

计算流体力学

第三章:有限差分方法

宁方飞 能源与动力工程学院

本章主要内容



- 有限差分的基本概念
- 有限差分格式的构造

• 编程作业: 迁移方程的离散及求解



• 有限差分: Finite difference

$$u_{x} = \frac{\partial u}{\partial x} = \lim_{\Delta x \to 0} \frac{u(x + \Delta x) - u(x)}{\Delta x}$$

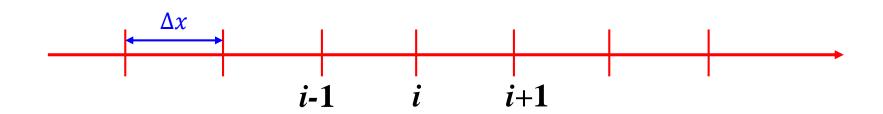
$$\frac{\partial u}{\partial x} \approx \frac{u(x + \Delta x) - u(x)}{\Delta x}$$

$$u_x = \frac{\partial u}{\partial x} = \lim_{\Delta x \to 0} \frac{u(x) - u(x - \Delta x)}{\Delta x}$$

$$\frac{\partial u}{\partial x} \approx \frac{u(x) - u(x - \Delta x)}{\Delta x}$$

$$u_{x} = \frac{\partial u}{\partial x} = \lim_{\Delta x \to 0} \frac{u(x + \Delta x) - u(x - \Delta x)}{2\Delta x}$$

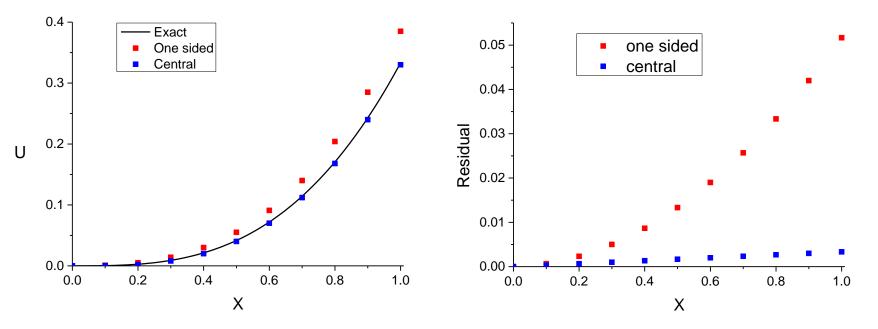
$$\frac{\partial u}{\partial x} \approx \frac{u(x + \Delta x) - u(x - \Delta x)}{2\Delta x}$$





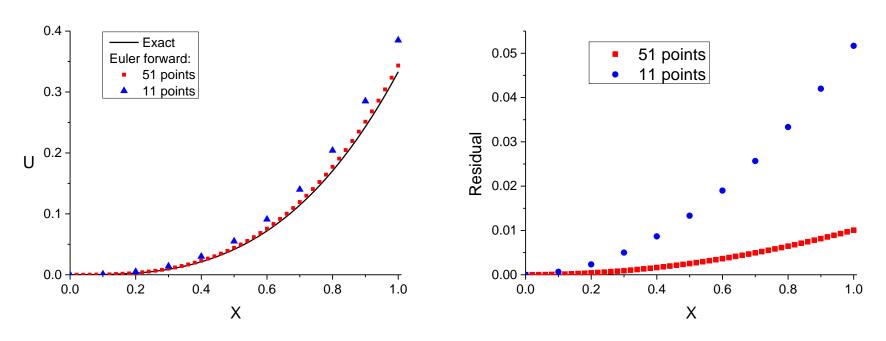
$$\frac{du}{dx} = x^2, \qquad u(0) = 0$$

$$i-1 \qquad i \qquad i+1$$



不同格式间的比较





单侧差分、不同离散点数的比较

不同差分格式、离散点 Δx 对精度都有影响,如何评价它们?



$$u_x = \frac{\partial u}{\partial x} = \lim_{\Delta x \to 0} \frac{u(x + \Delta x) - u(x)}{\Delta x}$$

$$u_{x} = \frac{\partial u}{\partial x} = \lim_{\Delta x \to 0} \frac{u(x + \Delta x) - u(x)}{\Delta x} \qquad u_{x} = \frac{\partial u}{\partial x} = \lim_{\Delta x \to 0} \frac{u(x + \Delta x) - u(x - \Delta x)}{2\Delta x}$$

$$u(x + \Delta x) = u(x) + \Delta x \frac{\partial u}{\partial x} + \frac{(\Delta x)^2}{2} \frac{\partial^2 u}{\partial x^2} + \frac{(\Delta x)^3}{6} \frac{\partial^3 u}{\partial x^3} + \cdots$$

$$u(x - \Delta x) = u(x) - \Delta x \frac{\partial u}{\partial x} + \frac{(\Delta x)^2}{2} \frac{\partial^2 u}{\partial x^2} - \frac{(\Delta x)^3}{6} \frac{\partial^3 u}{\partial x^3} + \cdots$$

$$\frac{\partial u}{\partial x} = \frac{u(x + \Delta x) - u(x)}{\Delta x}$$

$$\frac{\partial u}{\partial x} = \frac{u(x + \Delta x) - u(x)}{\Delta x} \qquad -\frac{\Delta x}{2} \frac{\partial^2 u}{\partial x^2} - \frac{(\Delta x)^2}{6} \frac{\partial^3 u}{\partial x^3} - \cdots$$

$$\frac{\partial u}{\partial x} = \frac{u(x + \Delta x) - u(x - \Delta x)}{2\Delta x} - 0 - \frac{(\Delta x)^2}{6} \frac{\partial^3 u}{\partial x^3} - \cdots$$

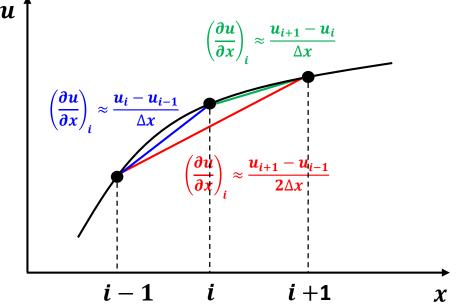
$$-\frac{(\Delta x)^2}{6}\frac{\partial^3 u}{\partial x^3}-\cdots$$

> 截断误差



$$\frac{\partial u}{\partial x} = \frac{u(x + \Delta x) - u(x)}{\Delta x} \qquad -\frac{\Delta x}{2} \frac{\partial^2 u}{\partial x^2} - \frac{(\Delta x)^2}{6} \frac{\partial^3 u}{\partial x^3} - \dots = \frac{u(x + \Delta x) - u(x)}{\Delta x} \qquad + O(\Delta x)$$

$$\frac{\partial u}{\partial x} = \frac{u(x + \Delta x) - u(x - \Delta x)}{2\Delta x} - 0 \qquad -\frac{(\Delta x)^2}{6} \frac{\partial^3 u}{\partial x^3} - \dots = \frac{u(x + \Delta x) - u(x - \Delta x)}{2\Delta x} + O((\Delta x)^2)$$



当 Δx → **0**时,能更快收敛至 $\frac{\partial u}{\partial x}$

• 截断误差 $O((\Delta x^n))$,n称为差分格式的精度阶数,n越大,则差分格式的精度越高



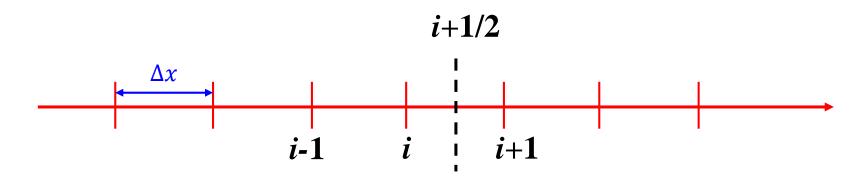
• 一维一阶导数的差分离散

$$\left(\frac{\partial u}{\partial x}\right)_{i} = \frac{u_{i+1} - u_{i}}{\Delta x} + O(\Delta x)$$
 Euler 前差

$$\left(\frac{\partial u}{\partial x}\right)_{i} = \frac{u_{i} - u_{i-1}}{\Delta x} + O(\Delta x)$$
 Euler 后差

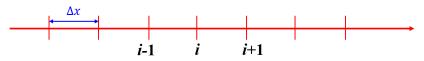
$$\left(\frac{\partial u}{\partial x}\right)_{i} = \frac{u_{i+1} - u_{i-1}}{2\Delta x} + O((\Delta x)^{2}) \quad \Leftrightarrow \triangle \stackrel{\text{def}}{=} \Omega$$

$$\left(\frac{\partial u}{\partial x}\right)_{i+1/2} = \frac{u_{i+1} - u_{i}}{\Delta x} + O((\Delta x)^{2})$$





• 一维一阶导数的差分离散



待定系数法:

$$\left(\frac{\partial u}{\partial x}\right)_{i+1} \approx (au_{i+1} + bu_i + cu_{i-1})/\Delta x$$

$$u_{i} = u_{i+1} - \Delta x \left(\frac{\partial u}{\partial x}\right)_{i+1} + \frac{(\Delta x)^{2}}{2} \left(\frac{\partial^{2} u}{\partial x^{2}}\right)_{i+1} - \frac{(\Delta x)^{3}}{6} \left(\frac{\partial^{3} u}{\partial x^{3}}\right)_{i+1} + \cdots$$

$$u_{i-1} = u_{i+1} - 2\Delta x \left(\frac{\partial u}{\partial x}\right)_{i+1} + \frac{(2\Delta x)^2}{2} \left(\frac{\partial^2 u}{\partial x^2}\right)_{i+1} - \frac{(2\Delta x)^3}{6} \left(\frac{\partial^3 u}{\partial x^3}\right)_{i+1} + \cdots$$

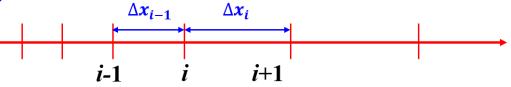
$$a+b+c=0$$
 $-b-2c=1$ $\frac{1}{2}b+2c=0$ $a=\frac{3}{2}$ $b=-2$ $c=\frac{1}{2}$

$$\left(\frac{\partial u}{\partial x}\right)_{i+1} = \frac{3u_{i+1} - 4u_i + u_{i-1}}{2\Delta x} + O((\Delta x)^2)$$
 三点Euler后差



• 一维一阶导数的差分离散

非均匀网格点分布:



$$\left(\frac{\partial u}{\partial x}\right)_{i} \approx (au_{i+1} + bu_{i} + cu_{i-1})/(\Delta x_{i-1} + \Delta x_{i})$$

$$u_{i+1} = u_i + \Delta x_i \left(\frac{\partial u}{\partial x}\right)_i + \frac{(\Delta x_i)^2}{2} \left(\frac{\partial^2 u}{\partial x^2}\right)_i + \frac{(\Delta x_i)^3}{6} \left(\frac{\partial^3 u}{\partial x^3}\right)_i + \cdots$$

$$u_{i-1} = u_i - \Delta x_{i-1} \left(\frac{\partial u}{\partial x} \right)_i + \frac{(\Delta x_{i-1})^2}{2} \left(\frac{\partial^2 u}{\partial x^2} \right)_i - \frac{(\Delta x_{i-1})^3}{6} \left(\frac{\partial^3 u}{\partial x^3} \right)_i + \cdots$$

$$a = \frac{\Delta x_{i-1}}{\Delta x_i} \qquad b = \frac{\Delta x_i}{\Delta x_{i-1}} - \frac{\Delta x_{i-1}}{\Delta x_i} \qquad c = -\frac{\Delta x_i}{\Delta x_{i-1}} \qquad r = \frac{\Delta x_i}{\Delta x_{i-1}}$$

$$c = -\frac{\Delta x_i}{\Delta x_{i-1}}$$

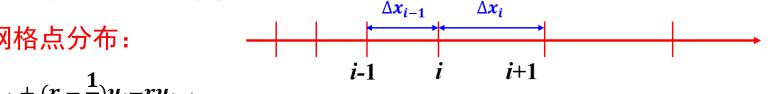
$$r = \frac{\Delta x_i}{\Delta x_{i-1}}$$

$$\left(\frac{\partial u}{\partial x}\right)_{i} = \frac{\frac{1}{r}u_{i+1} + (r - \frac{1}{r})u_{i} - ru_{i-1}}{\Delta x_{i-1} + \Delta x_{i}} + O((\Delta x)^{2})$$



• 一维一阶导数的差分离散

非均匀网格点分布:



$$\left(\frac{\partial u}{\partial x}\right)_{i} = \frac{\frac{1}{r}u_{i+1} + (r - \frac{1}{r})u_{i} - ru_{i-1}}{\Delta x_{i-1} + \Delta x_{i}} + O((\Delta x)^{2})$$

如果忽略网格点间距的变化:

$$r = 1$$

$$\left(\frac{\partial u}{\partial x}\right)_{i} = \frac{u_{i+1} - u_{i-1}}{\Delta x_{i-1} + \Delta x_{i}} + O((\Delta x)^{?})$$

$$\left(\frac{\partial u}{\partial x}\right)_{i} = \frac{u_{i+1} - u_{i-1}}{\Delta x_{i-1} + \Delta x_{i}} - \frac{\Delta x_{i} - \Delta x_{i-1}}{2} \left(\frac{\partial^{2} u}{\partial x^{2}}\right)_{i} + \cdots$$

$$\text{ Might}$$

$$u_{i+1} = u_i + \Delta x_i \left(\frac{\partial u}{\partial x}\right)_i + \left(\frac{(\Delta x_i)^2}{2}\left(\frac{\partial^2 u}{\partial x^2}\right)_i + \cdots\right)$$

$$u_{i-1} = u_i - \Delta x_{i-1} \left(\frac{\partial u}{\partial x} \right)_i + \frac{(\Delta x_{i-1})^2}{2} \left(\frac{\partial^2 u}{\partial x^2} \right)_i + \cdots$$

- 如果 $(\Delta x_i \Delta x_{i-1}) \sim \Delta x$,则格式只有一
- 当 $(\Delta x_i \Delta x_{i-1}) \sim O(\Delta x^2)$ 时,格式才有 二阶精度



• 一维一阶导数的差分离散

多项式法:

Step1: 给定N阶多项式

$$u(x) = a_0 + a_1(x - x_i) + a_2(x - x_i)^2 + a_3(x - x_i)^3 + \cdots$$

Step2: N阶多项式有N+1个待定系数,用多项式拟合N+1个离散点处u的值,可确定多项式系数

$$x = x_{i-k}, x_{i-k+1}, \cdots, x_i, \cdots, x_{i+m-1}, x_{i+m}$$

$$u = u_{i-k}, u_{i-k+1}, \cdots, u_i, \cdots, u_{i+m-1}, u_{i+m}$$
 代入多项式

Step3: 在i点的N阶精度的一阶差分为 $u'(x_i) = a_1$

在i点的N-1阶精度的二阶差分为 $u''(x_i) = 2a_2$



• 紧致差分离散

$$\frac{\partial u}{\partial x}\Big|_{i} = \frac{u_{i+1} - u_{i-1}}{2\Delta x} + O((\Delta x)^{2}) \qquad \frac{u_{i+1} - u_{i-1}}{2\Delta x} = \left(\frac{\partial u}{\partial x}\right)_{i} + O((\Delta x)^{2})$$

$$\frac{u_{i+1} - u_{i-1}}{2\Delta x} = a(u_{x})_{i+1} + b(u_{x})_{i} + c(u_{x})_{i-1} + O((\Delta x)^{n})$$

$$\frac{u_{i+1} - u_{i-1}}{2\Delta x} = (u_{x})_{i} \qquad + \frac{\Delta x^{2}}{3!}(u_{xxx})_{i} \qquad + \frac{\Delta x^{4}}{5!}(u_{5x})_{i} + O((\Delta x)^{6})$$

$$(u_{x})_{i+1} = (u_{x})_{i} + \Delta x(u_{xx})_{i} + \frac{\Delta x^{2}}{2}(u_{xxx})_{i} + \frac{\Delta x^{3}}{3!}(u_{4x})_{i} + \frac{\Delta x^{4}}{4!}(u_{5x})_{i} + O((\Delta x)^{5})$$

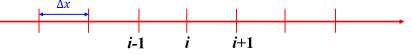
$$(u_{x})_{i-1} = (u_{x})_{i} - \Delta x(u_{xx})_{i} + \frac{\Delta x^{2}}{2}(u_{xxx})_{i} - \frac{\Delta x^{3}}{3!}(u_{4x})_{i} + \frac{\Delta x^{4}}{4!}(u_{5x})_{i} + O((\Delta x)^{5})$$

$$a + b + c = 1 \quad a = c \quad a + c = 1/3$$

$$\frac{1}{6}[(u_{x})_{i+1} + 4(u_{x})_{i} + (u_{x})_{i-1}] = \frac{u_{i+1} - u_{i-1}}{2\Delta x} + O((\Delta x)^{4})$$



• 紧致差分离散



$$u(x) = a_0 + a_1(x - x_i) + a_2(x - x_i)^2 + a_3(x - x_i)^3 + a_4(x - x_i)^4$$
 $u'(x) = a_1 + 2a_2(x - x_i) + 3a_3(x - x_i)^2 + 4a_4(x - x_i)^3$
 $x = x_{i-1}, x_i, x_{i+1}: u = u_{i-1}, u_i, u_{i+1}$
 $x = x_{i-1}, x_{i+1}: u' = u'_{i-1}, u'_{i+1}$

代入以上两式

$$\left(\frac{\partial u}{\partial x}\right)_{i} = a_{1} = -\frac{1}{4} \left(\frac{\partial u}{\partial x}\right)_{i-1} - \frac{1}{4} \left(\frac{\partial u}{\partial x}\right)_{i+1} + \frac{3}{4} \frac{u_{i+1} - u_{i-1}}{\Delta x}$$

$$\frac{1}{6}[(u_x)_{i+1} + 4(u_x)_i + (u_x)_{i-1}] = \frac{u_{i+1} - u_{i-1}}{2\Delta x} + O((\Delta x)^4)$$

具有4阶精度,一样的结果



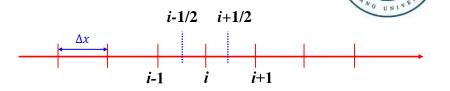
• 紧致差分离散

$$\frac{1}{6}[(u_x)_{i+1}+4(u_x)_i+(u_x)_{i-1}]=\frac{u_{i+1}-u_{i-1}}{2\Delta x}+O((\Delta x)^4)$$

$$\begin{bmatrix} & \ddots & & & \\ & 1 & 4 & 1 & & \\ & & 1 & 4 & 1 & \\ & & & 1 & 4 & 1 \\ & & & & \ddots & \end{bmatrix} \cdot \begin{bmatrix} \vdots \\ (u_x)_{i-1} \\ (u_x)_i \\ (u_x)_{i+1} \\ \vdots \end{bmatrix} = \frac{3}{\Delta x} \begin{bmatrix} \vdots \\ u_i - u_{i-2} \\ u_{i+1} - u_{i-1} \\ u_{i+2} - u_i \\ \vdots \end{bmatrix}$$

- 采用紧致格式需要隐式求解
- 任意一处离散点的u值的变化会影响所有的导数值
- 紧致格式可以用更少的点获得高阶精度





$$\left(\frac{\partial^2 u}{\partial x^2}\right)_{i} = \frac{(u_x)_{i+1/2} - (u_x)_{i-1/2}}{\Delta x} = \frac{u_{i+1} - 2u_i + u_{i-1}}{\Delta x^2} + O((\Delta x)^2)$$

$$\left(\frac{\partial^3 u}{\partial x^3}\right)_i = \frac{(u_{xx})_{i+1} - (u_{xx})_{i-1}}{2\Delta x} = \frac{u_{i+2} - 2u_{i+1} + 2u_{i-1} - u_{i-2}}{2\Delta x^3} + O((\Delta x)^2)$$

$$\left(\frac{\partial^3 u}{\partial x^3}\right)_{i+1/2} = \frac{(u_{xx})_{i+1} - (u_{xx})_i}{\Delta x} = \frac{u_{i+2} - 3u_{i+1} + 3u_i - u_{i-1}}{\Delta x^3} + O((\Delta x)^2)$$

$$\left(\frac{\partial^4 u}{\partial x^4}\right)_i = \frac{(u_{xxx})_{i+1/2} - (u_{xxx})_{i-1/2}}{\Delta x} = \frac{u_{i+2} - 4u_{i+1} + 6u_i - 4u_{i-1} + u_{i-2}}{\Delta x^4} + O((\Delta x)^2)$$



• 迁移方程的有限差分格式

$$\frac{\partial u}{\partial t} + a \frac{\partial u}{\partial x} = 0$$

计算域: $x_B \leq x \leq x_T$

初始条件: t=0时, $u(x,0)=u^0(x)$

边界条件: $x = x_U$ 时, $u(x_U, t) = g(t)$ $\begin{cases} a > 0: x_U = x_B \\ a < 0: x_U = x_T \end{cases}$

空间离散:

$$(u_t)_i = -a \frac{u_{i+1} - u_{i-1}}{2\Delta x}$$

$$(u_t)_i = -a \frac{u_i - u_{i-1}}{\Delta x}$$

$$(u_t)_i = -a \frac{3u_i - 4u_{i-1} + u_{i-2}}{2\Delta x}$$

中心差分

Euler后差

三点Euler后差

• 迁移方程的有限差分格式

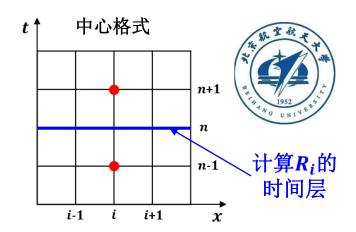
空间离散: $(u_t)_i = R_i$

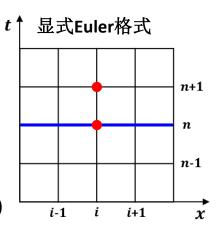
时间离散: 将时刻表示为: $t = n\Delta t$

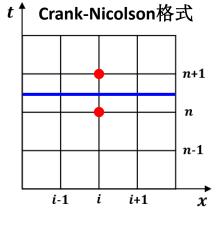
$$\frac{u_i^{n+1}-u_i^{n-1}}{2\wedge t}=R_i^n$$

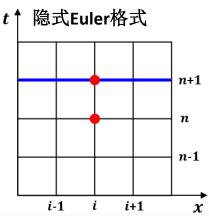
$$\frac{u_i^{n+1} - u_i^n}{\Delta t} = R_i^n, R_i^{n+1}, or \frac{1}{2} (R_i^n + R_i^{n+1})$$

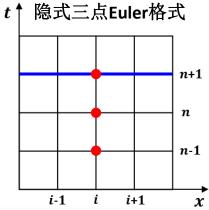
$$\frac{3u_i^{n+1} - 4u_i^n + u_i^{n-1}}{2\Delta t} = R_i^n \quad or \quad R_i^{n+1}$$







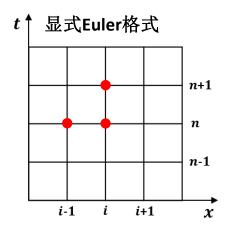






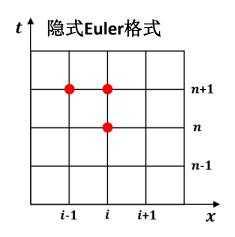
• 迁移方程的有限差分格式

更多离散格式:



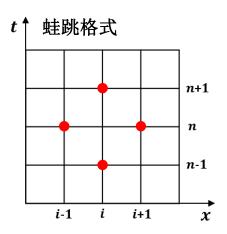
$$\frac{u_i^{n+1}-u_i^n}{\wedge t}=-a\frac{u_i^n-u_{i-1}^n}{\wedge x}$$

$$u_i^{n+1} = u_i^n - a \frac{\Delta t}{\Delta x} (u_i^n - u_{i-1}^n)$$



$$\frac{u_i^{n+1}-u_i^n}{\Delta t}=-a\frac{u_i^{n+1}-u_{i-1}^{n+1}}{\Delta x}$$

$$u_i^{n+1} = u_i^n - a \frac{\Delta t}{\Delta x} (u_i^{n+1} - u_{i-1}^{n+1})$$



$$\frac{u_i^{n+1} - u_i^n}{\Delta t} = -a \frac{u_i^{n+1} - u_{i-1}^{n+1}}{\Delta x} \qquad \frac{u_i^{n+1} - u_i^{n-1}}{2\Delta t} = -a \frac{u_{i+1}^n - u_{i-1}^n}{2\Delta x}$$

$$u_i^{n+1} = u_i^n - a \frac{\Delta t}{\Delta x} (u_i^{n+1} - u_{i-1}^{n+1}) \qquad u_i^{n+1} = u_i^{n-1} - a \frac{\Delta t}{\Delta x} (u_{i+1}^n - u_{i-1}^n)$$

列出的这些格式不一定都能给出正确的结果!



• 迁移方程的有限差分格式

Lax-Wendroff格式:

$$u_i^{n+1} = u_i^n + \Delta t \left(\frac{\partial u}{\partial t}\right)_i^n + \frac{\Delta t^2}{2} \left(\frac{\partial^2 u}{\partial t^2}\right)_i^n + O(\Delta t^3)$$
$$\left(\frac{\partial u}{\partial t}\right)_i = -a \frac{u_{i+1} - u_{i-1}}{2\Delta x} + O((\Delta x)^2)$$

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial}{\partial t} \left(\frac{\partial u}{\partial t} \right) = -a \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial t} \right) = a^2 \frac{\partial^2 u}{\partial x^2} = a^2 \frac{u_{i+1} - 2u_i + u_{i-1}}{\Delta x^2} + O((\Delta x)^2)$$

$$u_i^{n+1} = u_i^n - \frac{1}{2} \frac{a \Delta t}{\Delta x} (u_{i+1}^n - u_{i-1}^n) + \frac{1}{2} \left(\frac{a \Delta t}{\Delta x} \right)^2 (u_{i+1}^n - 2u_i^n + u_{i-1}^n)$$

$$u_i^{n+1} = u_i^n - \frac{\sigma}{2} (u_{i+1}^n - u_{i-1}^n) + \frac{\sigma^2}{2} (u_{i+1}^n - 2u_i^n + u_{i-1}^n) \qquad \sigma = \frac{a\Delta t}{\Delta x}$$

编程作业一: 迁移方程的离散及求解

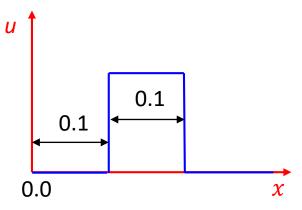


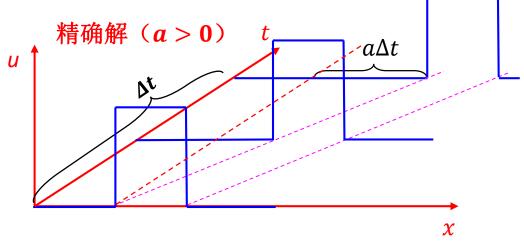
• 算例具体要求:

- 空间计算域: [0.0-1.0], 用至少100个点离散

- 对于显式格式,建议组合参量: $\left|a\frac{\Delta t}{\Delta x}\right| < 1$,隐式格式无限制

- 初场可以如下:





- 沿时间推进求解至时刻t = 0.7/a, (a > 0时)
- 完成不少于3种迁移方程离散格式的推导;
- 对每种离散格式,进行编程求解;
- 分析每种格式的结果、收敛性、精度等,完成分析报告,报告后 附源程序。



• 高维空间导数的差分离散

$$u(x + \Delta x, y + \Delta y) = u(x, y) + \left(\Delta x \frac{\partial}{\partial x} + \Delta y \frac{\partial}{\partial y}\right) u + \frac{1}{2} \left(\Delta x \frac{\partial}{\partial x} + \Delta y \frac{\partial}{\partial y}\right)^2 u + \frac{1}{3!} \left(\Delta x \frac{\partial}{\partial x} + \Delta y \frac{\partial}{\partial y}\right)^3 u + \cdots$$

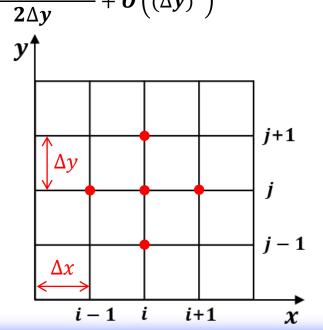
$$\left(\frac{\partial u}{\partial x}\right)_{i,i} = \frac{u_{i+1,j} - u_{i,j}}{\Delta x} + O(\Delta x); \quad \frac{u_{i,j} - u_{i-1,j}}{\Delta x} + O(\Delta x); \quad \frac{u_{i+1,j} - u_{i-1,j}}{2\Delta x} + O\left((\Delta x)^2\right)$$

$$\left(\frac{\partial u}{\partial y}\right)_{i,j} = \frac{u_{i,j+1} - u_{i,j}}{\Delta y} + O(\Delta y); \quad \frac{u_{i,j} - u_{i,j-1}}{\Delta y} + O(\Delta y); \quad \frac{u_{i,j+1} - u_{i,j-1}}{2\Delta y} + O(\Delta y)^2$$

$$\left(\frac{\partial^2 u}{\partial x^2}\right)_{i,i} = \frac{u_{i+1,j} - 2u_{i,j} + u_{i-1,j}}{\Delta x^2} + O\left((\Delta x)^2\right)$$

$$\left(\frac{\partial^2 u}{\partial y^2}\right)_{i,j} = \frac{u_{i,j+1} - 2u_{i,j} + u_{i,j-1}}{\Delta y^2} + O\left((\Delta y)^2\right)$$

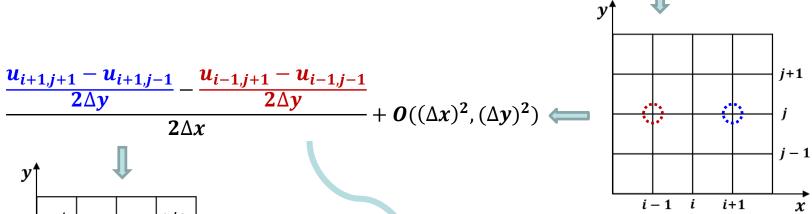
在笛卡尔网格下,函数的空间导数可视 为沿各自坐标方向的一维导数,所以可 用一维差分格式进行离散





• 二维空间导数的差分离散

$$\left(\frac{\partial^2 u}{\partial x \partial y}\right)_{i,j} = \left[\frac{\partial}{\partial x}(u_y)\right]_{i,j} = \left[\frac{\partial}{\partial y}(u_x)\right]_{i,j} = \mathbf{e.\,g.} \quad \frac{\left(\mathbf{u}_y\right)_{i+1,j} - \left(\mathbf{u}_y\right)_{i-1,j}}{2\Delta x} + O((\Delta x)^2)$$



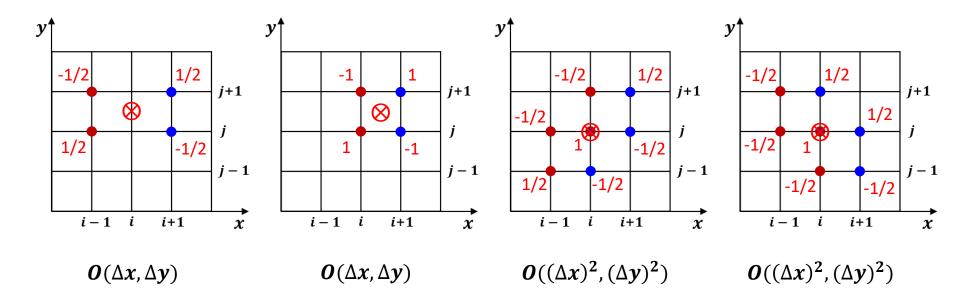
$$-1/4$$
 $1/4$ $j+1$ j
 $1/4$ $j-1/4$ $j-1/4$

$$\frac{u_{i+1,j+1} - u_{i+1,j-1} - u_{i-1,j+1} + u_{i-1,j-1}}{4\Delta x \Delta y} + O((\Delta x)^2, (\Delta y)^2)$$



• 二维空间导数的差分离散

$$\left(\frac{\partial^2 u}{\partial x \partial y}\right)_{i,j} = \left[\frac{\partial}{\partial x}(u_y)\right]_{i,j} = \left[\frac{\partial}{\partial y}(u_x)\right]_{i,j}$$

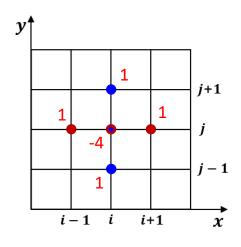




• Laplace算子的差分离散

$$\nabla^2 u = \overrightarrow{\nabla} \cdot \left(\overrightarrow{\nabla} u \right) = u_{xx} + u_{yy}$$

$$\left(\frac{\partial^2 u}{\partial x^2} \right)_{i,j} = \frac{u_{i+1,j} - 2u_{i,j} + u_{i-1,j}}{\Delta x^2} + O\left((\Delta x)^2 \right)$$



$$\left(\frac{\partial^2 u}{\partial y^2}\right)_{i,j} = \frac{u_{i,j+1} - 2u_{i,j} + u_{i,j-1}}{\Delta y^2} + O\left((\Delta y)^2\right)$$

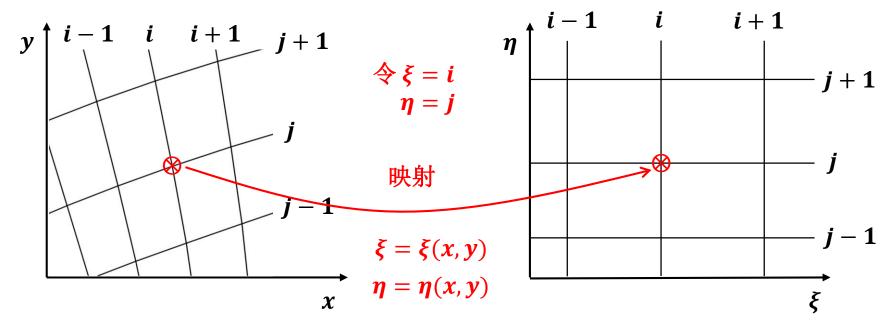
$$\left(\nabla^2 u \right)_{i,j} = \frac{u_{i+1,j} - 2u_{i,j} + u_{i-1,j}}{\Delta x^2} + \frac{u_{i,j+1} - 2u_{i,j} + u_{i,j-1}}{\Delta y^2} + O\left((\Delta x)^2, (\Delta y)^2 \right)$$

非线性扩散算子:

$$\vec{\nabla}(\mu\vec{\nabla}u)_{i,j} = \left[\frac{\partial}{\partial x}\left(\mu\frac{\partial u}{\partial x}\right)\right]_{i,j} + \left[\frac{\partial}{\partial y}\left(\mu\frac{\partial u}{\partial y}\right)\right]_{i,j} \qquad \mu = \mu(x,y)$$



• 曲线网格系统下的差分离散



$$\frac{\partial u}{\partial x} = \frac{\partial u}{\partial \xi} \frac{\partial \xi}{\partial x} + \frac{\partial u}{\partial \eta} \frac{\partial \eta}{\partial x}$$

以中心差分为例:

$$\left(\frac{\partial u}{\partial x}\right)_{i,j} = \frac{1}{2}(u_{i+1,j} - u_{i-1,j})(\xi_x)_{i,j} + \frac{1}{2}(u_{i,j+1} - u_{i,j-1})(\eta_x)_{i,j}$$

同样:
$$\left(\frac{\partial u}{\partial y}\right)_{i,j} = \frac{1}{2}(u_{i+1,j} - u_{i-1,j})(\xi_y)_{i,j} + \frac{1}{2}(u_{i,j+1} - u_{i,j-1})(\eta_y)_{i,j}$$

• 曲线网格系统下的差分离散

$$y \downarrow i-1 \qquad i \qquad i+1 \qquad j+1 \qquad \xi = \xi(x,y) \qquad \eta \qquad \qquad i-1 \qquad i \qquad i+1 \qquad \qquad j+1 \qquad \qquad j-1 \qquad \qquad j+1 \qquad$$

$$\begin{cases} \xi = \xi(x, y) \\ \eta = \eta(x, y) \end{cases} \stackrel{\mathbf{x} \to \mathbf{x} \to \mathbf{x}}{\mathbf{x} \to \mathbf{x}} \begin{cases} d\xi = \xi_x dx + \xi_y dy \\ d\eta = \eta_x dx + \eta_y dy \end{cases} \longrightarrow \begin{bmatrix} d\xi \\ d\eta \end{bmatrix} = \begin{bmatrix} \xi_x & \xi_y \\ \eta_x & \eta_y \end{bmatrix} \begin{bmatrix} dx \\ dy \end{bmatrix}$$

$$\begin{cases} x = x(\xi, \eta) \\ y = y(\xi, \eta) \end{cases} \qquad \begin{cases} dx = x_\xi d\xi + x_\eta d\eta \\ dy = y_\xi d\xi + y_\eta d\eta \end{cases} \longrightarrow \begin{bmatrix} dx \\ dy \end{bmatrix} = \begin{bmatrix} x_\xi & x_\eta \\ y_\xi & y_\eta \end{bmatrix} \begin{bmatrix} d\xi \\ d\eta \end{bmatrix}$$

$$\begin{bmatrix} \xi_x & \xi_y \\ \eta_x & \eta_y \end{bmatrix} = \begin{bmatrix} y_\eta & -x_\eta \\ -y_\xi & x_\xi \end{bmatrix} / (x_\xi y_\eta - x_\eta y_\xi) \qquad \qquad \begin{bmatrix} \xi_x & \xi_y \\ \eta_x & \eta_y \end{bmatrix} = \begin{bmatrix} x_\xi & x_\eta \\ y_\xi & y_\eta \end{bmatrix}^{-1}$$

$$(x_{\xi})_{i,j} = \frac{1}{2}(x_{i+1,j} - x_{i-1,j})$$

$$(x_{\xi})_{i,j} = \frac{1}{2}(x_{i+1,j} - x_{i-1,j})$$

$$(x_{\xi})_{i+\frac{1}{2}j} = \cdots \qquad (x_{\xi})_{i,j+\frac{1}{2}} = \cdots$$

$$(x_{\eta})_{i,j} = \frac{1}{2}(x_{i,j+1} - x_{i,j-1})$$

• 差分格式的守恒性概念

 $i-1 \qquad i \qquad i+1$ $i-\frac{1}{2} \qquad i+\frac{1}{2}$

格点*i*的 控制体

考虑如下守恒律:

$$\frac{\partial u}{\partial t} + \frac{\partial f}{\partial x} = 0$$

对空间导数作如下离散:

$$\left(\frac{\partial f}{\partial x}\right)_{i} = \frac{f_{i+\frac{1}{2}} - f_{i-\frac{1}{2}}}{x_{i+\frac{1}{2}} - x_{i-\frac{1}{2}}}$$

称之为守恒差分格式。Why?

控制方程的守恒性



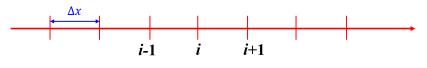
离散方程的守恒性



数值解的守恒性



一维Euler方程的差分离散



$$\frac{\partial U}{\partial t} + \frac{\partial F}{\partial x} = 0 \qquad U = \begin{bmatrix} \rho \\ \rho u \\ E \end{bmatrix} \qquad F = \begin{bmatrix} \rho u \\ \rho u^2 + p \\ \rho u \end{bmatrix}$$

$$F = \begin{bmatrix} \rho u \\ \rho u^2 + p \\ \rho u H \end{bmatrix}$$

守恒差分格式:

$$\left(\frac{\partial U}{\partial t}\right)_{i} + \frac{\widehat{\overline{F}}_{i+\frac{1}{2}} - \widehat{\overline{F}}_{i-\frac{1}{2}}}{\Delta x} = 0$$

上式不能视为对空间导数的中心差分,实际上 依据Lagrange中值定理, $(\hat{F}_{i+\frac{1}{2}} - \hat{F}_{i-\frac{1}{2}})/\Delta x = or \neq \frac{\partial F}{\partial x}$ 仅与 \hat{F}_1 的通量计算格式有关!

Lagrange中值定理:

f(x)在[x_B, x_T]上连续,在(x_B, x_T) 内可导,则在开区间 (x_B, x_T) 内至 少有一点 x_0 ,使下式成立:

$$f'(x_0) = \frac{f(x_T) - f(x_B)}{x_T - x_B}$$



• 二维Euler方程的差分离散

$$\frac{\partial U}{\partial t} + \frac{\partial F}{\partial x} + \frac{\partial G}{\partial y} = 0 \qquad U = \begin{bmatrix} \rho \\ \rho u \\ \rho v \\ E \end{bmatrix} \qquad F = \begin{bmatrix} \rho u \\ \rho u^2 + p \\ \rho uv \\ \rho uH \end{bmatrix} \qquad G = \begin{bmatrix} \rho v \\ \rho uv \\ \rho vH \end{bmatrix}$$

$$\frac{\partial U}{\partial t} + \frac{\partial F}{\partial \xi} \xi_x + \frac{\partial F}{\partial \eta} \eta_x + \frac{\partial G}{\partial \xi} \xi_y + \frac{\partial G}{\partial \eta} \eta_y = 0 \qquad \frac{1}{J} \frac{\partial U}{\partial t} + \frac{\partial}{\partial \xi} \frac{1}{J} \left(F \xi_x + G \xi_y \right) + \frac{\partial}{\partial \eta} \frac{1}{J} \left(F \eta_x + G \eta_y \right) = 0$$

$$\widehat{\overline{F}} = \frac{1}{J} (F\xi_x + G\xi_y) = \frac{1}{J} \begin{bmatrix} \rho U_{\xi} \\ \rho U_{\xi} u + p\xi_x \\ \rho U_{\xi} v + p\xi_y \\ \rho U_{\xi} H \end{bmatrix} \qquad U_{\xi} = u\xi_x + v\xi_y$$

$$\frac{1}{J} \frac{\partial U}{\partial t} + \frac{\partial \widehat{\overline{F}}}{\partial \xi} + \frac{\partial \widehat{\overline{G}}}{\partial n} = 0$$

(未考虑变形网格)

 $\widehat{\overline{G}} = \frac{1}{J}(F\eta_x + G\eta_y) = \frac{1}{J} \begin{bmatrix} \rho U_{\eta} \\ \rho U_{\eta} u + p\eta_x \\ \rho U_{\eta} v + p\eta_y \\ \rho U_{\eta} H \end{bmatrix} \quad U_{\eta} = u\eta_x + v\eta_y$

守恒差分格式?