

Advanced Fluid Mechanics

Homework 5

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$$\frac{u}{U} = 1 - \exp\left(-\frac{v_0 y}{\nu}\right) \quad (1.1)$$

边界层位移厚度

$$\delta^* = \int_0^\infty \frac{U - u}{U} dy = \int_0^\infty \exp\left(-\frac{v_0 y}{\nu}\right) dy = \frac{\nu}{v_0}. \quad (1.2)$$

粘性系数

$$\mu = \nu \rho, \quad (1.3)$$

表面摩擦力

$$\tau_0 = \mu \frac{du}{dy} \Big|_{y=0} = \rho U v_0. \quad (1.4)$$

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先统一一下符号标记

$$x = z, \quad (2.1)$$

$$\sigma = r, \quad (2.2)$$

$$v = u_\theta. \quad (2.3)$$

由 *Vortical Flows* 中的 (6.1.8b), 对于无粘情况,

$$\frac{D}{Dt} \left(\frac{\omega_\theta}{r} \right) = \frac{1}{r^4} \frac{\partial C^2}{\partial z} = \frac{2C}{r^4} \frac{\partial C}{\partial z}, \quad (2.4)$$

其中

$$C = vr. \quad (2.5)$$

又根据 (6.1.4a)

$$\omega_r = -\frac{1}{r} \frac{\partial C}{\partial z}, \quad (2.6)$$

$$\frac{D}{Dt} \left(\frac{\omega_\theta}{r} \right) = -\frac{2v\omega_r}{r^2}. \quad (2.7)$$

□

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在轴对称假设下, 流体的动能为

$$T = \pi \rho \int_{-\infty}^{+\infty} \int_0^{+\infty} (\psi \omega_\theta + V_\theta^2 r) dr dz. \quad (3.1)$$

对于 Hill 球涡,

$$V_\theta = 0, \quad (3.2)$$

$$\omega_\theta(r, z) = \begin{cases} -\frac{15}{2} \frac{U}{a^2} r, & R = \sqrt{r^2 + z^2} < a, \\ 0, & R = \sqrt{r^2 + z^2} > a, \end{cases} \quad (3.3)$$

$$\psi(r, z) = \begin{cases} -\frac{3}{4} U r^2 \left(1 - \frac{R^2}{a^2} \right) - \frac{3}{4} U r^2 \left(1 - \frac{R^2}{a^2} \right) - \frac{1}{2} U r^2, & R = \sqrt{r^2 + z^2} < a, \\ \frac{1}{2} U r^2 \left(1 - \frac{a^3}{R^3} \right) - \frac{3}{4} U r^2 \left(1 - \frac{R^2}{a^2} \right) - \frac{1}{2} U r^2, & R = \sqrt{r^2 + z^2} > a. \end{cases} \quad (3.4)$$

积分可得

$$T = \pi \rho \int_{-a}^a \int_0^{+\sqrt{a^2 - z^2}} (\psi \omega_\theta) dr dz = \frac{10}{7} \pi \rho U^2 a^3. \quad (3.5)$$

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在极坐标下, 方程

$$\nabla^2 \psi = c\psi \quad (4.1)$$

写为

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \psi}{\partial \theta^2} = c\psi. \quad (4.2)$$

把式子

$$\psi^{(i)} = c J_1(kr) \sin \theta \quad (4.3)$$

代入式 (4.2), 并假设 $c < 0$, $\sin \theta \neq 0$, $x = kr$ 可得

$$x^2 J_1'' + x J_1' - \left(\frac{c}{k^2} x^2 + 1 \right) J_1 = 0. \quad (4.4)$$

当 $k = \sqrt{-c}$ 时, 式 (4.3) 正好是 $m = 1$ 的贝塞尔方程的解.

在极坐标下, 流函数和速度的关系为

$$v_\theta^{(i)} = -\frac{\partial \psi^{(i)}}{\partial r} = -ck J_1'(kr) \sin \theta, \quad v_r^{(i)} = \frac{1}{r} \frac{\partial \psi^{(i)}}{\partial \theta} = \frac{c}{r} J_1(kr) \cos \theta. \quad (4.5)$$

$v_r = 0$ 时

$$J_1(kr_0) = 0. \quad (4.6)$$

假设 kr_0 是 J_1 的第一个零点, 则

$$v_\theta^{(i)} = -ck J_1'(kr_0) \sin \theta. \quad (4.7)$$

均匀来流的流函数

$$\psi^{(o)} = U_\infty \left(1 - \frac{a^2}{r^2} \right) r \sin \theta, \quad a = r_0. \quad (4.8)$$

速度

$$v_\theta^{(o)} = -\frac{\partial \psi^{(o)}}{\partial r} = -U_\infty \left(1 + \frac{a^2}{r^2} \right) \sin \theta, \quad v_r^{(o)} = \frac{1}{r} \frac{\partial \psi^{(o)}}{\partial \theta} = U_\infty \left(1 - \frac{a^2}{r^2} \right) \cos \theta. \quad (4.9)$$

匹配

$$v_\theta^{(o)}(r = a) = v_\theta^{(i)}(r = a) \quad (4.10)$$

得

$$U_\infty = ck J_1'(kr_0)/2. \quad (4.11)$$

得内外流函数

$$\psi = \begin{cases} c J_1(kr) \sin \theta, & r < a, \\ \frac{1}{2} ck J_1'(kr_0) \left(1 - \frac{a^2}{r^2} \right) r \sin \theta, & r > a. \end{cases} \quad (4.12)$$

其中 $k = \sqrt{-c}$, $a = r_0$.

假设 $c = -1$, 对于一阶贝塞尔函数, $r_0 = 3.8317$, 去流函数等值面, 可得流线图, 如图 4.1 所示.

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如图 5.1 所示, 从右上角顺时针绕一圈, 箭头先从东北方向向北旋转, 然后返回为东北方向, 然后再重复一遍, 所以该奇点的指标为 0.

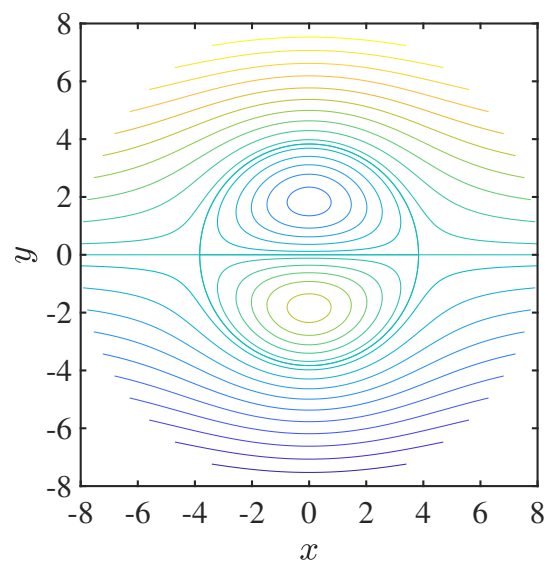


图 4.1. 流线图.

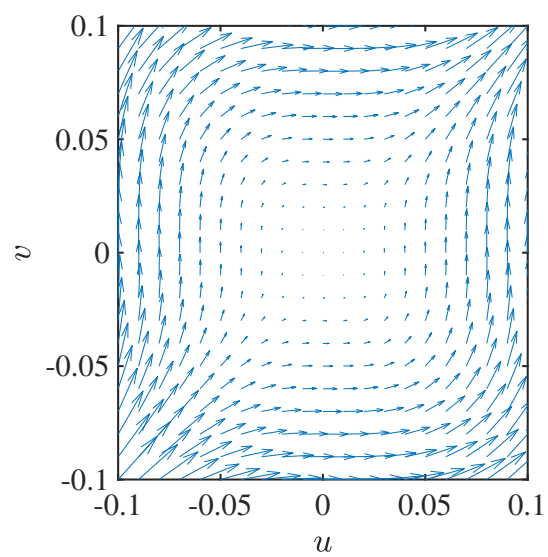


图 5.1. 流场矢量图.

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由于 $\left(\frac{\partial^2 w}{\partial z^2}\right)_0 = -\left[\frac{1}{h_1}\left(\frac{\partial^2 u}{\partial x \partial z}\right) + \frac{1}{h_2}\left(\frac{\partial^2 v}{\partial y \partial z}\right)\right]_0$, 判定流动分离的条件是: 在分离线上

$$\begin{cases} \left(\frac{\partial u}{\partial z}\right)_0 = 0, \\ \left(\frac{\partial^2 u}{\partial x \partial z}\right)_0 < 0, \\ \left(\frac{\partial^2 w}{\partial z^2}\right)_0 > 0 \text{ 或 } \left(\frac{1}{h_1}\frac{\partial^2 u}{\partial x \partial z} + \frac{1}{h_2}\frac{\partial^2 v}{\partial y \partial z}\right)_0 < 0. \end{cases} \quad (6.1)$$

对应于书上的

$$\tau_{x,x} > 0, \quad (6.2)$$

$$\tau_y = 0, \quad \tau_{y,x} = 0, \quad (6.3)$$

$$\tau_{y,y} < 0, \quad (6.4)$$

$$\nabla_\pi \cdot \boldsymbol{\tau} = \tau_{x,x} + \tau_{y,y} < 0. \quad (6.5)$$

注意两书的 x, y 标号是反的.

我从图书馆借了书, 书内容太多, 之前的部分我就不抄了.

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根据三角翼的升力公式

$$L \approx \rho U \int_W y \omega_x \, dS, \quad (7.1)$$

由于三角翼会卷起一个很强的涡, 并且是延着 x 方向, 即 ω_x 很大, 所以升力也很大.

如图 7.1 所示, 发动机的转子布满一圈三角形的叶片, 图示是正视图, 转动方向如图中箭头所示.

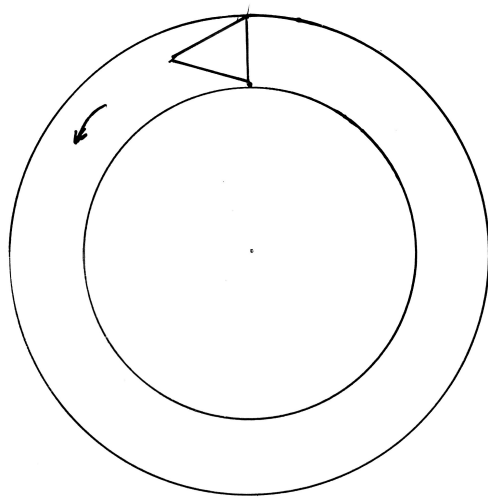


图 7.1. 发动机转子示意图.