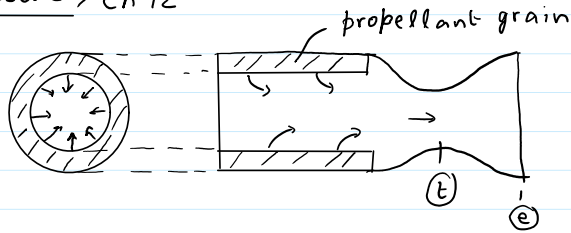


Solid Rocket Motors, Ch 12



Solid propellants

Homogeneous

fuel and oxidizer
in the same molecule

Ex: double-base propellant

(nitroglycerin-nitrocellulose)

Heterogeneous

Fuel and oxidizer
grossly mixed

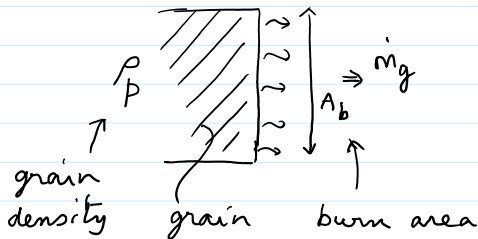
Ex: composite

(plastic-like binder: HTPB)

oxidizing crystals: AP, AN,
NP, KP, KN

Metal additives: Al, Mg, B

Solid propellants are observed to regress at a uniform rate during their burn (regression rate r)



$$\dot{m}_g = \rho_p A_b r$$

The regression rate is related to pressure as follows:

$$r = a p_o^n \quad (12.25)$$

n - pressure exponent (independent of temperature)

a - depends on initial temperature of the grain (T_p):

$$a = \frac{A}{T_i - T_p} \quad (12.26)$$

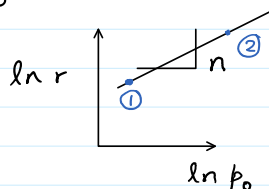
{ A , T_i are constants }

$$(12.25) \Rightarrow \ln r = \ln a + n \ln p_o$$

$$\ln r_2 = \ln a + n \ln p_{o2}$$

$$\ln r_1 = \ln a + n \ln p_{o1}$$

$$(\ln r_2 - \ln r_1) = n(\ln p_{o2} - \ln p_{o1})$$



$$\Rightarrow n = \frac{\ln(r_2/r_1)}{\ln(p_{02}/p_{01})}$$

Quasi-steady model

$$\dot{m}_g = \dot{m}_t$$

$$\dot{m}_g = \rho_p A_b r = \rho_p A_b a p_o^n$$

$$\dot{m}_t = \frac{p_0}{\sqrt{RT_0}} \sqrt{\gamma} \left(\frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{\gamma-1}} A_t \{M_t = 1\}$$

Equating \dot{m}_g to \dot{m}_f yields the steady state operating pressure of the rocket motor

$$p_0 = \left[\frac{A_b}{A_t} \frac{a_{p_p}}{\sqrt{\frac{\gamma}{RT_0} \left(\frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{\gamma-1}}}} \right]^{\frac{1}{1-\gamma}} \quad \text{See (12.29)}$$

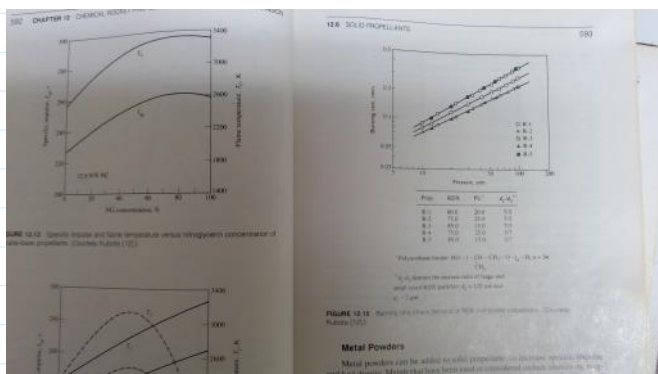
$$p_0 = \left(\frac{A_b}{A_t} a_p c^* \right)^{\frac{1}{1-n}}$$

Ex: Find $\left(\frac{A_b}{A_t}\right)$ required to reach a steady state

operating pressure of 14 MPa in a SRM with the following characteristics:

$$\rho_p = 1710 \text{ kg/m}^3, \gamma = 1.27, T_0 = 2220 \text{ K}, \bar{M} = 23,$$

R-4 propellant, see Fig 12.15 for burn data



At $p_{o1} = 10 \text{ atm}$, $r_1 = 0.07 \text{ cm/s}$

At $p_2 = 100 \text{ atm}$, $r_2 = 0.22 \text{ cm/s}$

$$\text{At } p_{02} = 100 \text{ atm}, r_2 = 0.22 \text{ cm/s}$$

$$n = \frac{\ln\left(\frac{r_2}{r_1}\right)}{\ln\left(\frac{p_{02}}{p_{01}}\right)} = \frac{\ln\left(\frac{0.22}{0.07}\right)}{\ln\left(\frac{100}{10}\right)} = 0.4973$$

$$a = \frac{r_1}{p_{01}^n} = \frac{(0.07)(10^{-2})}{[10(101325)]^{0.4973}} = 7.2186(10^{-7}) \left\{ \frac{\text{m/s}}{\text{Pa}^{0.4973}} \right\}$$

$$p_0 = 14 \text{ MPa}, \rho_p = 1710 \text{ kg/m}^3, \gamma = 1.27, R = \frac{\bar{R}}{\bar{M}} = 361.5 \frac{\text{J}}{\text{kg-K}}$$

$$c^* = \sqrt{\frac{1}{\gamma} R T_0 \left(\frac{\gamma+1}{2} \right)^{\frac{\gamma+1}{\gamma-1}}} = 1353.7 \text{ m/s}$$

$$p_0 = \left(\frac{A_b}{A_t} \alpha \rho_p c^* \right)^{\frac{1}{1-n}}$$

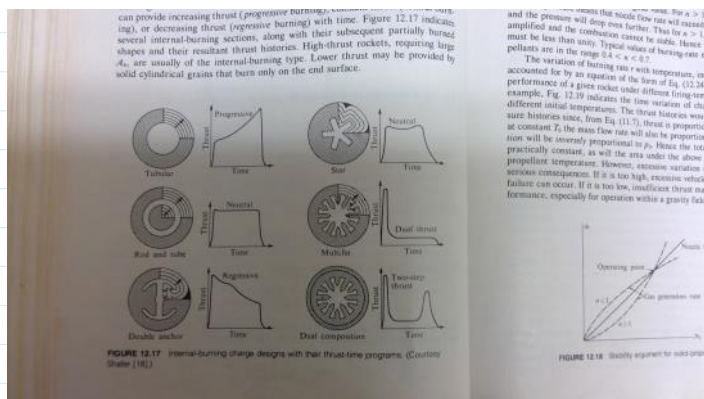
$$\Rightarrow \frac{A_b}{A_t} = \frac{p_0^{1-n}}{\alpha \rho_p c^*} = 2341$$

$$\{ r = a p_0^n = 7.2186(10^{-7}) [14(10^6)]^{0.4973} \}$$

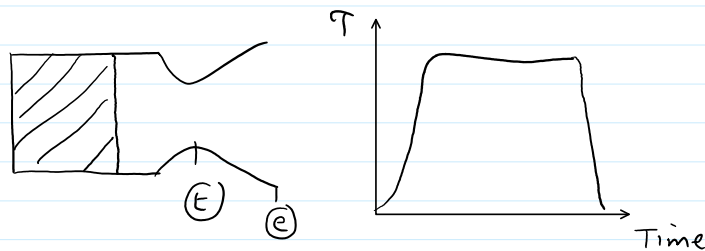
$\uparrow \frac{\text{m/s}}{\text{Pa}^{0.4973}}$

$$= 0.002584 \text{ m/s} = 0.258 \text{ cm/s} \}$$

See Fig 12.17 for schematics of several internal burning grains and associated thrust histories



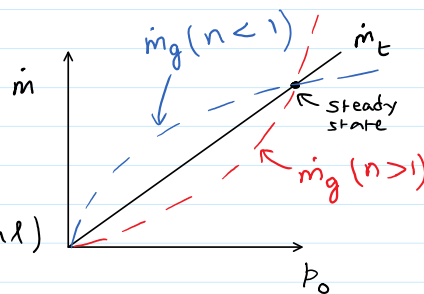
End burning grains burn neutrally.



Burning stability

$n > 1$: explosive behavior
 $n < 1$: stable behavior

$0.4 < n < 0.7$ (typical)



The sensitivity of burning rate to initial propellant temperature is defined as

$$\pi_r \stackrel{\text{def}}{=} \frac{1}{r} \left(\frac{\partial r}{\partial T_p} \right)_{P_0} = \left[\frac{\partial}{\partial T_p} (\ln r) \right]_{P_0} \quad (12.30)$$

$$r = a P_0^n, \quad \ln r = \ln a + n \ln P_0$$

$$\pi_r = \left[\frac{\partial (\ln r)}{\partial T_p} \right]_{P_0} = \frac{d(\ln a)}{dT_p} + 0$$

$$\text{Since } a = \frac{A}{T_1 - T_p}, \quad \ln a = \ln A - \ln(T_1 - T_p)$$

$$\frac{d(\ln a)}{dT_p} = + \frac{1}{T_1 - T_p}, \quad (\text{or}) \quad \frac{1}{a} \frac{da}{dT_p} = \frac{1}{T_1 - T_p}$$

$$\pi_r = \frac{1}{T_1 - T_p} \quad (12.31)$$

$$0.002 \text{ K}^{-1} < \pi_r < 0.008 \text{ K}^{-1} \quad (\text{typical } \pi_r \text{ range})$$

The sensitivity of p_0 to initial grain temperature is defined as

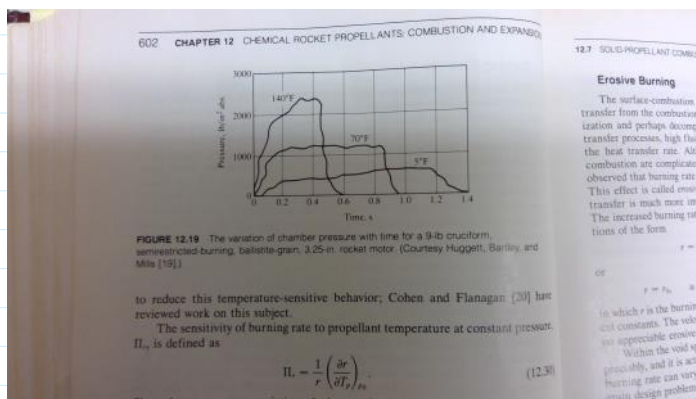
$$\Pi_p = \frac{1}{p_0} \left(\frac{\partial p_0}{\partial T_p} \right)_{\frac{A_b}{A_t}, T_0} = \left[\frac{\partial}{\partial T_p} (\ln p_0) \right]_{\frac{A_b}{A_t}, T_0}$$

From $p_0 = \left(\frac{A_b}{A_t} a p_p c^* \right)^{\frac{1}{1-n}}$, and observing only 'a' is initial temperature dependent,

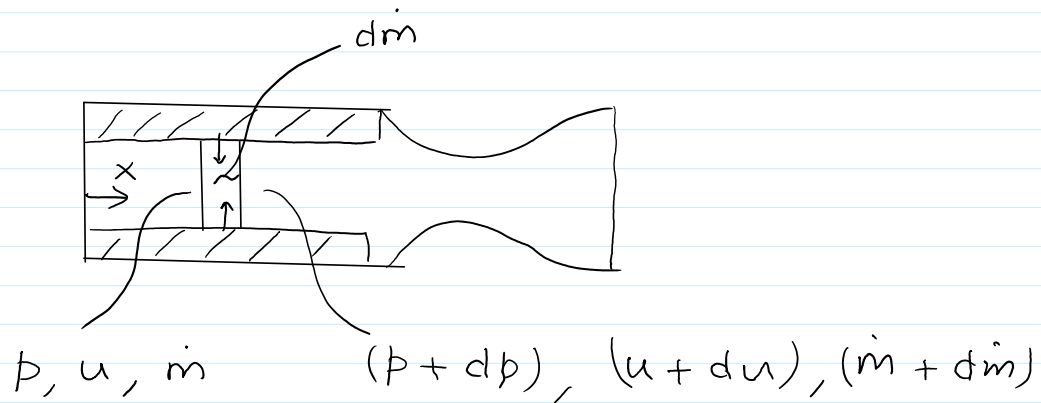
$$\left[\frac{\partial (\ln p_0)}{\partial T_p} \right]_{\frac{A_b}{A_t}, T_0} = \frac{1}{1-n} \frac{d(\ln a)}{dT_p} = \frac{1}{1-n} \Pi_r \quad (12.32)$$

$$\Pi_p = \frac{1}{1-n} \Pi_r = \frac{1}{(1-n)(T_i - T_p)}$$

To lessen the sensitivity of p_0 to initial grain temperature, the propellant should have high T_i and low n . Fig 12.19 shows the impact of the initial grain temperature on p_0 .



Axial pressure variation



x - Momentum:

$$\cancel{\dot{m}u} + (d\dot{m})(0) + \cancel{pA} - (p + dp)A = (\dot{m} + d\dot{m})(u + du) \quad \text{neglect}$$

$$= \cancel{\dot{m}u} + u d\dot{m} + \dot{m} du + (d\dot{m})(du) \quad \text{2nd order}$$

$$\Rightarrow -A dp = u d\dot{m} + \dot{m} du = d(\dot{m}u)$$

$$\int_{p=p_1}^p A dp = \int_{\dot{m}u=0}^{(\dot{m}u)_x} d(\dot{m}u)$$

$$\Rightarrow A(p_1 - p) = (\dot{m}u)_x$$

$$\Rightarrow p = p_1 - \frac{(\dot{m}u)_x}{A} \quad (12.35)$$

Since $u = \frac{\dot{m}}{\rho A}$,

$$p = p_1 - \frac{1}{\rho} \left(\frac{\dot{m}}{A} \right)^2 \quad (12.36)$$

If c is the circumference of the grain, the mass flow

rate at any station x is

$$\dot{m} = \rho_p (cx) r \quad \{A_b = cx\}$$

$$p = p_1 - \frac{R_T}{p} \left(\frac{\rho_p c x r}{A} \right)^2$$

$$\Rightarrow p = \frac{p_1}{2} \left[1 + \sqrt{1 - 4 R_T \left(\frac{\rho_p c r x}{p_1 A} \right)^2} \right] \quad (12.37)$$