

计算流体力学作业 3

College of Engineering 2001111690 袁磊祺

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Cauchy 问题的解

利用特征线理论分析问题

$$\begin{cases} u_t + a(u)u_x = 0, & x \in \mathbb{R}, t > 0, \\ u(x, 0) = u_0(x), & x \in \mathbb{R}. \end{cases} \quad (1.1)$$

并给出 (光滑的) 解.

解: $u_t + a(u)u_x = 0, u(x, 0) = u_0(x)$. 问题转化为

$$\begin{cases} \frac{dx}{dt} = a(u), & \frac{du}{dt} = 0, \\ x(0) = x_0. & u(0) = u_0(x_0). \end{cases} \quad (1.2)$$

由上述 ODE 初值问题得

$$u(x, t) = u_0(x_0), \quad (1.3)$$

$$x = x_0 + a(u_0(x_0))t = x_0 + a(u(x, t))t, \quad (1.4)$$

x_0 依赖于给定的点 (x, t) ,

$$u(x, t) = u_0(x - a(u_0(x_0))t) = u_0(x - a(u_0(x_0(x, t)))t) = u_0(x - a(u(x, t))t). \quad (1.5)$$

由式 (1.1) 得

$$u_t = u'_0(x_0) \frac{\partial x_0}{\partial t}, \quad u_x = u'_0(x_0) \frac{\partial x_0}{\partial x} \quad (1.6)$$

将式 (1.4) 的第一等号两端分别对 t 和 x 求导, 得

$$a(u_0(x_0)) + [1 + a' \cdot u'_0(x_0) \cdot t] \frac{\partial x_0}{\partial t} = 0 \quad (1.7)$$

$$(1 + a' u'_0(x_0) t) \frac{\partial x_0}{\partial x} = 1 \quad (1.8)$$

从式 (1.7) and (1.8) 得

$$\frac{\partial x_0}{\partial t} = -\frac{a(u_0(x_0))}{1 + (a' u'_0)_{x_0} t}, \quad \frac{\partial x_0}{\partial x} = \frac{1}{1 + (a' u'_0)_{x_0} t}. \quad (1.9)$$

将其代入式 (1.6) 知

$$u_t + a(u)u_x = 0. \quad (1.10)$$

$t = 0$ 时

$$u(0) = u_0(x_0), \quad (1.11)$$

所以式 (1.5) 满足式 (1.1).

无黏 Burgers 方程的定解问题

$$\begin{cases} u_t + (0.5u^2)_x = 0, \\ u_0(x) = \cos(\pi x), \quad x \in [-1, 1]. \end{cases} \quad (2.1)$$

此时 $a(u) = u$, 爆破点

$$t^* = -\frac{1}{a' u'_0} = \frac{1}{\pi \sin(\pi x_0)}, \quad (2.2)$$

只在 $x > 0$ 的部分会出现爆破, 最快达到爆破的点为 $x_0 = 0.5$, 经历时间 $t_0^* = \frac{1}{\pi}$.

初始图像如图 2.1 所示, 当其运动到爆破的时候如图 2.2 所示, 可以发现右边的线已经和 x 轴垂直了.

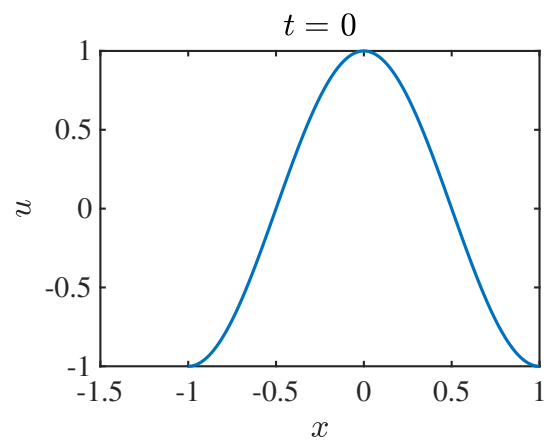


图 2.1. $t = 0$.

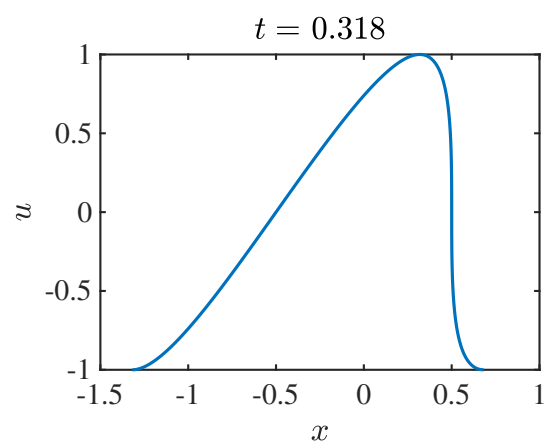


图 2.2. $t = \frac{1}{\pi}$. 爆破时刻.

Burgers 方程 Riemann 问题的弱解

弱解满足的方程为

$$\int_0^{+\infty} \int_{-\infty}^{+\infty} [\phi_t \mathbf{U} + \phi_x \mathbf{F}(\mathbf{U})] dx dt = - \int_{-\infty}^{+\infty} \phi(x, 0) \mathbf{U}(x, 0) dx. \quad (3.1)$$

对 Burgers 方程有

$$\int_0^{+\infty} \int_{-\infty}^{+\infty} \left[\phi_t u + \phi_x \frac{1}{2} u^2 \right] dx dt = - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx. \quad (3.2)$$

激波弱解

$$u(x, t) = \begin{cases} u_L, & x < st, \\ u_R, & x > st. \end{cases} \quad (3.3)$$

其中

$$s = \frac{u_L + u_R}{2}. \quad (3.4)$$

$$\int_0^{+\infty} \int_{-\infty}^{+\infty} \left[\phi_t u + \phi_x \frac{1}{2} u^2 \right] dx dt \quad (3.5)$$

$$= \int_0^{+\infty} \int_{-\infty}^{+\infty} \phi_t u dx dt + \int_0^{+\infty} \int_{-\infty}^{+\infty} \phi_x \frac{1}{2} u^2 dx dt \quad (3.6)$$

$$= - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx + \int_0^{+\infty} \phi(x, x/s) (u_L - u_R) dx \quad (3.7)$$

$$+ \frac{1}{2} \int_0^{+\infty} \phi(st, t) (u_L^2 - u_R^2) dt \quad (3.8)$$

$$= - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx + s \int_0^{+\infty} \phi(x, x/s) (u_L - u_R) d(x/s) \quad (3.9)$$

$$+ \frac{1}{2} \int_0^{+\infty} \phi(st, t) (u_L^2 - u_R^2) dt \quad (3.10)$$

$$= - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx. \quad (3.11)$$

$u_L < u_R$ 弱解

$$u(x, t) = \begin{cases} u_L, & x < st, \\ u_R, & x > st. \end{cases} \quad (3.12)$$

其中

$$s = \frac{u_L + u_R}{2}. \quad (3.13)$$

同样的

$$\int_0^{+\infty} \int_{-\infty}^{+\infty} \left[\phi_t u + \phi_x \frac{1}{2} u^2 \right] dx dt \quad (3.14)$$

$$= \int_0^{+\infty} \int_{-\infty}^{+\infty} \phi_t u dx dt + \int_0^{+\infty} \int_{-\infty}^{+\infty} \phi_x \frac{1}{2} u^2 dx dt \quad (3.15)$$

$$= - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx + \int_0^{+\infty} \phi(x, x/s) (u_L - u_R) dx \quad (3.16)$$

$$+ \frac{1}{2} \int_0^{+\infty} \phi(st, t) (u_L^2 - u_R^2) dt \quad (3.17)$$

$$= - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx + s \int_0^{+\infty} \phi(x, x/s) (u_L - u_R) d(x/s) \quad (3.18)$$

$$+ \frac{1}{2} \int_0^{+\infty} \phi(st, t) (u_L^2 - u_R^2) dt \quad (3.19)$$

$$= - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx. \quad (3.20)$$

稀疏波弱解

$$u(x, t) = \begin{cases} u_L, & x < s_m t \\ u_m, & s_m t \leq x \leq u_m t \\ \frac{x}{t}, & u_m t \leq x \leq u_R t \\ u_R, & x > u_R t \end{cases} \quad (3.21)$$

也是一个弱解, 其中 $u_m \in [u_L, u_R]$ 为任意, $s_m = \frac{u_L + u_m}{2}$. 不妨设 $u_L > 0$, $u_R > 0$.

$$\int_0^{+\infty} \int_{-\infty}^{+\infty} \left[\phi_t u + \phi_x \frac{1}{2} u^2 \right] dx dt \quad (3.22)$$

$$= - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx + \int_0^{+\infty} \left(-u_L \phi \left(x, \frac{x}{s_m} \right) + u_m \phi \left(x, \frac{x}{s_m} \right) \right. \quad (3.23)$$

$$\left. -u_m \phi \left(x, \frac{x}{u_m} \right) + \int_{\frac{x}{u_R}}^{\frac{x}{u_m}} \phi_t \frac{x}{t} dt + u_R \phi \left(x, \frac{x}{u_R} \right) \right) dx \quad (3.24)$$

$$+ \int_0^{+\infty} \frac{1}{2} \left(-u_R^2 \phi(u_R t, t) + \int_{u_m t}^{u_R t} \frac{x^2}{t^2} \phi_x(x, t) dx + u_m^2 \phi(u_m t, t) + (u_L^2 - u_m^2) \phi(s_m t, t) \right) dt. \quad (3.25)$$

$$= - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx + \int_0^{+\infty} \left(-u_m \phi \left(x, \frac{x}{u_m} \right) + \int_{\frac{x}{u_R}}^{\frac{x}{u_m}} \phi_t \frac{x}{t} dt + u_R \phi \left(x, \frac{x}{u_R} \right) \right) dx \quad (3.26)$$

$$+ \int_0^{+\infty} \frac{1}{2} \left(-u_R^2 \phi(u_R t, t) + \int_{u_m t}^{u_R t} \frac{x^2}{t^2} \phi_x dx + u_m^2 \phi(u_m t, t) \right) dt. \quad (3.27)$$

其中

$$\int_{\frac{x}{u_R}}^{\frac{x}{u_m}} \phi_t \frac{x}{t} dt = \left(\frac{x}{t} \phi \right) \Big|_{t=\frac{x}{u_R}}^{t=\frac{x}{u_m}} + \int_{\frac{x}{u_R}}^{\frac{x}{u_m}} \frac{x}{t^2} \phi dt \quad (3.28)$$

$$= u_m \phi \left(x, \frac{x}{u_m} \right) - u_R \phi \left(x, \frac{x}{u_R} \right) + \int_{\frac{x}{u_R}}^{\frac{x}{u_m}} \frac{x}{t^2} \phi dt, \quad (3.29)$$

$$\int_{u_m t}^{u_R t} \frac{x^2}{t^2} \phi_x dx = \frac{x^2}{t^2} \phi \Big|_{x=u_m t}^{x=u_R t} - \int_{u_m t}^{u_R t} \frac{2x}{t^2} \phi dx \quad (3.30)$$

$$= u_R^2 \phi(u_R t, t) - u_m^2 \phi(u_m t, t) - \int_{u_m t}^{u_R t} \frac{2x}{t^2} \phi dx. \quad (3.31)$$

所以

$$\int_0^{+\infty} \int_{-\infty}^{+\infty} \left[\phi_t u + \phi_x \frac{1}{2} u^2 \right] dx dt \quad (3.32)$$

$$= - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx + \int_0^{+\infty} \int_{\frac{x}{u_R}}^{\frac{x}{u_m}} \frac{x}{t^2} \phi dt dx - \int_0^{+\infty} \int_{u_m t}^{u_R t} \frac{x}{t^2} \phi dx dt \quad (3.33)$$

$$= - \int_{-\infty}^{+\infty} \phi(x, 0) u(x, 0) dx. \quad (3.34)$$

□

线化气体动力学方程

线化气体动力学方程组

$$\begin{cases} \frac{\partial \rho}{\partial t} + \rho_0 \frac{\partial u}{\partial x} = 0, \\ \frac{\partial u}{\partial t} + \frac{a_0^2}{\rho_0} \frac{\partial \rho}{\partial x} = 0. \end{cases} \quad (4.1)$$

一般初值问题的解

$$\mathbf{U}_t + \mathbf{A} \mathbf{U}_x = 0, \quad \mathbf{U} = \begin{pmatrix} \rho \\ u \end{pmatrix}, \quad \mathbf{A} = \begin{pmatrix} 0 & \rho_0 \\ a_0^2 / \rho_0 & 0 \end{pmatrix} \quad (4.2)$$

其中

$$\mathbf{R}^{(1)} = \begin{pmatrix} \rho_0 \\ -a_0 \end{pmatrix}, \quad \mathbf{R}^{(2)} = \begin{pmatrix} \rho_0 \\ a_0 \end{pmatrix}, \quad \begin{matrix} \lambda_1 = -a_0 \\ \lambda_2 = a_0 \end{matrix} \quad (4.3)$$

$$\mathbf{R}^{-1} = \frac{1}{2\rho_0 a_0} \begin{pmatrix} a_0 & -\rho_0 \\ a_0 & \rho_0 \end{pmatrix} \quad (4.4)$$

为了给出初值问题的解, 将它转化为相应的规范形

$$\begin{cases} \mathbf{W}_t + \mathbf{\Lambda} \mathbf{W}_x = 0 \\ \mathbf{W}(x, 0) = \mathbf{W}^{(0)} = \mathbf{R}^{-1} \mathbf{U}^{(0)}(x) \end{cases} \quad (4.5)$$

其中

$$\mathbf{\Lambda} = \begin{pmatrix} -a_0 & 0 \\ 0 & a_0 \end{pmatrix}. \quad (4.6)$$

该问题的解是

$$\mathbf{W}(x, t) = \left(W_1^{(0)}(x - \lambda_1 t), W_2^{(0)}(x - \lambda_2 t), \dots, W_m^{(0)}(x - \lambda_m t) \right)^T, \quad (4.7)$$

现在返回到表示式 $\mathbf{U} = \mathbf{RW}$, 即

$$\mathbf{U}(x, t) = \sum_{i=1}^m W_i(x, t) \mathbf{R}^{(i)} = \sum_{i=1}^m W_i^{(0)}(x - \lambda_i t) \mathbf{R}^{(i)}. \quad (4.8)$$

Riemann 问题的解

令 (U_R 类似处理)

$$\mathbf{U}_L = \begin{pmatrix} \rho_L \\ u_L \end{pmatrix} = \alpha_1 \begin{pmatrix} \rho_0 \\ -a_0 \end{pmatrix} + \alpha_2 \begin{pmatrix} \rho_0 \\ a_0 \end{pmatrix} \quad (4.9)$$

不难得

$$\alpha_1 = \frac{a_0 \rho_L - \rho_0 u_L}{2a_0 \rho_0}, \quad \alpha_2 = \frac{a_0 \rho_L + \rho_0 u_L}{2a_0 \rho_0} \quad (4.10)$$

$$\beta_1 = \frac{a_0 \rho_R - \rho_0 u_R}{2a_0 \rho_0}, \quad \beta_2 = \frac{a_0 \rho_R + \rho_0 u_R}{2a_0 \rho_0} \quad (4.11)$$

在星形域 ($\lambda_1 t < x < \lambda_2 t$) 内的解

$$\mathbf{U}^* = \begin{pmatrix} \rho^* \\ u^* \end{pmatrix} = \beta_1 \begin{pmatrix} \rho_0 \\ -a_0 \end{pmatrix} + \alpha_2 \begin{pmatrix} \rho_0 \\ a_0 \end{pmatrix} = \begin{pmatrix} \frac{1}{2}(\rho_L + \rho_R) - \frac{1}{2}(u_R - u_L)\rho_0/a_0 \\ \frac{1}{2}(u_L + u_R) - \frac{1}{2}(\rho_R - \rho_L)a_0/\rho_0 \end{pmatrix}. \quad (4.12)$$

$x < \lambda_1 t$ 的解为 \mathbf{U}_L , $x > \lambda_2 t$ 的解为 \mathbf{U}_R .