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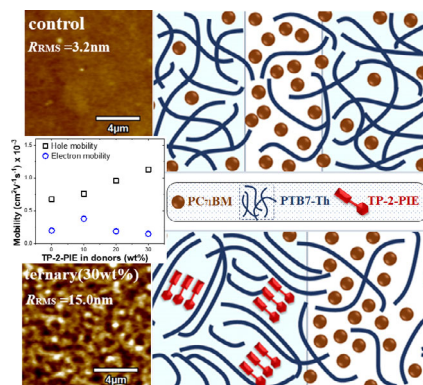
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## RESEARCH ARTICLES

H. Li, H. Gong, M. Sun, H. Zhang, C. Ji,  
C. Liang, F. You, X. Jing, X. Kong,\*  
Z. He\* ..... 2200101

**Tuning Molecular Interaction in Polymer  
Solar Cells via a Multifunctional Discotic  
Component to Enhance Photovoltaic  
Response**



Using a discotic liquid crystal material (TP-2-PIE) having a self-organized property as a third component is able to alter Flory-Huggins molecular interaction parameters in donors and enhance phase separations (poly[4,8-bis(5-(2-ethylhexyl)thiophen-2-yl)benzo[1,2-b:4,5-b']dithiophene-co-3-fluorothieno [3,4-b]-thiophene-2-carboxylate] [PTB7-Th] with TP-2-PIE) and acceptor ([6,6]-phenyl-C<sub>71</sub>-butyric acid methyl ester [PC<sub>71</sub>BM]). This can tune the charge carrier mobility and density of defect states in the active layer and improve device performance.

Q1

# Tuning Molecular Interaction in Polymer Solar Cells via a Multifunctional Discotic Component to Enhance Photovoltaic Response

Han Li, Hongkang Gong, Mengjie Sun, Huimin Zhang, Chao Ji, Chunjun Liang, Fangtian You, Xiping Jing, Xiangfei Kong,\* and Zhiqun He\*

**A donor–acceptor (D–A)-type discotic organic material (TP-2-PIE) having a self-organization ability is selected to be the third component blending with poly[4,8-bis(5-(2-ethylhexyl)thiophen-2-yl)benzo[1,2-b:4,5-b']dithiophene-co-3-fluorothieno [3,4-b]-thiophene-2-carboxylate] (PTB7-Th): [6,6]-phenyl-C<sub>71</sub>-butyric acid methyl ester (PC<sub>71</sub>BM) in ternary bulk-heterojunction polymer solar cells. The complementary absorption and energy transfer between TP-2-PIE and PTB7-Th contributes to the enhancement in photocurrent generation, improving the short-circuit current. In addition, TP-2-PIE alters the molecular interaction leading to an enhanced phase separation, which promotes carrier transport with a good balance and minimizes the density of trap states simultaneously. In this way, the overall photovoltaic performances are markedly enhanced at an optimal condition of 10 wt% TP-2-PIE in donors with a 12.6% efficiency increase. Current understanding on the functionality of this third component in ternary solar cells may be able to guide future device development.**

## 1. Introduction

Device engineering of polymer solar cells (PSCs) have made impressive progress in many aspects in recent years.<sup>[1–8]</sup> Techniques including synthesizing novel materials having better properties, matching optical absorption to maximize light harvesting, manipulating morphology of the active layer, and/or

employing ternary blends for fine tune device performance.<sup>[9–12]</sup> Functional organic electronic or semiconducting materials are among the most promising classes being developed due to their advantages in low-cost fabrication, light weight, and potential flexibility. Organic electronic materials, in particular, those having liquid crystalline (LC) phases, play a crucial role in tuning photovoltaic device performance or exploring the device physics. Although rod-like, calamitic LCs have been substantially exploited in display applications, whereas disc-like or discotic LC (DLC) finds more interest in the field of organic optoelectronic applications. The DLC materials possess many advantages, such as their ease to self-organize into  $\pi$ – $\pi$ -stacked columnar assemblies via strong molecular interactions; their charge carrier transport can

be many orders of magnitude higher in a particular direction, for example, an excellent exciton diffusion length ( $\approx 70$  nm) and a very high charge carrier mobility ( $\approx 1$  cm<sup>2</sup> V<sup>−1</sup> s<sup>−1</sup>) in hexahexylthiotriphenylene<sup>[13]</sup> due to a much stronger intra-column molecular interaction than that of intercolumns.<sup>[13–15]</sup> In addition, DLCs display self-healing of structural defects owing to their LC nature.<sup>[16]</sup>

Coincidentally, the early organic solar cell with a power conversion efficiency (PCE) about 1% reported by Tang in 1986 was prepared by a heterojunction structure consisting of two discotic molecules, copper phthalocyanine, and perylene tetracarboxylic derivative.<sup>[17]</sup> The pioneer work carried out using a well-defined hexa-peri-hexabenzocoronene (HBC)-based DLC to prepare heterojunction photovoltaic device (PCE  $\approx 2\%$ ) in combination with perylene diimide was reported by Schmidt-Mende et al. in 2001.<sup>[18]</sup> DLC hexaacetoxytriphenylene has been effectively used as an additive at 3 wt% in P3HT:PC61BM-based bulk-heterojunction (BHJ) solar cells leading to an enhancement in PCE from 3.03% to 3.97%.<sup>[19]</sup> However, the early LC materials (either rod-like or discotic) have a limited optical absorption due to relative simple core structure. They are not ideal components to directly use either as a donor or as an acceptor. A breakthrough was made with a number of small molecules as a third component. In 2015, a highly crystallized small organic molecule (*p*-DTS-(FBTTH<sub>2</sub>)<sub>2</sub>) as a third component was found to induce a  $\pi$ – $\pi$  stacking of poly[4,8-bis(5-(2-ethylhexyl)thiophen-2-yl)]

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benzo[1,2-b:4,5-b']dithiophene-co-3-fluorothieno [3,4-b]thiophene-2-carboxylate] (PTB7-Th) perpendicular to the substrate, with a face-on orientation, and achieved a PCE up to 10.5%.<sup>[20]</sup> Similarly, a nematic LC benzodithiophene terthiophene rhodanine (BTR) material was developed by Sun et al. in 2015, which has a broad enough optical absorption and is able to form highly oriented structures via solvent annealing. The photovoltaic device prepared with BTR as a donor and [6,6]-phenyl-C<sub>71</sub>-butyric acid methyl ester (PC<sub>71</sub>BM) as an acceptor achieved a PCE up to 9.3%.<sup>[21]</sup> Later, Ma et al. in 2017 employed BTR as a third component at 10 wt%, which raised the PCE of a PTB7-Th:PC<sub>71</sub>BM-based device up to 10.83%.<sup>[22]</sup> Only 5 wt% of BTR as a third component was able to boost PCE of PM6:Y6-based device to 16.6%.<sup>[23]</sup> More recently, by introducing highly crystalline molecules (IDIC as a third component and DTT as an additive) into a novel PTQ10: BTP-FTh blend, a record high efficiency of 19.05% can be realized.<sup>[24]</sup>

With the progressive enhancement in PCE of BHJ PSCs, however, the low crystallinity of the donor polymer, the difficulty in donor/acceptor phase separation, as well as the weak light absorption in active layers leave a legacy of difficulty in PCE improvement.<sup>[25]</sup> By optimization of the active-layer morphology, as well as broadening of light absorption, the short-circuit density ( $J_{sc}$ ) and fill factor ( $FF$ ) of PSCs could be increased.<sup>[26–28]</sup> Solvent or small molecular additive in an active layer has been a common method, in addition to thermal annealing, to manipulate morphology or crystallinity of the active layers.<sup>[29,30]</sup> Liquid crystal components have been regarded as a very promising selection in tuning the morphology of the active layers, in particular, the phase separation of donor and acceptor blends, thanks to its natural feature of self-assembly.<sup>[19,26,29]</sup>

In this work, a DLC TP-2-PIE<sup>[31]</sup> was selected as a third component in combination with PTB7-Th as a donor and PC<sub>71</sub>BM as an acceptor to fabricate PSCs. TP-2-PIE was particularly interesting as it has both electron donating and accepting units in the molecular structure. This investigation revealed that with the TP-2-PIE component, the photovoltaic performance of the devices was substantially enhanced. The functionalities of the TP-2-PIE component and physics of the device have been studied.

## 2. Results and Discussion

### 2.1. Polymer Photovoltaic Devices Tuning by a Discotic Liquid Crystal Component

A BHJ polymer photovoltaic cell was initially fabricated with an inverted configuration: ITO/ZnO/active layer/MoO<sub>3</sub>/Ag as seen in **Figure 1a**, where the active layer consisted of a PTB7-Th polymer as a donor and fullerene PC<sub>71</sub>BM as an acceptor, which was blended at a weight ratio of 1:1.5. Their molecular structures are shown in **Figure 2**.

Current density–voltage ( $J$ – $V$ ) characteristics of the devices were measured under AM 1.5 irradiation at 100 mW cm<sup>−2</sup> light intensity in ambient conditions as shown in **Figure 3a** and the photovoltaic parameters obtained are listed in **Table 1**. The device exhibits a moderate performance with a short-circuit current ( $J_{sc}$ ) of 17.37 mA cm<sup>−2</sup>, an open-circuit voltage ( $V_{oc}$ ) of 0.81 V, and a  $FF$  of 63.1%. The PCE of the device is about

8.76%. This work is attempting to explore how to further improve the device performance.

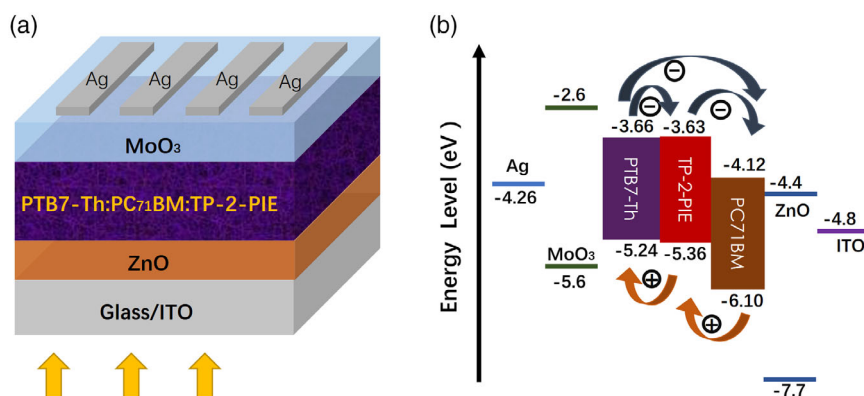
In a BHJ PSC, the selection of organic materials can be a crucial step determining the basic level of photovoltaic performance attainable. Further improvement by matching the material properties and/or tuning device structure can be achieved. The  $J_{sc}$ ,  $V_{oc}$ , and  $FF$  of the devices, for example, are closely dependent on the optical density for photo-harvesting, the energy level of the donor and acceptor materials that affects the driving force for exciton disassociation, and the morphology for charge carrier transport and extraction at the electrodes. Although some device parameters can be intrinsic to the materials used to a great extent, optimization of the device structure may also improve the device performance. A number of approaches can be employed to optimize the active layer or device structures. In this work, we consider a ternary strategy to tune the active layer to enhance photovoltaic performance. Therefore, a discotic liquid crystal compound, TP-2-PIE (see **Figure 2**), having both electron donor (D) and acceptor (A) units of a donor–acceptor (D–A) structure, has been evaluated for its suitability to use as a third component based on the following considerations: 1) light harvesting; 2) energy levels matching; 3) energy or charge transfer; and 4) self-organization.

Optical absorption spectra of the pristine and blended materials were measured as shown in **Figure 3b**. It was found that the TP-2-PIE specimen demonstrated a strong visible absorption with a multiple peak in the range of 430–550 nm, but the PTB7-Th specimen absorbed at a longer wavelength, in between 600 and 750 nm. That is, the two absorptions from TP-2-PIE and PTB7-Th are complementary well to each other. The addition of the TP-2-PIE component into the active layer could substantially enhance its absorption at short wavelength range (see **Figure 3b**). This would benefit the light harvesting in the active layer. Optical bandgap of the TP-2-PIE can be estimated through a Tauc Plot, which is 1.74 eV as seen **Figure S1** in Supporting Information.

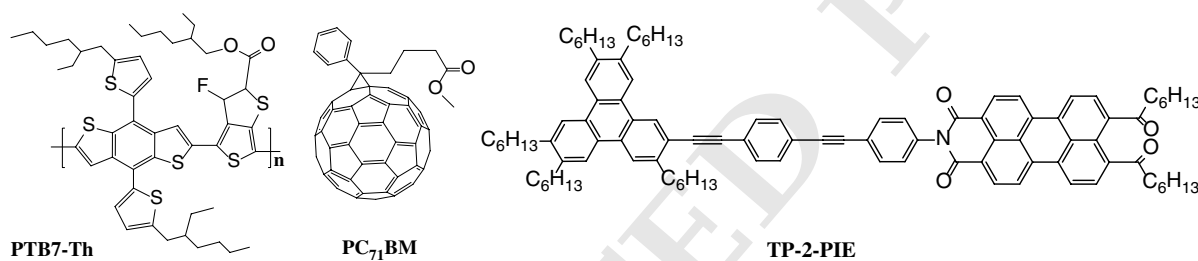
The energy levels of the TP-2-PIE, such as the highest occupied molecular orbital (HOMO) level and the lowest unoccupied molecular orbital (LUMO) level of TP-2-PIE, have been reported in previous work, to be −5.36 and −3.63 eV, respectively, measured by a cyclic voltammetry method.<sup>[31]</sup> An energy diagram including all functional layers is sketched in **Figure 1b**. It can be seen that the energy levels of TP-2-PIE match well with PTB7-Th. If PTB7-Th and PC<sub>71</sub>BM are used as donor and acceptor in the active layer, TP-2-PIE is expected to function as an extra donor according to its energy levels.

What is particularly interesting with the TP-2-PIE material is that this molecule has a triphenylene–perylene dyad structure linked through an ethynylphenyl bridge. However, the electron-donating (D) triphenylene moiety and the electron-accepting (A) perylene unit are decoupled in the TP-2-PIE dyad. According to the molecular simulations, the triphenylene moiety determines HOMO and the perylene monoimide diester unit governs the LUMO level of the TP-2-PIE dyad.<sup>[31]</sup> The D–A-type TP-2-PIE molecule does not respond to the light absorption and excitation as a whole. Instead, the two units are able to respond to the light individually as if they were independent molecules. Nevertheless, the D–A structure ensures the material has dual functionalities, either as a donor or as an acceptor depending on the counterpart species. The electron-rich conjugated bridge might promote intramolecular

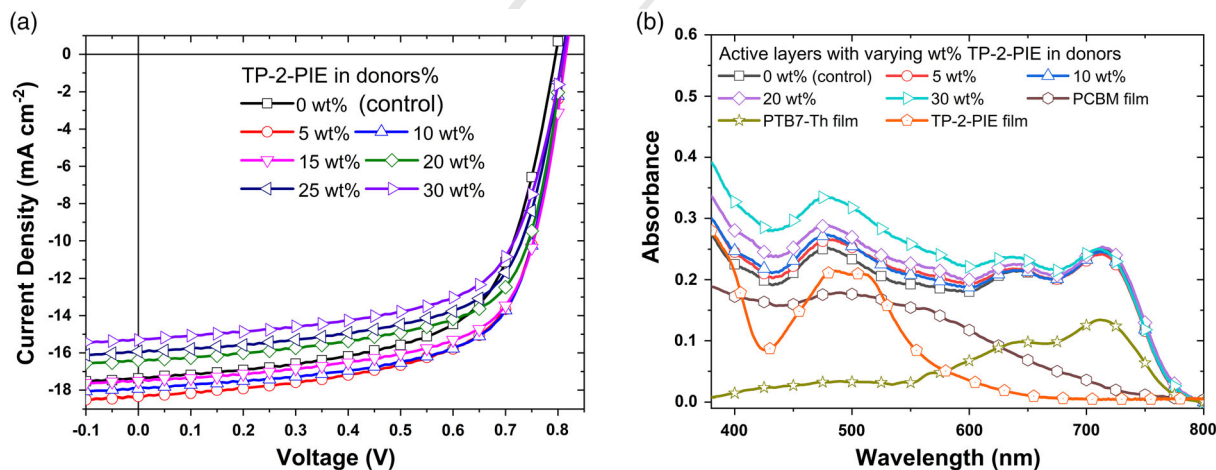




**Figure 1.** a) A schematic sketch of the device structure and b) an energy diagram of the device layers, where the values according to Kong et al. and Hofinger et al.<sup>[31,63,64]</sup> The carrier transport directions are indicated in arrows.



**Figure 2.** Chemical structures of poly[4,8-bis(5-(2-ethylhexyl)thiophen-2-yl)benzo[1,2-b:4,5-b']dithiophene-co-3-fluorothieno[3,4-b]thiophene-2-carboxylate] (PTB7-Th), [6,6]-phenyl-C<sub>71</sub>-butyric acid methyl ester (PC<sub>71</sub>BM), and TP-2-PIE.



**Figure 3.** a) J–V characteristics measured under illumination of the devices, and b) UV-vis absorption spectra of the active layers and pristine specimens.

1 charge transfer within the molecular skeleton and promote charge-  
2 separated states of the molecules.<sup>[32]</sup>

3 In addition, the DLC TP-2-PIE has an advantage for molecular  
4 orientation. A homeotropic alignment is frequently a preferred  
5 orientation.<sup>[15,33,34]</sup> This has the potential to induce an ordered  
6 structure in the aggregate, which may be beneficial for the charge  
7 carrier transporting properties.<sup>[15,31,33]</sup>

8 Based on the earlier considerations, a series of ternary photo-  
9 voltaic devices were fabricated using TP-2-PIE as the third  
10 component to mix with PTB7-Th to form a composite donor

(PTB7-Th:TP-2-PIE). The active layers then become 1  
PTB7-Th<sub>1-x</sub>:TP-2-PIE<sub>x</sub>:PC<sub>71</sub>BM, where the fraction *x* was taken 2  
from 0, 5, 10, 15, 20, 25, and 30 wt%, respectively. The total weight 3  
ratio between donors and acceptors was kept unchanged at 1:1.5. 4

J–V characteristics and the photovoltaic parameters of the 5  
devices were measured as given in Figure 3a and Table 1. 6  
About 20 cells were repeatedly fabricated and measured at each 7  
condition to ensure reproducibility. The statistical data can be 8  
found in Figure S2, Supporting Information. It is clear that 9  
the optimized device was the one with an active layer consisting 10

**Table 1.** Photovoltaic parameters derived from  $J$ - $V$  measurements of devices with different weight ratio of TP-2-PIE in donors.

PTB7-Th:TP-2-PIE: PC <sub>71</sub> BM (weight ratio) <sup>a)</sup>	$V_{OC}$ [V]	$J_{SC}$ [mA cm <sup>-2</sup> ]	$FF$ [%]	$PCE$ [%]	$R_S$ [Ω]	$R_{SH} \times 10^2$ [Ω]
1:0:1.5 (control)	0.81	17.37	63.1	8.76	6.23	4.84
0.95:0.05:1.5	0.81	18.32	66.1	9.80	5.82	6.32
0.90:0.10:1.5	0.81	17.92	68.0	9.86	5.71	6.97
0.85:0.15:1.5	0.81	17.53	66.9	9.61	4.79	6.77
0.80:0.20:1.5	0.81	16.39	67.1	8.90	4.92	6.78
0.75:0.25:1.5	0.81	15.92	65.9	8.50	5.91	6.41
0.70:0.30:1.5	0.81	15.32	64.7	8.02	6.66	5.80

<sup>a)</sup>PTB7-Th: poly[4,8-bis(5-(2-ethylhexyl)thiophen-2-yl)benzo[1,2-b:4,5-b']dithiophene-co-3-fluorothieno [3,4-b]-thiophene-2-carboxylate]; PC<sub>71</sub>BM: [6,6]-phenyl-C<sub>71</sub>-butyric acid methyl ester; FF: fill factor.

TP-2-PIE in donor were examined (see **Figure 4a**) (EQEs of the whole set of devices can be found in Figure S3, Supporting Information.) There is a distinct increase in photocurrent between 430 and 510 nm spectral ranges from the device having TP-2-PIE component. Obviously, the complementary absorption of TP-2-PIE component contributed to the enhancement of photo-harvesting in this range (Figure 3b). The integrating current ( $J_{int}$ ) can be further determined to be  $J_{int} = 17.13 \text{ mA cm}^{-2}$  for the control and  $17.64 \text{ mA cm}^{-2}$  for the device having 10 wt% of TP-2-PIE in donor, which are comparable to the  $J_{SC}$  measured from the devices (Table 1). The better photo-harvesting increased the photocurrent output from the device having 10 wt% of TP-2-PIE in donor.

Organic photovoltaic device performance is closely related to exciton formation, dissociation, and charge-separation processes. Photoluminescence (PL) can be an indication to probe a photo-excited exciton-formation process. PL spectra from varying fractions of PTB7-Th:TP-2-PIE blends were measured by exciting at  $\lambda_{ex} = 475 \text{ nm}$  (Figure 4b). It can be clearly seen that the pristine TP-2-PIE film emits around 677 nm, while PL from pristine PTB7-Th film was quite different from that from the TP-2-PIE film, in which the latter exhibited a broad emission around  $\approx 837 \text{ nm}$  with a shoulder near  $\approx 765 \text{ nm}$ . The center emission was about 160 nm away from the PL emission of the TP-2-PIE specimen.

From Figure 4b, it was found that the intensities of the PL were progressively enhanced as the TP-2-PIE component was added gradually into the PTB7-Th matrix, and the emission profiles from the blends remained similar to that from the pristine PTB7-Th. This indicates that the TP-2-PIE component had no contribution to the PL spectra of the blends, even though the excitation wavelength (475 nm) was close to the maximum absorption for the TP-2-PIE (a minimum absorption for PTB7-Th). Only PTB7-Th emission can be observed in the PTB7-Th:TP-2-PIE-mixed specimens. However, the peak positions shift gradually to the blue as the fraction of TP-2-PIE increased. It peaks at 826 nm for the specimen with 30 wt% TP-2-PIE in donor with about a 10 nm shift (see Figure 4b).

If comparing the absorption spectrum for the PTB7-Th and the PL spectrum for the TP-2-PIE shown in Figure 4c, it was discovered that the two are heavily overlapped in the region from 600 to 770 nm. This implies that an energy transfer channel between the TP-2-PIE and PTB7-Th components could be opened, which allows transferring the excitation energy in TP-2-PIE to the adjacent PTB7-Th species to enhance PL emission from PTB7-Th. The time-resolved PL (TRPL) spectra of the blends were also evaluated and shown in Figure S4, Supporting Information, which can be fitted by a double exponential decay process with a fast ( $\tau_1$ ) and a slow ( $\tau_2$ ) lifetime as listed in Table S1, Supporting Information. The average decay times ( $\tau_{avg}$ ) of the blends increased progressively from 0.87 to 1.31 ns as the TP-2-PIE concentrations were increased. This also supports the energy transfer mechanism between the two donors.

The missing of TP-2-PIE emission, the enhancement of PTB7-Th emission, and the extended PL decay times when increasing TP-2-PIE concentration in the blends all form evidence for the energy transfer between the TP-2-PIE and PTB7-Th. The existence of such energy transfer would lead to extra exciton

1 of PTB7-Th:TP-2-PIE:PC<sub>71</sub>BM at a ratio of 0.9:0.1:1.5, which  
2 achieved  $V_{OC} = 0.81 \text{ V}$ ,  $J_{SC} = 17.92 \text{ mA cm}^{-2}$ , and  $FF = 68\%$ ,  
3 along with a 12.5% increment in PCE in comparison with that  
4 from the control device. The best PCE in the ternary device  
5 approached 9.86%.

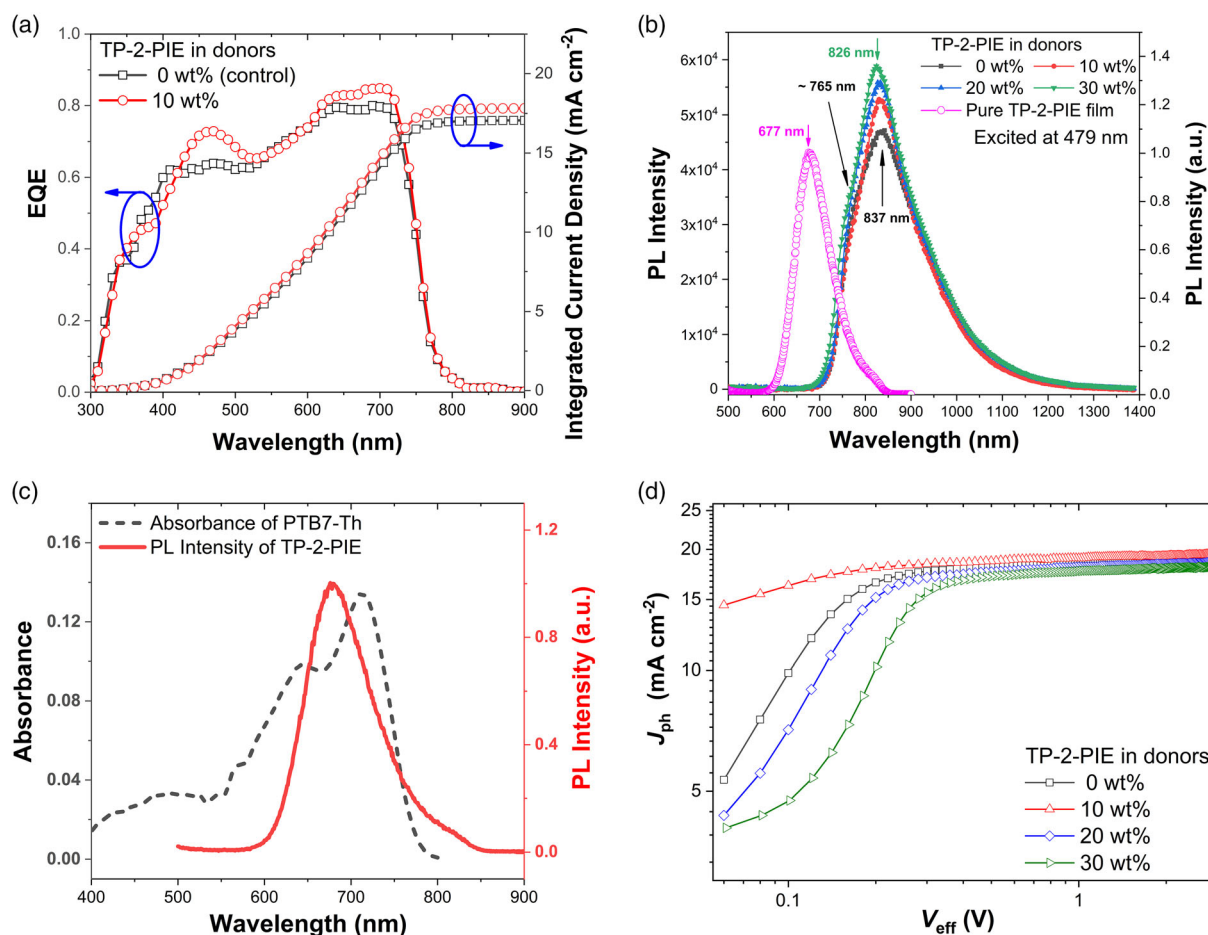
6 It is interesting to note that the addition of TP-2-PIE compo-  
7 nent has no influence on  $V_{OC}$ . The similarity in  $V_{OC}$  of the devi-  
8 ces is consistent with the similarity in energy levels between the  
9 two donors.<sup>[35]</sup> However, the addition of TP-2-PIE component in  
10 the donor does have a strong impact on the  $J_{SC}$ ,  $FF$ , as well as  
11 series and shunt resistances ( $R_S$  and  $R_{SH}$ ) of the devices to some  
12 extent (see Table 1). A small fraction of TP-2-PIE can give rise to  
13 an abrupt increase in both  $J_{SC}$  and  $FF$ . However, the perfor-  
14 mance of the device deteriorated if the fraction of TP-2-PIE  
15 was further increased.

16 A quick way to analyze the change in  $FF$  is to evaluate the  $R_S$   
17 and  $R_{SH}$  variation of the devices, which can be estimated from  
18 the reciprocal of the slope for  $J$ - $V$  curves.<sup>[36]</sup> It was revealed from  
19 Table 1 that the  $R_S$  of the devices decreased gradually upon the  
20 addition of TP-2-PIE component, but it was increased again if the  
21 fraction of TP-2-PIE was further increased (see Table 1). The low-  
22 est  $R_S$  value was  $4.79 \Omega$  obtained in 15 wt% of TP-2-PIE compo-  
23 nent in donor. At the same time, a substantial rise in  $R_{SH}$  of the  
24 devices was found as the TP-2-PIE component was added. An  
25 optimized value of  $R_{SH} = 6.97 \times 10^2 \Omega$  was obtained at a rela-  
26 tively low fraction of 10 wt% TP-2-PIE in donor. The reduction  
27 of  $R_S$  and the increase in  $R_{SH}$  contributed to the improvement in  
28  $FF$ .<sup>[37]</sup> Overall, the optimized  $FF$  was found in 10 wt% of TP-2-  
29 PIE component in donor with over 7% improvement. The reduc-  
30 tion in  $R_S$  indicates a better contact at the electrode of the devices  
31 if the TP-2-PIE component was added in donors. The in-depth  
32 physical mechanism in relation with the TP-2-PIE component  
33 and the device physics is subject to a thorough investigation.

## 34 2.2. Excitation Generation, Dissociation, and Energy Transfer 35 Processes in the Devices

36 To understand the variation in photocurrent, external quantum  
37 efficiencies (EQEs) of the control and the device having 10 wt% of





**Figure 4.** a) External quantum efficiency (EQE) spectra of the cells having different fractions of TP-2-PIE as indicated; b) steady-state photoluminescence (PL) spectra of PTB7-Th:TP-2-PIE films; c) comparison of PL spectrum of a TP-2-PIE film and the absorption spectrum of a PTB7-Th film; d) photocurrent density versus effective voltage characteristics of ternary devices having different fraction of TP-2-PIE additive.

1 formation in the donor layer, which could potentially enhance  
2 photocurrent generation in the devices.

3 The carrier generation and recombination processes of the  
4 active layers were analyzed through dark and illuminated  $J$ - $V$   
5 characteristics of the devices. In photocurrent generation pro-  
6 cess, the generation rate ( $G$ ) of charge carriers can be evaluated  
7 experimentally. The photocurrent ( $J_{ph}$ ) can be obtained by  
8  $J_{ph} = J_L - J_D$ , where  $J_L$  and  $J_D$  are the current densities measured  
9 under a solar simulator light illumination and in the dark,  
10 respectively.<sup>[38]</sup> If plotting  $J_{ph}$  against applied voltage ( $V$ ), a com-  
11 pensation voltage  $V_0$  can be determined at  $J_{ph} = 0$ .<sup>[39]</sup> An effective  
12 voltage ( $V_{eff}$ ) can therefore be obtained through  $V_{eff} = V_0 - V$ . A  
13 relation between  $J_{ph}$  versus  $V_{eff}$  is represented in Figure 4d.

14 It can be seen from Figure 4d that as  $V_{eff}$  increased,  $J_{ph}$  grad-  
15 ually saturated, showing that the bounded  $e$ - $h$  pairs (the excitons)  
16 were separated into free charge carriers. A saturation current  
17 density ( $J_{sat}$ ) was reached at a relatively high  $V_{eff}$  ( $V_{eff} > 2$  V).  
18 In this work,  $J_{sat}$  values were taken at  $V_{eff} = 3$  V and are listed  
19 in Table 2. It was interesting to find that the  $J_{sat}$  values are  
20 20.10 and 19.25 mA cm<sup>-2</sup> for the devices with 5 and the 10 wt  
21 % TP-2-PIE in donors, which are substantially higher than the

**Table 2.** Maximum exciton generation rate ( $G_{max}$ ) and charge dissociation probabilities  $P(E, T)$  of devices with different weight ratio of TP-2-PIE in donors.

PTB7-Th:TP-2-PIE:PC <sub>71</sub> BM (weight ratio)	$J_{sat}$ at 3 V [mA cm <sup>-2</sup> ]	$G_{max} \times 10^{28}$ [m <sup>-3</sup> s <sup>-1</sup> ]	$P_1(E, T)^{a)}$ [%]	$P_2(E, T)^{b)}$ [%]
1.00:0.00:1.5	18.89	1.18	90.8	69.4
0.95:0.05:1.5	20.10	1.26	91.1	71.1
0.90:0.10:1.5	19.25	1.20	93.1	75.1
0.85:0.15:1.5	19.10	1.19	91.7	70.9
0.80:0.20:1.5	17.64	1.10	87.9	69.9
0.75:0.25:1.5	17.35	1.08	87.1	69.7
0.70:0.30:1.5	17.13	1.07	84.4	68.5

<sup>a)</sup>The dissociation probability at short-circuit condition; <sup>b)</sup>the dissociation probability at maximum power condition.

$J_{sat}$  value of 18.89 mA cm<sup>-2</sup> from the binary device without 1  
TP-2-PIE. However, the  $J_{sat}$  dropped if the TP-2-PIE component 2  
was over 20%. This indicates that only a small fraction of the 3

1 TP-2-PIE component present in the active layer could signifi-  
2 cantly enhance  $J_{\text{sat}}$ .

3 Supposing no recombination of free charge carriers if space-  
4 charge effects can be neglected, the  $J_{\text{sat}}$  could be transported and  
5 collected by the electrodes. The  $J_{\text{sat}}$  through the external circuit  
6 can be given by  $J_{\text{ph}} = qLG$ , where  $q$  is the elementary electric  
7 charge and  $L$  is the thickness of the active layer.<sup>[39,40]</sup> The maxi-  
8 mum exciton generation rate ( $G_{\text{max}}$ ) can be evaluated using  
9  $J_{\text{sat}} = qLG_{\text{max}}$  and the results are also listed in Table 2. It was  
10 found that in ternary devices with a relatively low fraction (5–  
11 15 wt%) of TP-2-PIE in donors, the  $G_{\text{max}}$  is higher than that  
12 for the control device. This indicates that the charge separations  
13 were enhanced.

14 Assuming that charge is fully separated at  $J_{\text{sat}}$ , a dissociation  
15 probability  $P(E, T) = J_{\text{ph}}/J_{\text{sat}}$  can be defined, where  $P_1(E, T)$  and  
16  $P_2(E, T)$  are the dissociation probabilities at a short-circuit condi-  
17 tion and at a maximum power condition, respectively.<sup>[39]</sup> The cal-  
18 culated values are listed in Table 2. The highest values for  
19  $P_1 = 93.1\%$  and  $P_2 = 75.1\%$  are from the ternary device having  
20 10 wt% TP-2-PIE in donors, indicating the most efficient condi-  
21 tion for charge separation and dissociation. This is in good agree-  
22 ment with what has been found that the device having 10 wt% of  
23 TP-2-PIE performed the best in photovoltaic response as seen in  
24 Table 1.

25 The results indicate that only a small fraction of TP-2-PIE hav-  
26 ing an extended absorption along with its energy transfer is  
27 enough to enhance charge dissociation, which explains well  
28 the enhancement of photocurrent in photovoltaic devices.

### 29 2.3. The Device Physics in Relation to the Molecular 30 Interaction, Morphology, and Charge Carrier Transport of the 31 Active Layers

32 The charge generation and separation in a donor-/acceptor-  
33 blended organic polymer matrix are generally complicated,  
34 which can be closely related to the morphology of the blended  
35 system, in particular, the formation of two-component-blended  
36 domains. The molecular interactions can be a key factor, which  
37 may alter the miscibility of the components and hence the phase  
38 separation.

39 In the current investigation, surface morphologies of the  
40 blended active layers with different concentration of the TP-2-  
41 PIE component were explored using atomic force microscopy  
42 (AFM) as shown in Figure 5. It was interesting to observe from  
43 the topographic images that the introduction of TP-2-PIE com-  
44 ponent into the initial binary (PTB7-Th:PC<sub>71</sub>BM) film is able  
45 to change the surface morphology distinctly. The surface of  
46 the control film was relatively smooth, but granule-like structures  
47 were developed as the TP-2-PIE fraction was increased. To quan-  
48 tify such a change, root-mean-square roughness ( $R_{\text{RMS}}$ ) of the  
49 active layers was evaluated, which appeared to increase gradually  
50 from 3.2 nm for the control specimen up to 15.0 nm for the layer  
51 having 30 wt% TP-2-PIE in donors. That is, the surface became  
52 rougher with structure developed as the TP-2-PIE fraction was  
53 increased. Obviously, the changes in surface morphology were  
54 induced by the introduction of LC TP-2-PIE component. The  
55 changes in surface morphology might reflect a change of mor-  
56 phology in the bulk, such as molecular aggregation or phase

separation. It is therefore useful to analyze the molecular inter-  
action to see how these blended domains formed.

Due to the PTB7-Th polymer being almost amorphous in  
nature, no distinct crystalline diffractions can be observed in  
wide-angle X-Ray diffraction. A direct observation of changes  
in structural order is difficult. Nevertheless, with the help of a  
Flory–Huggins model,<sup>[41–43]</sup> it is possible to evaluate the molec-  
ular interaction and hence the miscibility in multicomponents.  
Therefore, the phase separation of a blended system can be  
analyzed.

According to Flory–Huggins theory, an interaction parameter  
( $\chi$ ), relating a thermodynamic driving force to the morphology  
formations in an amorphous polymeric system, is a particularly  
important value in the control of miscibility, phase separation,  
and morphological stability.<sup>[44,45]</sup> This could help us to investi-  
gate the role of the TP-2-PIE molecule on the microstructural  
morphology and the phase separation of the donor/acceptor  
blends.<sup>[46]</sup>

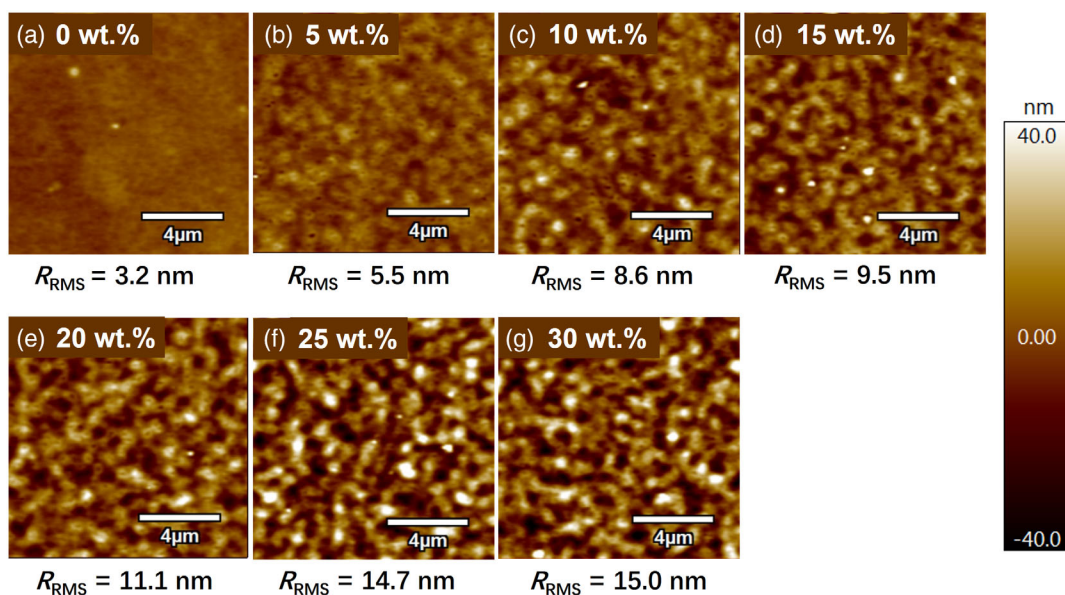
The Flory–Huggins interaction parameter  $\chi$  can be related  
to the cohesive energy density of the components for an amor-  
phous polymer–small molecule blend system through  
Equation (1).<sup>[47,48]</sup>

$$\chi_{ij} = \frac{V_0}{RT} (\delta_i - \delta_j)^2 + 0.34 \quad (1) \quad \text{Q10}$$

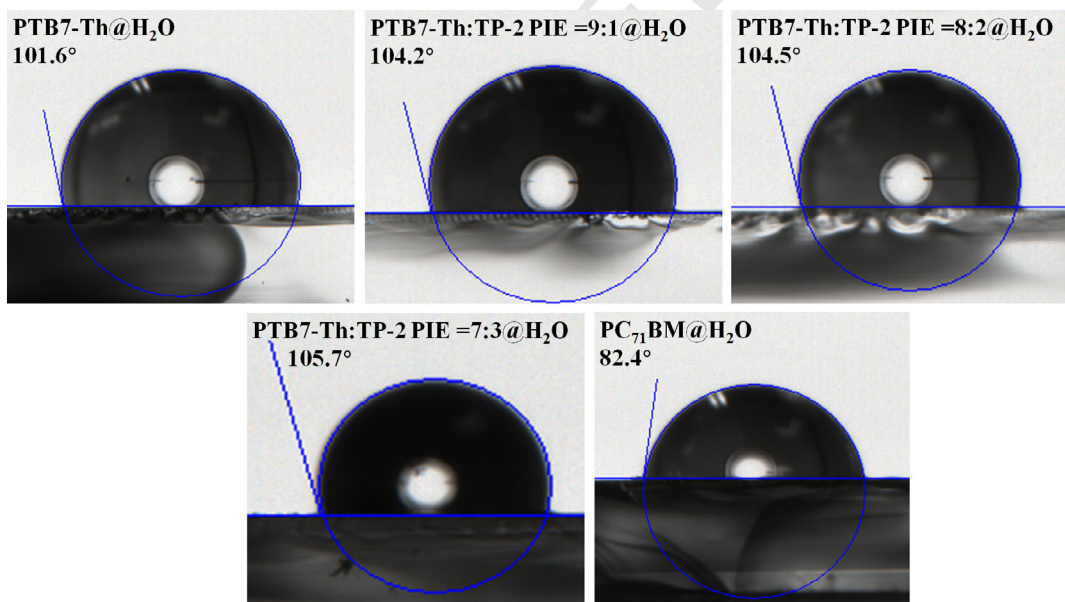
where  $\chi_{ij}$  is the interaction parameter for the component  $i$  and  $j$ ,  
respectively,  $V_0$  is the geometric mean of the polymer segment  
molar volume,  $R$  is the gas constant,  $T$  is the absolute tempera-  
ture in K, and the  $\delta_i$  and  $\delta_j$  are the Hildebrand solubility param-  
eters for the components  $i$  and  $j$ , respectively. The squared  
difference  $(\delta_i - \delta_j)^2$  between components is a measure of the  
increase in energy on mixing. And, 0.34 represents a correction  
factor for entropic contributions in a polymer solution system.

In this work, the solubility parameters were estimated through  
surface energy measurement as  $\delta$  value is the square root of the  
surface energy ( $\gamma_s$ ) ( $\delta = K\gamma_s^{1/2}$ ) and  $K$  is a universal proportional-  
ity constant, which can be determined by measuring a known  
material.<sup>[49]</sup> The surface energy  $\gamma_s$  can be obtained experimentally  
by measuring the contact angles as described in Supporting  
Information. The contact angles for films of the pristine or  
blended donor(s), as well as the acceptor, are shown in  
Figure 6 (and Table S2, Supporting Information). The  $\gamma_s$  values  
were calculated and tabulated in Table 3. Therefore, the param-  
eter  $\chi$  values were analyzed based on the mixed donor (PTB7-Th  
with TP-2-PIE) and the acceptor (PC<sub>71</sub>BM) components, and the  
resulted can be found in Table 3.

It can be seen clearly that the addition of TP-2-PIE into the  
PTB7-Th matrix led to a gradual increase in the surface contact  
angles to the testing liquid. Hence, the  $\gamma_s$  value decreased grad-  
ually from 29.05 mN m<sup>−1</sup> for the pristine PTB7-Th layer down to  
about 24.24 mN m<sup>−1</sup> for the donor layer having 30 wt% of TP-2-  
PIE component. This reduced the  $\delta$  values accordingly. The  $\gamma_s$   
value for the donor layer having 30 wt% of TP-2-PIE component  
is much lower than that for the PC<sub>71</sub>BM (41.58 mN m<sup>−1</sup>). As a  
result, the differences in solubility parameters ( $\Delta\delta$ ) between the  
donor and acceptor increased, that is, from  $\Delta\delta = 3.34$  to 5.05, as  
the fraction of TP-2-PIE component increased. Consequently, an  
increase of TP-2-PIE concentration substantially enhanced the



**Figure 5.** Tapping-mode atomic force microscopy (AFM) topography height images of a) PTB7-Th:PC<sub>71</sub>BM as-cast, b–g) PTB7-Th:PC<sub>71</sub>BM with different weight fractions of TP-2-PIE in donors as indicated. Root-mean-square roughness ( $R_{\text{RMS}}$ ) is marked below each graph.



**Figure 6.** Contact angles of different ratios of TP-2-PIE in donors and PCBM on H<sub>2</sub>O.

**Table 3.** The summarized surface free energy  $\gamma_s$ , solubility parameters  $\delta$ , and Flory–Huggins interaction parameters  $\chi$  of different ratios of TP-2-PIE in donors and PCBM.

<i>i</i> (weight ratio)	$\gamma_s$ [mN m <sup>-1</sup> ]	$\delta$ [MPa <sup>1/2</sup> ]	$\chi_{(i, \text{PC}_{71}\text{BM})}$
PTB7-Th:TP-2-PIE = 10:0	29.05	19.77	0.68
PTB7-Th:TP-2-PIE = 9:1	25.90	18.67	1.22
PTB7-Th:TP-2-PIE = 8:2	25.63	18.57	1.27
PTB7-Th:TP-2-PIE = 7:3	24.24	18.06	1.57
PC <sub>71</sub> BM	41.58	23.11	n/a

interaction parameter  $\chi_{\text{donor/acceptor}}$ . The  $\chi_{\text{PTB7-Th/PCBM}}$  took a value of 0.68 for the binary PTB7-Th:PC<sub>71</sub>BM pair; while the  $\chi_{\text{(PTB7-Th:TP-2-PIE = 7:3)/PCBM}}$  raised to 1.57 for the ternary mixed donor (PTB7-Th:TP-2-PIE = 7:3) and PC<sub>71</sub>BM pair.

Flory–Huggins interaction parameters  $\chi$  plays an important role in the morphology formation.<sup>[48,50,51]</sup> An increase in  $\chi$  value would lead to a decrease in miscibility of the components, which determines the driving force for phase separation. In comparison of the  $\chi$  values with the AFM images in Figure 5, it is found that a lower  $\chi$  value as in the control layer has a relatively good miscibility between the D–A components demonstrating a smooth



1 surface with less phase separation. A structured surface with  
2 higher roughness corresponds to a larger  $\chi$  value for a reduced  
3 miscibility in the active layer having a TP-2-PIE component,  
4 which displayed an enhancement in phase separation, or the  
5 phase-separated domain growing larger. The  $\chi$  value analyses  
6 agree well with the experimental observations.

7 Since a good miscibility with a low  $\chi$  value could lead to a low  
8 domain purity,<sup>[52]</sup> a higher  $\chi$  value is therefore beneficial for a  
9 higher degree of phase separation forming domains with a rela-  
10 tively high purity. Phase separation with high domain purity has  
11 been found to offer charge-transporting channels more effi-  
12 ciently, which is beneficial for the photovoltaic performance.<sup>[51]</sup>  
13 However, our ternary devices performed well only when the  
14 TP-2-PIE component was less than 20% (see Table 1).  
15 Although a further increase in TP-2-PIE concentration to over  
16 25 wt% enhanced the phase separation, the device performance  
17 deteriorated. For example, the active layer having 30 wt% of  
18 TP-2-PIE component was highly phase separated (Figure 5g),  
19 and the device efficiency was substantially reduced (Table 1).  
20 Further charge-transporting processes were investigated to  
21 understand this.

Q12 22 Single-carrier hole-only (ITO/PEDOT:PSS/Active layer/Au)  
23 and electron-only (ITO/ZnO/Active layer/Al) devices having dif-  
24 ferent fractions of TP-2-PIE component in the active layer were  
25 fabricated.  $J$ - $V$  characteristics of these ternary devices were plot-  
26 ted on a double logarithmic scale as shown in Figure 7a,b. The  
27 hole or electron mobility in actively layers can be estimated based  
28 on the space-charge limited current (SCLC) method similar to  
29 our previous work.<sup>[15,53]</sup>

30 It can be seen from Figure 7a,b that  $J$  increased linearly  
31 against  $V$  (i.e.,  $J \propto V$ ) in all devices at low voltage bias regions,  
32 indicating that an ohmic contact was fulfilled. As the bias voltage  
33 increased in hole-only devices (see Figure 7a, the current went  
34 into a trap-filled limited region first, followed by the third region  
35 having a quadratic voltage dependence ( $J \propto V^n$ ,  $n = 2$ ), that is, an  
36 SCLC was reached. However, in the electron-only devices, no  
37 trap-filled limited current can be observed in Figure 7b. As  
38 the bias voltage increased, current went directly into a trap-free  
39 SCLC region with a quadratic voltage dependence ( $J \propto V^2$ ). In  
40 the SCLC region, the charge carrier mobility can be obtained  
41 through the Mott–Gurney equation as seen in Equation (2),  
42 which is combined with the Poole–Frenkel expression [ $\mu = \mu_0$   
43  $\exp(\gamma\sqrt{F})$ ] to take account of the field dependency and the imper-  
44 fect quadratic relation.<sup>[15,38]</sup>

$$J_{\text{SCLC}} = \frac{9 \epsilon_0 \epsilon_r \mu_0}{8 L} \exp(0.891 \gamma \sqrt{F}) F^2 \quad (2)$$

45 In Equation (2),  $F$  is the applied electric field defined as  $F = V/L$   
46 and  $L$  is the thickness of the active layer, which is 100 nm in  
47 this work;  $\epsilon$  and  $\epsilon_0$  are the relative and vacuum permittivity of the  
48 materials, respectively, and we take  $\epsilon_r = 3$  for organic materials in  
49 this work;  $\mu_0$  is the carrier mobility at zero electric field; and  $\gamma$  is a  
50 constant. The  $\gamma$  and  $\mu_0$  can be obtained from the slope and the  
51 intercept of  $\ln(J/F^2)$  versus  $F^{1/2}$  plot in the high-field region (see  
52 Figure S5, Supporting Information) so that the values of hole  
53 mobility ( $\mu_h$ ) and electron mobility ( $\mu_e$ ) can be calculated at an  
54 electric field of  $10^6 \text{ V m}^{-1}$ .<sup>[15]</sup> The balance of the hole and electron  
55 transport can be evaluated via a ratio of  $\mu_e/\mu_h$  calculated for the

corresponding devices. These results are listed in Table 4. As the  
trap-filled limited current occurred solely in the hole-only devi-  
ces, the densities of trap states ( $n_t$ ) were calculated only for holes  
following a relation of  $V_{\text{TFL}} = q n_t L^2 / 2 \epsilon \epsilon_0$ , where  $V_{\text{TFL}}$  is the trap-  
filled limited voltage and  $q$  is the elementary charge. These  
results are also listed in Table 4.

It is interesting to see from the Table 4 that the TP-2-PIE com-  
ponent had a relatively weak effect on the  $\mu_e$  of the device. The  $\mu_e$   
showed a slight increase in the device having 10 wt% of TP-2-PIE  
in donor with a  $\mu_e$  value of  $3.8 \times 10^{-4}$ . However, the TP-2-PIE  
component was able to tune  $\mu_h$  substantially. The  $\mu_h$  increased  
gradually as the TP-2-PIE concentrations were increased. The  
 $\mu_h$  values enhanced monotonically from  $6.8 \times 10^{-4} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$   
for the control device to  $1.1 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  for the device hav-  
ing 30 wt% of TP-2-PIE in donor.

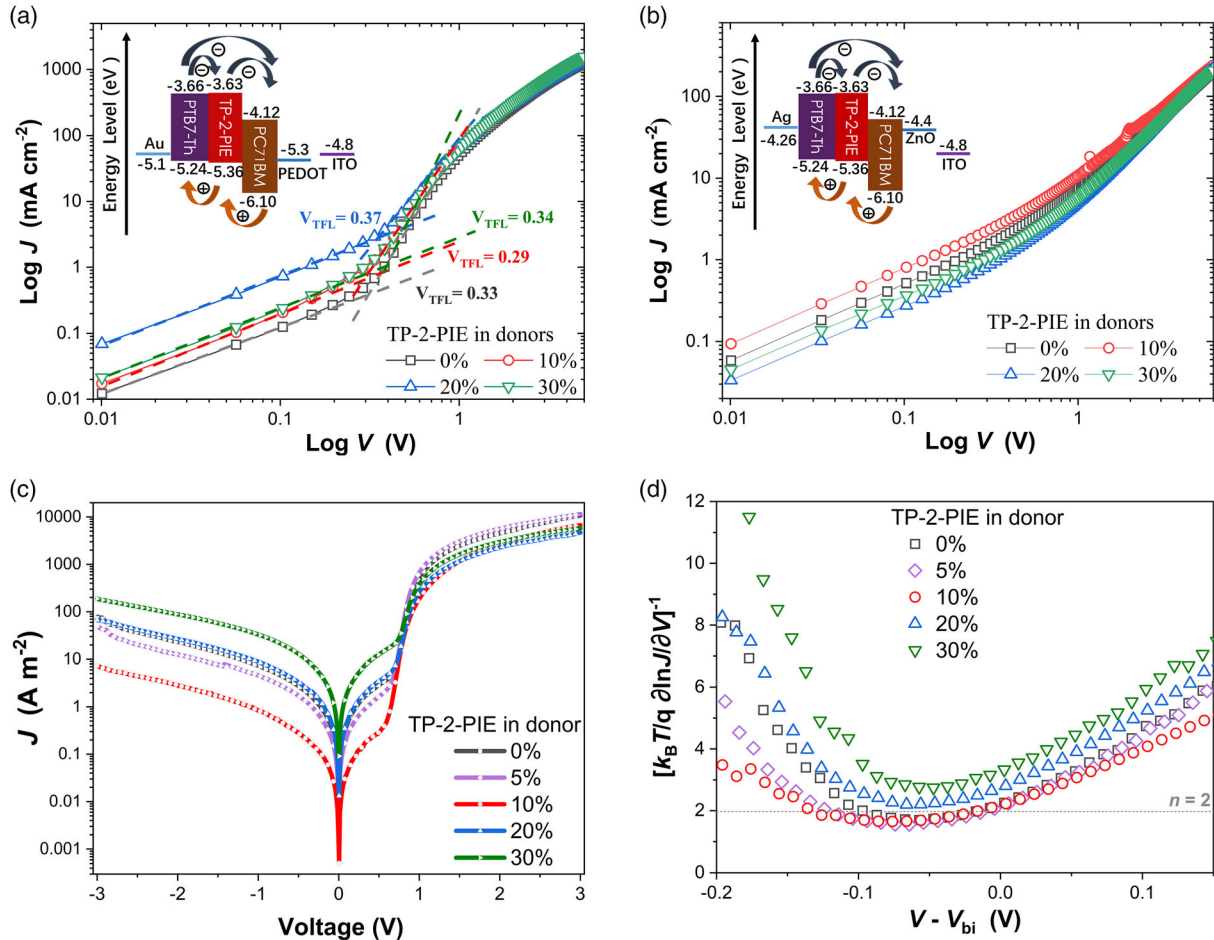
The increase in charge carrier mobilities is in good agreement  
with the change in Flory–Huggins interaction parameters  $\chi$  of  
the D–A pairs (see Table 3). It is very interesting that the gradual  
increase in  $\mu_h$  was coincident with the increase in  $\chi$  values with a  
decreased miscibility and an enhanced phase separation and/or  
domain purity.

It is known that the DLC TP-2-PIE is able to crystallize at room  
temperature and forms LC phase at elevated temperatures. More  
importantly, such a DLC molecule possesses a self-organization  
property, which can be orientated face-on based on molecular  
 $\pi$ - $\pi$  stacking.<sup>[20,31,33]</sup> This may induce a higher degree of struc-  
tural order in organic polymer systems.<sup>[41,54]</sup>

Based on the earlier analyses, the changing in morphology  
may be explained as follows. To simplify the complicated ternary  
system, we start from two mixed phases for the control (binary)  
specimen: one phase was dominated by PTB7-Th polymer  
(donor) and the other phase was dominated by PC<sub>71</sub>BM  
(acceptor). Both phases are mixed with D/A components, but  
with different purities. This can be illustrated in Figure 8.

From the molecular structures, it was expected that the TP-2-  
PIE component would mix better with polymer PTB7-Th than  
with fullerene PC<sub>71</sub>BM. Due to the Flory–Huggins interaction  
parameter that became larger when the fraction of TP-2-PIE com-  
ponent was increased, the D–A phase separation was promoted.  
A highly order TP-2-PIE component would induce more ordered  
structure in the donor-dominated domain. It not only developed  
ordered structures and resulted in phase separation, but might  
also enhance the purity in donor-dominated PTB7-Th:TP-2-PIE  
phase as shown in Figure 8. A phase-separated morphology con-  
sisting of higher degree of order in structure with a good domain  
purity is expected to offer a good charge transport channel and  
improve the charge carrier mobility.<sup>[51]</sup> Coincidentally, in our  
experiments, the  $\mu_h$  of the active layer was gradually increased  
as the TP-2-PIE concentration increased. However, the  $\mu_e$  of  
the active layer was less affected, which increased slightly as  
the TP-2-PIE concentration increased but then reduced again  
if the fraction of TP-2-PIE component was further increased.  
This indicates that the TP-2-PIE component main functioned  
in the donor-dominated phase.

In this work, although the  $\mu_h$  of the device was substantially  
increased due to the addition of TP-2-PIE component, the pho-  
tovoltaic performance was first improved and then deteriorated  
as the fraction of TP-2-PIE component increased. This can be  
understood by the balance between the hole and electron



**Figure 7.**  $J$ - $V$  characteristics of a) hole-only and b) electron-only devices plotted on a double logarithmic scale; c) dark  $J$ - $V$  and d)  $n$ - $V$  characteristics are obtained from the photovoltaic devices. (The TP-2-PIE fractions indicated are weight fractions).

**Table 4.** Hole ( $\mu_h$ ) and electron ( $\mu_e$ ) mobilities measured by space-charge limited current (SCLC) method from devices having different fraction of TP-2-PIE in donors.

PTB7-Th: TP-2-PIE: PC <sub>71</sub> BM (weight ratio)	$\mu_h$ [cm <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup> ]	$\mu_e$ [cm <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup> ]	$\mu_h/\mu_e$	$n_t^{a)}$ [cm <sup>-3</sup> ]
1.0:0:1.5	$6.8 \times 10^{-4}$	$2.0 \times 10^{-4}$	3.4	$1.1 \times 10^{16}$
0.9:0.1:1.5	$7.6 \times 10^{-4}$	$3.8 \times 10^{-4}$	2.0	$9.6 \times 10^{15}$
0.8:0.2:1.5	$9.6 \times 10^{-4}$	$1.9 \times 10^{-4}$	5.1	$1.2 \times 10^{16}$
0.7:0.3:1.5	$1.1 \times 10^{-3}$	$1.5 \times 10^{-4}$	7.3	$1.1 \times 10^{16}$

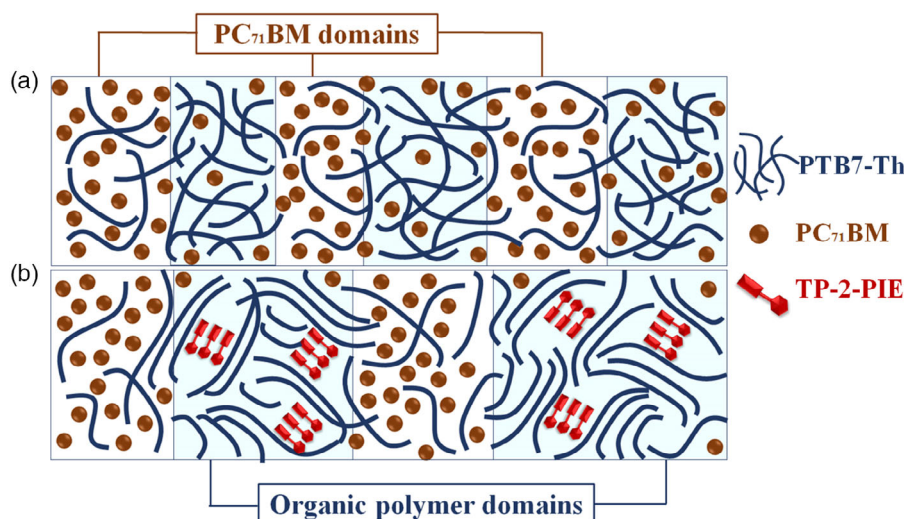
<sup>a)</sup>Density of trap states was calculated from the hole-only devices.

mobilities. It can be seen from Table 4 that the device with 10 wt% TP-2-PIE in donor demonstrated the lowest  $\mu_h/\mu_e$  ratio of 2.0. It means that a symmetrical electric field distribution could be established in the active layer to reach a transport balance. It is the balance in charge carrier mobility, which contributes to the increase in  $J_{sc}$  and  $FF$  of the device.<sup>[55,56]</sup> The unbalanced carrier mobilities would result in a space charge

accumulation, which would reduce the  $J_{sc}$  and enhanced the recombination at the interface leading to low  $FF$ .<sup>[56]</sup>

Further semiconducting physical analysis was carried out through the dark  $J$  -  $V$  characteristics of the ternary devices, which is plotted in a semilogarithmic scale as depicted in Figure 7c. It can be seen from Figure 7c that the current density fell into three distinct regions. At low voltage bias, a leakage current can be seen, which was substantially reduced when a small fraction of TP-2-PIE component was added into the active layer. However, further addition of TP-2-PIE component increased the leakage current again. The minimum leakage current was achieved in device having 10 wt% of TP-2-PIE in donor. As the bias voltage increased, but was still below the built-in voltage ( $V_{bi}$ ), the  $J$ - $V$  curve went into the second region, where the current is diffusion dominated and has an exponential dependence on voltage, which can be described using a classical Shockley diode equation (Equation (3)) given as follows.<sup>[57]</sup>

$$J = J_0 \left[ \exp \left( \frac{qV}{nk_B T} \right) - 1 \right] \quad (3)$$



**Figure 8.** A schematic diagram of donor-acceptor phase separation induced by TP-2-PIE component, which induced enhanced order in organic polymer-dominated phase with enlarged area.

1 where  $J_0$  is the saturation current density,  $n$  is an empirical ide-  
2 ality factor,  $k_B$  is the Boltzmann constant, and  $T$  is the tempera-  
3 ture. In this work,  $n$  values were obtained by fitting experimental  
4 results using Equation (4). The results were plotted against volt-  
5 age as shown in Figure 7d.

$$n = \left( \frac{k_B T}{q} \frac{\partial \ln J}{\partial V} \right)^{-1} \quad (4)$$

From this investigation, it was found that the TP-2-PIE com-  
ponent has multiple functionalities in the ternary active layer. 2  
Tuning the TP-2-PIE concentration could alter the molecular 3  
interaction and hence the phase separation and/or domain 4  
purity, which further changed the charge transport mobility 5  
and the density of trap states. A good phase separation with 6  
higher degree of structural order and a reduced density of trap 7  
states could enhance charge carrier mobility with a higher extent 8  
for hole mobility. Therefore, a mobility balance can be achieved 9  
at an optimized condition having 10 wt% of TP-2-PIE compo- 10  
nent. At the optimized condition, the ideality factor of the device 11  
was also reduced. 12

The molecular interaction, good phase separation or domain 13  
purity, the balance transport, and the reduced density of trap 14  
states ultimately enhance the photovoltaic performance of the 15  
ternary devices. 16

### 3. Conclusion

In this work, a D-A discotic liquid crystal organic TP-2-PIE was 18  
elaborately selected as a third component to fabricate the ternary 19  
PSCs by blending with PTB7-Th:PC<sub>71</sub>BM active layers. The TP-2- 20  
PIE molecule has both electron donating and accepting moieties 21  
to interact with other components as well as the ability to self- 22  
organize into a highly ordered or orientated molecular assembly 23  
to tune the morphology. 24

It was found that the TP-2-PIE component had no effect on 25  
 $V_{OC}$  on devices using the PTB7-Th:PC<sub>71</sub>BM pair, but substan- 26  
tially changed the  $J_{SC}$  and  $FF$  of the devices. This investigation 27  
discovered that the TP-2-PIE component has multiple function- 28  
alities in the ternary device: 1) the complementary absorption by 29  
TP-2-PIE and energy transfer between the TP-2-PIE and PTB7-Th 30  
granted the mixed donors with an extended photo-harvesting and 31  
extra exciton formation. The photocurrent and hence the  $J_{SC}$  of 32  
the devices can be enhanced; and 2) The TP-2-PIE component 33  
altered the molecular interaction quantified by the Flory- 34

6 For organic semiconductors,  $n$  value contains important infor-  
7 mation, such as recombination and carrier transport processes.  
8 For a classical semiconducting device of an ideal  $p-n$  junction  
9 diode without traps and carrier recombination, the  $n$  is expected  
10 to be equal to unity. However, if a trap-assisted recombination  
11 process is presented, the  $n$  may deviate from 1 and take a value  
12 of 2.<sup>[57]</sup> It has been reported that the situation in a BHJ organic  
13 solar cell can be more complicated, such as for polymer:fullerene  
14 organic solar cells, and ideality factors were typically 1.3–2.0 as  
15 reported.<sup>[58–60]</sup> The deviation of the ideality factors from unity  
16 was observed in devices with traps free for both hole (donor)  
17 and electron (acceptor). Therefore, the nonideality of the dark  
18 current in BHJ organic solar cells cannot be used as an evidence  
19 for the presence of trap-assisted recombination. It can be the  
20 transport-dominating constituent in the donor:acceptor blend.<sup>[61]</sup>

21 It was found in this work that the addition of a small fraction  
22 of TP-2-PIE component into active is able to reduce the  $n$  value  
23 slightly from  $n = 1.71$  for the control device down to  $n = 1.56$  and  
24 1.65 for the devices having 5 and 10 wt% of TP-2-PIE component  
25 in donor. However, with further increasing in TP-2-PIE concen-  
26 tration, the  $n$  value would increase to levels over 2. These results  
27 are in good agreement with the calculated density of trap states  
28 ( $n_t$ ) as seen in Table 4. It can be clearly seen that the improve-  
29 ment in ideality factor of devices was correlated to the decrease in  
30 the calculated density of trap states. The confirmed recombina-  
31 tion in these ternary devices is a trap-assisted process.



Huggins interaction parameter leading to a better phase separation with higher domain purity in the active layers. This resulted in an enhancement of charge carrier transport with a good balance and minimized density of trap states. The ideality factor  $n$  of the devices was also lowered. These improved the overall performance of the photovoltaic devices and markedly enhanced the  $J_{SC}$  and  $FF$  of the devices.

The optimal device with 10 wt% TP-2-PIE in donors achieved a 12.6% improvement in PCE as compared with the control (binary) device. This investigation offers a better understanding with regards to the functionalities of the third component in ternary BHJ organic solar cells, which may act as a guide for further device development.

## 4. Experimental Section

**Materials:** A liquid crystal compound TP-2-PIE was synthesized and reported elsewhere.<sup>[31]</sup> PTB7-Th and PC<sub>71</sub>BM were purchased from 1-Material Inc. These materials were used in active layers with chemical structures which are shown in Figure 2. The 1,8-diiodooctane was obtained from TCI Inc. Silver metallic particles (99.99%) and MoO<sub>3</sub> (99.95%) were from ZhongNuo Advanced Material (Beijing) Technology Co. and Sigma-Aldrich, respectively. Patterned indium tin oxide (ITO)-coated glasses were purchased from South China Science & Technology Company Limited, which has a sheet resistance of 15  $\Omega$  square<sup>-1</sup>. Organic solvents (AR) and other reagents were purchased from Acros if not specified.

ZnO precursor was prepared by dissolving zinc acetate dihydrate (purchased from Aladdin Inc.) into 2-methoxyethanol solvent, which was purchased from TCI Inc. Precursor for active layer (active solution) was made by co-dissolving pristine PTB7-Th or blended with TP-2-PIE as donors (10 mg mL<sup>-1</sup>) and PC<sub>71</sub>BM as an acceptor (15 mg mL<sup>-1</sup>) into dichlorobenzene at a weight ratio of 1:1.5. The 1,8-diiodooctane (3 vol%) was also added to the active solution as a solvent additive.

**Device Fabrication:** Device fabrication procedures were as follows: ITO substrates were ultrasonically cleaned first with detergent for 24 h, which were then washed with deionized water and acetone for 20 min as each. After drying with N<sub>2</sub> flow, the cleaned ITO substrates were treated in O<sub>2</sub> plasma environment for 20 min before use. First of all, a thin layer of ZnO, the electron transporting layer, was prepared on top of the ITO substrates by spin-coating the ZnO precursor at a spinning rate of 3000 rpm for 40 s. It was followed by thermal annealing at 200 °C in air for 40 min.<sup>[62]</sup> This was then transferred into a nitrogen-filled glove box to prepare the active layer. Active layer was prepared by spin-coating the active solution at 1200 rpm for 60 s on the top of ZnO layer to form a layer having an area of 0.038 cm<sup>2</sup>. Hole transporting layer of molybdenum trioxide (MoO<sub>3</sub>) was deposited using a vacuum evaporator and finally a silver electrode was deposited similarly at a pressure of  $1 \times 10^{-4}$  Pa.

**Materials and Device Measurements:** Thickness of the active layers was measured using a surface profile (Ambios XP-2), which was about 100 nm. UV-vis-near-infrared (NIR) spectrophotometers (Shimadzu UV-310PS and UV2600) had been used to measure the absorption spectra of the film specimen. X-Ray diffraction patterns were taken by a Bruker D8 advance diffractometer. Steady-state and TRPL spectra were measured using a PL spectrometer (Edinburgh Instruments FLS 1000) using a NIR photomultiplier tube (Hamamatsu R5509-73 PMT) as detector with liquid nitrogen cooling for 2 h before measurement. The emissions were corrected for the instrumental spectral response. The transient emission was excited with a 406 nm laser diode. The contact angles were measured optically through a static sessile-drop method using an automatic goniometer and surface energies were calculated using a droplet analyzer (Guangdong Hokuto Instrument Co. & Ltd., CA200).

Current-voltage characteristics of the devices were measured using a source measure unit (Agilent B2902A). Photovoltaic responses were obtained under a standard illumination of AM 1.5G solar irradiation at

100 mW cm<sup>-2</sup> from a solar simulator (Abet Sun 2000). EQEs of the solar cells were evaluated through a solar cell quantum efficiency measurement system (Zolix Solar Cell Scan 100).

## Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

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## Conflict of Interest

The authors declare no conflict of interest.

## Data Availability Statement

The data that support the findings of this study are available from the corresponding author upon reasonable request.

## Keywords

charge carrier mobility, discotic, energy transfer, Flory–Huggins interaction parameter, phase separation, photovoltaics, polymer solar cells

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