

1. SCALING INTO DIMENSIONAL VARIABLES

This note aims to show how to formally display Benney-Luke equations (BLE) (1.5) with dimensional variables. We start from the potential flow equations with dimensional variables (t, x, y, z) :

PFE

$$(1.1a) \quad \Delta\phi + \partial_{zz}\phi = 0 \quad \text{in } \Omega \subset \mathbb{R}^3,$$

$$(1.1b) \quad \partial_t\eta + \nabla\phi \cdot \nabla\eta - \partial_z\phi = 0 \quad \text{at } z = h,$$

$$(1.1c) \quad \partial_t\phi + \frac{1}{2}|\nabla\phi|^2 + \frac{1}{2}(\partial_z\phi)^2 + g\eta = 0 \quad \text{at } z = h,$$

$$(1.1d) \quad \mathbf{n} \cdot \nabla\phi = 0 \quad \text{on } \partial\Omega_b,$$

$$(1.1e) \quad \partial_z\phi = 0 \quad \text{at } z = 0,$$

where g is the gravitational constant, $9.8m/s^2$. Note that unit of t and x, y, z, η are a second and a meter, respectively.

Now, we formally derive Benney-Luke equations from (PFE) (1.1). First, let us introduce non-dimensional variables

scale_ptob

$$(1.2) \quad \hat{x} = \frac{\sqrt{\mu}}{H_0}x, \quad \hat{y} = \frac{\sqrt{\mu}}{H_0}y, \quad \hat{z} = \frac{1}{H_0}z, \quad \hat{t} = \frac{\sqrt{gH_0\mu}}{H_0}t, \quad \hat{\eta} = \frac{1}{\epsilon H_0}\eta, \quad \hat{\phi} = \frac{\sqrt{\mu}}{\epsilon H_0\sqrt{\epsilon H_0}}\phi.$$

where H_0 is rest depth, the amplitude parameter $\epsilon = a/H_0 \ll 1$, and the dispersion parameter $\mu = (H_0/\lambda)^2 \ll 1$. After that, let Φ denote the seabed potential at $\hat{z} = 0$,

$$(1.3) \quad \Phi(\hat{x}, \hat{y}, \hat{t}) := \hat{\phi}(\hat{x}, \hat{y}, 0, \hat{t}).$$

And then, $\hat{\phi}$ is assumed to have Taylor expansion in terms of z as

taylor

$$(1.4) \quad \hat{\phi}(\hat{t}, \hat{x}, \hat{y}, \hat{z}) = \Phi(\hat{t}, \hat{x}, \hat{y}) - \frac{\mu}{2}\hat{z}^2\hat{\Delta}\Phi + \frac{\mu^2}{24}\hat{\Delta}^2\hat{z}^4\Phi + O(\mu^3\hat{z}^3).$$

(For more detail, [1]) Substituting (1.2) and (1.4) into (PFE) (1.1) and truncating terms up to order $O(\epsilon^2\mu, \epsilon^3)$ yields Benney-Luke equations (BLE) with non-dimensional variables $(\hat{t}, \hat{x}, \hat{y}) \in \mathbb{R}^3$:

BLE

BLE1

$$(1.5a) \quad \partial_{\hat{t}}\Phi - \frac{\mu}{2}\partial_{\hat{t}}\hat{\nabla}^2\Phi + \frac{\epsilon}{2}|\hat{\nabla}\Phi|^2 + \hat{\eta} = 0 \quad \text{in } \Omega,$$

BLE2

$$(1.5b) \quad \partial_{\hat{t}}\hat{\eta} - \frac{\mu}{2}\partial_{\hat{t}}\hat{\nabla}^2\hat{\eta} + \hat{\nabla} \cdot ((1 + \epsilon\hat{\eta})\hat{\nabla}\Phi) - \frac{2}{3}\mu\hat{\nabla}^4\Phi = 0 \quad \text{in } \Omega,$$

bd1

$$(1.5c) \quad \mathbf{n} \cdot \hat{\nabla}\Phi = 0 \quad \text{on } \partial\Omega,$$

bd2

$$(1.5d) \quad \mathbf{n} \cdot \hat{\nabla}(\hat{\nabla}^2\Phi) = 0 \quad \text{on } \partial\Omega.$$

As a result, after solving (BLE) (1.5), the BLE solutions are formally displayed with dimensional variables by substituting (scale_ptob) (1.2) into the solutions.

REFERENCES

- [1] Bokhove, Onno and Kalogirou, Anna, *Variational water wave modelling: from continuum to experiment*, (Bridges, T., Groves, M., and Nicholls, D.), LMS Lecture Note Series, Cambridge University Press, 426, 226–260 2016. 1

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