

# Numerical Methods for Solving PDE's Arising in Option Pricing

## *An Overview*

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# Overview

- Review of Black-Scholes PDE for European Options
- Conversion to the Standard Diffusion Equation
- Finite Difference and Finite Elements Methods
- Pricing American options
- Methods for Exotic Options

# The Black-Scholes PDE

Recall from Lecture 5 that we modeled the stock price  $S$  as an Itô process:

$$dS = \mu S \, dt + \sigma S \, dz \quad (1)$$

If  $f(S, t)$  is a function of the stock price, then Itô's Lemma tells us

$$df = \sigma S \frac{\partial f}{\partial S} dz + \left( \mu S \frac{\partial f}{\partial S} + \frac{1}{2} \sigma^2 S^2 \frac{\partial^2 f}{\partial S^2} + \frac{\partial f}{\partial t} \right) dt. \quad (2)$$

(See Lecture 6.)

# The Black-Scholes PDE

The random component  $dz$  in these equations can be eliminated using a simple hedging argument. We obtain the famed **Black-Scholes partial differential equation**:

$$\boxed{\frac{\partial f}{\partial t} + rS \frac{\partial f}{\partial S} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 f}{\partial S^2} - rf = 0.} \quad (3)$$

For now, we will consider a European call option to keep things simple.

# The Black-Scholes PDE

To investigate the existence of solution to (3), we rewrite the equation in operator notation:

$$\left( \frac{\partial}{\partial t} - r + rS \frac{\partial}{\partial S} + \frac{1}{2} \sigma^2 S^2 \frac{\partial^2}{\partial S^2} \right) C = 0 \quad (4)$$

An operator of this form is called **parabolic**.

PDE theory asserts that in order for this problem to have a unique solution it is sufficient to know the value of the solution at one point in time and at  $S = 0$  and  $S = \infty$ . These are called **boundary conditions**.

# The Black-Scholes PDE

- Usually know value of solution at  $t = 0$ —an *initial* condition. In the Black-Scholes case, however, we know the *final* value of the solution—for a European option the payoff function at time  $t = T$ . (Black-Scholes PDE is “backwards” wrt time.)
- We know that if  $S = 0$ , the option is worthless, so  $C(0, t) = 0$  for  $0 \leq t \leq T$ .
- We also know that if  $S \rightarrow \infty$ , exercise is almost certain. Thus  $C(S, t) \sim S$  as  $S \rightarrow \infty$ .

# The Heat Equation

- The prototypical example of a parabolic PDE is the **heat (or diffusion) equation**.
- The (one-dimensional) heat equation models the heat flow in a uniform medium (i.e., the thermal conductivity does not change with location) over time.
- The heat equation has been extensively studied. If we can transform the Black-Scholes equation into the heat equation, we can appeal to standard results and methods to study the solution to the Black-Scholes equation.

# The Heat Equation

The heat equation on the domain  $(x, t) \in (-\infty, \infty) \times [0, \infty)$  is given by

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \quad (5)$$

$$u(x, 0) = u_0(x), \quad -\infty < x < \infty \quad (6)$$

$$u(x, t) \sim F(x, t), \quad x \rightarrow \infty \quad (7)$$

$$u(x, t) \sim G(x, t), \quad x \rightarrow -\infty \quad (8)$$



# The Heat Equation

The general solution to the heat equation

$$u(x, t) = \frac{1}{2\sqrt{t\pi}} \int_{-\infty}^{\infty} u_0(s) e^{-\frac{(x-s)^2}{4t}} dt \quad (9)$$

$$= u_0(x) * K(x, t), \quad (10)$$

where  $K(x, t)$  is the **heat kernel** or fundamental solution:

$$K(x, t) = \begin{cases} (4\pi t)^{-1/2} e^{-x^2/4t}, & t > 0 \\ 0 & t \leq 0 \end{cases}. \quad (11)$$

and “\*” represents convolution.

# Conversion to the Heat Equation

To solve the Black-Scholes equation analytically, we convert it to a forward diffusion equation using the change of variables

$$S = Ke^x, \quad t = T - 2\tau/\sigma^2, \quad C = Kv(x, \tau) \quad (12)$$

Notice we are effectively now working with  $\log S$  instead of  $S$ . Also notice that  $x$  and  $v$  will be dimensionless.

# Conversion to the Heat Equation

This reduces the Black-Scholes equation to the equation

$$\frac{\partial v}{\partial \tau} = \frac{\partial^2 v}{\partial x^2} + (k - 1) \frac{dv}{dx} - kv, \quad (13)$$

where  $k = r / \frac{1}{2} \sigma^2$  is a dimensionless parameter. The boundary condition becomes

$$v(x, 0) = \max(e^x - 1, 0) \quad (14)$$

# Conversion to the Heat Equation

We can eliminate the  $\frac{dv}{dx}$  term in (13) using an “integrating factor”

$$v(x, \tau) = e^{\alpha x + \beta \tau} u(x, \tau), \quad (15)$$

where  $\alpha = -0.5(k - 1)$  and  $\beta = -0.25(k + 1)^2$ .

# Conversion to the Heat Equation

This reduces (13) to the diffusion equation

$$\frac{\partial u}{\partial \tau} = \frac{\partial^2 u}{\partial x^2} \quad (16)$$

with boundary conditions

$$u(x, 0) = u_0(x) \quad (17)$$

$$u(0, \tau) \rightarrow 0 \quad \text{as } x \rightarrow -\infty \quad (18)$$

$$u(x, \tau) \rightarrow g(x, \tau) \quad \text{as } x \rightarrow \infty, \quad (19)$$

# Conversion to the Heat Equation

where the boundary functions are

$$u_0(x) = \max(e^{\frac{1}{2}(k+1)x} - e^{\frac{1}{2}(k-1)x}, 0) \quad (20)$$

$$g(x, \tau) = e^{\frac{1}{2}(k+1)x + \frac{1}{4}(k+1)^2\tau} \quad (21)$$

We can now find the solution  $u(x, t)$  in terms of the fundamental solution (9). Some messy integration there, but technique is the same as that used in Lecture 6.

# Conversion to the Heat Equation

To find the solution to the Black-Scholes equation, we simply transform our solution  $u(x, t)$  back to the original variables to yield the solution to the Black-Scholes equation:

$$C(S, t) = S\phi(d_1) - Ke^{-r(T-t)}\phi(d_2) \quad (22)$$

where  $\phi(x)$  is the CDF for the normal distribution and  $d_1$  and  $d_2$  are the constants introduced in Lecture 6.

# Numerical Methods for PDE's

Two major methods:

- Finite Differences—approximate *derivatives*
- Finite Elements—approximate *integrals*

Both methods approximate a PDE by a matrix problem. For most of the problems we will consider here, both give the same matrix problem.



# Finite Difference Method (FDM)

Approximate derivatives by difference quotients. More than one approximation.

In one dimension:

- Forward

$$f'(t) \approx \frac{f(t + \delta t) - f(t)}{\delta t} \quad (23)$$

- Centered

$$f'(t) \approx \frac{f(t + \delta t) - f(t - \delta t)}{2\delta t} \quad (24)$$

- Backward

$$f'(t) \approx \frac{f(t) - f(t - \delta t)}{\delta t} \quad (25)$$

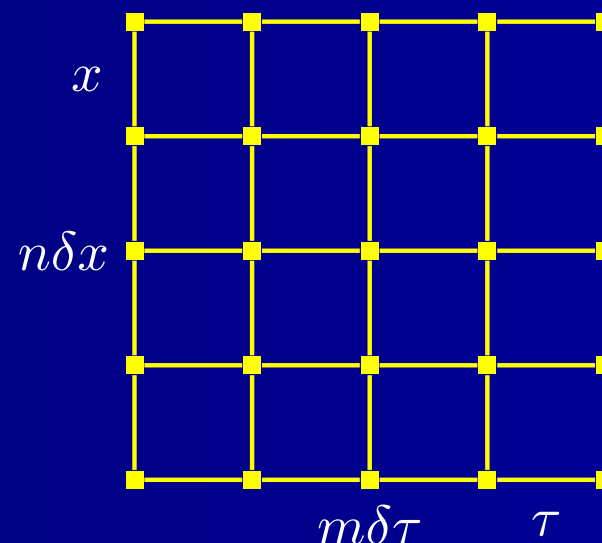
# Finite Differences Method (FDM)

By repeated use of first-order formulas can get formulas for higher order derivatives. We will use a centered difference for  $\frac{\partial^2 u}{\partial x^2}$ :

$$\frac{\partial^2 u}{\partial x^2}(x, t) \approx \frac{u(x + \delta x, t) - 2u(x, t) + u(x - \delta x, t)}{(\delta x)^2} \quad (26)$$

# Creating a Mesh

- Notation:  $u_n^m = u(n\delta x, m\delta t)$
- Issue: discretize  $S$  in untransformed Black-Scholes, or discretize  $\log S$  in transformed equation?
- Stock price has lognormal increments, so  $\log S$  models the underlying Brownian motion better.
- Approx. infinite mesh by finite one. Needs some care.



# Explicit Finite Differences

Discretize  $\frac{\partial u}{\partial t}$  by forward difference,  $\frac{\partial^2 u}{\partial x^2}$  by (26):

$$\frac{u(x, \tau + \delta\tau) - u(x, \tau)}{\delta\tau} \approx \frac{u(x + \delta x, \tau) - 2 * u(x, \tau) + u(x - \delta x, \tau)}{(\delta x)^2} \quad (27)$$

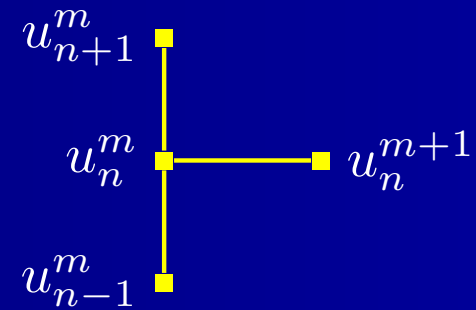
In terms of our mesh (see the previous slide) this equation says that

$$u_n^{m+1} \approx \alpha u_{n+1}^m + (1 - 2\alpha)u_n^m + \alpha u_{n-1}^m \quad (28)$$

where  $\alpha = \delta\tau / (\delta x)^2$ .

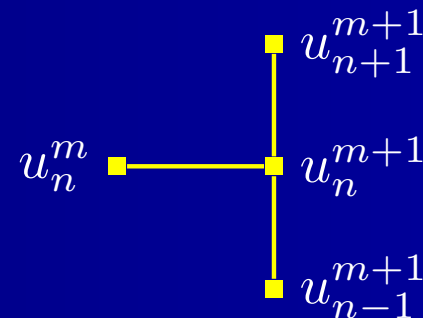
# Explicit Finite Differences

- Notice that  $u_n^{m+1}$  depends on the known quantities  $u_{n-1}^m$ ,  $u_n^m$ , and  $u_{n+1}^m$ . Hence the name “explicit”.
- Notice structure is similar to a trinomial tree (going backwards in time). Up prob. is  $\alpha$ , down prob. is  $\alpha$ , stay put prob. is  $(1 - 2\alpha)$ .
- $\alpha$  is a stability parameter: need  $0 < \alpha \leq \frac{1}{2}$  for to keep rounding errors from growing each step.
- Severely limits the spatial resolution of explicit finite differences.



# Implicit Finite Differences

- Use backward approximation for  $\frac{\partial u}{\partial t}$ .
- Now  $u_{n-1}^{m+1}$ ,  $u_n^{m+1}$ , and  $u_{n+1}^{m+1}$  all depend on  $u_n^m$  in an **implicit** manner.
- Requires us to solve a tridiagonal linear system to find values at the next time step.
- Always stable, for any value of  $\alpha > 0$ .



# Crank-Nicolson Method

- Average explicit and implicit methods
- Stable for all  $\alpha > 0$
- Converges faster than explicit or implicit methods  
( $O((\delta\tau)^2)$  vs.  $O(\delta\tau)$ )
- Again need to solve a linear system.

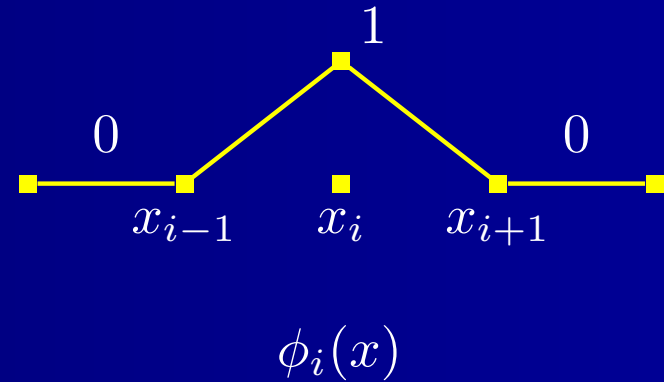
# Finite Element Method (FEM)

- Suppose the solution  $u$  to the heat equation lies in some space (e.g.  $L^2([0, 1])$ ).
- Approximate  $u$  by  $\tilde{u}$  in some finite-dimensional subspace of the original space
- Connected to the idea of a **weak** solution to a PDE: if  $L$  is a diff. operator and  $f$  is a “right-hand side”, find  $u$  s.t.  $\langle Lu, v \rangle = \langle f, v \rangle$  for all **test functions**  $v$ . Here  $\langle \circ, \star \rangle$  denotes integration wrt to some measure.



# Finite Element Method (FEM)

- Discretize spatial domain  $\{x_0, \dots, x_n\}$ .
- A typical test function (a.k.a **shape function**) is shown at right.
- Using  $n - 1$  shape functions, our approximation to  $u$  is



$$u_n(x) = \sum_{i=1}^{n-1} \alpha_i \phi_i(x) \quad (29)$$

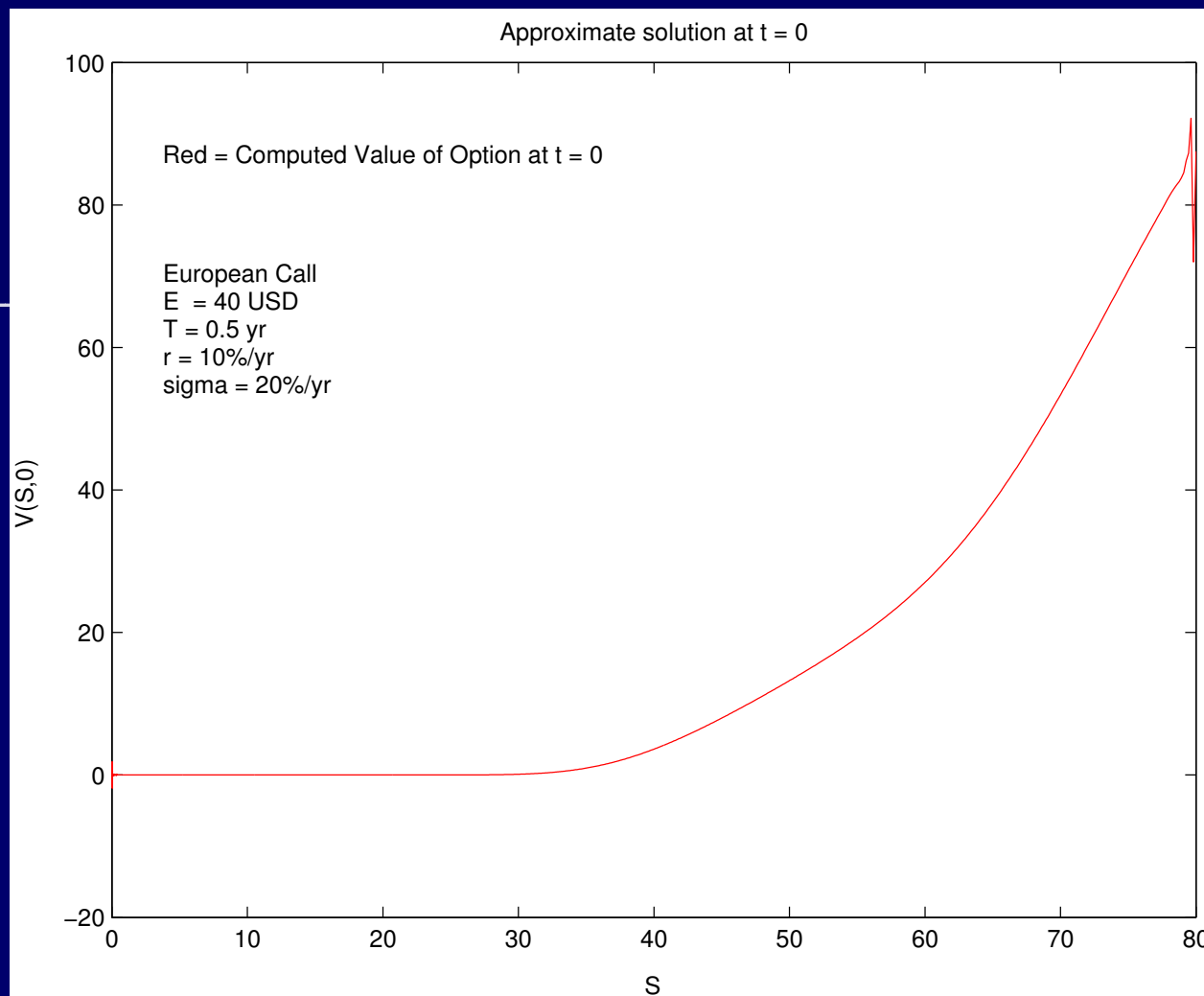
where the  $\alpha_i$ 's are coefficients to be determined by plugging into weak form of PDE.

# Pricing European Calls and Puts

Convert to heat equation, solve, convert back.

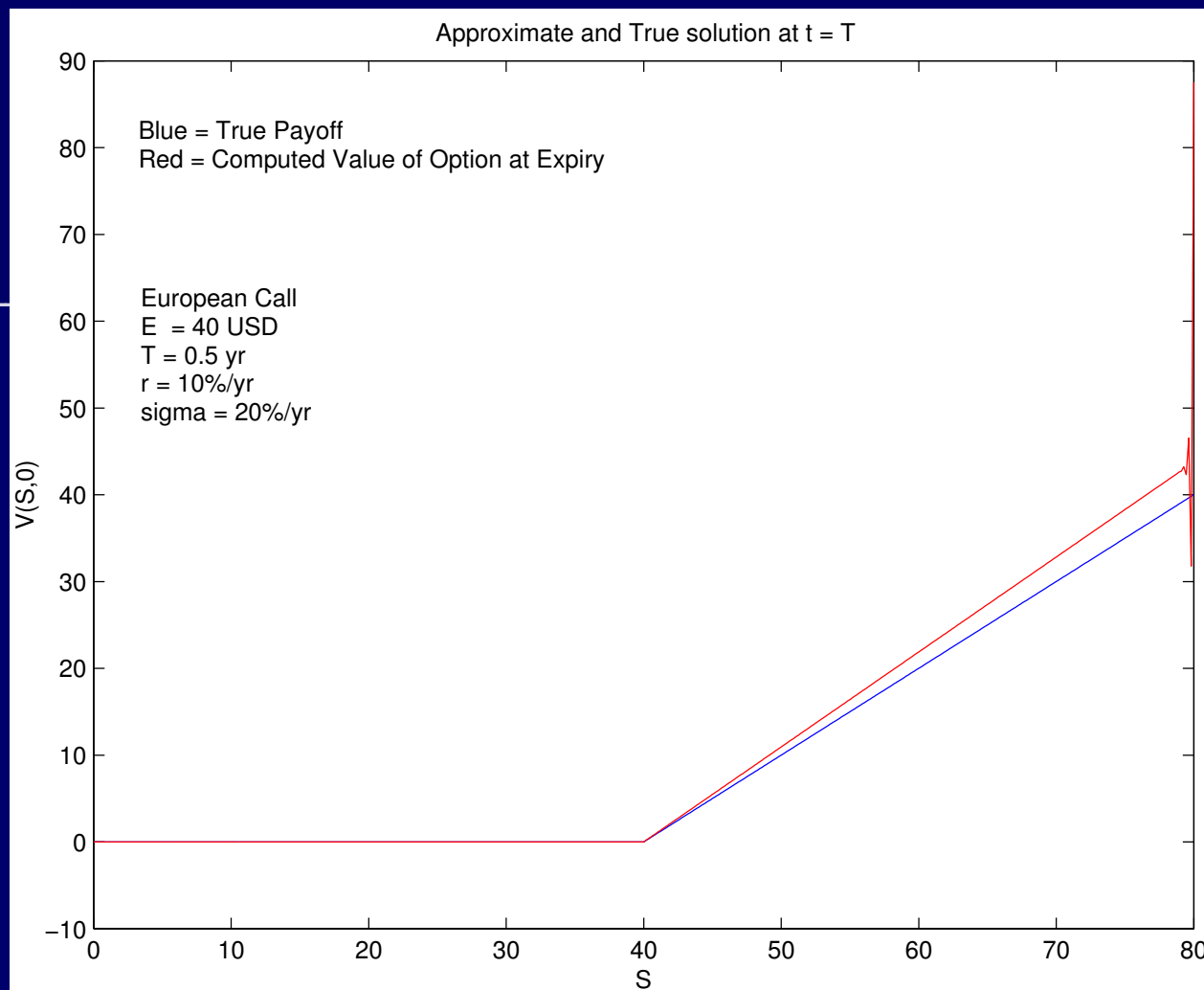
- Could use FDM or FEM in space
- Crank-Nicholson in time
- Large sparse symmetric linear system—use an iterative method (such as conjugate gradients with a good preconditioner) for best results.

Example on next slide.



Price at time  $t = 0$  of a 6-month European call with strike price 40 USD. (See Hull, p249).

The mesh consisted of 5000 points in the  $x$ -direction and 60 points in the  $\tau$  direction. Total MATLAB time was approx. 10 minutes on a (loaded) P4 1.5 GHz with 512 MB RAM. Note the endpoint instability due to the truncation of the infinite domain.



Price at expiry of a the 6-month European call mentioned in the previous slide. (See Hull, p249).

# Pricing American Options

- American options are potentially worth more than their European cousins. Cannot blindly apply the above methods to American options!
- When exercise early is not optimal, the American option should behave just like the European option, i.e., it should satisfy the Black-Scholes equation.
- When early exercise is optimal, arbitrage argument for Black-Scholes only tells us that the return for the portfolio cannot exceed the return on a risk-free deposit.

# Pricing American Options

Thus an American option satisfies a Black-Scholes inequality:

$$\frac{\partial f}{\partial t} + rS \frac{\partial f}{\partial S} + \frac{1}{2} \sigma^2 S^2 \frac{\partial^2 f}{\partial S^2} - rf \leq 0. \quad (30)$$

# Pricing American Options

- By “no arbitrage”, value of option cannot fall below intrinsic value (payoff)
- Thus must have  $P(S, t) \geq \max(K - S, 0)$  for a put.
- When early exercise is optimal, option value is the payoff value.
- Problem: *a priori* we don't know where the exercise boundary is!
- This is an example of a **free boundary problem**—we have a boundary condition, but the location of the boundary must be determined as part of the solution.

# Free Boundary Problems

## Example: The Stefan Problem

- Apply heat to a block of ice at one end. The ice melts, but not uniformly.
- Under some simplifying assumptions, the diffusion equation still holds for the liquid part.
- Boundary between water and ice is unknown *a priori*, changes over time.
- Modeled like the heat equation, but with conditions on the free boundary.



# Pricing American Options

- We make the same change of variables as in the European case to convert the Black-Scholes equation into a diffusion equation.

- Boundary functions are now

$$u_0(x) = \max(e^{\frac{1}{2}(k-1)x} - e^{\frac{1}{2}(k+1)x}, 0) \quad (31)$$

$$g(x, \tau) \geq e^{\frac{1}{4}((k-1)^2 + 4k)\tau} u(x, 0) \quad (32)$$

- Could probably apply numerical methods here, but there is a better way!

# Linear Complementarity Form

- Essentially a “weak” formulation of the problem—recall the construction of the finite element method
- LC form of Black-Scholes for American put:

$$\left( \frac{\partial u}{\partial \tau} - \frac{\partial^2 u}{\partial x^2} \right) (u(x, \tau) - g(x, \tau)) = 0 \quad (33)$$

$$\left( \frac{\partial}{\partial u} \right) \tau - \frac{\partial^2 u}{\partial x^2} \geq 0 \quad (u(x, \tau) - g(x, \tau)) \geq 0 \quad (34)$$

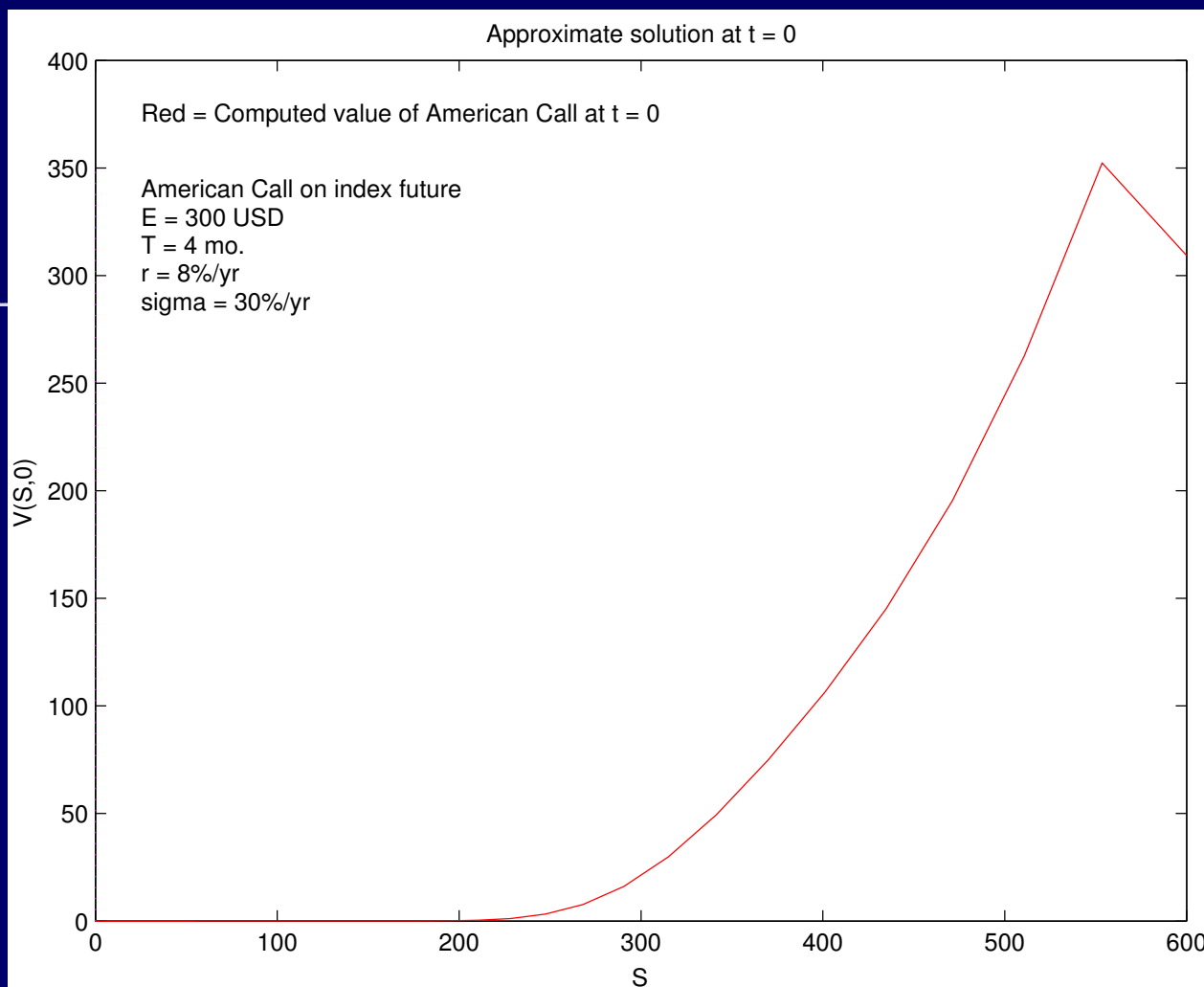
- Advantage: no explicit mention of free boundary!

# Pricing American Options

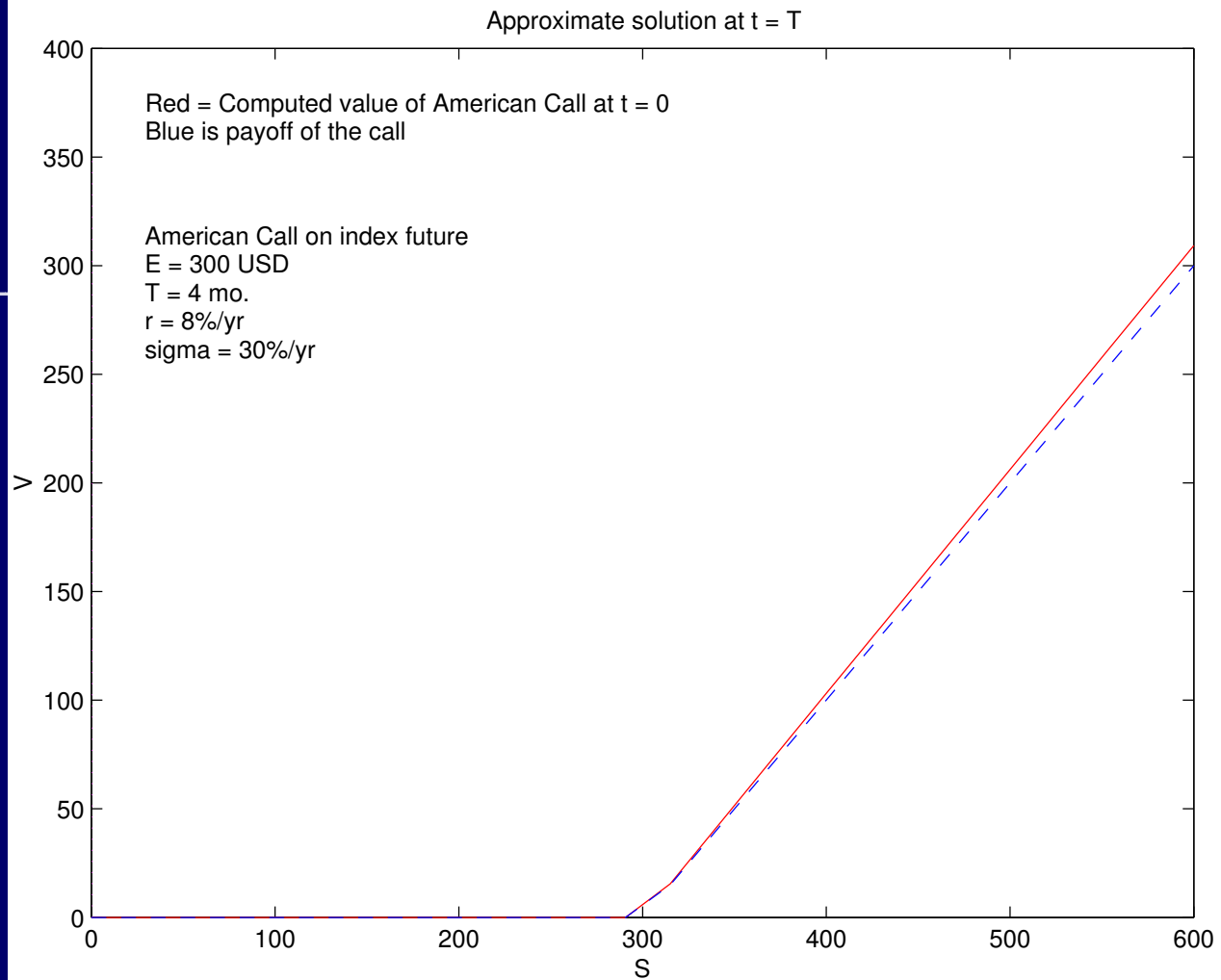
- Discretize LC form using FDM or FEM with Crank-Nicolson.
- Need to solve a constrained matrix problem of the form

$$Ax \geq b, \quad x \geq c, \quad (x - c)^T (Ax - b) = 0$$

- **Projected Successive Over Relaxation** algorithm is the standard tool for this.
- Solve for free boundary after the fact using its definition.



Price at time  $t = 0$  of a 4-month American call on index futures with strike price 300 USD. (See Hull, p400). The mesh consisted of 250 points in the  $x$ -direction and 120 points in the  $\tau$  direction. Total MATLAB time was approx. 10 minutes on a (loaded) P4 1.5 GHz with 512 MB RAM. Note the endpoint instability due to the truncation of the infinite domain.



Price at time  $t = T$  of same option.

# Methods for Other Types of Options

- European/American call/put with constant dividend yield. Slight change to above method, very easy.
- European/American Cash-or-nothing call. Different payoff function, but an identical methodology can be used. Must be careful about the jump discontinuity in the payoff.

# Methods for Other Types of Options

- **Other vanilla barrier options.** Same change of variables, extra boundary condition. Use the “method of images” with the above methodology.
- **Asian options.** Need a little bit of everything!
- **Loopback options.** For European, use the “method of images”. For American, use LC form.

# Methods for Other Types of Options

- Russian options (lookback with infinite time horizon). LC form seems to be a good strategy in most setups.
- Passport options (call on the balance of a trading account—see Wilmott's web site).  
Hamilton-Jacobi-Bellman equations, multiple free boundaries...
- Parisian options (barrier, but barrier only active if price remains beyond the barrier for prespecified time). Wilmott has a solution using basic methods we described above.



# Questions for Further Study

Some questions for further study.

Conservation Law for American Options?

# References

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