# Part IB — Methods Example Sheet 3

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- We are given

$$\underbrace{\ddot{\theta} + 2p\dot{\theta} + (p^2 + q^2)}_{f}\theta = f(t)$$

with  $\theta(0) = \dot{\theta}(0) = 0$ , and  $p > 0, q \neq 0$ .

Want to find G such that  $\mathcal{L}G = \delta(t - \tau)$ , so that for each value of  $\tau$ , the Green's function will solve the homogeneous equation LG = 0 whenever  $t \neq \tau$ .

We construct G for  $0 \le t < \tau$  as a general solution of the homogeneous equation, so that  $G = Ay_1(t) + By_2(t)$ . Have

$$G(t,\tau) = \Theta(t-\tau)e^{-p(t-\tau)}[C(\tau)\cos(q(t-\tau)) + D(\tau)\sin(q(t-\tau))]$$

where  $\Theta$  is the Heaviside step function. Continuity demands  $G(\tau,\tau)=0$  so  $C(\tau)=0$ . The jump condition (with  $\alpha(\tau)=1$ ) enforces  $D(\tau)=\frac{1}{q}$ . Therefore the Green's function is

$$G(t,\tau) = \Theta(t-\tau) \frac{1}{q} e^{-p(t-\tau)} \sin(q(t-\tau))$$

and the general solution to  $\mathcal{L}\theta = f(t)$  obeying  $\theta(0) = \dot{\theta}(0) = 0$  is

$$\theta(t) = \frac{1}{q} \int_0^t e^{-p(t-\tau)} \sin(q(t-\tau)) f(\tau) d\tau$$

 Next we use Fourier Transforms. Taking the Fourier transform of the differential equation gives

$$(i\omega)^2\tilde{\theta} + 2ip\omega\tilde{\theta} + (p^2 + q^2)\tilde{\theta} = \tilde{f}$$

and so

$$\tilde{\theta} = \frac{\tilde{f}}{-\omega^2 + 2ip\omega + (p^2 + q^2)} =: \tilde{R}(\omega)\tilde{f}(\omega)$$

which solves the equation algebraically in Fourier space. Note that

$$\begin{split} \tilde{R}(\omega) &= \frac{1}{-\omega^2 + 2ip\omega + (p^2 + q^2)} \\ &= \frac{1}{(i\omega + p)^2 - (qi)^2} \\ &= \frac{1}{2qi} \left[ \frac{1}{i\omega + p - qi} - \frac{1}{i\omega + p + qi} \right] \end{split}$$

To solve for  $\theta$  in real space we take the inverse Fourier transform to find

$$\begin{aligned} \theta(t) &= \int_0^t R(t-u)f(u) \, \mathrm{d}u \\ &= \int_0^t \left[ \frac{1}{2\pi} \int_{-\infty}^\infty \tilde{R}(\omega) e^{i\omega(t-u)} \, \mathrm{d}\omega \right] f(u) \, \mathrm{d}u \\ &= \frac{1}{q} \int_0^t \underbrace{\left[ \frac{1}{4i\pi} \int_{-\infty}^\infty e^{i\omega(t-u)} \left[ \frac{1}{i\omega + p - qi} - \frac{1}{i\omega + p + qi} \right] \, \mathrm{d}\omega \right]}_{(*)} f(u) \, \mathrm{d}u \end{aligned}$$

which agrees with the first result, if we can show that

$$(*) = e^{-p(t-u)}\sin[q(t-u)]$$

but not sure how to evaluate the integral.

The general homogeneous solution is  $c_1 \sinh \lambda x + c_2 \cosh \lambda x$  so we can take  $y_1(x) = \sinh \lambda x$  and  $y_2(x) = \sinh[\lambda(1-x)]$  as our homogeneous solutions satisfying the boundary conditions at x=0 and x=1 respectively. Then

$$G(x;\varepsilon) = \begin{cases} A(\varepsilon) \sinh(\lambda x) & \text{when } 0 \le x < \varepsilon \\ B(\varepsilon) \sinh[\lambda(1-x)] & \text{when } \varepsilon < x \le 1 \end{cases}$$

Applying the continuity condition we get

$$A \sinh \lambda \varepsilon = B \sinh[\lambda(1-\varepsilon)]$$

while the jump condition gives

$$B(-\lambda \cosh[\lambda(1-\varepsilon)]) - A\lambda \cosh(\lambda \varepsilon) = -1$$

Solving the for A and B gives us the Green's function

$$G(x;\varepsilon) = -\frac{1}{\lambda \sinh \lambda} \left[ \Theta(\varepsilon - x) \sinh[\lambda (1 - \varepsilon)] \sinh \lambda x + \Theta(x - \varepsilon) \sinh(\varepsilon \lambda) \sinh[\lambda (1 - x)] \right]$$

Using this Green's function we are immediately able to write down the complete solution as

$$y = -\frac{1}{\lambda \sinh \lambda} \left\{ \sinh \lambda x \int_{x}^{1} f(\varepsilon) \sinh \lambda (1 - \varepsilon) d\varepsilon + \sinh \lambda (1 - x) \int_{0}^{x} f(\varepsilon) \sinh \lambda \varepsilon d\varepsilon \right\}$$

Use the substitution y = z/x, (quotient rule speeds thing up)

$$L_x y = -\frac{1}{x} \frac{\mathrm{d}^2 z}{\mathrm{d}x^2} + \frac{z}{x}$$

So

$$y = \frac{1}{x} (c_1 \cosh x + c_2 \sinh x)$$
$$y = \frac{1}{x} (c'_1 e^x + c'_2 e^{-x})$$

For solutions that are (a) bounded as  $x \to 0$ , we have  $y = A\frac{1}{x}\sinh x$ . For solutions that are (b) bounded as  $x \to \infty$ , have  $y = B\frac{1}{x}e^{-x}$ .

Green's function satisfying  $L_xG = \delta(x - \xi)$  and both conditions (a), (b) is given by

$$G(x;xi) = \begin{cases} A(\xi) \frac{\sinh x}{x} & \text{if } 0 < x < \xi \\ B(\xi) \frac{e^{-x}}{x} & \text{if } \xi < x \end{cases}$$

Continuity  $\Rightarrow$ 

$$A \sinh \xi = Be^{-\xi}$$

Jump  $(\alpha = -1) \Rightarrow$ 

$$Be^{\xi}(\xi+1) + A(\xi\cosh\xi - \sinh\xi) = \xi^2$$

Solving,

$$A = \xi e^{-\xi}$$
  $B = \xi \sinh \xi$ 

Next, to solve

$$L_x y(x) = \begin{cases} 1 & \text{if } 0 \le x \le R \\ 0 & \text{if } x > R \end{cases}$$

Using this G,

$$y = x^{-1}e^{-x} \int_0^x \xi \sinh \xi \, d\xi + x^{-1} \sinh x \int_x^R \xi e^{-xi} \, d\xi$$

For  $0 \le x \le R$ , solve. (things combine into a sinh) For x > R note the RHS integral vanishes, so solve.

Can use a similar argument to the one in lectures to show that  $G(t,\tau)=0$  whenever  $t\in[0,\tau]$ . Following the usual procedure we get

$$G(t,\tau) = \Theta(t-\tau)[A(\tau)\cos[k(t-\tau)] + B(\tau)\sin[k(t-\tau)] + C(\tau)(t-\tau) + D(\tau)]$$

Continuity demands  $G(\tau,\tau) = G'(\tau,\tau) = G''(\tau,\tau) = 0$ , yielding

$$A(\tau) + D(\tau) = 0$$
  

$$kB(\tau) + C(\tau) = 0$$
  

$$A(\tau) = 0$$

respectively. The third equation also implies  $D(\tau)=0,$  and the jump condition on G''' gives

$$-k^3B(\tau) = 1$$

 $\Rightarrow B(\tau)=-k^{-3}\Rightarrow C(\tau)=k^{-2},$  showing the Green's function is indeed what is required.

Solution to  $\mathcal{L}y = e^{-t}$  is given by

$$y(t) = \int_0^t [k^{-2}(t-\tau) - k^{-3}\sin k(t-\tau)]e^{-\tau} d\tau$$

Under the substitution  $y = \phi[x]$ , hence  $\frac{dx}{dy} = \frac{1}{|\phi'(x)|}$  (monotone increasing) the result becomes easy to show

$$\begin{split} \int_a^b f(x) \delta[\phi(x)] \; \mathrm{d}x &= \int_{\phi(a)}^{\phi(b)} f(\phi^{-1}(y)) \frac{1}{|\phi'(\phi^{-1}(y))|} \delta(y) \mathrm{d}y \\ &= f(\phi^{-1}(0)) \frac{1}{|\phi'(\phi^{-1}(0))|} \\ &= \frac{f(c)}{|\phi'(c)|} \quad \text{as } \phi^{-1}(0) = c \end{split}$$

If monotone decreasing,  $\frac{\mathrm{d}x}{\mathrm{d}y} = -\frac{1}{|\phi'(x)|}$ , but  $\phi(b) < \phi(a)$ , so  $\int_{\phi(a)}^{\phi(b)} = -\int_{\phi(b)}^{\phi(a)}$  and the result holds.

Hence for a general  $\phi(x)$  with simple zeros at  $c_1, c_2, c_3 \cdots, c_N$  we have

$$\int_{a}^{b} f(x)\delta[\phi(x)] dx = \sum_{i=1}^{N} \frac{f(c_i)}{|\phi'(c_i)|}$$
 (\*)

Next for any  $a \in \mathbb{R} \setminus \{0\}$ 

$$\int_{\mathbb{R}} \delta(at)\phi(t) dt = \frac{1}{|a|} \int_{\mathbb{R}} \delta(y)\phi(y \ a) dy = \frac{1}{|a|}\phi(0)$$

so we may write  $\delta(at) = \delta(t)/|a|$ .

Lastly, we apply (\*) for f(x) = |x| and  $\phi(x) = x^2 - a^2$ , which has simple zeros at x = a and x = -a. Get

$$\int_{-\infty}^{\infty} |x| \delta(x^2 - a^2) \, dx = \frac{f(a)}{|\phi'(a)|} + \frac{f(-a)}{|\phi'(-a)|}$$
$$= \frac{|a|}{2|a|} + \frac{|a|}{2|a|}$$
$$= 1$$

as required

(i) f(x) = 1, |x| < c.

$$\begin{split} \tilde{f}(k) &= \int_{-c}^{c} e^{-ikx} \, \mathrm{d}x \\ &= \frac{i}{k} \left[ e^{-ikx} \right]_{-c}^{c} \\ &= \frac{i}{k} (-2i \sin kc) \\ &= \frac{2 \sin kc}{k} \end{split}$$

(ii)  $\mathcal{F}[e^{iax}f(x)] = \tilde{f}(k-a)$ 

we have for  $f(x) = e^{iax}$  and using our previous answer,

$$\tilde{f}(k) = \frac{2\sin[(k-a)c]}{k-a}$$

(iii)  $f(x) = \sin(ax)$ 

Noting that this is the imaginary part of what we've just done (but not seeing how to make use of that... so stuck with calculating ) :

No idiot, use  $\sin(ax) = \frac{1}{2i(e^{iax} - e^{-iax})}$  and linearity of Fourier Transform...

$$\begin{split} \tilde{f}(k) &= \int_{-c}^{c} \sin(ax) e^{-ikx} \, \mathrm{d}x \\ &= \frac{i}{k} \left[ \sin(ax) e^{-ikx} \right]_{-c}^{c} - \frac{ia}{k} \int_{-c}^{c} \cos(ax) e^{-ikx} \, \mathrm{d}x \\ &= \frac{i}{k} 2 \sin(ax) \cos(kc) - \frac{ia}{k} \left( \frac{i}{k} \left[ \cos(ax) e^{-ikx} \right]_{-c}^{c} + \frac{ia}{k} \int_{-c}^{c} \sin(ax) e^{-ikx} \, \mathrm{d}x \right) \\ &= \frac{i}{k} 2 \sin(ax) \cos(kc) + \frac{a}{k^2} \left( -2 \cos(ax) \sin(kc) + a\tilde{f}(k) \right) \end{split}$$

Thus multiplying upon rearranging, we get

$$\tilde{f}(k) = \frac{2i\left(-k\cos[kc]\sin[ac] + a\cos[ac]\sin[kc]\right)}{a^2 - k^2}$$

(iv) Next, by the differentiating property of Fourier Transforms, know that

$$\mathcal{F}[a\cos(ax)] = ik\tilde{f}(k)$$

where  $\tilde{f}(k)$  is the FT of  $f(x) = \sin(ax)$ . Then, by scaling,

$$\mathcal{F}[\cos(ax)] = |a|ik\tilde{f}(k)$$

Thus the Fourier transform of cos(ax) is given by

$$\tilde{f}(k) = \frac{-2ka\left(-k\cos[kc]\sin[ac] + a\cos[ac]\sin[kc]\right)}{a^2 - k^2}$$

Have that

$$f(x) = \begin{cases} 0 & \text{if } x < 0 \\ e^{-x} & \text{if } x > 0 \end{cases}$$

The Fourier transform is given by

$$\begin{split} \tilde{f}(k) &= \int_0^\infty e^{-(ik+1)x} \, \mathrm{d}x \\ &= \frac{1}{-ik-1} \left[ e^{-(ik+1)x} \right]_0^\infty \\ &= \frac{1}{1+ik} \\ &= \frac{1-ik}{1+k^2} \end{split}$$

The inverse Fourier transform is given by

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ikx} \tilde{f}(k) \, dk$$

Thus

$$f(0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1 - ik}{1 + k^2} dk$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{1 + k^2} dk - \frac{i}{2\pi} \int_{-\infty}^{\infty} \frac{k}{1 + k^2} dk$$

$$= \frac{1}{2\pi} \left[ \arctan k \right]_{-\infty}^{\infty} - \frac{i}{4\pi} \underbrace{\left[ \log(1 + k^2) \right]_{-\infty}^{\infty}}_{=0}$$

$$= \frac{1}{2\pi} \left( \frac{\pi}{2} + \frac{\pi}{2} \right)$$

$$= \frac{1}{2}$$

Wait. Note that

This is easy. Define  $g(x)=e^{-x^2}$ , calculate  $\tilde{g}(k)$ , then do rephasing and scaling in any order. Is good practice. Should get

$$\tilde{f}(k) = e^{-ik\mu} \frac{1}{n} \sqrt{\pi} e^{-k^2/4n^2}$$

$$f(x) = e^{-n^2(x-\mu)^2}$$

$$f(x + \mu) = e^{-n^2 x^2}$$
$$f'(x + \mu) = -2n^2 x e^{-n^2 x^2}$$

Calculating the Fourier Transform of the last function gives

$$\tilde{f}(k) = -2n^2 \int_{-\infty}^{\infty} x e^{-n^2 x^2} e^{-ikx} dx$$

Let us suppose we have N measurements of a function h(t), where N is even, with constant sampling interval  $\Delta$ , ie. we have the set of measurements

$$h_m = h(t_m), \quad t_m = m\Delta, \ m = 0, 1, \dots, N-1$$

Parseval's theorem for DFT is

$$\sum_{m=0}^{N-1} |h(t_m)|^2 = \frac{1}{N} \sum_{n=0}^{N-1} |\tilde{h}_d(f_n)|^2 \qquad (*)$$

where  $h_d$  is the Discrete Fourier transform

$$\tilde{h}_d(f_n) = \sum_{n=0}^{N-1} h_m \exp\left[\frac{-2\pi i}{N} mn\right]$$

and

$$f_n = \frac{n}{N\Lambda}, n = -N/2, \cdots, N/2,$$

To prove (\*), consider a fixed  $h(t_m)$  on the LSH. Applying the inversion formula, and making a Riemann approximation to the integral, we obtain

$$h(t_m) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i\omega t_m} \tilde{h}(\omega) d\omega$$

$$= \int_{-\infty}^{\infty} e^{2\pi i f t_m} h(f) df \qquad (2\pi f = \omega)$$

$$\approx \frac{\Delta}{\Delta N} \sum_{n=0}^{N-1} \tilde{h}_d(f_n) e^{2\pi i f_n t_m}$$

$$= \frac{1}{N} \sum_{n=0}^{N-1} \tilde{h}_d(f_n) \exp\left[\frac{2\pi i}{N} mn\right]$$

$$= \frac{1}{N} \sum_{n=0}^{N-1} \tilde{h}_d(f_n) \kappa^{mn}$$

where  $\kappa = e^{2\pi i/N}$  is an Nth root of unity. Now

$$|h(t_m)|^2 = (h(t_m))^* h(t_m) = \frac{1}{N^2} \sum_{p,q=0}^{N-1} \tilde{h}_d^*(f_p) \kappa^{-mp} \tilde{h}_d(f_q) \kappa^{mq}$$

$$= \frac{1}{N^2} \sum_{p,q=0}^{N-1} \tilde{h}_d^*(f_p) \tilde{h}_d(f_q) \delta_{pq}$$

$$= \frac{1}{N^2} \sum_{p=0}^{N-1} \tilde{h}_d^*(f_p) \tilde{h}_d(f_p)$$

$$= \frac{1}{N} \sum_{p=0}^{N-1} |h_d(f_p)|^2$$

noting the independence from m. Hence, summing over N on both sides gives (\*) as required, as the right hand side just gains a factor of N.

Thus the Fourier transform of cos(x) (Q7 (iv)) is given by

$$\begin{split} \tilde{f}(k) &= \frac{-2k \left(-k \cos[k\frac{\pi}{2}] \sin[\frac{\pi}{2}] + \cos[\frac{\pi}{2}] \sin[k\frac{\pi}{2}]\right)}{1 - k^2} \\ &= \frac{2}{1 - k^2} \cos\left(\frac{k\pi}{2}\right) \end{split}$$

And the Fourier transform of the derivative  $-\sin(x)$  is given by

$$\tilde{f}(k) = -i\frac{2k}{1 - k^2}\cos\left(\frac{k\pi}{2}\right)$$

Inverse Fourier transform of  $\tilde{f}(k)$  is

$$f(x) = \frac{1}{2\pi} \int_{-1}^{1} e^{ikx} e^{k} - e^{-k} dk$$

$$= \frac{1}{\pi} \int_{-1}^{1} e^{ikx} \sinh k dk$$

$$= \frac{1}{\pi} \left[ -\frac{i}{x} e^{ikx} \sinh k \right]_{-1}^{1} + \frac{i}{x\pi} \int_{-1}^{1} e^{ikx} \cosh k dk$$

$$= \frac{-i}{\pi x} \left( e^{ix} \sinh 1 + e^{-ix} \sinh 1 \right) + \frac{i}{x\pi} \left( \left[ -\frac{i}{x} e^{ikx} \cosh k \right]_{-1}^{1} + \frac{i}{x} \int_{-1}^{1} e^{ikx} \sinh k dk \right)$$

$$= \frac{-2i}{\pi x} \left( \cos x \sinh 1 \right) + \frac{i}{x\pi} \left( -\frac{2i}{x} i \sin x \cosh 1 + \frac{i}{x} f(x) \right)$$

$$\Rightarrow \pi x^{2} f(x) = -2ix \cos x \sinh 1 + 2i \sin x \cosh 1 - f(x)$$

Hence rearranging gives

$$f(x) = \frac{2i}{\pi(1+x^2)}(\cosh 1\sin x - x\cos x \sinh 1)$$

(a long algebraic exercise. good for me?)

Next, take the Fourier transform of  $\nabla^2(x,y) = 0$  with respect to x;

$$\int_{-\infty}^{\infty} \left( \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} \right) e^{-ikx} \, dx = 0$$

$$\Rightarrow 0 = \int_{-\infty}^{\infty} \frac{\partial^2 \phi}{\partial x^2} e^{-ikx} \, dx + \int_{-\infty}^{\infty} \frac{\partial^2 \phi}{\partial y^2} e^{-ikx} \, dx$$

$$= -k^2 \tilde{\phi}(k) + \frac{\partial^2}{\partial y^2} \int_{-\infty}^{\infty} \phi e^{-ikx} \, dx$$

$$= -k^2 \tilde{\phi}(k) + \frac{\partial^2 \tilde{\phi}(k)}{\partial y^2}$$

Solving this second order ODE in k gives

$$\tilde{\phi}(k,y) = A(k)e^{ky} + B(k)e^{-ky}$$

We have the boundary conditions  $\tilde{\phi}(k,0) = \tilde{f}(k)$ , and  $\tilde{\phi}(k,1) = 0$ . This gives

$$A + B = e^k - e^{-k}$$
$$Ae^k + Be^{-k} = 0$$

respectively. Solving gives  $A = -e^{-k}$ ,  $B = e^k$ . Thus, we take the inverse FT to finish.

$$\phi(x,y) = -\frac{1}{2\pi} \int_{-1}^{1} e^{k(y-1)} - e^{-k(y-1)} dk$$