

# Statistical mechanical models for stabiliser codes subject to correlated noise

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THE UNIVERSITY OF  
SYDNEY



EQUIS

# Statistical mechanical mapping

Stabiliser code  
& iid noise model



Disordered statistical  
mechanical model

Error-correcting  
threshold



Quenched phase  
transition

Maximum likelihood  
decoding



Calculating partition  
functions/free energy  
differences

Allows us to reappropriate techniques for studying stat. mech. systems to study quantum codes.

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# Our results

- Generalise mapping to arbitrary codes and correlated noise
- Show how to apply our mapping to circuit noise via the history code
- Numerically demonstrate that mild correlations can lower the threshold of the toric code considerably
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# Independent case

Let  $\llbracket A, B \rrbracket$  be the scalar commutator of two Paulis,  $AB =: \llbracket A, B \rrbracket BA$ .

For a stabiliser code generated by  $\{S_k\}_k$ , and an error Pauli  $E$ ,

$$H_E(\vec{s}) := - \sum_i \sum_{\sigma \in \mathcal{P}_i} J_i(\sigma) \llbracket \sigma, E \rrbracket \prod_{k: \llbracket \sigma, S_k \rrbracket = -1} s_k$$

for some couplings  $J_i(\sigma) \rightarrow \mathbb{R}$ .

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Suppose we have an independent error model

$$\Pr(E) = \prod_i p_i(E_i).$$

Nishimori condition: 
$$\beta J_i(\sigma) = \frac{1}{4} \sum_{\tau \in \mathcal{P}} \log p_i(\tau) \llbracket \sigma, \tau \rrbracket ,$$

Using the Fourier-like orthogonality relation  $\frac{1}{4} \sum_{\sigma} \llbracket \sigma, \tau \rrbracket = \delta_{\tau, I}$  we get that

$$-\beta H_E(\vec{1}) = \sum_i \sum_{\tau} \log p_i(\tau) \frac{1}{4} \sum_{\sigma, \tau} \llbracket \sigma, E\tau \rrbracket = \sum_i \log p_i(E_i) = \log \Pr(E).$$

The previous symmetry gives us that

$$Z_E = \sum_{\vec{s}} e^{-\beta H_E(\vec{s})} = \sum_S e^{-\beta H_{ES}(\vec{1})} = \sum_S \Pr(ES) = \Pr(\bar{E}).$$

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# Error correction threshold as a quenched phase transition

Consider the free energy cost of a logical operator  $L$ ,

$$\Delta_E(L) = -\frac{1}{\beta} \log Z_{EL} + \frac{1}{\beta} Z_E = \frac{1}{\beta} \log \frac{\Pr(E)}{\Pr(EL)}.$$

For a fixed error  $E$

	Quantum code	Stat. mech. system
Below threshold	$\{\Pr(EL)\}_L$ peaked	$\Delta(L) \rightarrow \infty$ (in mean)
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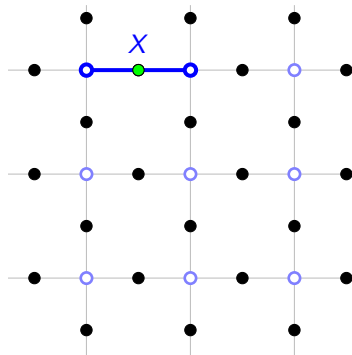
# Independent examples

## Toric code

Bit-flip  $\rightarrow$  Random-bond Ising<sup>1</sup>

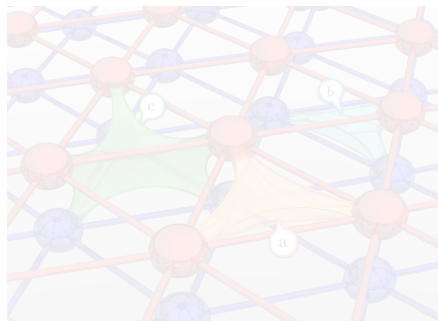
Indep.  $X&Z \rightarrow 2\times\text{RBIM}$  (uncoupled)

Depolarising  $\rightarrow 2\times\text{RBIM}$  (coupled)<sup>2</sup>



## Color code

Depolarising  $\rightarrow$  Eight-vertex model<sup>2</sup>



<sup>1</sup>Dennis et.al., JMP 2002, doi:10/cs2mtf, arXiv:quant-ph/0110143

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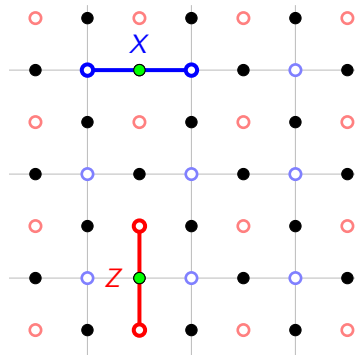
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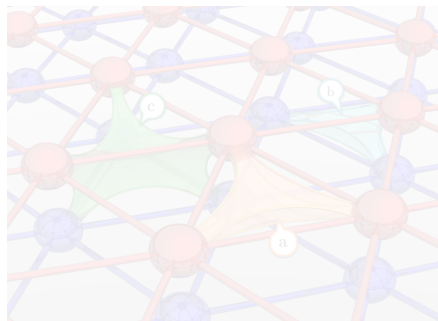
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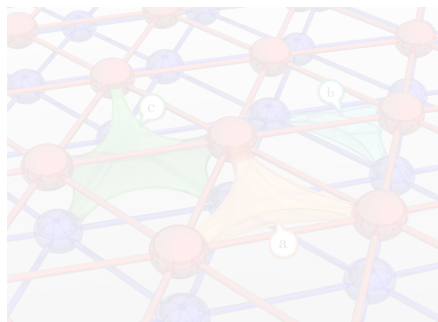
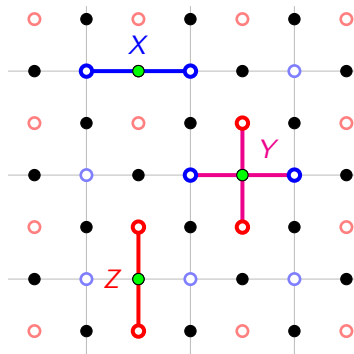
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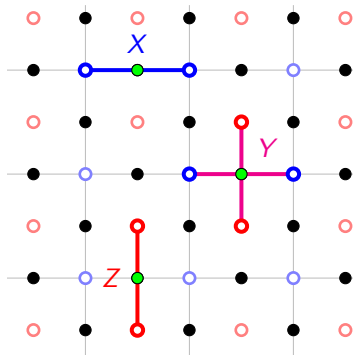
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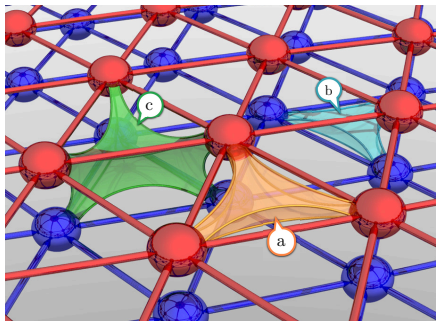
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# Correlated case

The key point independence gave us was the ability to factor our noise model

$$\Pr(E) = \prod_i p_i(E_i).$$

We can generalise this to correlated models:

## Factored distribution

An error model factors over regions  $\{R_j\}_j$  if there exist  $\phi_j : \mathcal{P}_{R_j} \rightarrow \mathbb{R}$  such that

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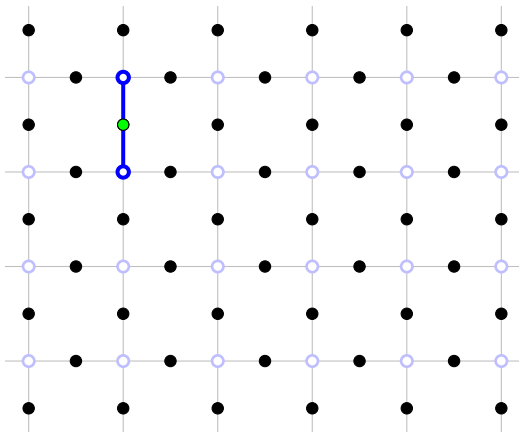
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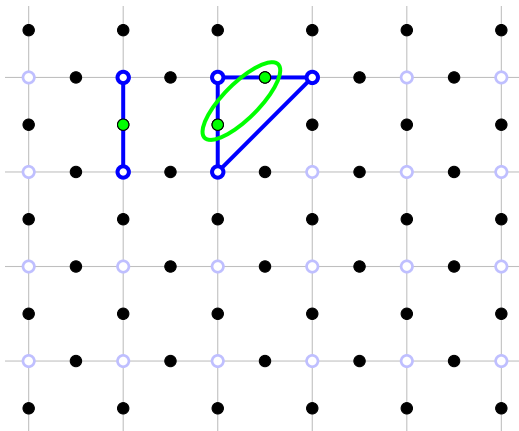
## Toric code with correlated bit-flips

Correlations induce longer-range interactions



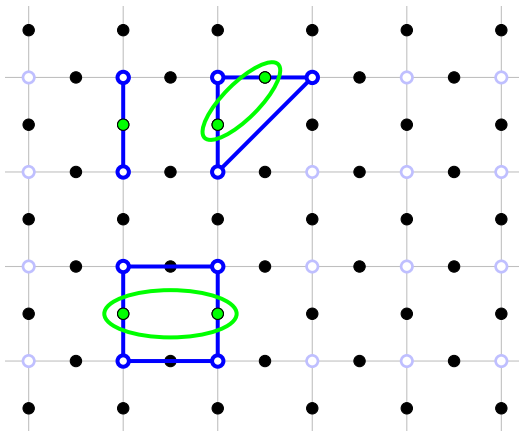
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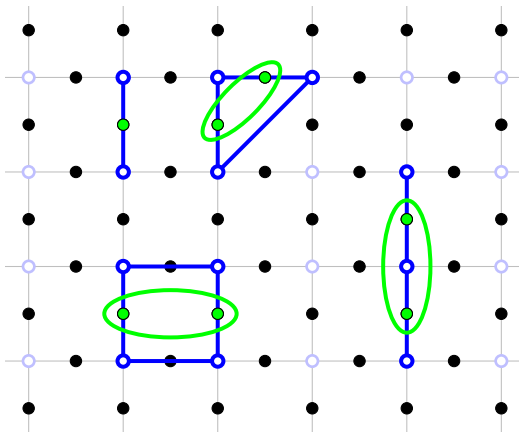
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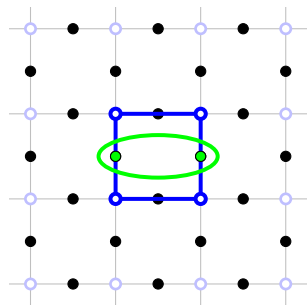
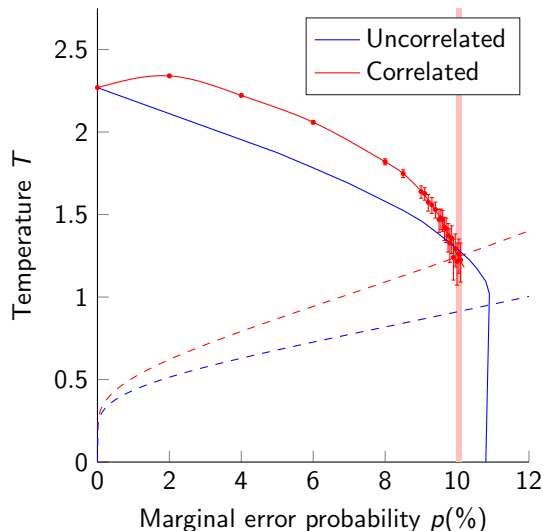


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# Monte Carlo simulations



## Thresholds

Indep.<sup>1</sup>:  $p_{th} = 10.94(2)\%$

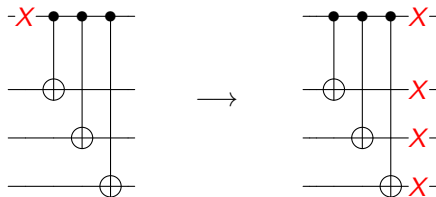
Corr.:  $p_{th} = 10.04(6)\%$

<sup>1</sup>Dennis et.al., JMP 2002, doi:10/cs2mtf, arXiv:quant-ph/0110143

# Circuit noise

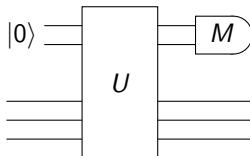
Noise followed by ideal measurements is unrealistic. In reality, circuits will be faulty.

Applying measurement circuits will tend to spread around and correlate noise:



# Circuit noise

We will consider measurement circuits of the form



where  $U$  is a Clifford and  $M$  is a Pauli.

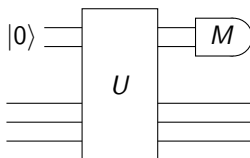
For convenience we only consider independent noise on each circuit. We also will push noise through until after the unitary:





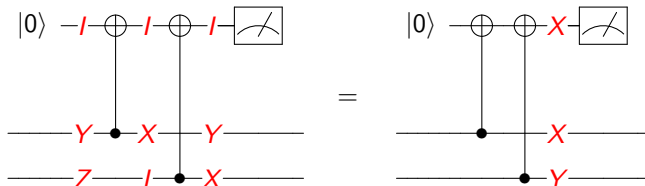
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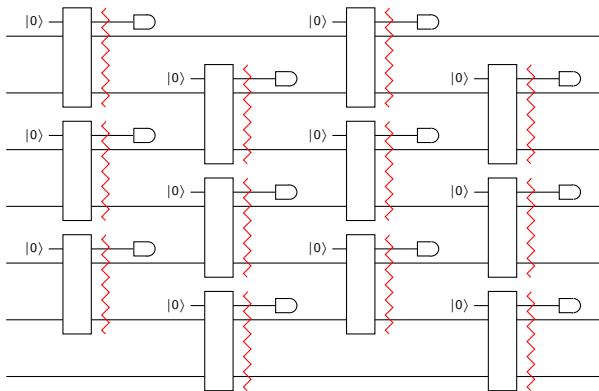


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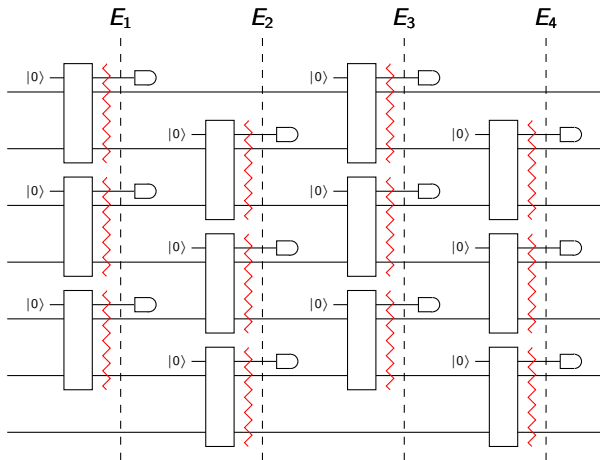


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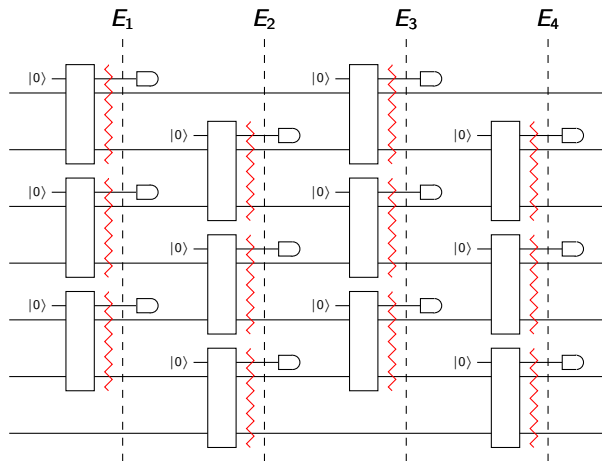
The distribution of the error history  $(E_1, \dots, E_T)$  factors.

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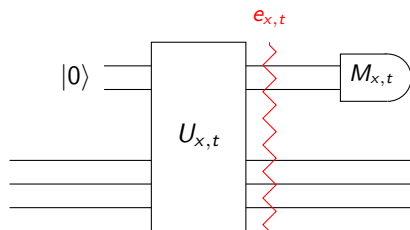
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# Factorising circuit noise



$(E_t)_t$  and  $(e_{x,t})_{x,t}$  are 1-to-1, so

$$\Pr\left((E_t)_t\right) = \Pr\left((e_{x,t})_{x,t}\right).$$

If  $e_{x,t} \sim p_{x,t}$  independently, then

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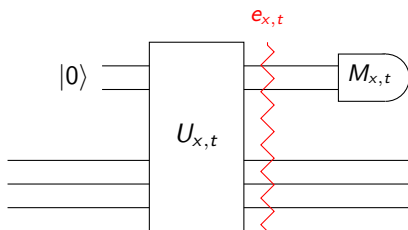
In terms of the accumulated errors:

$$E_t|_{R_{x,t}^+} = e_{x,t} \left( U_{x,t} E_{t-1} |_{R_{x,t}^-} U_{x,t}^\dagger \right)$$

This implies the distribution of error histories factor over  $\{R_{x,t}^+ \cup R_{x,t}^-\}_{x,t}$ :

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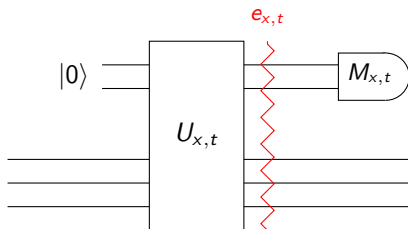
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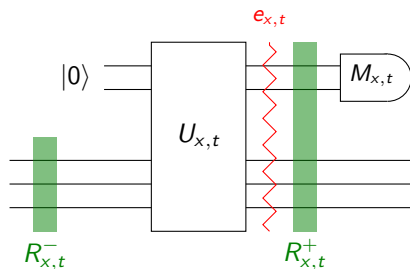
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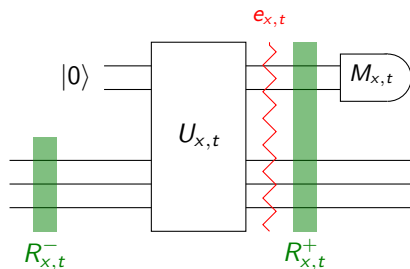
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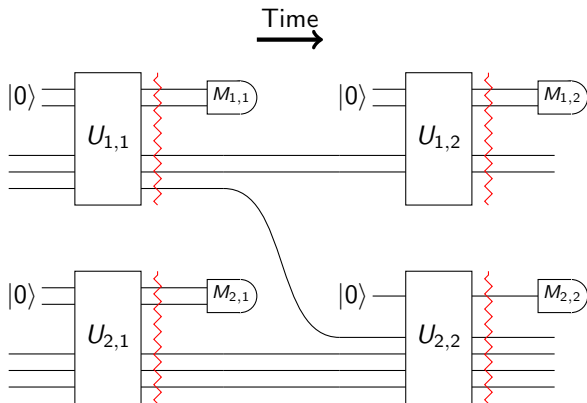
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$$E_t|_{R_{x,t}^+} = e_{x,t} \left( U_{x,t} E_{t-1}|_{R_{x,t}^-} U_{x,t}^\dagger \right)$$

This implies the distribution of error histories factor over  $\{R_{x,t}^+ \cup R_{x,t}^-\}_{x,t}$ :

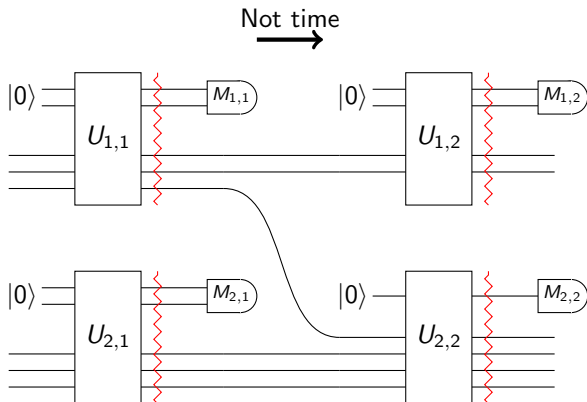
$$\Pr\left((E_t)_t\right) = \prod_{x,t} p_{x,t} \left( E_t|_{R_{x,t}^+} \left( U_{x,t} E_{t-1}^\dagger|_{R_{x,t}^-} U_{x,t}^\dagger \right) \right)$$

# History code



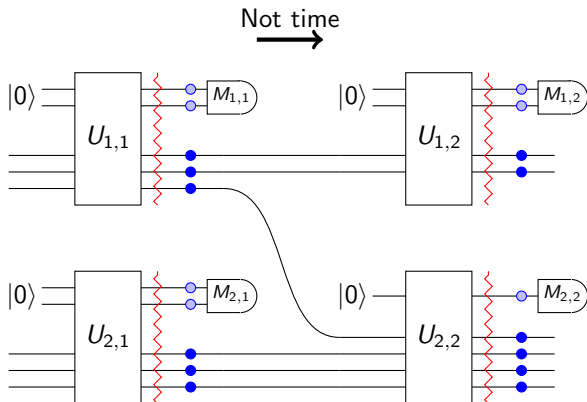
- The error model is given by error history, which factor.
- Stabilisers are the measurements  $(M_{x,t})_{x,t}$ , and stabilisers are final time-slice.
- Logicals are logicals at final time-slice.

# History code



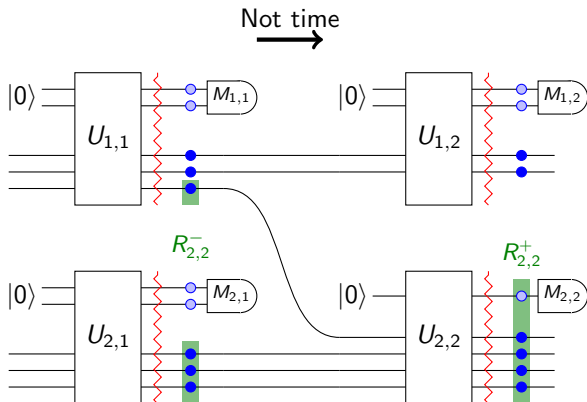
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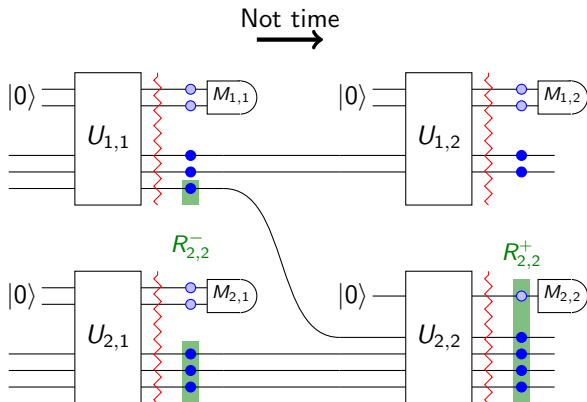
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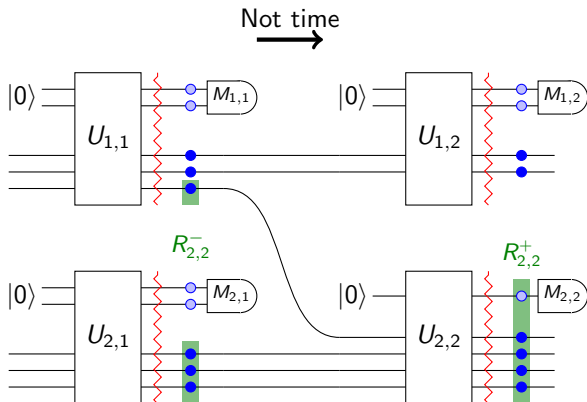
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# Conclusions and further work

- Extended the stat. mech. mapping to correlated models
  - Can apply stat. mech. mapping to circuit noise via the history code
  - Stat. mech. mapping gives tensor network ML decoders
- 
- Can we apply this to experimentally relevant correlated models?
  - Can we use the decoders to understand to better understand the connection between correlation and the threshold (ongoing work).