$$
$$\Rightarrow \mu = n\pi$$

Thus infinite positive eigenvalues  $\lambda_n = n^2 \pi^2$ 

The associated eigenvectors are thus proportional to  $e^{-2x}\sin(n\pi x)$ .

Eigenvectors associated with distinct eigenvalues are indeed orthogonal on the interval, if the weight function  $e^{4x}$  is correctly included in the inner product integral  $I_{mn}$ ,  $(m \neq n)$  defined as

$$I_{mn} = \int_0^1 e^{4x} Y_n(x) Y_m(x) dx$$

where  $Y_n$  are  $Y_m$  are normalized eigenfunction with distinct eigenvalues  $\lambda_n=n^2\pi^2$  and  $\lambda_m=m^2\pi^2$ 

(i) Using L to denote the operator,

$$L := \frac{\mathrm{d}}{\mathrm{d}x} \left( x \frac{\mathrm{d}u}{\mathrm{d}x} \right)$$

the Sturm-Liouville form is

$$Lu = -\lambda xu, \quad 0 < x < 1$$

hence with weight function x.

We seek a linear substitution to turn this into Bessel's equation of order zero. We cannot make a substitution for u as the linearity of the operator makes this redundant; after trying a few things, we see that  $x=\frac{z}{\sqrt{\lambda}}$  is the way forward. Turns it into

$$\frac{\mathrm{d}}{\mathrm{d}z} \left( z \frac{\mathrm{d}}{\mathrm{d}z} \right) y = -zy$$

which is Bessel's equation, and we are told the general solution is of the form

$$y(z) = AJ_0(z) + B[R(z) + J_0(z)\log(z)]$$

where A, B are constants and R(z) is a 'regular function'.

Now we have u(x) bounded as  $x \to 0$ , same must be true for y(z) as  $z \to 0$ . From the series definition of  $J_0$ , we know that  $J_0(0) = 1$ . Hence as  $\log(z) \to -\infty$  as  $z \to 0$ , we have B = 0, concluding that  $y(z) = AJ_0(z)$ , ie.

$$u(x) = AJ_0(\sqrt{\lambda}x)$$

and  $u(1) = 0 \Rightarrow AJ_0(\sqrt{\lambda}) = 0$ , so  $\sqrt{\lambda} = j_0$  for  $n = 1, 2, \cdots$ .

Thus the operator L has eigenfunctions  $u_n(x) = J_0(j_n x)$  with eigenvalues  $\lambda_n = j_n^2$ .

(ii) Acting the operator L on its eigenfunction  $J_0(\alpha x)$  given

$$\frac{\mathrm{d}}{\mathrm{d}x} \left( x \frac{\mathrm{d}}{\mathrm{d}x} \right) J_0(\alpha x) = -\alpha^2 x J_0(\alpha x)$$

Multiplying by  $J_0(\beta x)$  and integrating gives

$$\int_0^1 J_0(\beta x) \frac{\mathrm{d}}{\mathrm{d}x} \left( x \frac{\mathrm{d}}{\mathrm{d}x} \right) J_0(\alpha x) \, \mathrm{d}x = -\alpha^2 \int_0^1 x J_0(\alpha x) J_0(\beta x) \, \mathrm{d}x$$

. .

Next part: setting  $\alpha = j_n$ ,  $\beta = j_m$ , we note that  $J_0(j_n) = J_0(j_m) = 0$ , thus the identity follows.

Next: note that our first result is only valid for when  $\beta \neq \alpha$ . So we set  $\alpha = j_n$ ,  $\beta = j_n + \varepsilon$ , and the result should pop out.

(iii) Summarising what we have so far.

$$L := \frac{\mathrm{d}}{\mathrm{d}x} \left( x \frac{\mathrm{d}}{\mathrm{d}x} \right) \qquad \text{with weight } w = x$$

$$\lambda_n = j_n^2, \qquad u_n(x) = J_0(j_n x)$$

orth. relation: 
$$\int_0^1 x J_0(j_n x) J_0(j_m x) dx = \frac{1}{2} \left[ J_0'(j_n) \right]^2 \delta_{mn}$$

To solve

$$Lu + \tilde{\lambda}xu = xf(x) \tag{*}$$

Seek eigenfunction expansions

$$u = \sum_{n=1}^{\infty} a_n J_0(j_n x)$$
  $f(x) = \sum_{n=1}^{\infty} b_n J_0(j_n x)$ 

Substitute into (\*)

$$\sum_{n=1}^{\infty} a_n \underbrace{L[J_0(j_n x)]}_{-j_n^2 x J_0(j_n x)} + \tilde{\lambda} x \sum_{n=1}^{\infty} a_n J_0(j_n x) = x \sum_{n=1}^{\infty} b_n J_0(j_n x)$$

Comparing coefficients (note how the x makes this easy)

$$-a_n j_n^2 + \tilde{\lambda} a_n = b_n$$

$$\Rightarrow a_n = \frac{b_n}{\tilde{\lambda} - j_n^2}$$

Noting that  $\tilde{\lambda}$  is not an eigenvalue, ie.  $\tilde{\lambda} \neq \lambda_n = j_n^2$ .

To find the eigenfunction expansion of u it remains to find  $b_n$  st.

$$f(x) = \sum_{n=1}^{\infty} b_n J_0(j_n x)$$

Multiply by  $xJ_0(j_nx)$  and integrate, thus

$$\int_0^1 x f(x) J_0(j_m x) \, \mathrm{d}x = \sum_{n=1}^\infty b_n \int_0^1 x J_0(j_n x) J_0(j_m x) \, \mathrm{d}x$$
$$= \sum_{n=1}^\infty b_n \frac{1}{2} [J_0'(j_n)]^2 \delta_{mn}$$
$$= b_m \frac{1}{2} [J_0'(j_m)]^2$$

ie.

$$b_n = \frac{2\int_0^1 t f(t) J_0(j_m t) dt}{[J_0 1(j_n)]^2}$$

Hence

$$u(x) = 2\sum_{n=1}^{\infty} \frac{\int_0^1 t f(t) J_0(j_m t) dt}{[J_0 1(j_n)]^2 (\tilde{\lambda} - j_n^2)} J_0(j_n x)$$