

# Generation of milliCharged particles at the LHC

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## 1 Introduction

We discuss the MC generation of milliCharged particles ( $\chi$ ) in  $pp$  collisions at the LHC in the central rapidity region. Except for Drell-Yan (DY), this is done in a few steps

1. Generate in various ways SM particles ( $A$ ) that can decay as  $A \rightarrow X\chi^+\chi^-$
2. Keep track of the cross-section for SM production of  $A$ .
3. Calculate the branching ratio for  $A \rightarrow X\chi^+\chi^-$  for milliCharge  $Q = 1$ . For different charges, it is easy to scale the branching ratio as  $Q^2$ .
4. Generate the  $A \rightarrow X\chi^+\chi^-$  decay with proper kinematics.
5. Save the “event” for further processing in a root file.

**Note: the plots and discussions in Sections 5.1, 5.2, 5.3, 5.4, 6 are for the Run2 configuration  $\sqrt{s} = 13$  TeV. The modifications for the Run3 configurations  $\sqrt{s} = 13.6$  TeV are discussed in the Appendix.**

## 2 Special (simplest) case: Drell-Yan

The Drell-Yan process  $pp \rightarrow \chi^+\chi^-$  is generated with the external model interface in MadGraph. Note that we cannot use SM  $pp \rightarrow \ell^+\ell^-$  with a different lepton mass because the couplings of the  $Z$  to  $\chi$ s and leptons are different (in our model  $\chi$ s are weak isospin singlets[1][2]).

Technically, this was achieved in MadGraph by simply modifying the mass and couplings of the SM electron. The couplings are  $\varepsilon e$  and  $-\varepsilon e \tan \theta_W$  to the  $\gamma$  and  $Z_0$ , respectively, where  $\varepsilon = Q/e$ .

The  $pp \rightarrow \chi^+\chi^-$  process is generated with up to 1 extra parton. This is not very relevant at lower masses since a large majority of events don’t include an extra parton. At higher masses, however, it can increase cross section by  $O(20-30\%)$ . The MadGraph parameter `xqcut` is set to 0, but there still seems to be a 20 GeV cut on the  $p_T$  of the extra parton.

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\*Run 2 version of the note only

We can't find a parameter to control this, but there appears to be a  $\chi^+\chi^-$  invariant mass cut of 2 GeV. Correspondingly, the cross section plateaus below  $m_\chi = 1$  GeV. This shouldn't matter, however, as production from  $J/\psi$ ,  $\phi$ ,  $\omega$ ,  $\rho$ , etc. is expected to dominate at low masses.

A plot of DY cross section as a function of mass is shown along with the other production modes in Fig. 8. Note that DY is handled specially here: while for the other modes we require the mother particle to have  $|\eta| < 2$ , this does not make sense for DY (unless there is an extra parton produced, the  $\eta$  of the ‘‘mother’’ is  $\pm\infty$ ). Instead we have required at least one mCP to have  $|\eta| < 1$ . This threshold was chosen so that the probability of an mCP pointing in the vicinity of the detector is similar between DY and the other modes, in order to facilitate comparison of cross sections.

### 3 List of non-DY processes

The main processes for  $A \rightarrow e^+e^-X$  at the LHC are:

- $\pi^0 \rightarrow e^+e^-\gamma$  (BR=1.17%, Dalitz decay)
- $\eta \rightarrow e^+e^-\gamma$  (BR=0.7%, Dalitz decay)
- $\eta' \rightarrow e^+e^-\gamma$  (BR=5e-4, Dalitz decay)
- $\omega \rightarrow \pi^0 e^+e^-$  (BR=8e-4, Dalitz decay)
- $\eta' \rightarrow \omega e^+e^-$  (BR=2e-4, Dalitz decay)
- $\eta' \rightarrow \pi^+\pi^-e^+e^-$  (BR=2e-3, skip, see comment in Section 6)
- $V \rightarrow e^+e^-$  ( $V = \text{onia}, \phi, \rho, \omega$ )

Note that if we are only interested in  $\chi$  masses above 100 MeV, Dalitz decays of  $\pi^0$  and decays of  $\eta'$  into  $\omega$  do not contribute. In all cases the BR for  $A \rightarrow \chi^+\chi^-X$  can be obtained by rescaling  $A \rightarrow e^+e^-X$  by a factor of  $Q^2$  times an additional mass-dependent factor. In general this factor consists of a phase space piece  $\sqrt{1 - (2m_\chi/m_A)^2}$  plus an additional piece that arises from the matrix element.

Note also that among these processes, the paper [2] that proposed the milliquan experiment only included onia. Furthermore it is not clear that it included the  $B \rightarrow \psi$  process, it is not clear that it included  $\psi(2S)$ ,  $\Upsilon(2S)$ , and  $\Upsilon(3S)$ , and it imposed requirements  $p_T > 6.5$  GeV (2 GeV) on  $J/\psi$  and  $\Upsilon(1S)$ , respectively. The Letter of Intent [3] also mentions  $\Upsilon(2S)$  and  $\Upsilon(3S)$ , and it is possible that the latest NSF proposal included even more processes. In any case: here we make an attempt to include as many processes as possible and to remove  $p_T$  requirements.

## 4 Branching Ratios of non-DY processes (Q=1).

### 4.1 Dalitz BR

The partial width for  $A \rightarrow e^+e^-\gamma$  can be written as[4, 5]:

$$\frac{d\Gamma}{dq^2} = \frac{2\alpha}{3\pi q^2} \left(1 + \frac{2m_e^2}{q^2}\right) \sqrt{1 - \frac{4m_e^2}{q^2}} \left(1 - \frac{q^2}{m_A^2}\right)^3 |F(q^2)|^2 \Gamma(A \rightarrow \gamma\gamma) \quad (1)$$

where  $m_A$  is the mass of  $A$ ,  $q^2$  is the mass-squared of the  $e^+e^-$  pair, and  $F(q^2)$  is a form factor. Note the  $1/q^2$  factor that results in a highly non-uniform invariant mass distribution sharply peaked at low values. The form factor is such that  $F(0) = 1$  and for pions is usually parametrized near  $q^2 = 0$  as  $F(q^2) = 1 + a \frac{q^2}{m_\pi^2}$  with  $a \approx 0.03$ . The form factor can also be estimated in the Vector Dominance Model (VDM) as

$$|F(q^2)|^2 = \frac{m_\rho^4 + m_\rho^2 \Gamma_\rho^2}{(m_\rho^2 - q^2)^2 + m_\rho^2 \Gamma_\rho^2} \quad (2)$$

where  $m_\rho$  and  $\Gamma_\rho$  are the mass and width of the  $\rho$  meson. The VDM model assumes that the decay proceeds through  $\pi^0 \rightarrow \gamma V^*$ ,  $V^* \rightarrow e^+e^-$  and  $V = \rho$  or  $\omega$ ; equation 2 neglects the difference between  $\rho$  and  $\omega$ .

In the case of  $A \rightarrow e^+e^-X$ , when  $X$  is not  $\gamma$ , the partial width can be written as

$$\begin{aligned} \frac{d\Gamma}{dq^2} = & \frac{\alpha}{3\pi q^2} \left(1 + \frac{2m_e^2}{q^2}\right) \sqrt{1 - \frac{4m_e^2}{q^2}} \cdot \\ & \left[ \left(1 + \frac{q^2}{m_A^2 - m_X^2}\right)^2 - \frac{4m_A^2 q^2}{(m_A^2 - m_X^2)^2} \right]^{3/2} |F_{AX}(q^2)|^2 \Gamma(A \rightarrow X\gamma) \end{aligned} \quad (3)$$

where  $m_X$  is the mass of  $X$ , and the transition form factor  $F_{AX}$  can also be approximated as in equation 2.

For milliCharged particles, Dalitz decays branching ratios can be obtained by integrating equation 1 or 3 from  $q^2 = 4m_\chi^2$  to the kinematical limit, substituting  $m_\chi$  for  $m_e$ , and rescaling by  $Q^2$ . Some numerical results for  $Q = 1$  are given in Table 1. Note the sharp drop in branching ratios with mass. This is due to the  $1/q^2$  behavior. The calculations with the electron and muon masses are in excellent agreement with the PDG.

### 4.2 Vector meson branching ratios

At lowest order the SM decay rate for  $V \rightarrow \ell\ell$  is given by the Van Royen-Weisskopf formula[6, 7]:

$$\Gamma(V \rightarrow \ell\ell) = 4\pi\alpha^2 \frac{f_V^2}{m_V} Q_q^2 (1 - 4x_\ell^2)^{1/2} (1 + 2x_\ell^2) \quad (4)$$

where  $f_V$  is the vector decay constant,  $m_V$  is the vector mass,  $Q_q$  is the charge of the quark that makes up the meson,  $x_\ell = m_\ell/m_V$ , and  $m_\ell$  is the lepton mass.

$m_\chi$ (MeV)	$\pi^0 \rightarrow \chi\chi\gamma$	$\eta \rightarrow \chi\chi\gamma$	$\eta' \rightarrow \chi\chi\gamma$	$\eta' \rightarrow \chi\chi\omega$	$\omega \rightarrow \chi\chi\pi^0$
0.511 ( $=m_e$ ) PDG for $ee$	1.17 e-2 (1.17 $\pm$ 0.04)e-2	6.6 e-3 (6.9 $\pm$ 0.4)e-4	4.6 e-4 (4.7 $\pm$ 0.3)e-4	1.8 e-4 (2.0 $\pm$ 0.4)e-4	7.6 e-4 (7.7 $\pm$ 0.6)e-4
10	2.8 e-3	2.9 e-3	2.5 e-4	5.7 e-5	3.7 e-4
30	3.5 e-4	1.6 e-3	1.8 e-4	1.7 e-5	2.3 e-4
50	1.2 e-5	1.0 e-3	1.4 e-4	4.3 e-6	1.6 e-4
60	2.7 e-7	8.2 e-4	1.3 e-4	1.7 e-6	1.4 e-4
90		4.3 e-4	1.0 e-4		9.2 e-5
105.7 ( $=m_\mu$ ) PDG for $\mu\mu$		3.0 e-4 (3.1 $\pm$ 0.4) e-4	9.2 e-5 (1.1 $\pm$ 0.3) e-4		7.4 e-5 (1.3 $\pm$ 0.2) e-4
150		8.9 e-5	6.8 e-5		3.7 e-5
200		1.2 e-5	4.8 e-5		1.5 e-5
250		1.0 e-7	3.2 e-5		3.6 e-6
400			5.6 e-7		

Table 1: Branching ratios for different Dalitz decay modes as a function of  $m_\chi$  for  $Q = 1$  calculated based on equations 1, 2, and 3. When possible we compare with the values from the 2019 PDG.

Decay mode	PDG BR value	Calculation
$\Upsilon(1S) \rightarrow \tau\tau$	(2.6 $\pm$ 0.1)e-2	(2.4 $\pm$ 0.1)e-2
$\psi(2S) \rightarrow \tau\tau$	(3.1 $\pm$ 0.4)e-3	(3.1 $\pm$ 0.0)e-3
$\phi \rightarrow \mu\mu$	(2.9 $\pm$ 0.2)e-4	(3.0 $\pm$ 0.0)e-4
$\omega \rightarrow \mu\mu$	(7.4 $\pm$ 1.8)e-5	(7.3 $\pm$ 0.2)e-5
$\rho \rightarrow \mu\mu$	(4.6 $\pm$ 0.3)e-5	(4.7 $\pm$ 0.1)e-5

Table 2: Comparison of the PDG and calculated values of the  $V \rightarrow \ell\ell$  branching ratios. The quoted uncertainty on the calculation reflects the uncertainty on the  $V \rightarrow ee$  branching ratio.

Thus the ratio of BR for  $V \rightarrow \chi\chi$  to  $V \rightarrow ee$  is given by

$$\frac{\Gamma(V \rightarrow \chi\chi)}{\Gamma(V \rightarrow ee)} = Q^2 \frac{(1 - 4x_\chi^2)^{1/2}(1 + 2x_\chi^2)}{(1 - 4x_\ell^2)^{1/2}(1 + 2x_\ell^2)} \quad (5)$$

where  $x_\chi = m_\chi/m_V$ , and  $m_\chi$  is the mass of  $\chi$ .

As a sanity check, we use equation 5 to predict several  $V \rightarrow \ell\ell$  branching ratios and compare them with their PDG values, see Table 2. To be fair the tests with  $\ell = \mu$  are not very powerful since  $x_\mu^2 \approx x_e^2 \approx 0$ .

## 5 Generation of SM particles decaying into $\chi^+\chi^-$

The key features of our approach are the following

- Use theory or published data, or some MC to generate  $p_T$  distributions for SM particles saved as histograms in ROOT files (Drell Yan is an exception, as discussed previously).

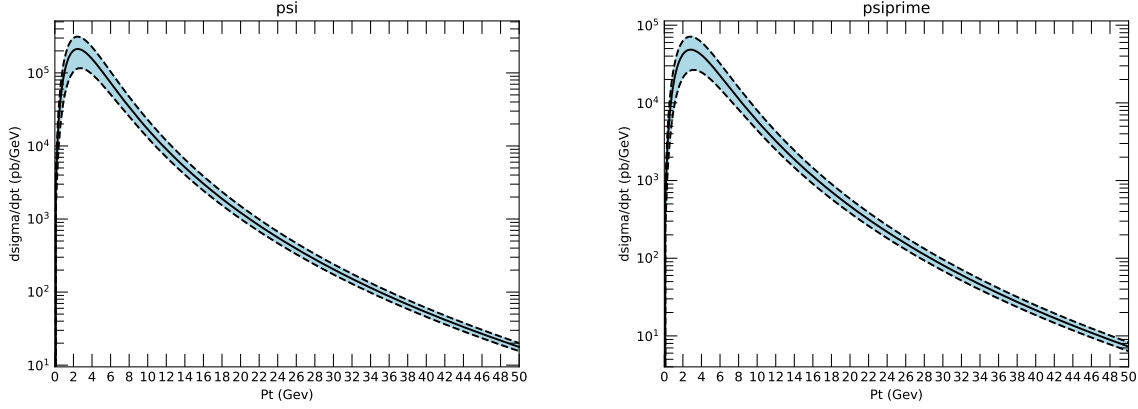


Figure 1: Transverse momentum distributions of  $J/\psi$  (left) and  $\psi'$  from bottom quark decays. Note: this is from a single  $b$ , multiply by two to include  $\bar{b}$ . The units are pb/GeV and the distributions are integrated over  $|\eta| < 1.0$ .

- Sample the ROOT histograms to generate SM particles of a given  $p_T$ .
- Pick azimuthal angles  $\phi$  and pseudorapity  $\eta$  in a limited range, matched to the acceptance of milliquan. The  $\eta$  distribution is assumed to be flat, see comments in Section 6.
- Decay the SM particles into milliCharged particles.
- When possible, keep track of theoretical uncertainties.
- In general it is sufficient to generate SM particles at low and moderate  $p_T$  since that is where the cross-section is largest.

## 5.1 $J/\psi$ and $\psi'$ from b-decays

We use the tool available in

<http://www.lpthe.jussieu.fr/~cacciari/fonll/fonllform.html>

to generate histograms of  $p_T$  distributions (cross-sections) for charmonium from bottom decays, including theoretical uncertainties[9, 10]. See Figure 1

## 5.2 Direct onia production

### 5.2.1 Direct bottonium

There have been many measurement of the  $p_T$  spectra of  $\Upsilon$  in  $pp$  collisions at the LHC by CMS [11, 12, 13, 14], Atlas [15, 16], and LHCb [17, 18, 19, 20, 21]. The LHCb measurements are in the forward region. The only measurement at 13 TeV in the central region is from CMS [14]. Unfortunately, it is limited to  $p_T > 20$  GeV.

Due to the lack of 13 TeV data, initially we planned to use theoretical predictions as a basis of the  $\Upsilon$  event generation. We contacted the theorists [22] that provided the state-of-the-art calculations used to confront the data in Reference [14]. We asked them to extend

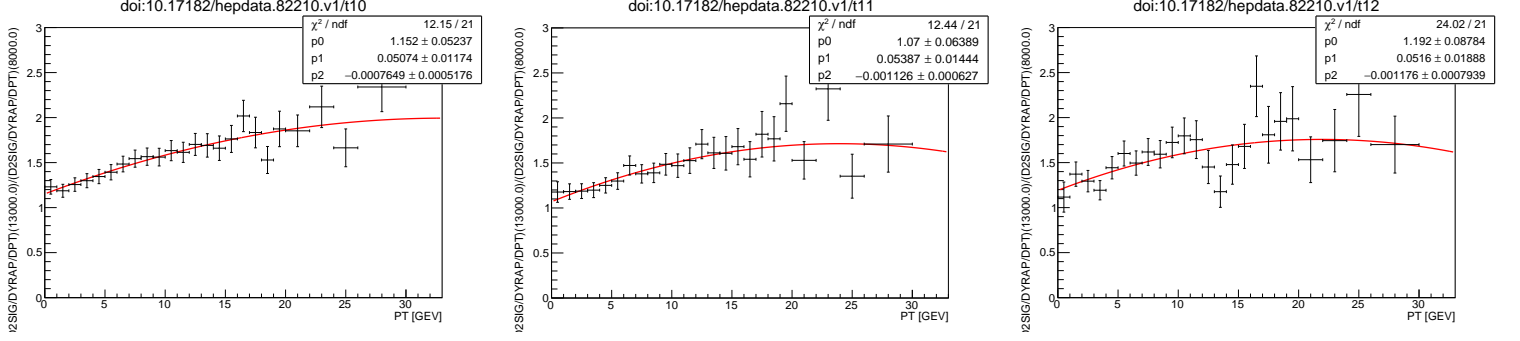


Figure 2: Ratio of 13 to 7 GeV  $\Upsilon$  cross-section for  $2.0 < y < 2.5$  from LHCb [17]. From left to right: 1S, 2S, 3S. The quadratic fits are ours.

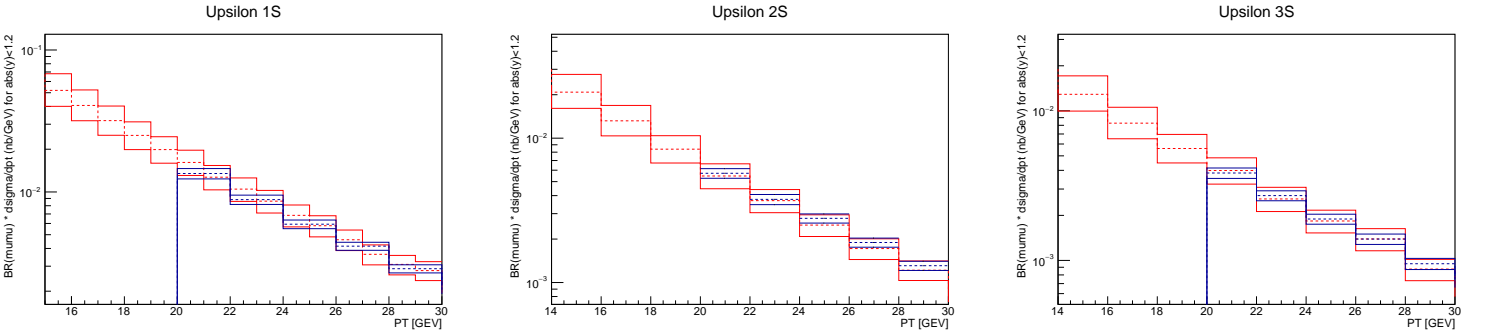


Figure 3: Comparison of the rescaled 7 GeV Atlas  $\Upsilon$  spectra (red) with the 13 GeV CMS spectra (blue) in the neighborhood of 20 GeV, where the matching of the two spectra takes place. From left to right: 1S, 2S, 3S. The dashed lines represent the central values, the solid lines cover the uncertainty range. This is  $\mathcal{B}(\Upsilon \rightarrow \mu\mu) \cdot d\sigma/dp_T$  in nb/GeV integrated over  $|\eta| < 1.2$ .

their predictions to lower  $p_T$ , unfortunately they claim that these are unreliable below 15 GeV.

As a result we decided to use 7 TeV data for  $p_T < 20$  GeV and the CMS 13 TeV data at higher  $p_T$ . A key ingredient is the ratio of 13 and 7 GeV  $\Upsilon$  production cross-sections. These have been measured for  $p_T > 20$  GeV and  $|\eta| < 1.2$  by CMS, see Figure 2 of Reference [14]. The ratios are about 1.7 at  $p_T = 20$  GeV, irrespective of  $\Upsilon$  state (1S, 2S, or 3S), and increase slowly to about 2 at  $p_T = 40$  GeV. The ratios have also been measured by LHCb [17] all the way down to zero  $p_T$  for  $2.0 < y < 2.5$ , see Figure 2. The LHCb ratios in the 20-30 GeV region measured at slightly higher  $y$  are in agreement with the more central ratios measured by CMS.

Consequently, we rescale the **measured** 7 GeV **central** low  $p_T$   $\Upsilon$  spectra to 13 TeV using the curves of Figure 2; we combine these with the **measured** 13 TeV **central** high  $p_T$  spectra to obtain an inclusive 13 TeV spectra. The 7 TeV data is from Atlas [16], since it happens to be in a more convenient format than the equivalent CMS data. On the other hand the 13 TeV spectra are from CMS [14], since there is no Atlas measurement at 13 TeV.

We demonstrate in Figure 3 that the matching of the Atlas and CMS cross-sections works

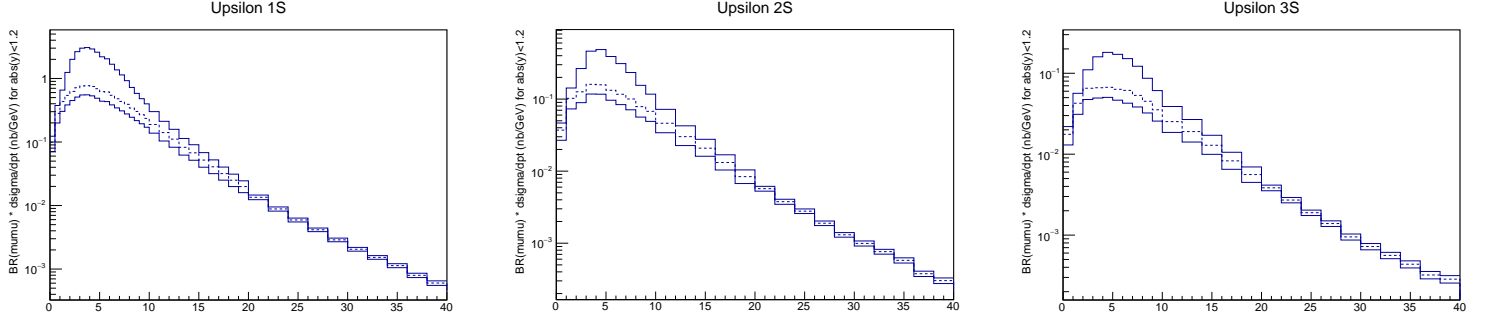


Figure 4: Combined ATLAS 7 TeV, CMS 13 TeV  $\Upsilon$  spectra. From left to right: 1S, 2S, 3S. The dashed line represent the central value, the solid lines cover the uncertainty range. This is  $\mathcal{B}(\Upsilon \rightarrow \mu\mu) \cdot d\sigma/dp_T$  in nb/GeV integrated over  $|\eta| < 1.2$ .

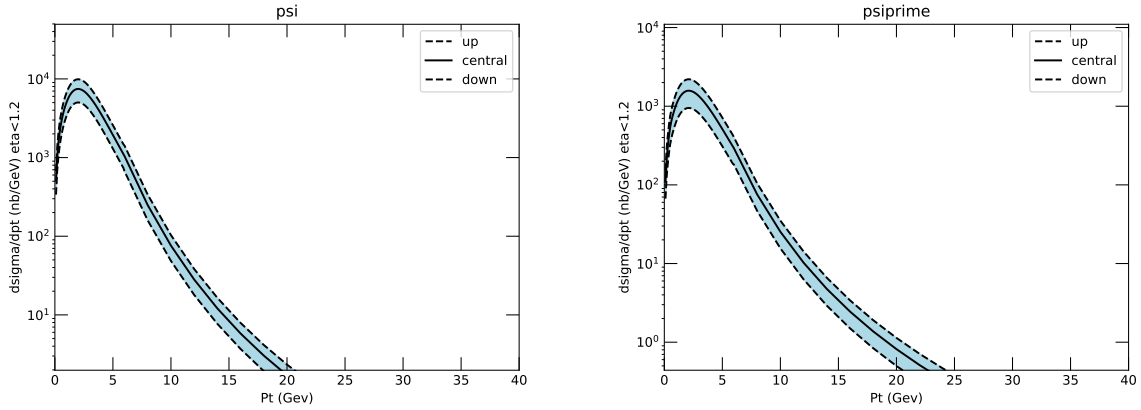


Figure 5: Transverse momentum distributions of  $J/\psi$  (left) and  $\psi'$  from direct production. The units are nb/GeV and the distributions are integrated over  $|\eta| < 1.2$ .

well. The combined spectra to be used in the event generation are in Figure 4.

### 5.2.2 Direct charmonium

We take the charmonium spectra from theory [23, 25, 24] see Figure 5.

## 5.3 $J/\psi$ sanity check

Most of the  $J/\psi$  cross-section is from direct production, see Figure 6, left. As a sanity check we also show the fractions of  $J/\psi$  and  $\psi'$  from b decays measured at lower energy, see Figure 6, right[26].

## 5.4 $\pi^0$ , $\eta$ , $\eta'$ , $\phi$ , $\rho$ , and $\omega$

We generate these from PYTHIA. The measurement of the  $\pi^\pm p_T$  spectrum from CMS[27] is in good agreement with PYTHIA 8 Minimum Bias at low momentum, which gives us some

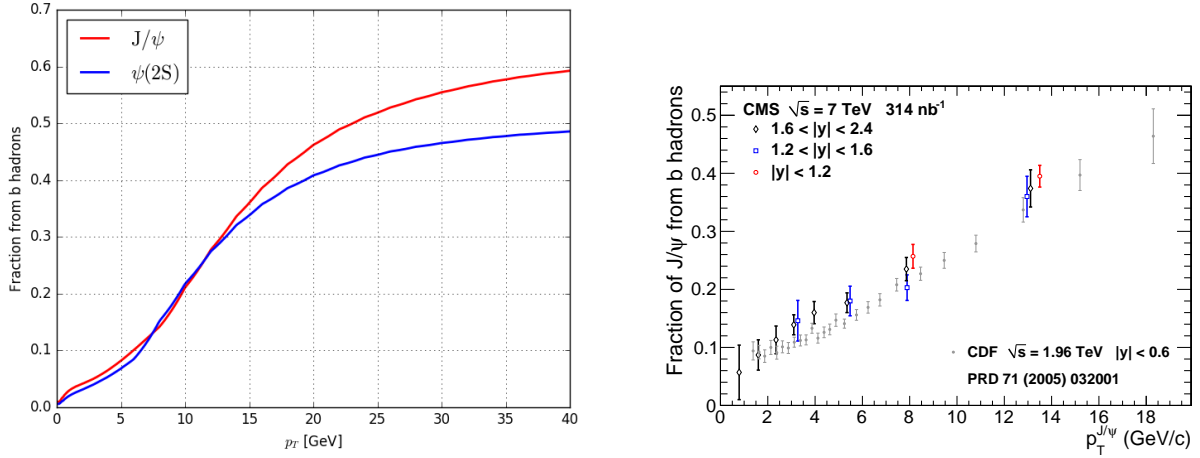


Figure 6: Fraction of charmonium from b decays according to the calculations used in this note (left), and the corresponding measurement ( $J/\psi$  only) at lower energy (right). The left plot is for  $|\eta| < 1.2$ .

confidence in PYTHIA MinBias. However, we compare PYTHIA to actual measurements of specific particles where possible, to guide the choice of tune and any corrections if necessary.

For  $\pi_0$ 's and  $\eta$ 's, ALICE measures the differential cross section at 8 TeV and compares to various PYTHIA tunes [31]. The PYTHIA 8 Monash2013 tune [32] predicts  $\pi_0$ 's well and  $\eta$ 's well until low  $p_T$  ( $< 3$  GeV), where it overpredicts. So we use the Monash2013 tune and scale down low- $p_T$   $\eta$ 's by a function that grows from 1.0 at 3 GeV to 2.0 at 0.5 GeV.

For  $\rho$ 's, ALICE measures the differential cross section at 2.76 TeV [33]. Again Monash2013 agrees well, except at very low  $p_T$  ( $< 1$  GeV). So we use this tune, and scale low- $p_T$   $\rho$ 's down by a function that grows from 1.0 at 1 GeV to 2.0 at 0.5 GeV.

For  $\omega$ 's, ALICE measures the differential cross section at 7 TeV [34]. Good agreement with Monash2013 is seen everywhere, but the measurement only goes down in  $p_T$  to 2 GeV. Since  $\omega$  production should be similar to  $\rho$  production, we use the same low- $p_T$  reweighting scheme as described above.

For  $\phi$ 's, ATLAS measures the differential cross section at 7 TeV [35]. Good agreement is seen with the PYTHIA 6 DW tune [36], while the PYTHIA 8 tunes all underpredict. So for  $\phi$  mesons only, we use the PYTHIA 6 DW tune.

To account for observed disagreements with data and variation between tunes, we apply a 30% uncertainty on the production rate for each of these light mesons, uncorrelated across different meson types.

We do not attempt to use QCD  $2 \rightarrow 2$  at very low  $p_T$  since the process is infrared divergent. PYTHIA `SoftQcd:nonDiffraction` includes all hard QCD processes[28] so in principle this is all that is needed. However, we run out of statistics at high  $p_T$ . So at high  $p_T$  we stitch together the Min Bias distributions with distributions obtained from QCD  $2 \rightarrow 2$  at moderate  $p_T$ .

The QCD samples are “ $p_T$ -binned”, using the `PhaseSpace:pTHat[Min/Max]` parameters in PYTHIA, in bins 15-30 GeV, 30-50 GeV, and 50-80 GeV. The 13 TeV MinBias cross-section



is taken to be  $80 \pm 10$  mb (80 mb is the value measured by ATLAS [29] and predicted by PYTHIA, and the 10 mb uncertainty is to account for a 70 mb measurement by CMS [30]). Then the stitching procedure is the following:

- The QCD samples are first normalized to their LO cross-sections.
- Next, we estimate a “QCD-MinBias scale factor” by integrating over some  $p_T$  region where the MinBias/QCD ratio is roughly flat.
- The QCD samples are all renormalized by this scale factor.
- The samples are stitched together by visually picking the  $p_T$  where the curves cross each other.

The resulting  $p_T$  curves are shown in Figure 7. It is not clear what kind of uncertainties we should assign. Eventually we’ll have to come up with something (50%?).

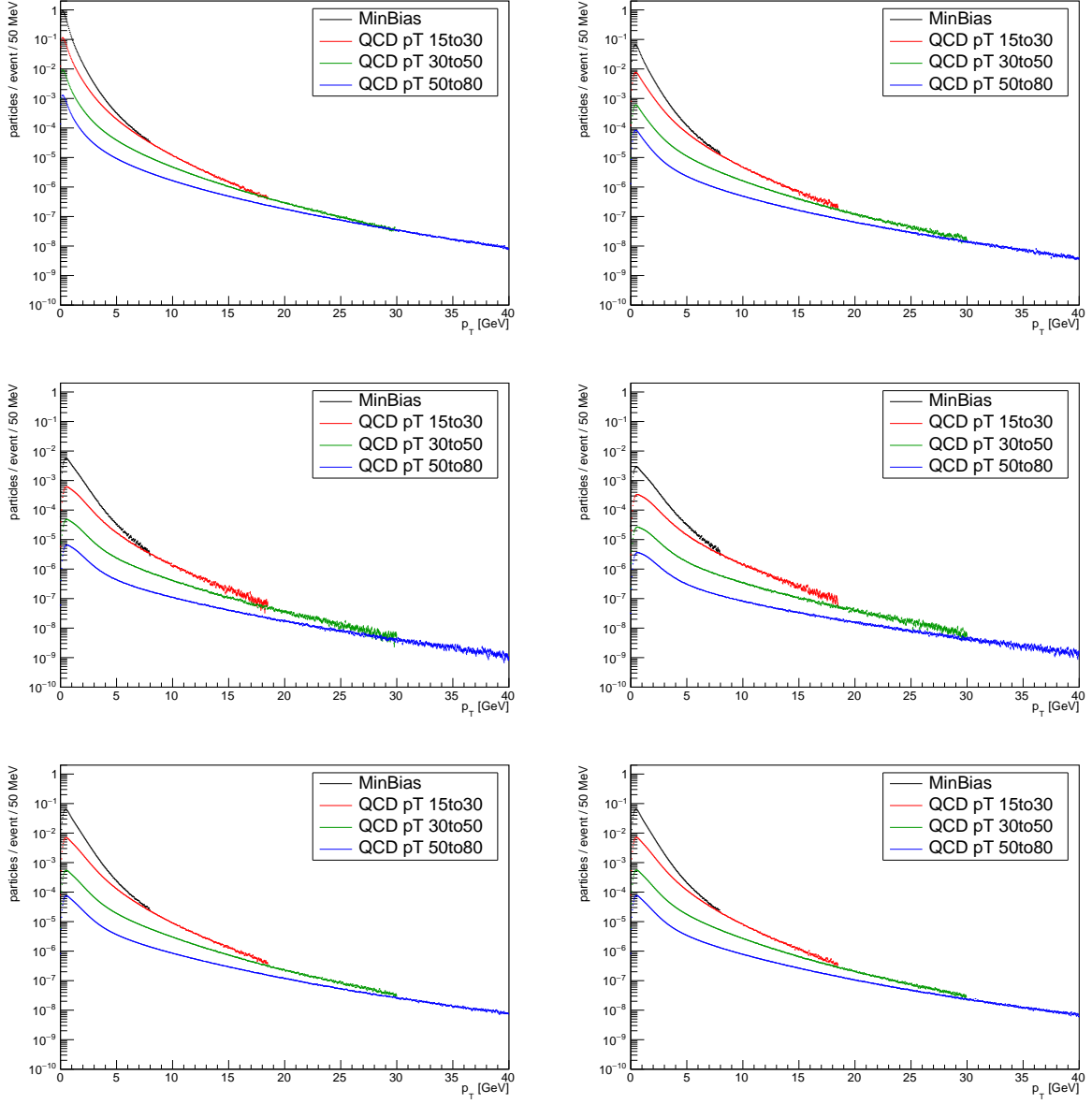


Figure 7: Transverse momentum distributions of  $\pi^0$ ,  $\eta$ ,  $\eta'$ ,  $\phi$ ,  $\rho$ , and  $\omega$ , top left to bottom right, for  $|\eta| < 1$ .

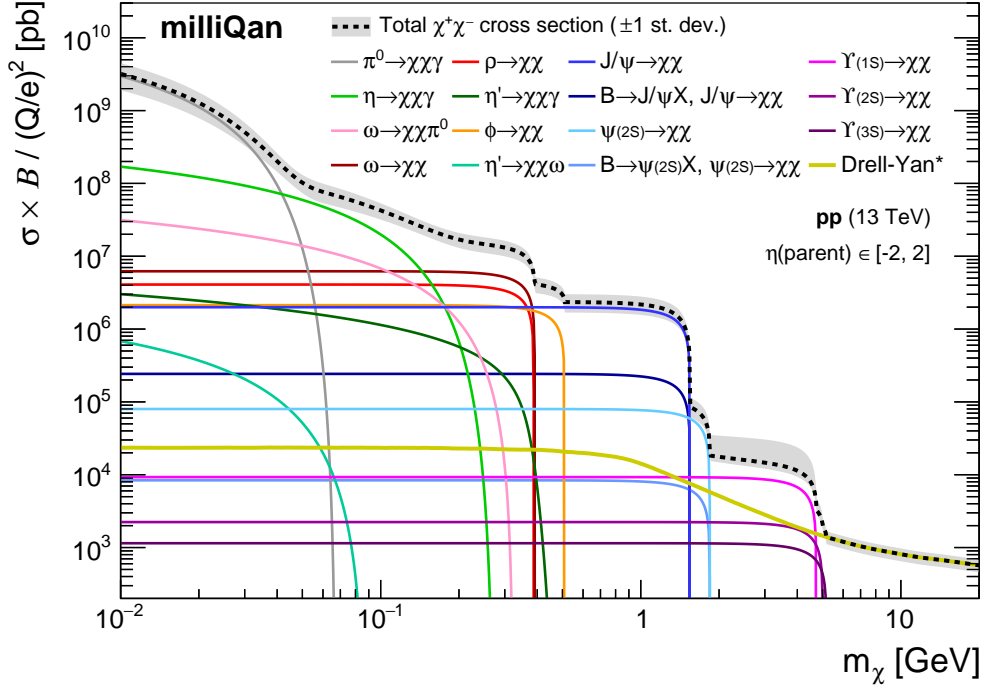


Figure 8: Cross-section times branching ratios for various decays into  $\chi^+\chi^-$  pairs. Note: this is the cross section for producing the mother particle times the branching ratio for the mother particle to decay into  $\chi^+\chi^-$ . Multiply by two to get the  $\chi^\pm$  production cross section. For non-DY modes, the mother particle must be in  $|\eta| < 2$ , but the daughter  $\chi^\pm$  could have  $|\eta| > 2$ . On the other hand some  $\chi$ s with  $|\eta| < 2$  can arise from mothers of  $|\eta| > 2$ . For DY, it doesn't make sense to use “mother  $\eta$ ” (it's almost always  $\pm\infty$ ), so the plotted cross section requires at least one mCP to have  $|\eta| < 1$ . This threshold was chosen so that the probability of an mCP pointing in the vicinity of the detector is similar to that from the other non-DY production modes, in order to facilitate comparison. The plateau in DY cross section below  $m_\chi = 1$  GeV is due to a 2 GeV  $\chi^+\chi^-$  invariant mass cut in MadGraph.

## 6 Putting cross sections and BR together

The total cross section for milliCharged particle production (excluding DY) is shown in Fig. 8. As mentioned in Section 3, we are skipping  $\eta' \rightarrow \pi^+\pi^-\chi^+\chi^-$ . Note that the other contributions from  $\eta'$  in the plot are negligible. Since for  $\chi = e$  the branching ratio of the skipped process is four times higher than that of  $\eta' \rightarrow e^+e^-\gamma$  (Dalitz decay), one worries that this process could actually be important. However, for reasonable  $\chi$  masses, say above 100 MeV, this process is suppressed by phase space with respect to Dalitz decay, e.g., from the PDG,  $\text{BR}(\eta' \rightarrow \mu^+\mu^-\gamma) \approx 1\text{e-}4$  while  $\text{BR}(\eta' \rightarrow \pi^+\pi^-\mu^+\mu^-) < 3\text{e-}5$ . So we can safely ignore it.

As mentioned in the caption of Figure 8 some of the  $\chi$ s in the central region (i.e.: potentially within the acceptance of milliQan) could result from decays of mother particles at higher  $|\eta|$ , especially in the cases where the mother particles have low momentum and the

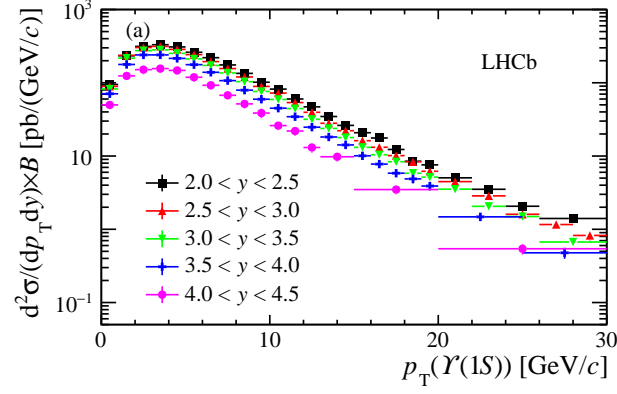


Figure 9:  $\Upsilon(1S)$  spectra in different rapidity ranges from LHCb [17].

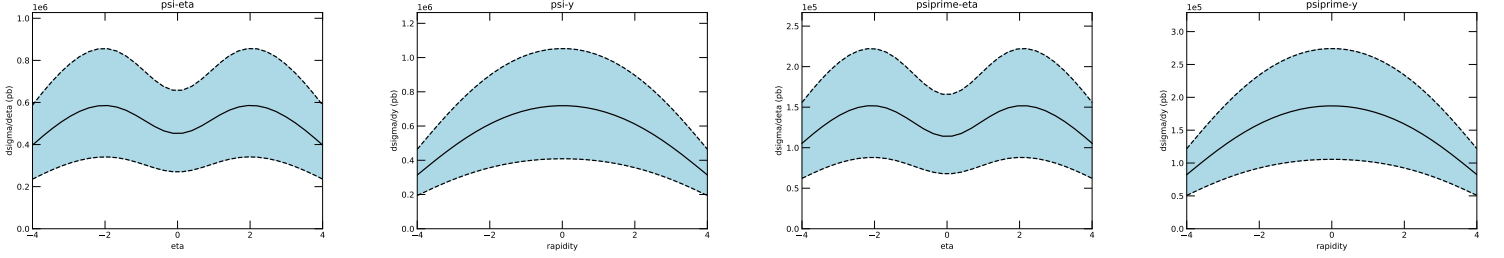


Figure 10: Pseudorapidity and rapidity distributions from  $\psi$  (left) and  $\psi(2S)$  (right) from b-decays[9, 10]. These are for  $p_T < 10$  GeV, i.e., in the low boost regime. Note that in this  $p_T$  range the theoretical uncertainties are large.

boost in the decay is small. Particle production is nearly flat in rapidity, see Figures 9, 10 for some examples. The event generation will then assume a flat pseudorapidity distribution. At the analysis level one will be able to keep track of the contributions from high  $|\eta|$  mothers, and assign a systematic if necessary.

## 7 Generation of the decays

We provide functions that take as input the lab frame 4-vector of either pseudoscalar ( $P$ ) or vector ( $V$ ) meson, and return the 4 vectors of the two  $\chi$ s from the decay. We assume that the  $V$ 's are unpolarized.

### 7.1 Generation of pseudoscalar Dalitz decays

The implementation goes as follows (this also works for the Dalitz decay of the  $\omega$ , which has  $J=1$ , as long as it is unpolarized):

- Rotate the 4-vector of  $P$  from the lab frame into frame  $S_1$  such that  $P$  is traveling in the  $z$ -direction.
- Boost along  $z$  into frame  $S_2$  where  $P$  is at rest.
- Pick a  $q^2$  according to equation 1.
- Generate a decay  $P \rightarrow X\gamma^*$  where  $\gamma^*$  is a particle of  $m^2 = q^2$ . The  $\gamma^*$  direction is random in  $\phi$  and random in  $\cos\theta$ .
- Rotate the  $\gamma^*$  4-vector into a frame  $S_3$  such that the  $\gamma^*$  is traveling in the  $z$ -direction.
- Boost along  $z$  into frame  $S_4$  where  $\gamma^*$  is at rest.
- Generate a decay  $\gamma^* \rightarrow \chi^+\chi^-$  such that the angle  $\phi$  of the  $\chi^+$  is random and  $\cos\theta$  is picked according to [8]

$$\frac{dN}{d\cos\theta} = 1 + \cos^2\theta + \frac{4m_\chi^2}{q^2} \sin^2\theta \quad (6)$$

- Set the 3-vector of the  $\chi^-$  to be back-to-back with the  $\chi^+$ .
- Boost the 4-vectors of the  $\chi$ s from  $S_4$  to  $S_3$ .
- Rotate the 4-vectors of the  $\chi$ s from  $S_3$  to  $S_2$ .
- Boost the 4-vectors of the  $\chi$ s from  $S_2$  to  $S_1$ .
- Rotate the 4-vectors of the  $\chi$ s from  $S_1$  to the lab frame.

### 7.2 Generation of vector decays

The procedure is the following:

- Rotate the 4-vector of  $V$  from the lab frame into frame  $S_1$  such that  $V$  is traveling in the  $z$ -direction.
- Boost along  $z$  into frame  $S_2$  where  $V$  is at rest.

- Generate a decay  $V \rightarrow \chi^+ \chi^-$  such that both the angle  $\phi$  and the  $\cos \theta$  of the  $\chi^+$  are random.
- Set the 3-vector of the  $\chi^-$  to be back-to-back with the  $\chi^+$ .
- Boost the 4-vectors of the  $\chi$ s from  $S_2$  to  $S_1$ .
- Rotate the 4-vectors of the  $\chi$ s from  $S_1$  to the lab frame.

## A Modifications for Run 3

### A.1 $J/\psi$ and $\psi'$ from b-decays - Run 3

We use the exact method described in Section 5.1, changing the center-of-mass energy from 13 to 13.6 TeV. Note that for the Run 2 configuration we had used the CTEQ6.6 pdfs, whereas for the Run 3 configuration we used the NNPDF30\_nlo\_as\_0118 pdfs. The reason for this is simply that the CTEQ6.6 pdf is not an option in the <http://www.lpthe.jussieu.fr/~cacciari/fonll/fonllform.html> tool for 13.6 TeV. Comparisons of Run 2 vs. Run 3 cross-sections as a function of  $P_T$  are shown in Figure 11

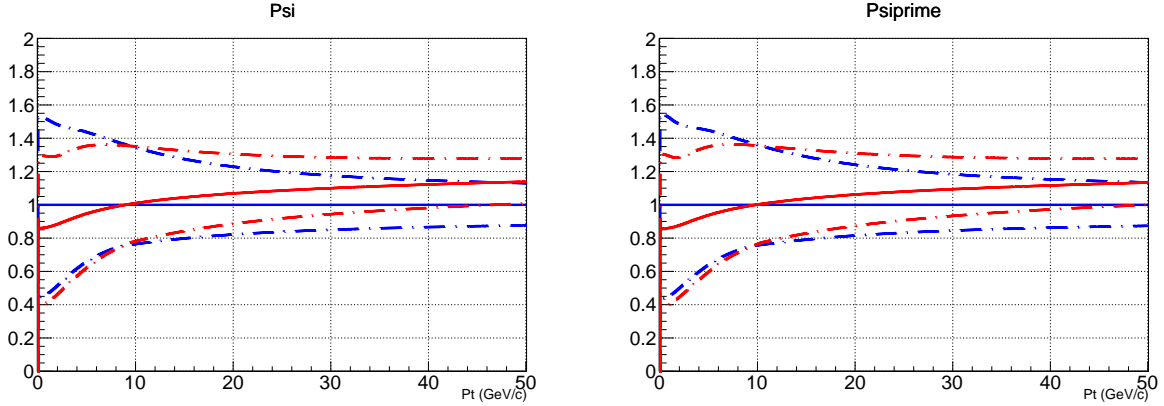


Figure 11: Run 2 vs. Run 3 comparisons of the  $J/\psi$  (left) and  $\psi'$  (right)  $P_T$  spectra from b-decays. The blue band is the Run2 variation as a function of  $P_T$  from the central value. (Dashed blue = up/down divided by central, Solid blue = central = 1.0). The red band is the ratio of the central Run 3 spectrum to the central Run2 spectrum. (Dashed red = up/down Run 3 divided by Run 2 central, Solid red = Run 3 central/Run 2 central) Note that the PDFs have changed between Run2 and Run3. See text for details.

## A.2 Direct bottonium - Run 3

The quadratic functions from Figure 2 were used to extrapolate the LHCb 7 TeV data to 13 TeV (see discussion in Section 5.2.1). We use these same functions to (linearly) further extrapolate the direct bottonium 13 TeV  $P_T$  spectra in Figure 4 up to 13.6 TeV.

**HOWEVER:** The  $P_T$  spectrum used in the generation of the events was within  $|y| < 1.2$ . The event generation assumed a flat rapidity spectrum, which is a good assumption, but it also assumed  $\eta = y$  which is not a good assumption at low  $P_T$ . (Suffices to say that at  $P_T = 0$  the polar angle  $\theta = \pm\pi$ ). If  $\frac{dN}{dy} = C$  is flat, then  $\frac{dN}{d\eta} = C \frac{d\eta}{dy}$ . Since the millikan detector is at  $\eta = 0.11$ , we now also rescale the bottonium spectra by an addotional factor of  $\frac{d\eta}{dy}$  evaluated at  $\eta = 0.11$ .

The factor  $\frac{d\eta}{dy}$  evaluated at  $\eta = 0.11$  depends on the mass and the  $p_T$ . This factor is shown in Fig. 12 for  $\Upsilon(1S)$  and  $J/\psi$  as a function of  $p_T$ . The total Run 2 to Run 3 scale factors are displayed in Figure 13.

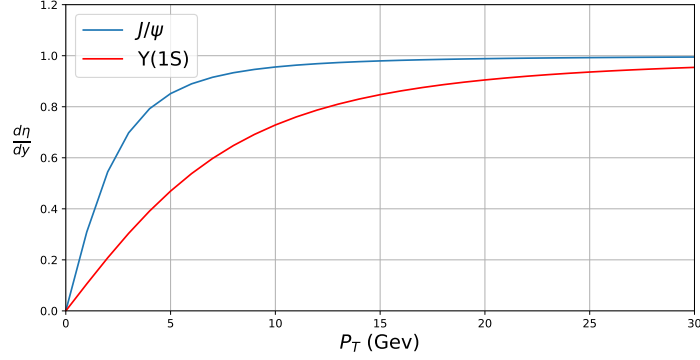


Figure 12:  $\frac{d\eta}{dy}$  evaluated at  $\eta = 0.11$  as a function of  $P_T$

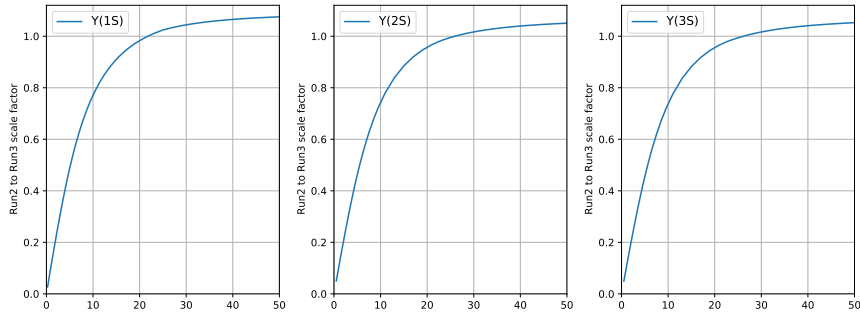


Figure 13: Run 2 to Run 3 scale factors for  $\Upsilon$  production. These are further extrapolations from the LHCb 7 to 13 TeV ratios of Figure 2 and include  $\frac{d\eta}{dy}$  evaluated at  $\eta = 0.11$  as explained in the text.



### A.3 Direct Charmonium - Run 3

As was done for Run 2 (Section 5.2.2) we take the charmonium spectra from theory. However in the process of updating to the Run 3 calculation (provided to us again by Yan-Qing Ma, yqma@pku.edu.cn) we found that due to a miscommunication the  $P_T$  spectra that we were provided were in a rapidity ( $y$ ) instead of pseudo-rapidity ( $\eta$ ) range. This then will necessitate a correction for  $\frac{d\eta}{dy}$ , which will suppress the effective cross-section at low  $P_T$ , as discussed in Section A.2. Furthermore, the  $P_T$  spectra were generated with two different calculations, one appropriate for low  $P_T$  and one appropriate for high  $P_T$ . We have now improved the matching between the low and high  $P_T$  spectra (Fig. 14). The overall changes from Run 2 are summarized in Fig.15

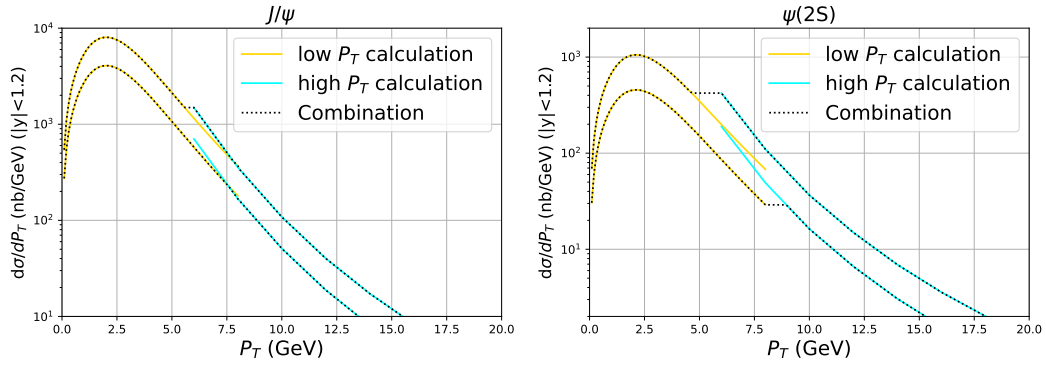


Figure 14: The direct charmonium spectra from theory. We show the theoretical calculations for high and low  $P_T$ , as well as our combinations.

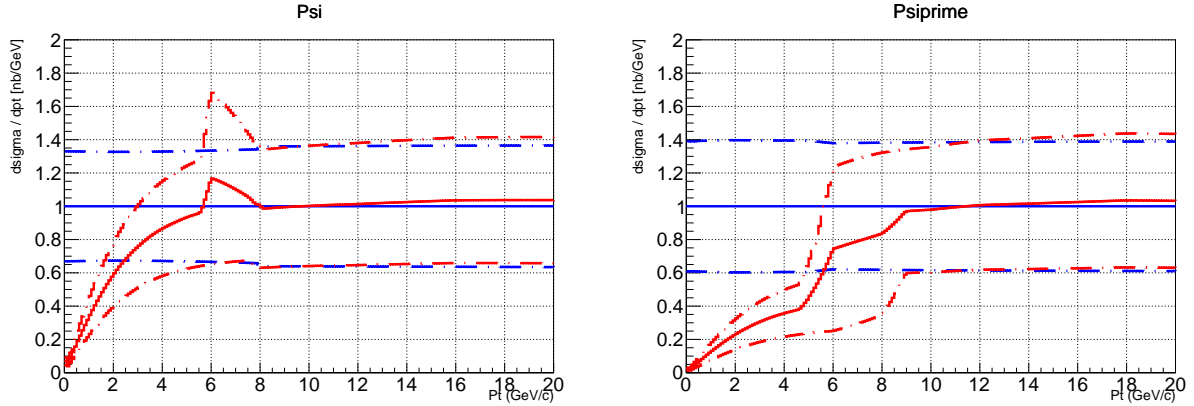


Figure 15: Run 2 vs. Run 3 comparisons of the direct  $J/\psi$  (left) and  $\psi'$  (right)  $P_T$  spectra. See Fig. 11 for an explanation of the colored bands.

#### A.4 $\pi^0$ , $\eta$ , $\eta'$ , $\phi$ , $\rho$ , and $\omega$ - Run 3

We have regenerated these from Pythia 8 and Pythia 6, as discussed in Section 5.4. In Fig. 16 we show a comparison between the Run 2 and the Run 3 spectra. We find that the Run 3 spectra are a few percent higher than the Run 2 spectra, except for the  $\phi$  spectrum. This spectrum is supposed to have been generated with Pythia 6, whereas all the other spectra are generated with Pythia 8. We have double-checked the current (Run 3) spectrum and we believe it to be correct. We do not understand why it differs so much from the previous (Run 2) spectrum. We suspect a bug in the Run 2 spectrum

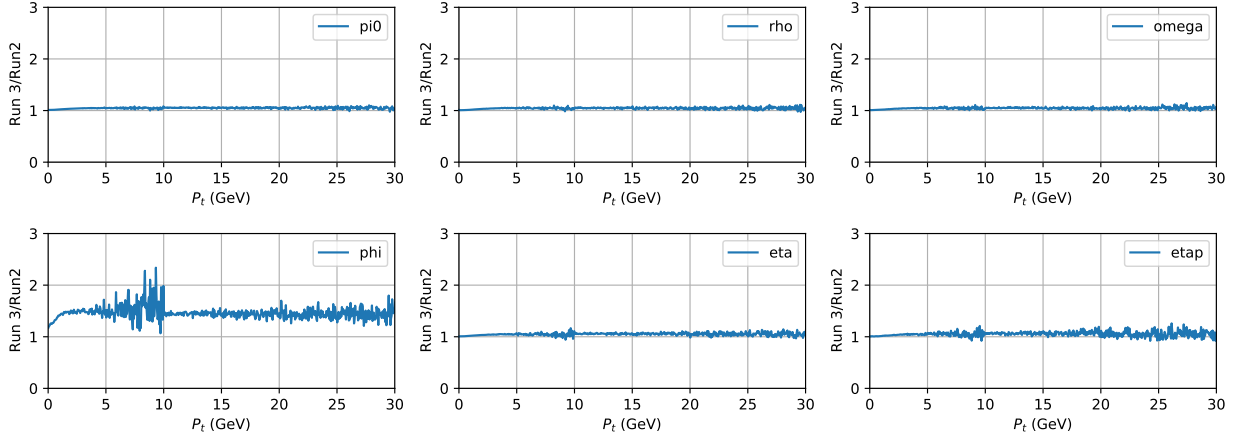


Figure 16: Run 2 vs. Run 3 ratio of the meson spectra. See text for detail and comments about the  $\phi$  spectrum.

## B Muons from the $pp$ interaction

These muons are used in the experiment as a control sample. Their generation was not described in the Run 2 version of the note. For completeness we describe it here.

### B.1 Muons from Heavy Flavor decays

These are  $b \rightarrow \mu$ ,  $b \rightarrow c \rightarrow \mu$ , and  $c \rightarrow \mu$ . As is the case for charmonium states from bottom decays, we obtain  $p_T$  distributions for these muons from a state of the art heavy flavor calculation[9, 10], see Fig. 17. The differences between the Run 2 and Run 3 distributions are minimal (Fig. 18).

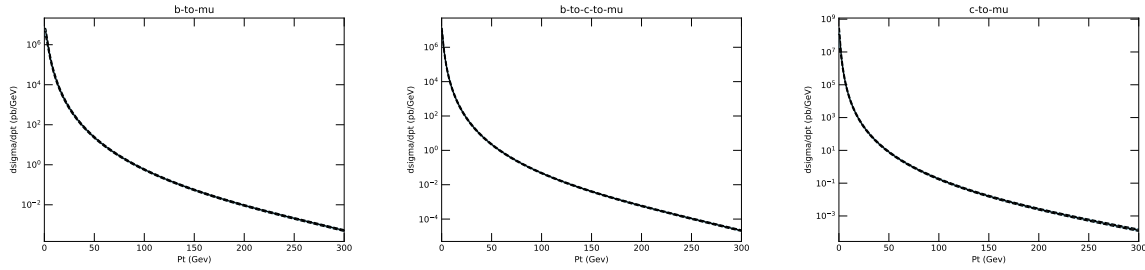


Figure 17: Run 3 transverse momentum distributions of muons from  $b \rightarrow \mu$  (left)  $b \rightarrow c \rightarrow \mu$  (center), and  $c \rightarrow \mu$ . Units are pb/GeV and the distributions are integrated over  $|\eta| < 1.0$ .

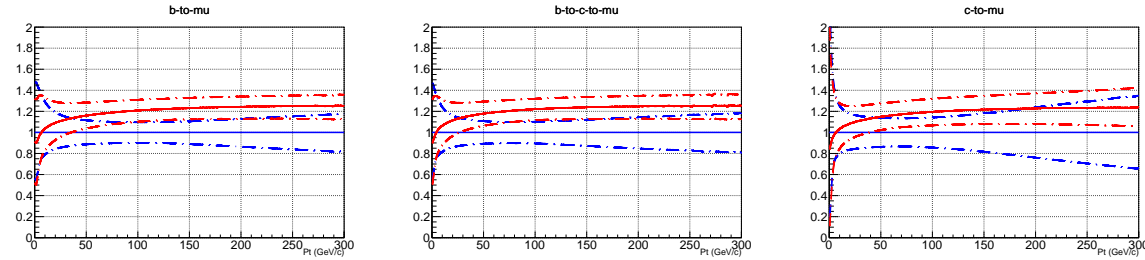


Figure 18: Run 2 vs. Run 3 comparisons of the  $P_T$  spectra for  $b \rightarrow \mu$  (left)  $b \rightarrow c \rightarrow \mu$  (center), and  $c \rightarrow \mu$  (right). The blue band is the Run2 variation as a function of  $P_T$  from the central value. (Dashed blue = up/down divided by central, Solid blue = central = 1.0). The red band is the ratio of the central Run 3 spectrum to the central Run2 spectrum. (Dashed red = up/down Run 3 divided by Run 2 central, Solid red = Run 3 central/Run 2 central) Note that the PDFs have changed between Run2 and Run3, see discussion in Section A.1.

### B.2 Muons from pions and kaons decays in flight

These are generated with Pythia 8 in the same process that collects the  $P_T$  spectra of mesons as described in Sections 5.4 and A.4. Decays are generated within a cylinder of radius 2 m and length  $\pm 4$  m. The radius of the front face of the CMS HB is  $r = 1.81$  m, and the front face of the CMS HE is at  $z = \pm 3.9$  m, so this choice of geometry is reasonable (also

accounting for the fact that pions and kaons will not interact as soon as they enter the HCAL).

The  $P_T$  spectrum of decayed muons is collected in a histograms of “muons per minimum bias event per 50 MeV” in  $|\eta| < 2$ . Before propagation, this spectrum is rescaled to  $|\eta| < 1$  through a multiplication by 1.13/2 where the factor of 1.13 was derived in Run2 from a MC study<sup>1</sup>. The decay-in-flight muon spectrum and the Run 3 to Run 2 comparison is shown in Fig. 19.

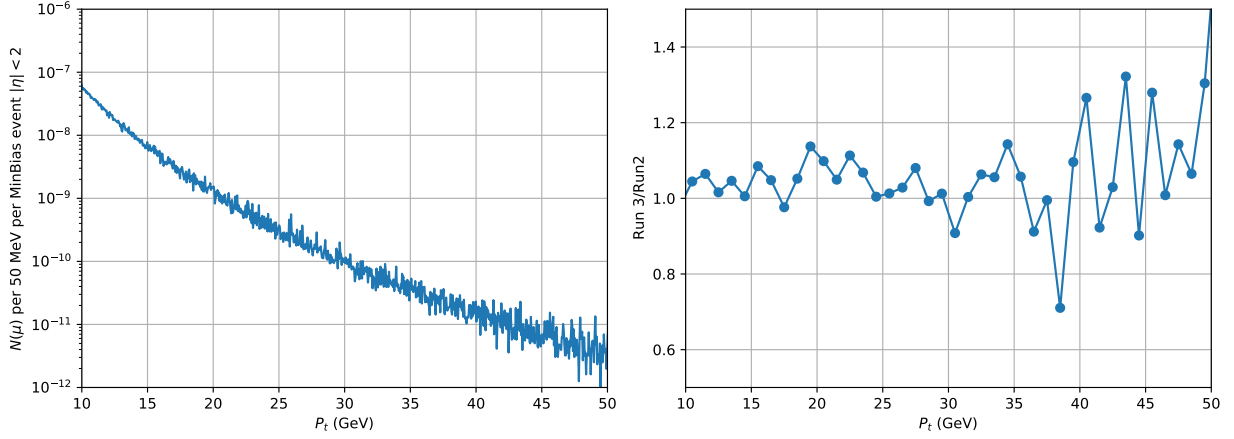


Figure 19: Run 3 transverse momentum spectrum of decay-in-flight muons (left), and Run 3 to Run 2 comparison.

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<sup>1</sup>Details of this study have been lost.

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