

The Harmonic Oscillator

Anthony Mezzacappa

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$$\begin{aligned} H &= \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2 \\ &= \frac{1}{2m} \left[p^2 + (m\omega x)^2 \right] \end{aligned}$$

This has the form $u^2 + v^2$, which can be written as

$$u^2 + v^2 = (iu + v)(-iu + v)$$

Classically, we are dealing with functions. Quantum mechanically we are dealing with operators, so we have to be careful about ordering.

The Hamiltonian operator is

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 \hat{x}^2$$

Define the operators

$$\begin{aligned} \hat{a} &\equiv \frac{1}{\sqrt{2m\omega}}(m\omega\hat{x} + i\hat{p}) \\ \hat{a}^\dagger &\equiv \frac{1}{\sqrt{2m\omega}}(m\omega\hat{x} - i\hat{p}) \end{aligned}$$

Now consider the product

$$\begin{aligned} \hat{a}^\dagger \hat{a} &= \frac{1}{\sqrt{2m\omega}}(m\omega\hat{x} - i\hat{p})(m\omega\hat{x} + i\hat{p}) \\ &= \frac{1}{\sqrt{2m\omega}} \{ m^2 \omega^2 \hat{x}^2 + im\omega[\hat{x}, \hat{p}] + \hat{p}^2 \} \end{aligned}$$

and the product

$$\begin{aligned}\hat{a}\hat{a}^\dagger &= \frac{1}{\sqrt{2m\omega}}(m\omega\hat{x} - i\hat{p})(m\omega\hat{x} + i\hat{p}) \\ &= \frac{1}{\sqrt{2m\omega}} \{m^2\omega^2\hat{x}^2 - im\omega[\hat{x}, \hat{p}] + \hat{p}^2\}\end{aligned}$$

The sum

$$\begin{aligned}\hat{a}^\dagger\hat{a} + \hat{a}\hat{a}^\dagger &= m\omega^2\hat{x}^2 + \frac{\hat{p}^2}{m\omega} \\ &= \frac{2}{\omega}\hat{H}\end{aligned}$$

Then

$$\hat{H} = \frac{1}{2}\omega(\hat{a}^\dagger\hat{a} + \hat{a}\hat{a}^\dagger)$$

Let's consider the commutator of the operators \hat{a} and \hat{a}^\dagger . ($\hbar = 1$)

$$\begin{aligned}[\hat{a}, \hat{a}^\dagger] &= \frac{1}{2m\omega} \{(m\omega\hat{x} + i\hat{p})(m\omega\hat{x} - i\hat{p}) - (m\omega\hat{x} - i\hat{p})(m\omega\hat{x} + i\hat{p})\} \\ &= \frac{1}{2m\omega} \{m^2\omega^2\hat{x}^2 - im\omega[\hat{x}, \hat{p}] + \hat{p}^2 - (m^2\omega^2\hat{x}^2 + im\omega[\hat{x}, \hat{p}] + \hat{p}^2)\} \\ &= -i(i) \\ &= +1\end{aligned}$$

That is

$$[\hat{a}, \hat{a}^\dagger] = 1 \quad \Rightarrow \quad \hat{a}\hat{a}^\dagger = 1 + \hat{a}^\dagger\hat{a}$$

and of course

$$\begin{aligned}[\hat{a}, \hat{a}] &= 0 \\ [\hat{a}^\dagger, \hat{a}^\dagger] &= 0\end{aligned}$$

Then

$$\begin{aligned}\hat{H} &= \frac{1}{2}\omega(\hat{a}^\dagger\hat{a} + \hat{a}\hat{a}^\dagger) \\ &= \frac{1}{2}\omega(\hat{a}^\dagger\hat{a} + 1 + \hat{a}^\dagger\hat{a}) \\ &= \omega(\hat{a}^\dagger\hat{a} + \frac{1}{2})\end{aligned}$$

Now look at

$$\begin{aligned}[\hat{H}, \hat{a}] &= \omega \{(\hat{a}^\dagger\hat{a} + \frac{1}{2})\hat{a} - \hat{a}(\hat{a}^\dagger\hat{a} + \frac{1}{2})\} \\ &= \omega(\hat{a}^\dagger\hat{a}\hat{a} - \hat{a}\hat{a}^\dagger\hat{a}) \\ &= \omega[\hat{a}^\dagger\hat{a}\hat{a} - (1 + \hat{a}^\dagger\hat{a})\hat{a}] \\ &= -\omega\hat{a}\end{aligned}$$

and

$$\begin{aligned}
[\hat{H}, \hat{a}^\dagger] &= \omega \left\{ (\hat{a}^\dagger \hat{a} + \tfrac{1}{2}) \hat{a}^\dagger - \hat{a}^\dagger (\hat{a}^\dagger \hat{a} + \tfrac{1}{2}) \right\} \\
&= \omega (\hat{a}^\dagger \hat{a} \hat{a}^\dagger - \hat{a}^\dagger \hat{a}^\dagger \hat{a}) \\
&= \omega [\hat{a}^\dagger \hat{a} \hat{a}^\dagger - \hat{a}^\dagger (\hat{a} \hat{a}^\dagger - 1)] \\
&= \omega \hat{a}^\dagger
\end{aligned}$$

Label the states by

$$|n\rangle \quad E_n = (n + \tfrac{1}{2})\omega \quad \hat{H} |n\rangle = (n + \tfrac{1}{2})\omega |n\rangle$$

and look at

$$\begin{aligned}
\hat{H} \hat{a} |n\rangle &= \omega (\hat{a}^\dagger \hat{a} + \tfrac{1}{2}) \hat{a} |n\rangle \\
&= \omega (\hat{a}^\dagger \hat{a} \hat{a} + \tfrac{1}{2} \hat{a}) |n\rangle \\
&= \omega [(\hat{a} \hat{a}^\dagger - 1) \hat{a} + \tfrac{1}{2} \hat{a}] |n\rangle \\
&= \omega \hat{a} (\hat{a}^\dagger \hat{a} + \tfrac{1}{2} - 1) |n\rangle \\
&= \hat{a} (\hat{H} - \omega) |n\rangle \\
&= \hat{a} [(n + \tfrac{1}{2})\omega - \omega] |n\rangle \\
&= (n - \tfrac{1}{2})\omega \hat{a} |n\rangle \\
&\Rightarrow \hat{a} |n\rangle \propto |n - 1\rangle
\end{aligned}$$

What is the constant of proportionality?

From

$$\hat{H} |n\rangle = (\hat{a}^\dagger \hat{a} + \tfrac{1}{2})\omega |n\rangle = (n + \tfrac{1}{2})\omega |n\rangle$$

we see that the number operator is

$$\hat{N} = \hat{a}^\dagger \hat{a}$$

Then

$$\langle n | \hat{a}^\dagger \hat{a} | n \rangle = n$$

But

$$\hat{a} |n\rangle = c |n - 1\rangle$$

and

$$\langle n | \hat{a}^\dagger = c^* \langle n - 1 |$$

Then

$$n = \langle n | \hat{a}^\dagger \hat{a} | n \rangle = |c|^2 \langle n - 1 | n - 1 \rangle = |c|^2$$

and

$$c = \sqrt{n}$$

That is

$$\hat{a} |n\rangle = \sqrt{n} |n-1\rangle$$

Now look at

$$\begin{aligned} \hat{H} \hat{a}^\dagger |n\rangle &= \omega(\hat{a}^\dagger \hat{a} + \tfrac{1}{2}) \hat{a}^\dagger |n\rangle \\ &= \omega(\hat{a}^\dagger \hat{a} + \tfrac{1}{2}) \hat{a}^\dagger |n\rangle \\ &= \omega(\hat{a}^\dagger \hat{a} \hat{a}^\dagger + \tfrac{1}{2} \hat{a}^\dagger) |n\rangle & \hat{a} \hat{a}^\dagger - \hat{a}^\dagger \hat{a} = 1 \\ &= \omega \hat{a}^\dagger [(1 + \hat{a}^\dagger \hat{a}) + \tfrac{1}{2}] |n\rangle \\ &= \hat{a}^\dagger (\hat{H} + \omega) |n\rangle \\ &= \hat{a}^\dagger (n + \tfrac{1}{2} + 1) \omega |n\rangle \\ &= (n + \tfrac{1}{2} + 1) \omega \hat{a}^\dagger |n\rangle \\ &\Rightarrow \hat{a}^\dagger |n\rangle \propto |n+1\rangle \end{aligned}$$

But

$$\langle n+1 | \hat{a}^\dagger \hat{a} | n+1 \rangle = n+1$$

and

$$n+1 = \langle n+1 | \hat{a}^\dagger \hat{a} | n+1 \rangle = |c|^2 \langle n+1 | n+1 \rangle = |c|^2$$

Then

$$\hat{a}^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle$$

One of the most important aspects of the spectrum of states of the quantized harmonic oscillator is that its energy levels are separated by a uniform amount: $\hbar\omega$. The state $|n\rangle$ has energy $n\hbar\omega$ and the state $|n+1\rangle$ has energy $(n+1)\hbar\omega$.

With

$$\hat{H} |n\rangle = (n + \tfrac{1}{2}) \hbar\omega |n\rangle$$

we have the vacuum state $|0\rangle$ with energy $\frac{1}{2}\hbar\omega$. Then we can think of the state $|n\rangle$ as a state of n quanta each of energy $\hbar\omega$ – i.e., we can think of it as a multiparticle state.

With this interpretation, the operator \hat{a}^\dagger is a creation operator that creates a quantum of energy $\hbar\omega$, whereas \hat{a} is a annihilation operator that annihilates a quantum of energy $\hbar\omega$.

Now we can generate all of the states beginning with the vacuum state defined by

$$\hat{a} |0\rangle = 0$$

Then

$$|n\rangle = \frac{1}{\sqrt{n!}} (\hat{a}^\dagger)^n |0\rangle$$

$\frac{1}{\sqrt{n!}}$ cancels out the $\sqrt{n!}$ that appears when the \hat{a}^\dagger act n times.

For example,

$$\begin{aligned} |2\rangle &= \frac{1}{\sqrt{2}} \hat{a}^\dagger (\hat{a}^\dagger |0\rangle) \\ &= \frac{1}{\sqrt{2}} \hat{a}^\dagger (\sqrt{1} |1\rangle) \\ &= \frac{1}{\sqrt{2}} \sqrt{1+1} |2\rangle \\ &= |2\rangle \end{aligned}$$