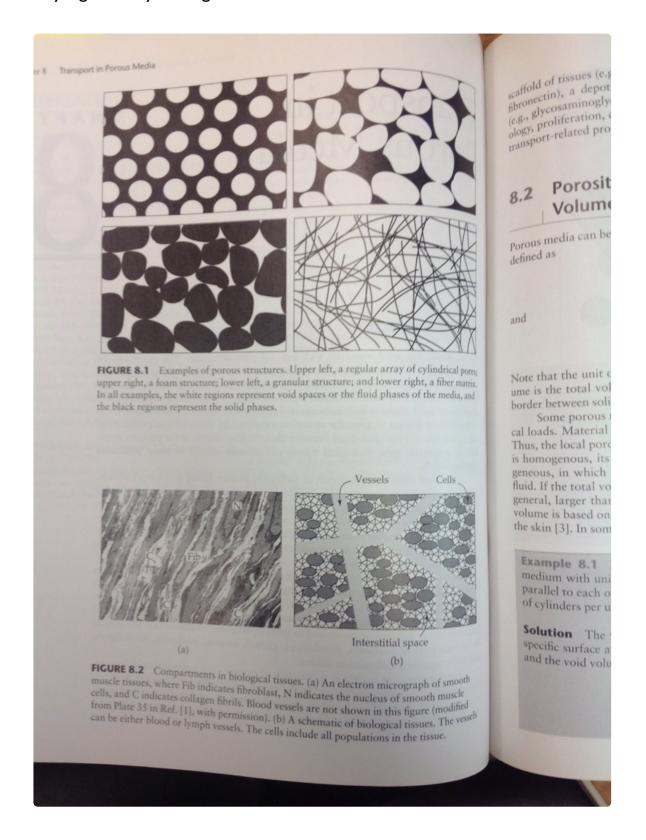
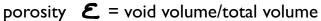
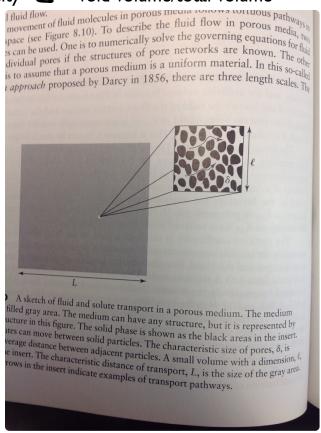
Porous media are solid materials with internal pore structures. Porous structures vary significantly among different media as shown below.



A regular array of cylindrical pores can be found in micro or nanofabricated materials. A foam structure is composed of a continuous solid phase with interconnected channels or isolated pores and is often observed as a sponge. A granular structure exhibited by a pile of sand, consists of solid particles and the void space beteen them. A fiber matrix is the primary structure in polymer gels. Biological tissues can contain seeal of these structures simultaneously. Definitions:

specific surface s = Total interface area/total volume





It is an average value and does not provide information on how different pores are connected or on how any pores are available for water and solute transport. Pores can further be subdivided into penetrable pores and isolated pores. As the name suggests, isolated pores are not available for solute transport. As shown in the figure above, the path between points A and B in a porous medium is tortuous depending on the nature of the medium. In other words, the length

traversed in going from point A to B is greater than the shortest length. One can define tortuosity T as $T = \binom{L}{L_{chart}}^2$

where L and Lshort refer to the actual path length and shortest pathlength respectively. One can see that tortuosity is always greater than one. Not all pores are accessible to solutes. Accessibility will depend on molecular properties of the solute. If the solute molecule is larger than the pore size, it was a solute.

properties of the solute. If the solute molecule is larger than the pore size, it will be inaccessible. The partition coefficient of a solute is the ratio of available volume and the total volume, i.e.

$$\phi = \frac{\text{Available Volume}}{E(\text{total Volume})}$$
 (2)

For a circular pore and a flexible random coil solute molecule such as protein, the partition coefficient is given by,

$$\phi_{\text{cylinder}} = 4 \sum_{n=1}^{\infty} \frac{1}{n^2} \exp\left[-\frac{2}{n} \lambda_c^2\right]$$
 (3)

$$\lambda_{c} = \frac{Rg}{R_{pore}} (4)$$

where Rg, the radius of gyration of protein is given by,

$$R_{3}^{2} = \frac{1}{2N^{2}} \sum_{i} \sum_{j} \pi_{ij}^{2} (5)$$

 \mathcal{P}_{ij} being the distance between residues i and j and N is the total number of residues. In eq. (3), of are the zeros of $\mathcal{J}(\mathcal{C}_n) = \mathcal{D}_j$ being Bessel function of zero order.

Flow through porous medium:

One can average all the properties over a length scale of I to obtain smooth volume averaged properties. Therefore, one can still consider the porous medium as a continuum over length scale I.

y = velocity of fluid volume averaged over

y = velocity of fluid volume averaged through pores.

Continuity:
$$\nabla \cdot \tilde{V} = 0 \quad (7)$$

$$\nabla \cdot \tilde{V} = 0 \quad (7)$$
Darcy's law:
$$\tilde{V} = -K \nabla P \quad (8)$$

$$K = \text{hydraulic conductivity}$$

$$V = -K \nabla P \quad (8)$$

Substituting in equation of continuity, one obtains,

For constant K (homogeneous medium), one obtains,

$$\nabla \cdot \nabla p = 0$$
 $\nabla p = 0$
(10) Laplace Equation

One dimensional flow in rectangular coordinates:

$$\frac{dp}{dx^{2}} = 0$$

$$\frac{dp}{dx} = c_{1} \quad p = c_{1} \times + c_{2}$$

$$x = 0 \quad p = p_{1}$$

$$x = 0 \quad p = p_{2}$$

$$c_{1} = \frac{(P_{2} - P_{1})}{R}$$

$$c_{1} = \frac{(P_{2} - P_{1})}{R}$$

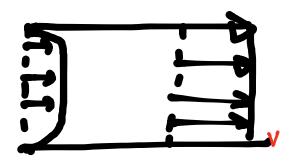
$$P = \frac{P_{1} - (P_{1} - \frac{P_{2}}{R})}{A} \left(\frac{\frac{X}{R}}{R}\right)$$

$$V_{X} = -\frac{K}{\frac{dP}{dx}} = -KC_{1}$$

$$V_{X} = \frac{K}{R} \frac{(P_{1} - P_{2})}{R} \quad \text{Flux}$$

Note that is constant. Therefore, $V = \frac{V}{E}$ is also a constant. In other

words, the resistance to flow near the walks much greater than the viscous resistance. Therefore, if you plot the velocity profile within a pore, you will obtain a flat velocity profile with a steep gradient near the wall. Therefore, Darcy's law does NOT satisfy NO SLIP condition. Darcy's law is valid only for dense porous materials or for low values of hydraulic conductivity K.



Darcy's law is applicable only when the porous medium is dense, I.e. The permeability is low. When the resistance of fluid flow is not large, the approximation of no slip at the wall breaks down. Consequently, one cannot use Darcy's law. The equation of motion has to be modified as

$$\mu \nabla \overset{2}{V} - \frac{1}{K} \overset{V}{\sim} - \nabla p = 0 \quad (11)$$
viscous resistance

This is known as Brinkman's equation.

Let us examine interstitial flow through cells. The two cell membranes are separated by h (see schematic below).

Lubrication Approximation

$$\mu \frac{\partial^2 y}{\partial y^2} - \frac{\mu}{R} y_x - \frac{\partial p}{\partial x} = 0 \quad (13)$$

œ۲,

$$\mu \frac{\partial^2 V_X}{\partial y^2} - \frac{\mu}{R} V_X = \frac{\partial p}{\partial x} = Constant$$
function of y
function of Y

$$\frac{d^2V_X}{dy^2} - \frac{1}{R}V_X = \frac{B}{\mu}$$

The solution to the above equation is

$$V = C_1 \sinh \left(\frac{y}{\sqrt{R}}\right)$$

$$+ C_2 \cosh \left(\frac{y}{\sqrt{R}}\right) (5)$$

$$- \frac{R}{\mu} B$$

$$B \cdot C : - y = 0 \quad dV_X = 0 \quad (6)$$

$$y = R \qquad \forall x = 0 \qquad (7)$$

$$y = R \qquad B \left[1 - \frac{\cosh(y)}{\cosh(x)} \right]$$

$$y = -\frac{R}{\mu} B \left[1 - \frac{\cosh(y)}{\cosh(x)} \right]$$

$$(8)$$

One can evaluate the volumetric flow rate per unit cross sectional area from

$$9 = \frac{1}{R} \int_{X}^{R} \sqrt{dy}$$

$$= -\frac{R}{\mu} B \left[1 - 2\sqrt{R} \cosh\left(\frac{R}{2\sqrt{R}}\right) \right]$$

The effective hydraulic conductivity is defined as the ratio of fluid flux to pressure gradient and is given by

$$K = \frac{1}{4} \left[\frac{2\sqrt{R} \tanh(\frac{R}{2\sqrt{R}})}{R} \right]$$
channel \(\frac{\pi}{4}\) \(\frac{(9)}{2}\)

