

## Problem Set 7

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### • Solving Ordinary Differential Equation

#### 1. Result

Run the code `solve_ODE.c`.

```
cindytsai@TURQUOISEA /cygdrive/d/GitHub/Computational_Physics/Assignment/ProblemSet_7
$ gcc -o a.out solve_ODE.c

cindytsai@TURQUOISEA /cygdrive/d/GitHub/Computational_Physics/Assignment/ProblemSet_7
$ ./a.out
```

| Step | x    | Exact_sol | Euler   | ModifiedEuler | ImprovedEuler | RungeKutta3rd |
|------|------|-----------|---------|---------------|---------------|---------------|
| 1    | 0.01 | 0.01000   | 0.01000 | 0.01000       | 0.01000       | 0.01000       |
| 2    | 0.02 | 0.02000   | 0.02000 | 0.02000       | 0.02000       | 0.02000       |
| 3    | 0.03 | 0.03001   | 0.03001 | 0.03001       | 0.03001       | 0.03001       |
| 4    | 0.04 | 0.04002   | 0.04001 | 0.04002       | 0.04002       | 0.04002       |
| 5    | 0.05 | 0.05004   | 0.05003 | 0.05004       | 0.05004       | 0.05004       |
| 6    | 0.06 | 0.06007   | 0.06006 | 0.06007       | 0.06007       | 0.06007       |
| 7    | 0.07 | 0.07011   | 0.07009 | 0.07011       | 0.07012       | 0.07011       |
| 8    | 0.08 | 0.08017   | 0.08014 | 0.08017       | 0.08017       | 0.08017       |
| 9    | 0.09 | 0.09024   | 0.09020 | 0.09024       | 0.09025       | 0.09024       |
| 10   | 0.10 | 0.10033   | 0.10029 | 0.10033       | 0.10034       | 0.10033       |
| 11   | 0.11 | 0.11045   | 0.11039 | 0.11044       | 0.11045       | 0.11045       |
| 12   | 0.12 | 0.12058   | 0.12051 | 0.12058       | 0.12058       | 0.12058       |
| 13   | 0.13 | 0.13074   | 0.13065 | 0.13074       | 0.13074       | 0.13074       |

#### 2. Discussion

Equation

$$\frac{dy}{dx} = y^2 + 1 \text{ and } y(0) = 0$$

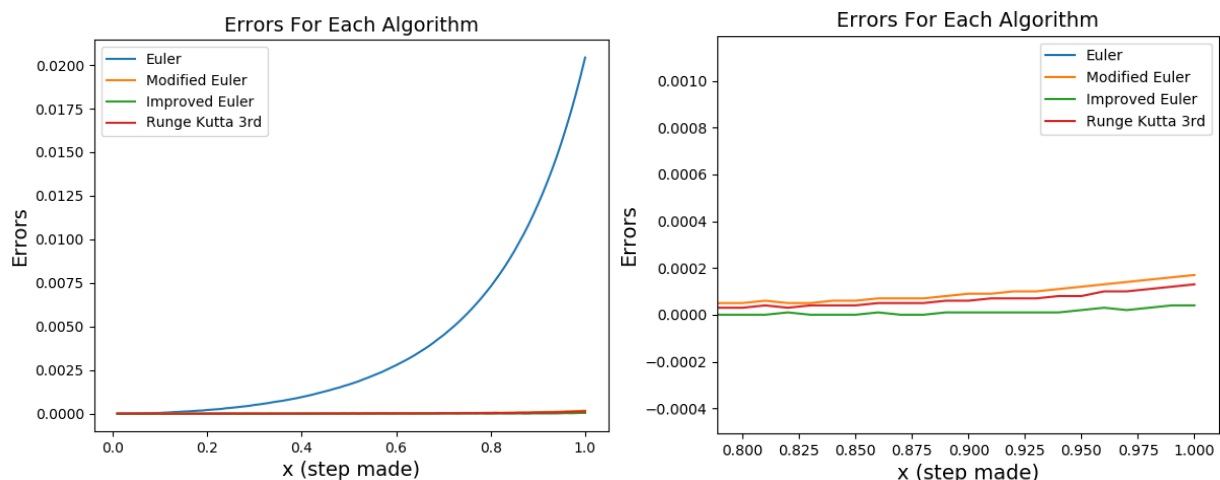
Gives

$$\int \frac{1}{y^2 + 1} dy = \int dx = x + \text{const.} = \tan^{-1} y \Rightarrow y = \tan x$$

And compare each algorithm's solution to the exact solution, where

$\text{error} = |\text{exact solution} - \text{algorithm}|$ , we have:

(In file `solve_ODE_errorplot.py`)



We can see that as each step is approximate by algorithms, the error accumulates.

So errors will grow in every algorithm. We would want to avoid using Euler

method for solving ODE, since the error boost up. It is interesting that the error size is

*Modified Euler > Runge Kutta 3rd > Improved Euler*

Improved Euler method is the same as Runge Kutta 2<sup>nd</sup> order method. I guess that different shape of the differential equations has different best approximation method.

- **Prove the 3<sup>rd</sup> order and 4<sup>th</sup> order Runge-Kutta Formulas**

**1. Prove of 3<sup>rd</sup> order Runge-Kutta Formula**

From

$$\frac{dy}{dx} = f(x, y) \Rightarrow y(x_0 + h) = y(x_0) + \int_{x_0}^{x_0+h} f(x, y(x)) dx$$

We have

$$y(x_0 + h) \cong y(x_0) + h[\alpha \cdot f(x_0, y_0) + \beta \cdot f(x_0 + \gamma \cdot h, y_0 + \delta \cdot hf(x_0, y_0))]$$

And let

$$y_0 = y(x_0) \text{ and } f(x_0, y_0) = f_0$$

Do the Taylor Expansion of  $f(x_0 + \dots, y_0 + \dots)$  and  $y(x_0 + h)$  to the 3<sup>rd</sup> order,

$$\begin{aligned} LHS &= y_0 + h \left\{ \alpha f_0 + \beta \left[ f_0 + \gamma h \frac{\partial f}{\partial x} + \delta h f_0 \frac{\partial f}{\partial y} + \gamma^2 h^2 \frac{\partial^2 f}{\partial x^2} + 2\gamma\delta h^2 f_0 \frac{\partial^2 f}{\partial x \partial y} + \delta^2 h^2 f_0^2 \frac{\partial^2 f}{\partial y^2} \right]_{x_0, y_0} \right\} \\ \Rightarrow LHS &\cong y_0 + h(\alpha + \beta)f_0 + \beta\gamma h^2 \frac{\partial f}{\partial x} + \beta\delta h^2 f_0 \frac{\partial f}{\partial y} + \frac{1}{2}\beta\gamma^2 h^3 \frac{\partial^2 f}{\partial x^2} + \frac{1}{2} \\ &\quad \cdot 2\beta\gamma\delta h^3 f_0 \frac{\partial^2 f}{\partial x \partial y} + \frac{1}{2}\beta\delta^2 h^3 f_0^2 \frac{\partial^2 f}{\partial y^2} \text{ at } (x_0, y_0) \end{aligned}$$

$$RHS \cong y_0 + h \frac{dy}{dx} \Big|_{x_0} + \frac{h^2}{2!} \frac{d^2 y}{dx^2} \Big|_{x_0} + \frac{h^3}{3!} \frac{d^3 y}{dx^3} \Big|_{x_0}$$

$$\Rightarrow RHS = y_0 + hf_0 + \frac{h^2}{2!} \left[ \frac{\partial f}{\partial x} + f_0 \frac{\partial f}{\partial y} \right] + \frac{h^3}{3!} \left[ \frac{\partial^2 f}{\partial x^2} + 2f_0 \frac{\partial^2 f}{\partial x \partial y} + f_0^2 \frac{\partial^2 f}{\partial y^2} \right]$$

where  $\frac{d^2 y}{dx^2} = \frac{df(x, y)}{dx} = \frac{\partial f}{\partial x} + f \frac{\partial f}{\partial y}$  and  $\frac{d^3 y}{dx^3} = \frac{d}{dx} \left[ \frac{\partial f}{\partial x} + f \frac{\partial f}{\partial y} \right] = \frac{\partial^2 f}{\partial x^2} + 2f \frac{\partial^2 f}{\partial x \partial y} +$

$f^2 \frac{\partial^2 f}{\partial y^2} + \frac{df}{dx} \frac{\partial f}{\partial y}$ . We neglect the last term  $\frac{df}{dx} \frac{\partial f}{\partial y}$ , since it's a higher order.

Then compare LHS and RHS term by term, we have

$$\alpha + \beta = 1$$

$$\beta\gamma = \frac{1}{2} \text{ and } \beta\delta = \frac{1}{2}$$

$$\beta\gamma^2 = \frac{1}{3} \text{ and } \beta\gamma\delta = \frac{1}{3} \text{ and } \beta\delta^2 = \frac{1}{3}$$

We can solve it and find out that:

$$\alpha = \frac{1}{4} \quad \beta = \frac{3}{4} \quad \gamma = \frac{2}{3} \quad \delta = \frac{2}{3}$$

Which is 3<sup>rd</sup> order Runge-Kutta Formula

$$y(x_0 + h) \cong y(x_0) + h \left[ \frac{1}{4} \cdot f(x_0, y_0) + \frac{3}{4} \cdot f(x_0 + \frac{2}{3} \cdot h, y_0 + \frac{2}{3} \cdot hf(x_0, y_0)) \right]$$

## 2. Prove of 4<sup>th</sup> order Runge-Kutta Formula

We know that the 4<sup>th</sup> order Runge-Kutta Formula looks just like Simpson's Rule when doing integration and take  $f(x, y) = f(x)$ , so we can assume that the equation is

$$y(x_0 + h) \cong y(x_0) + \frac{h}{6} [f_0 + 2f_1 + 2f_2 + f_3]$$

And

$$f_0 = f(x_0, y_0)$$

$$f_1 = f(x_0 + \frac{h}{2}, y_0 + \alpha hf_0)$$

$$f_2 = f(x_0 + \frac{h}{2}, y_0 + \beta hf_1)$$

$$f_3 = f(x_0 + h, y_0 + \gamma hf_2)$$

We assumed the formula to be like this, in the way let they will shrink to Simpson's Rule when we take  $f(x, y) = f(x)$ .

And again same as 3<sup>rd</sup> order Runge-Kutta Formula, we do the Taylor Expansion at  $(x_0, y_0)$  to find the coefficients.

$$f_1 \cong f_0 + \frac{h}{2} \frac{\partial f}{\partial x} + \alpha hf_0 \frac{\partial f}{\partial y} + \frac{1}{2!} \left( \frac{h^2}{4} \frac{\partial^2 f}{\partial x^2} + \frac{h}{2} \alpha hf_0 \frac{\partial^2 f}{\partial x \partial y} + \alpha^2 h^2 f_0^2 \frac{\partial^2 f}{\partial y^2} \right)$$

$$f_2 \cong f_0 + \frac{h}{2} \frac{\partial f}{\partial x} + \beta hf_1 \frac{\partial f}{\partial y} + \frac{1}{2!} \left( \frac{h^2}{4} \frac{\partial^2 f}{\partial x^2} + \frac{h}{2} \beta hf_1 \frac{\partial^2 f}{\partial x \partial y} + \beta^2 h^2 f_1^2 \frac{\partial^2 f}{\partial y^2} \right)$$

$$f_3 \cong f_0 + h \frac{\partial f}{\partial x} + \gamma hf_2 \frac{\partial f}{\partial y} + \frac{1}{2!} \left( h^2 \frac{\partial^2 f}{\partial x^2} + h\gamma hf_2 \frac{\partial^2 f}{\partial x \partial y} + \gamma^2 h^2 f_2^2 \frac{\partial^2 f}{\partial y^2} \right)$$

Since for  $f_0$ ,  $f_x$  and  $f_y \dots$  are all higher order terms, we neglect them in approximating  $f_2$  and  $f_3$ . Thus  $f_1$  and  $f_2$  inside  $f_2$  and  $f_3$  respectively are  $f_0$ .

Bring them back to  $y(x_0 + h)$  and compare to the Taylor Expansion of itself,

$$y(x_0 + h) \cong y(x_0) + \frac{h}{6} [f_0 + 2f_1 + 2f_2 + f_3]$$

$$y(x_0 + h) \cong y(x_0) + hf_0 + \frac{h^2}{2!} \left[ \frac{\partial f}{\partial x} + f_0 \frac{\partial f}{\partial y} \right] + \frac{h^3}{3!} \left[ \frac{\partial^2 f}{\partial x^2} + 2f_0 \frac{\partial^2 f}{\partial x \partial y} + f_0^2 \frac{\partial^2 f}{\partial y^2} \right]$$

Compare the term  $\frac{\partial f}{\partial y}$ , we have  $2\alpha + 2\beta + \gamma = 3$ .

Compare the term  $\frac{\partial f}{\partial x \partial y}$ , we have  $\alpha + \beta + \gamma = 2$ .

Compare the term  $\frac{\partial^2 f}{\partial y^2}$ , we have  $\alpha^2 + \beta^2 + \frac{\gamma^2}{2} = 1$

Solve these equation, finally we get

$$\alpha = \frac{1}{2} \quad \beta = \frac{1}{2} \quad \gamma = 1$$

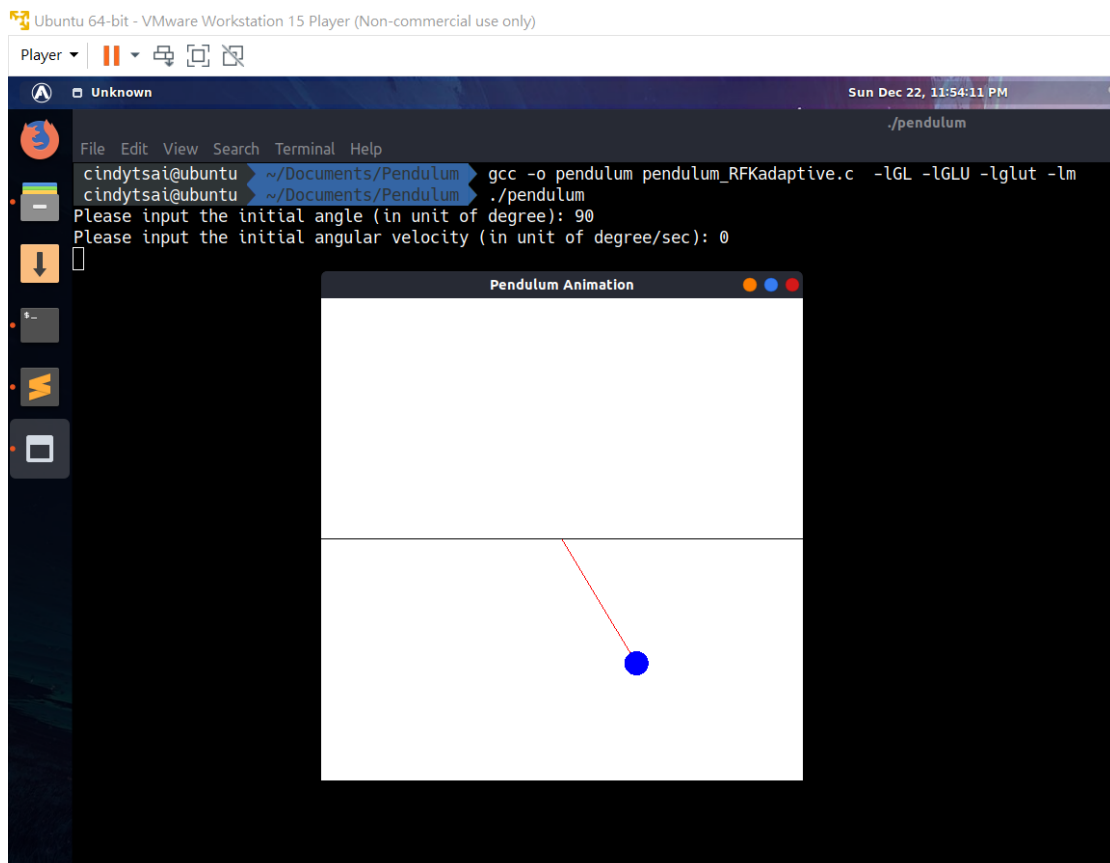
Which give us 4<sup>th</sup> order Runge-Kutta Formula.

- **Physical Pendulum**

- 1. Result**

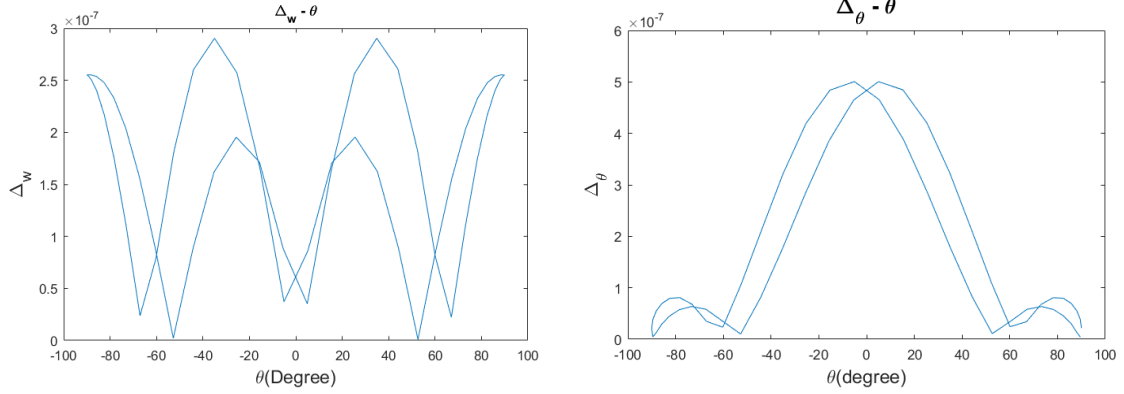
My settings on Cygwin for OpenGL doesn't seem to work, so I use VMware Virtual Machine with Ubuntu Linux instead.

Compile and run the code *pendulum\_RFKadaptive.c*.



## 2. Discussion

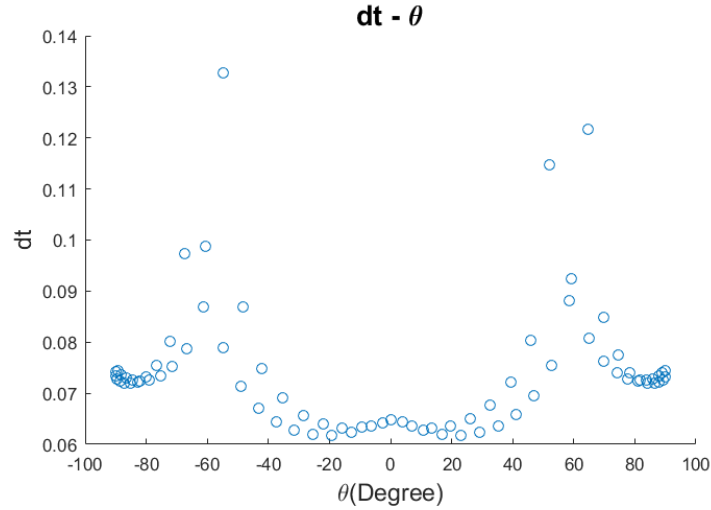
We can see that  $w$  and  $\theta$  are complementary, due to energy conservation.



We can see that we need adaptive step size, if we want them to meet an accuracy under  $\Delta_0 \approx 10^{-7}$ . The above diagram gives how much they vary under  $h$  and  $\frac{h}{2}$ .

If we define the truncation error as  $\Delta \equiv \left| w_{\frac{h}{2}} - w_h \right| + \left| \theta_{\frac{h}{2}} - \theta_h \right|$ , so that we

consider both affect from  $w$  and  $\theta$ , the two only physical parameter, we would have the adaptive step size like this:



We can see that if we add  $\Delta_w$  and  $\Delta_\theta$  together with respect to  $\theta$ , we would see that they are complement to the  $dt - \theta$  diagram, just as what we want the program should do.