

ABE 307

Date: Nov 5th, 2017

Boundary Layer Velocity Profile

Consider a flow of fluid over a flat plate. Far from the plate let us say that the flow is uniform with a constant velocity vo. This is called: Free Street Velocity.

No slip condition, plate is stationary.

So what is the fluid flow like over the plate?

Velocity increases from 0 to approx Voo from y=0 to y=8

Boundary Layer: The thin layer near the vicinity of the plate in which there is a velocity profile, after which the fluid flows with free stream velocity.

Why are we interested in velocity profile in boundary layer?

To be able to calculate the drag forces or shear forces Dimension of Problem? of fluid on the solid object.

Planar flow,

⇒ Assume steady state & incompressible fluid.

3D (Vx, Vy).

→ Assume steady state & incompressible fluid.

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- Line

x=0 K—1—>

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Equation of Continuity:

$$\frac{\partial VX}{\partial X} + \frac{\partial VY}{\partial Y} + \frac{\partial VZ}{\partial Z} = 0$$
Equation of Motion: These equation of motions are written to describe flow with

Equation of Motion: These equation of motions are written to describe flow within the boundary layer.

Equation of Motion for incompressible, Newtonian fluid (Navier-Stokes equation) 3 components in Cartesian

$$\alpha$$
-component $\left(\frac{3x}{3x} + \frac{3y}{3y} - \frac{1}{3}\frac{3P}{3x} + \frac{1}{9}\left(\frac{3x}{3x} + \frac{3^2yx}{3y^2}\right)\right)$

$$\lambda - \infty$$
 who want, $\lambda = -\frac{3\lambda}{3} + \frac{3\lambda}{3} + \frac{3\lambda}{3} = -\frac{3\lambda}{3} = -\frac{$

Order of Magnitude Analysis for Various Terms In Order to Simplify the Equation

order of magnitude. $\frac{9x}{50x} + \lambda \hat{A} \cdot \frac{9\hat{A}}{9x^{X}} = -\overline{1 \cdot 9b} + \overline{10} \left[\frac{9x}{9x^{X}} + \frac{9\hat{A}_{5}}{9x^{X}} \right]$ $\frac{9x}{9x^{X}} + \lambda \hat{A} \cdot \frac{9\hat{A}_{5}}{9x^{X}} = -\overline{1 \cdot 9b} + \overline{10} \left[\frac{9x}{9x^{X}} + \frac{9\hat{A}_{5}}{9x^{X}} \right]$ $\frac{9x}{9x^{X}} + \lambda \hat{A} \cdot \frac{9\hat{A}_{5}}{9x^{X}} = -\overline{1 \cdot 9b} + \overline{10} \left[\frac{9x}{9x^{X}} + \frac{9\hat{A}_{5}}{9x^{X}} \right]$ $\frac{9x}{9x^{X}} + \lambda \hat{A} \cdot \frac{9\hat{A}_{5}}{9x^{X}} = -\overline{1 \cdot 9b} + \overline{10} \left[\frac{9x}{9x^{X}} + \frac{9\hat{A}_{5}}{9x^{X}} \right]$ Sov 6 $\frac{9\times}{9\times}$ $\sqrt{\frac{r}{\Lambda\infty}}$ 3/x ~ 0 (8) L NX SVX ~ D(rog) $\frac{9x}{9x} + \frac{9\lambda}{9n\lambda} = 0$ $\frac{3y}{\sqrt{3}\sqrt{x}} \sim 0 \left(\frac{\sqrt{\infty}}{8}, \frac{\sqrt{\infty}}{8}\right)$ $\frac{89}{904} = \frac{89}{-900}$ ry = (3xx, 04 1 3 3 V (Voo 2/2) $=\frac{\sqrt{\infty}.8}{1}$ $\frac{3x^2}{3^2 \sqrt{x}} = \frac{3x}{3} \left(\frac{3x}{3 \sqrt{x}} \right)^2 0 \left(\frac{1}{1} \cdot \frac{\sqrt{\infty}}{\sqrt{x}} \right)$ $\sqrt[8]{50}\left(\frac{\sqrt{\infty}}{\sqrt{2}}\right)$ $\frac{345}{20x} = \frac{94}{9}\left(\frac{94}{90x}\right) \sim 0\left[\frac{8}{1}, \frac{8}{100}\right] \sim 0\left(\frac{85}{100}\right)$ -13P ~ 0 (Va)

$$V_{x} \frac{\partial x}{\partial x} + v_{y} \frac{\partial y}{\partial y} = \frac{1}{2} \frac{\partial P}{\partial x} + \frac{\partial P}{\partial x} \frac{\partial x}{\partial y} + \frac{\partial V_{x}}{\partial y}$$

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 $v_y \cdot \frac{\partial v_y}{\partial y} \sim 0 \left(\frac{8v_0}{2} \cdot \frac{8v_0}{8} \cdot \frac{1}{8}\right) \sim 0 \left(\frac{v_0 \cdot 8}{12}\right)$, $-\frac{1}{5} \frac{\partial P}{\partial y} \sim 0 \left(\frac{8v_0^2}{12}\right)$,

Since $\left(-\frac{1}{5} \cdot \frac{\partial P}{\partial y}\right) < < 0 \left(\frac{-1}{5} \cdot \frac{\partial P}{\partial x}\right)$ we can ignore, y-combonant of equation of use can ignore, y-combonant of equation of

so, our final equations that we need to 55 solve for getting velocity profile in Boundary layer one:
(b) $\frac{\partial Vx}{\partial x} + \frac{\partial Vy}{\partial y} = 0$ (continuity equation), Prandth Boundary $\frac{\partial Vx}{\partial x} + \frac{\partial Vy}{\partial y} = -\frac{1}{9} \cdot \frac{\partial P}{\partial x} + \frac{\partial P}{\partial y} \cdot \frac{\partial P}{$
Further simplification, eliminate 3P using Potential flow, aince outside the boundary layer, potential flow theory is applicable.
B.C $v_{x}=0$, $y=0$ & No slip condition. $v_{y}=0$, $y=0$ \tag{No slip condition.} $v_{y}=0$
$\frac{1}{2}g(vx^2+vy^2) + P = \omega t + \frac{1}{2}g(ve^2+vy^2) + \frac{1}{2}g(v$
Vx dvx + vy. dvx = vedve + nodvx

From Prandfly equations - we desive on the Von-Kasıman momentum integral equation.

Deriving the Von-Karman Momentum Integral Equation

The equation obtained for Boundary Layer:

$$v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} = v_e \frac{dv_e}{dx} + v \frac{\partial^2 v_x}{\partial y^2} - 4.4.11$$

Step 1: Transform this equation into PDE with single dependent variable

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Hint: Use equation of continuity. Leliminate vy.

$$\frac{\partial Vx}{\partial x} + \frac{\partial Vy}{\partial y} = 0 \implies Vy = -\int \frac{\partial Vx}{\partial x} \cdot \frac{\partial y}{\partial y}$$

$$Vx \cdot \frac{\partial Vx}{\partial x} - \int \frac{\partial Vx}{\partial x} \cdot \frac{\partial Vx}{\partial y} = Ve \cdot \frac{\partial Ve}{\partial x} + o \cdot \frac{\partial Vx}{\partial y^2}.$$

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This is now a PDE in single dependent variable that can be integrated to get Von-Karman equation.

Step 2: Multiply the new Eqn () with density
$$\rho$$
 and integrate from $y=0$ to $y=\infty$

$$\int_{0}^{\infty} \int_{0}^{\infty} \int_{0$$

Your equations now has four different terms to be integrated.

Step 3: Integrate Term IV using
$$\frac{\partial v_y}{\partial y} = 0$$
 at $y = \infty$

Such that $\frac{\partial V_y}{\partial y} = 0$

Why is this boundary condition true?

 $\frac{\partial V_x}{\partial y} = 0$; because outside boundary

dependent on y .

Step 4: Simplify Term II using integral by parts.

If $\frac{\partial v_y}{\partial y} = 0$

Use $dv = \frac{\partial v_y}{\partial y} dy$; $u = \int_0^y \frac{\partial v_y}{\partial x} dy$ between $y = 0$ and $y = \infty$

Integrate $\frac{\partial v_y}{\partial y} dy$; $\frac{\partial v_y}{\partial x} = 0$

Differentiate $\frac{\partial v_y}{\partial x} dy = 0$

Differentiate $\frac{\partial v_y}{\partial x} dy = 0$

The first $\frac{\partial v_y}{\partial x} dy$

Substitute the modified form of II and IV to get the new equation:

Take the term that has viscosity μ on one side and rest of the term on one side of equation.

$$\frac{\partial y}{\partial y} = \int ve. \frac{\partial v}{\partial x} \frac{\partial y}{\partial y} - \int vx. \frac{\partial x}{\partial x} \frac{\partial y}{\partial x}$$

$$+ \int ve \int \frac{\partial x}{\partial x} \frac{\partial y}{\partial y} - \int vx. \frac{\partial x}{\partial x} \frac{\partial y}{\partial y}$$

What does the term with viscosity represent?

Shear Stress

Now to the last equation, add and subtract the term
$$\rho \int_0^\infty v_x \frac{dv_e}{dx} dy$$
 $\frac{\partial Vx}{\partial y} = \frac{g}{g} = \frac{g}{g}$

Simplification now involves collection of terms together.

$$= \int_{0}^{\infty} (ve^{-vx}) \cdot \frac{dve}{dx} \cdot dy$$

$$+ \int_{0}^{\infty} \frac{d(vevx)}{dx} \cdot dy \cdot - \int_{0}^{\infty} \frac{d(wx)}{dx} \cdot dx$$

$$+ \int_{0}^{\infty} \frac{d(vevx)}{dx} \cdot dy + \int_{0}^{\infty} \frac{d(vevx)}{dx} \cdot dy$$

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Von-karman momentum integral equation.