

Tau Mass Reconstruction. One More Status Report

Claudio (with lots of help from Hualin)

10 Feb 2022

H \rightarrow $\tau\tau$ kinematics

There are up to 6 unknowns

$$\vec{a} = (x_1, \varphi_1, m_{\nu\nu 1}, x_2, \varphi_2, m_{\nu\nu 2})$$

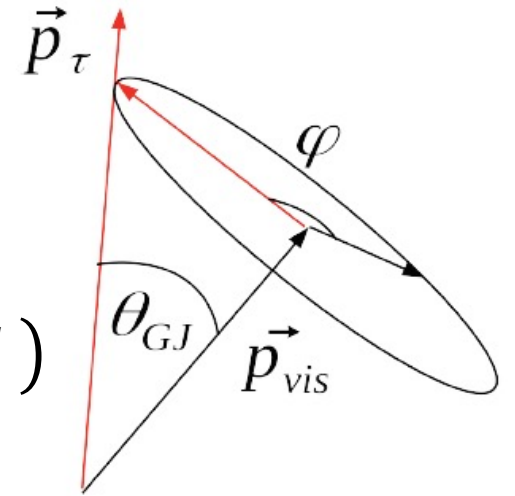
There is “event data”

$$\vec{y} = (E_{vis1}, \vec{p}_{vis1}, E_{vis2}, \vec{p}_{vis2}, \vec{E}_T^{miss})$$

$$\frac{E_{vis}}{E_\tau}$$

$\nu\nu$ inv. mass
(leptonic decays)

Note: θ_{GJ} can be calculated from
the other vars. (not immediately obvious)



SvFit has a likelihood, maximized as a function of m_{test}

$$\mathcal{L}(m_{test} | event\ data) = \int p(m_{\tau\tau} | \vec{y}, \vec{a}) \delta(m_{\tau\tau} - m_{test}) d\vec{a}$$

$$m_{\tau\tau}(\vec{y}, \vec{a})$$

$$p(m_{\tau\tau} | \vec{y}, \vec{a}) = ME \cdot TF \cdot REG$$

matrix
element

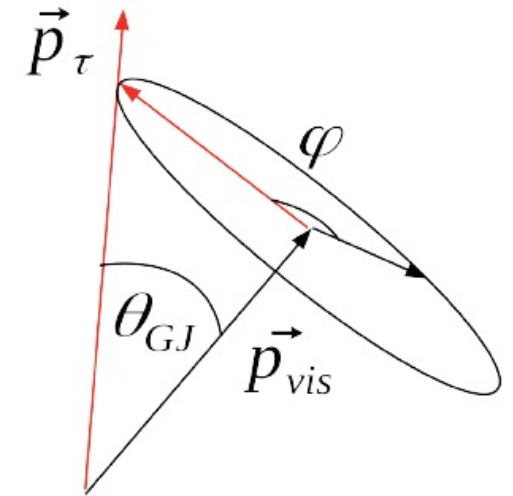
transfer
factor, ie,
detector effect

fudge
factor

Integral is complicated, done with MC methods, slow

FastMTT, simplification of SVFit (faster)

- Collinear approximation, $\theta_{GJ} = 0$
 - $\vec{a} = (x_1, \cancel{\phi_1}, m_{vv1}, x_2, \cancel{\phi_2}, m_{vv2})$
- Matrix elements = constant
- Only “Transfer Factor” is for E_T^{miss}



$$\mathcal{L}(m_{test} | event\ data) = MET_TF(\vec{E}_T^{miss\ reco}, \vec{E}_T^{miss\ hypo}) \cdot \int \delta(m_{\tau\tau} - m_{test}) d\vec{a} \quad \bullet\ REG$$

$$MET_TF = \mathcal{L}(\vec{E}_T^{miss\ reco} | \vec{E}_T^{miss\ hypo}) = \frac{1}{2\pi\sqrt{|V|}} \cdot \exp \left[-\frac{1}{2} \left((\vec{E}_T^{miss\ reco} - \vec{E}_T^{miss\ hypo})^T \cdot V^{-1} \cdot (\vec{E}_T^{miss\ reco} - \vec{E}_T^{miss\ hypo}) \right) \right]$$

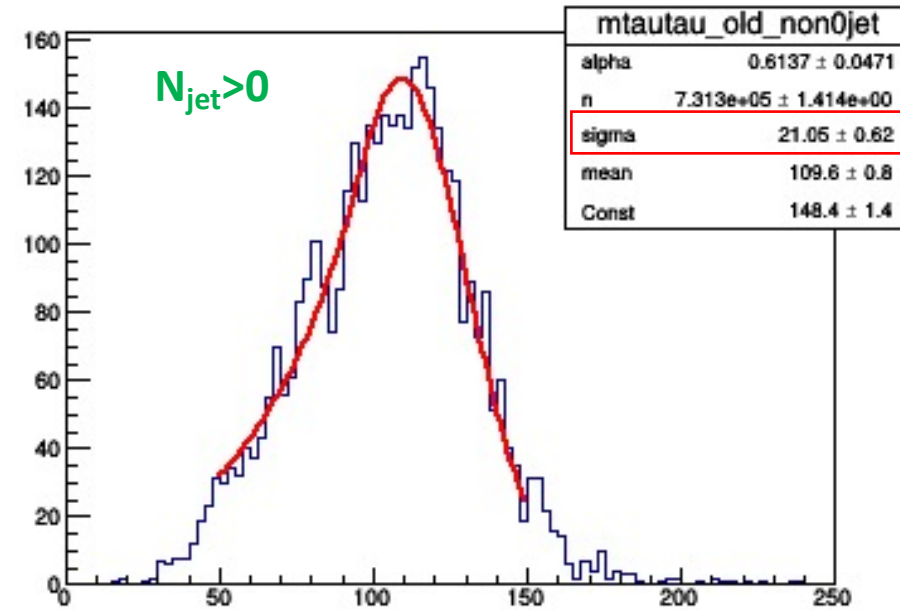
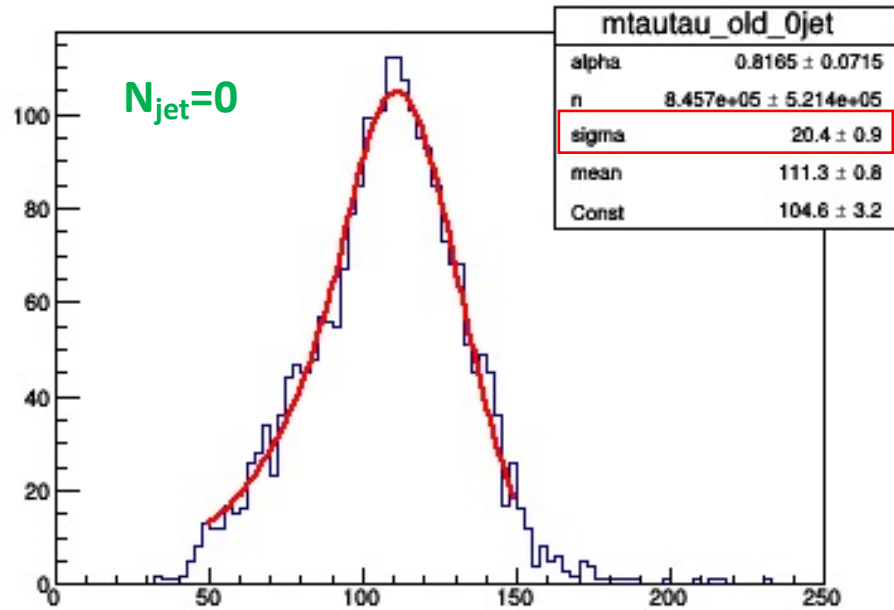
$$= \frac{1}{m_{test}^3}$$

(controls high tails)

depends on x_1 and x_2

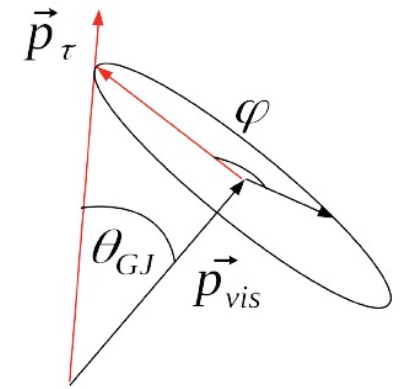
- The integral over the δ -function can be done analytically
 - There is a bit of fudging of the limits of integrations as well

FastMTT out of the box. Crystal Ball Fit

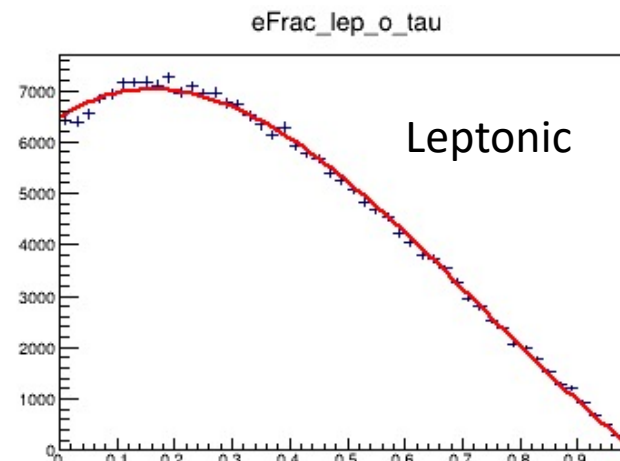
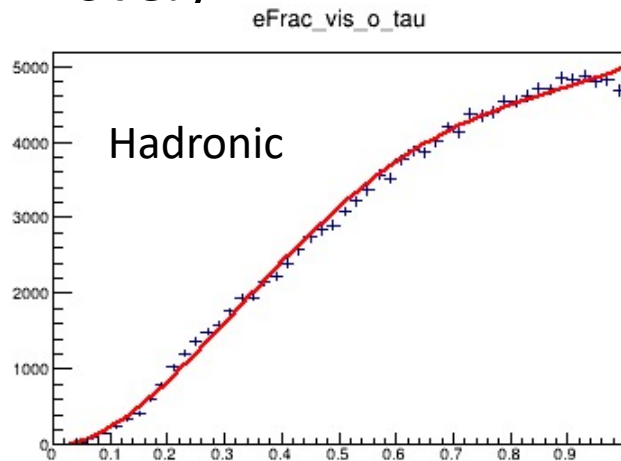


Should probably take the time to look at these for different final states!!!

Improvement of FastMTT (CC version)



- Super-Collinear approximation, $\theta_{GJ} = 0$, $\nu\nu$ also collinear
 - $\vec{a} = (x_1, \cancel{q_1}, \cancel{m_{\nu\nu 1}}, x_2, \cancel{q_2}, \cancel{m_{\nu\nu 2}})$
- Include expected x_i pdfs from MC (\sim matrix element)
 - In practice since the pdfs come from MC the lack of collinearity of the neutrinos in leptonic decays is actually included! (integrated over)



Fits to
“pol4”
(hadronic)

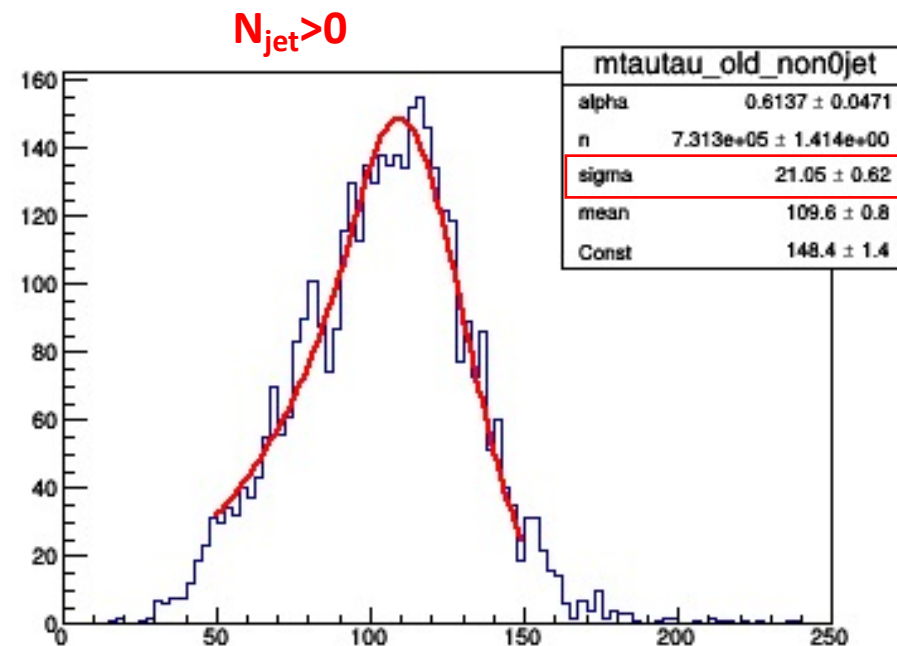
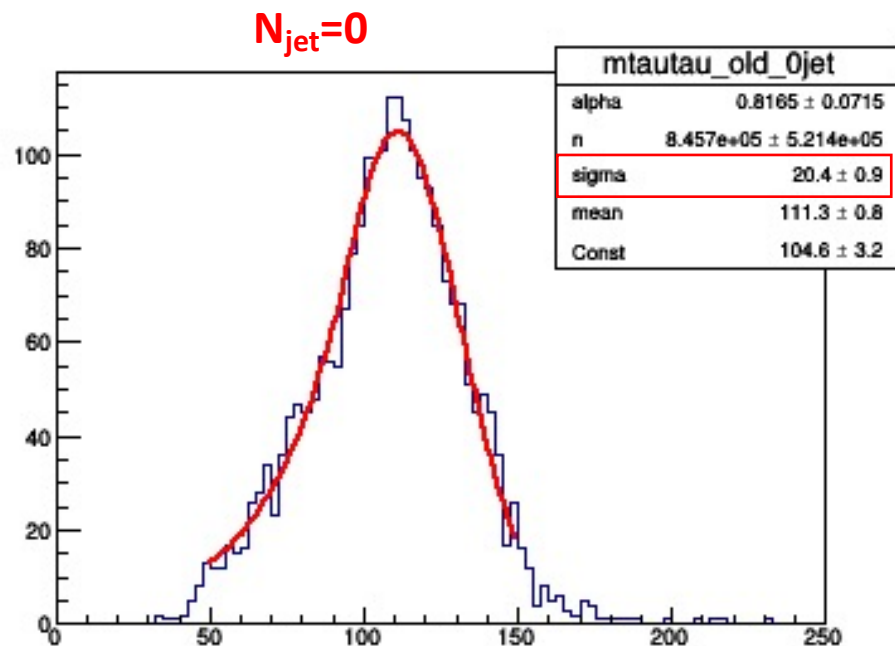
```
*****
Minimizer is Linear / Migrad
Chi2      =      201.181
NDf       =         44
p0        =      -20.8584  +/-   5.60061
p1        =      -103.46   +/-   181.624
p2        =      28912.2   +/-  1039.63
p3        =      -40754    +/-   1916.33
p4        =      16937.9   +/-   1087.52
```

“pol3”
(leptonic)

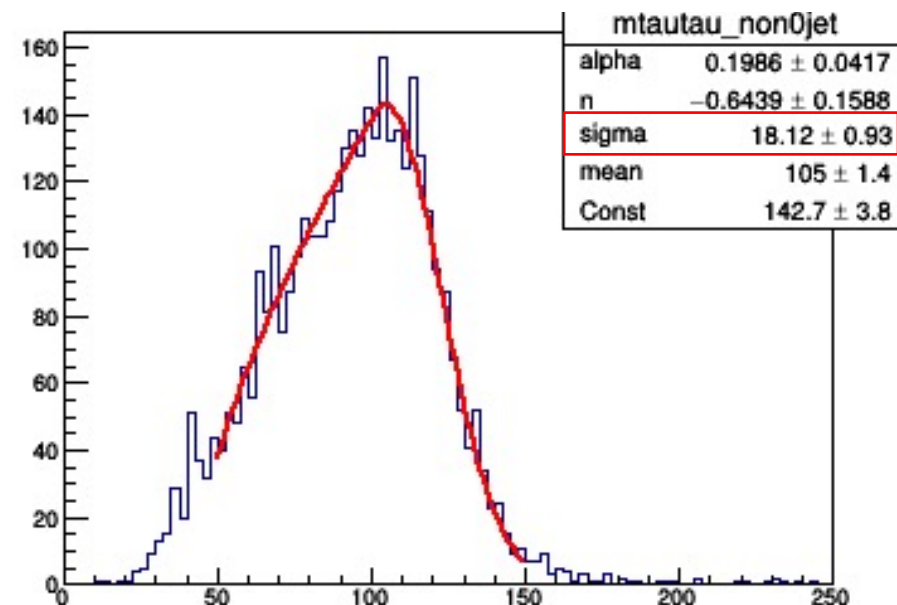
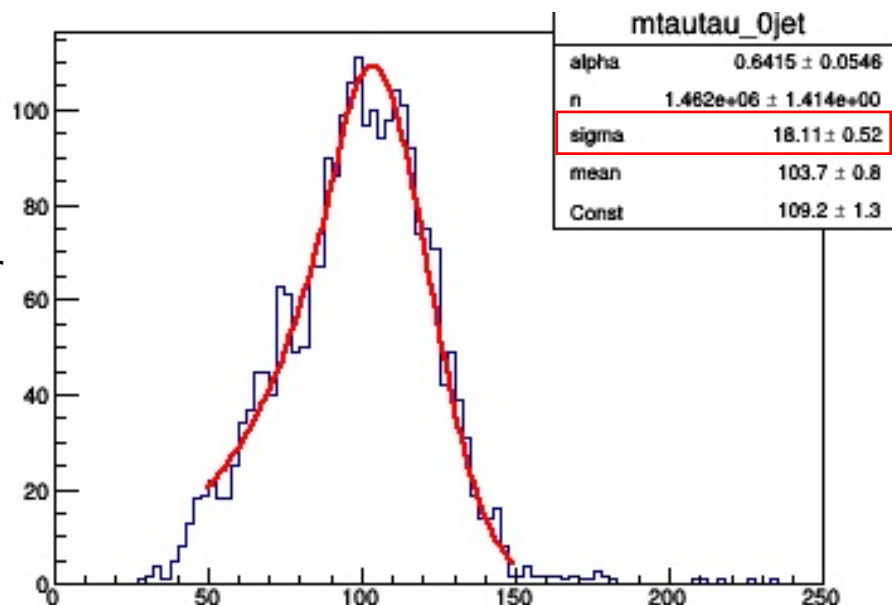
```
*****
Minimizer is Linear / Migrad
Chi2      =      118.649
NDf       =         46
p0        =      6462.7    +/-   42.8792
p1        =      7119.87   +/-   319.185
p2        =     -24764.4    +/-   642.921
p3        =      11189.4   +/-   371.598
```

- With polynomial functions the integral with the δ -function can be done analytically (still very fast!)
 - $\int p_1(x_1)p_2(x_2)\delta(m_{\tau\tau} - m_{test})dx_1dx_2$

FastMTT out
of the box



FastMTT with
matrix element

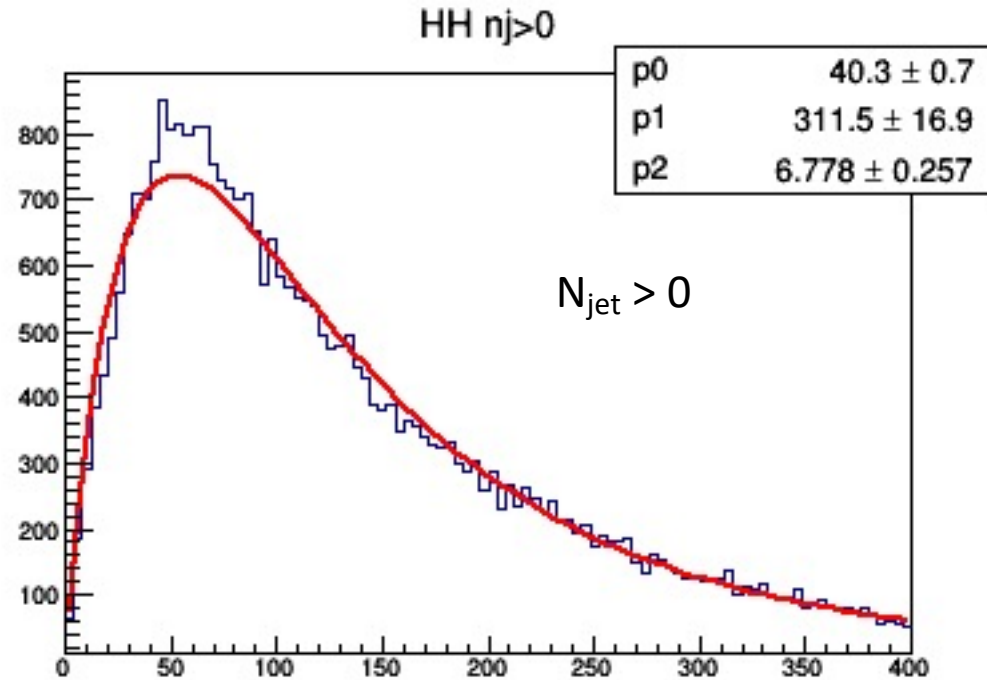
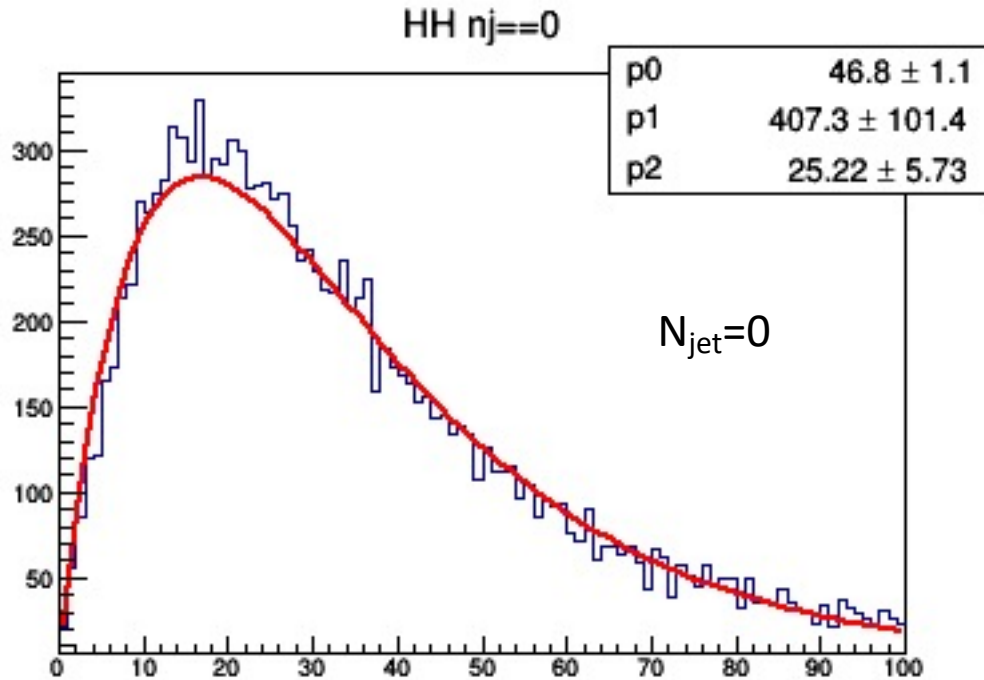


Same fudge factor
(not tuned)

~ 10% narrower

Would like to incorporate $P_T(HH)$ information

GEN expectations:



```
Double_t func1(Double_t *x, Double_t *par){  
    Double_t fitval;  
    Double_t pt = x[0];  
    Double_t p0 = par[0];  
    Double_t p1 = par[1];  
    Double_t p2 = par[2];  
    fitval = pt*p0*pow((1+pt/p1), -p2);  
    return fitval;  
}
```


Not so easy to do in the SvFit/FastMTT approach

$$\mathcal{L}(m_{test}, p_{T,test} | event\ data) = MET_TF(\vec{E}_T^{miss\ reco}, \vec{E}_T^{miss\ hypo}) \cdot \int \delta(m_{\tau\tau} - m_{test}) d\vec{a} \cdot \int \delta(p_{Tx,\tau\tau} - p_{Tx,test}) d\vec{a} \cdot \int \delta(p_{Ty,\tau\tau} - p_{Ty,test}) d\vec{a} \quad (8)$$

- The δ functions are nasty unless we remove the matrix element completely
- (There may be a way...need to talk to Hualin as one needs to run the actual FastMTT code, not my kludges)

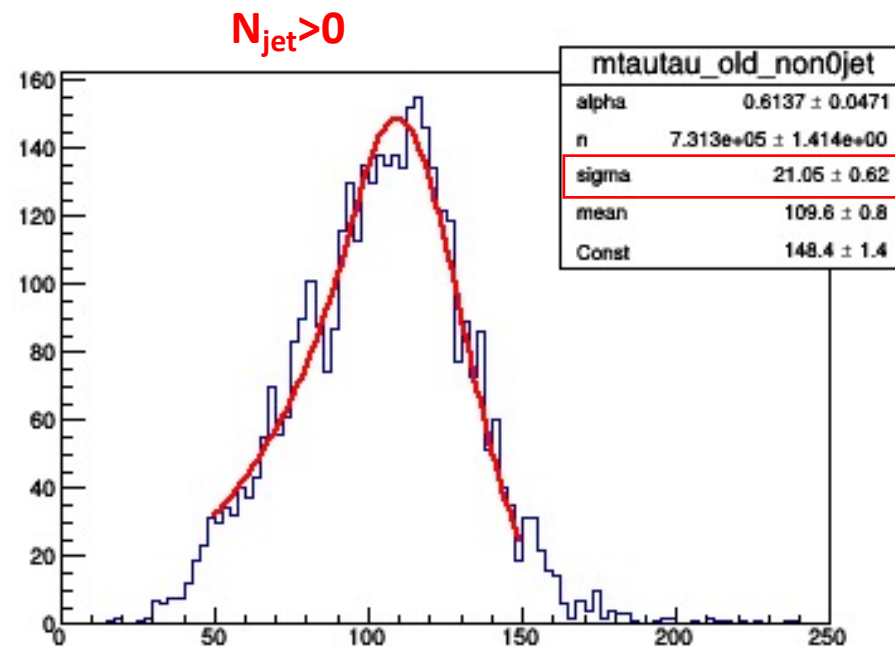
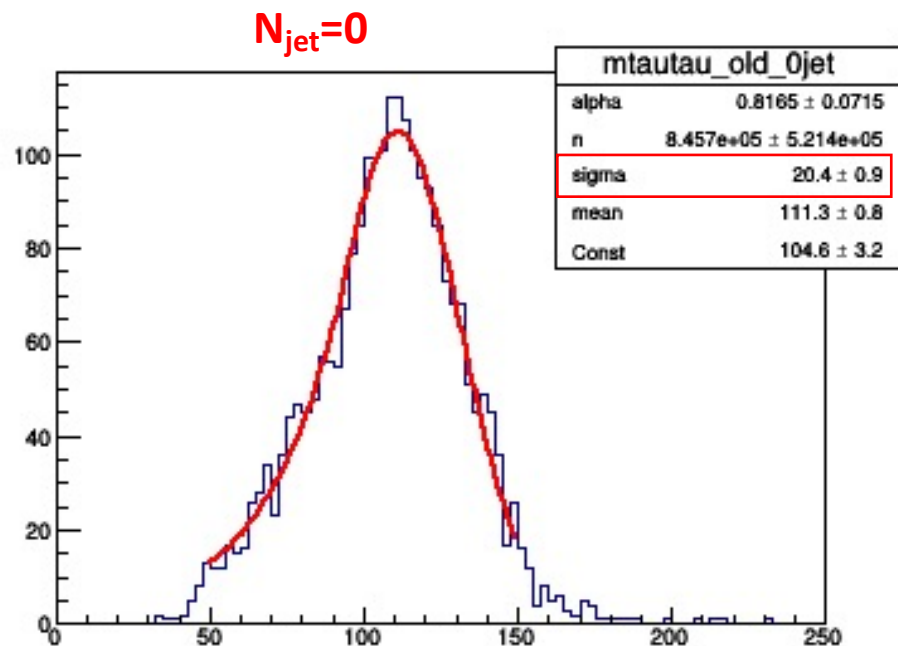
Instead, try a different approach, similar to Atlas

$$\mathcal{L} = MET_TF \cdot p(x_1) \cdot p(x_2) \cdot p(P_T) \cdot REG$$

Since $P_T = P_T(x_1, x_2)$ this is probably not entirely kosher. Not sure. But neither is *REG*...

- First, check this out without the P_T information

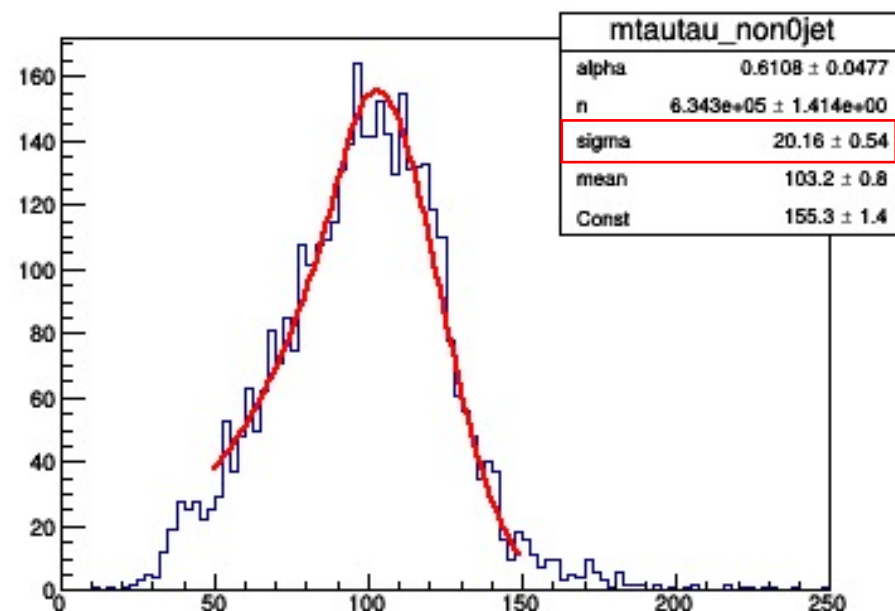
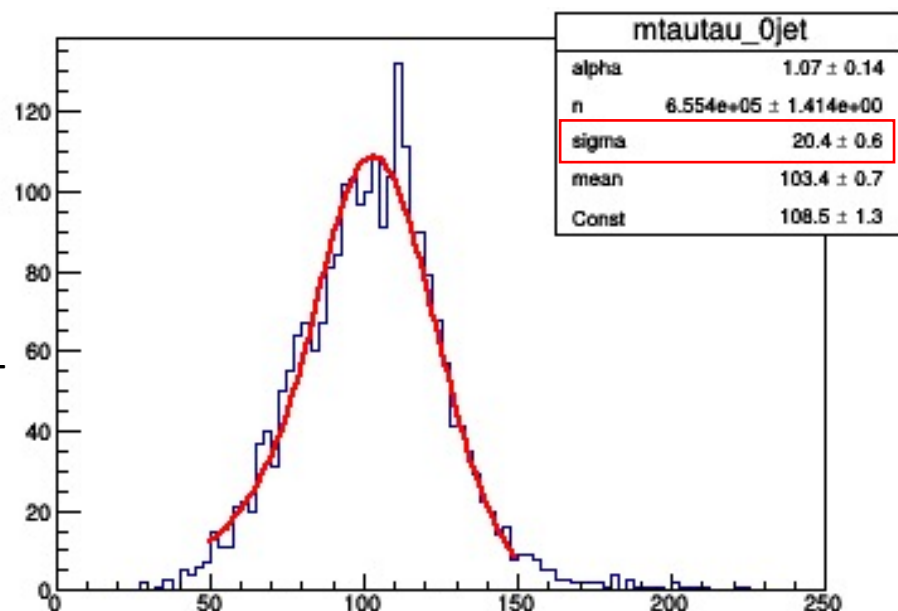
FastMTT out
of the box



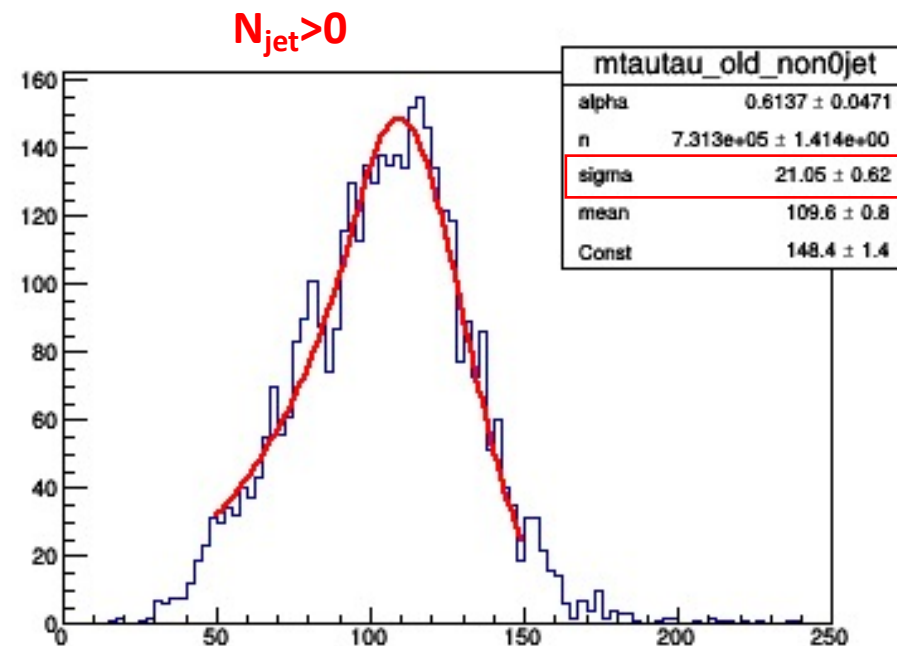
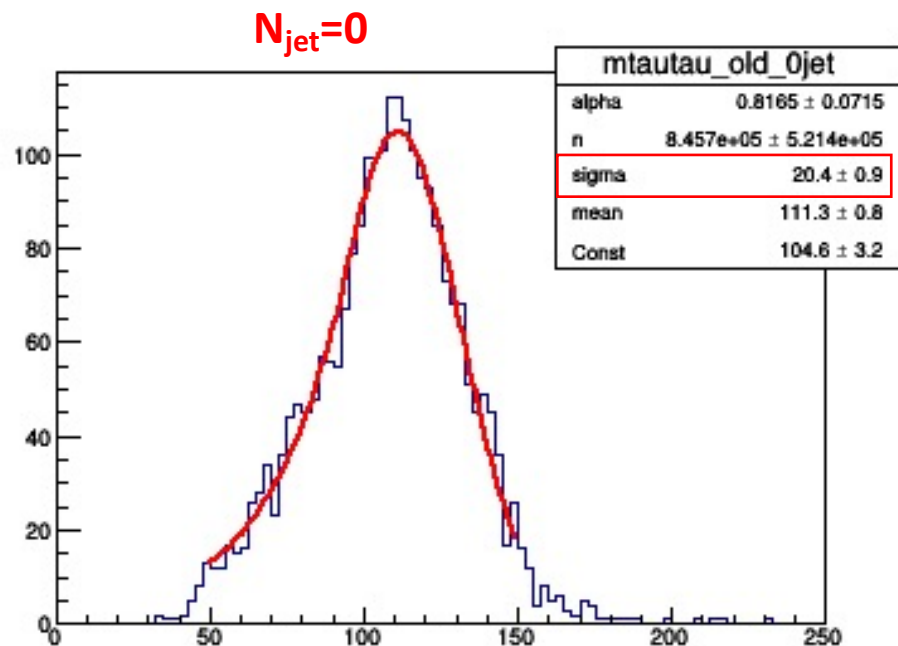
Atlas-like, no P_T

Same fudge factor
(not tuned)

~ same as FastMTT



FastMTT out
of the box

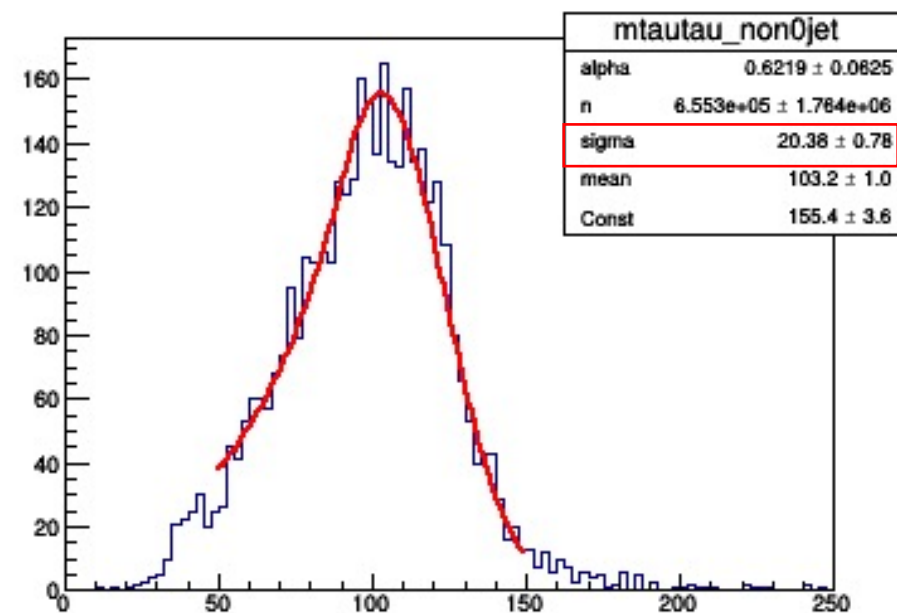
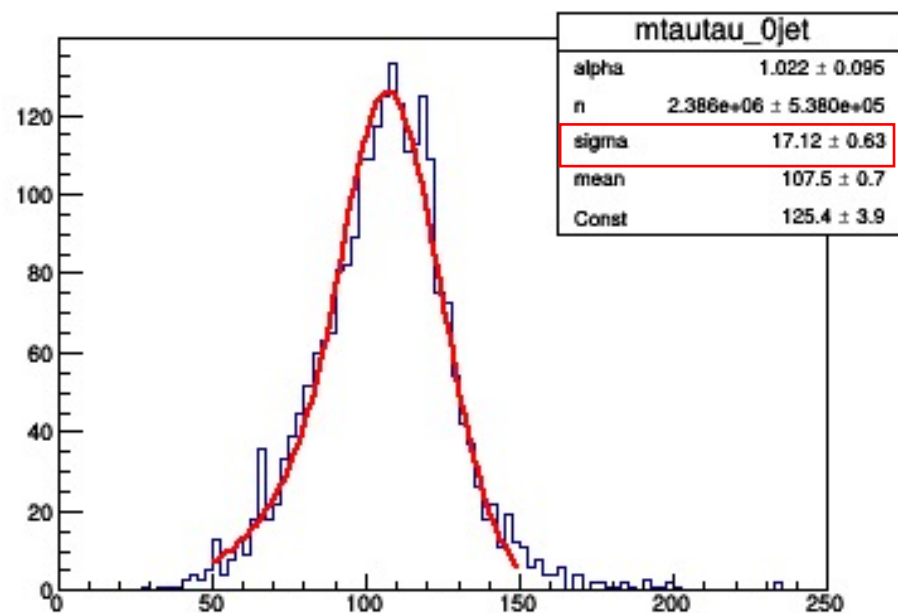


Atlas-like, with P_T

Same fudge factor
(not tuned)

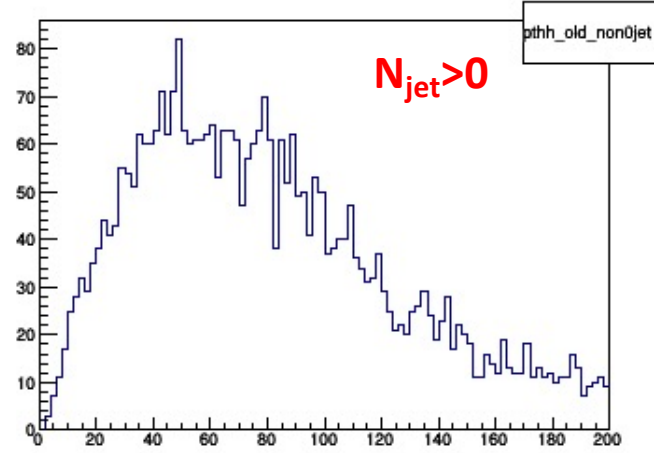
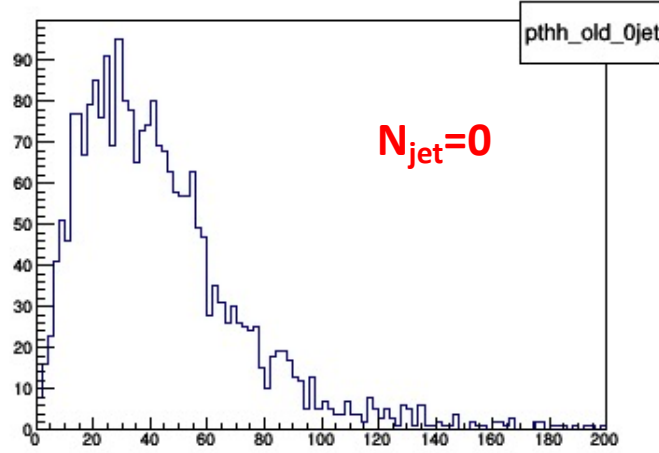
0 jets:
~ 15% narrower

> 0 jet
~ 3% narrower
(maybe)

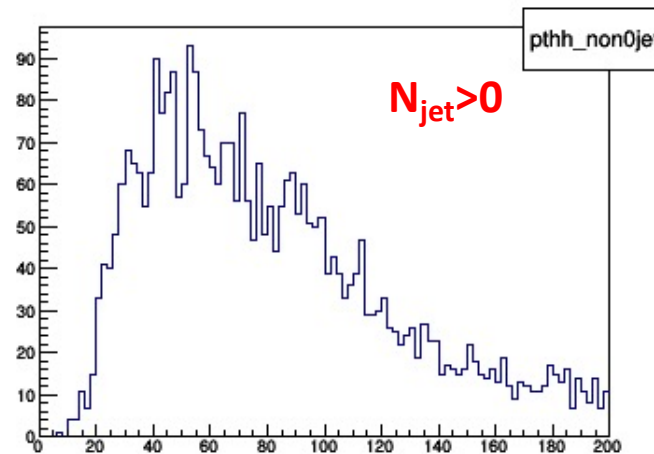
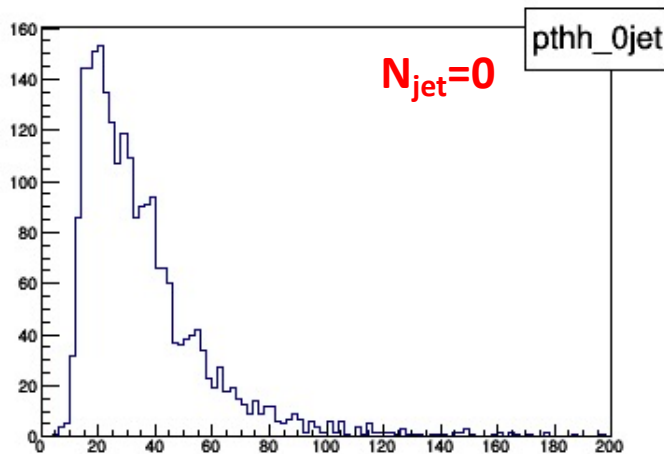


What happened to P_T ?

FastMTT out of the box



Atlas-like, with P_T



There is a depletion of events near zero in the Atlas-like (with P_T) algorithm.

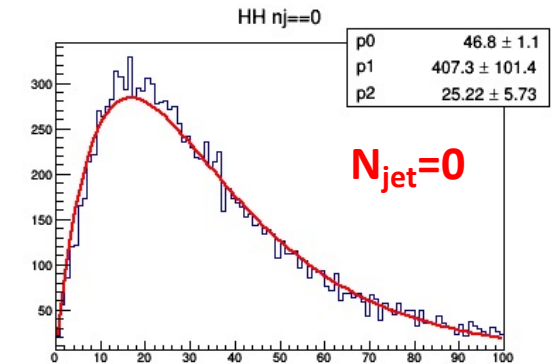
This is not unexpected.

There are large uncertainties.

The P_T pdfs go to zero at $P_T = 0$.

Because of the large uncertainties the penalty to move any given event away from zero is not so large.

Therefore, events are moved away from 0



Summary (1)

- It is possible to improve FastMTT a little bit by adding the matrix element.
- An Atlas-like reconstruction does just as well as FastMTT
- In the Atlas-like reconstruction it is “easy” to include the $P_T(\text{HH})$ information
- The additional information improves resolution in the $N_{\text{jet}} = 0$ sample by about 15%, but hardly at all in the $N_{\text{jet}} > 0$ case.
 - It may be improved by fine-tuning the fudge factor (eg: different fudge factors for different tau decays?)
- The $P_T(\text{HH})$ distribution becomes a little weird, but actually I think it makes sense.
- The inclusion of the $P_T(\text{HH})$ is probably not kosher
- I think there may be a way to put the $P_T(\text{HH})$ information in the FastMTT framework. Not sure
- Franny’s study seems to indicate that none of this matters much if at all 😂

Summary (2)

- The “FastMTT+ME” the “Atlas-like” algos are coded by me inside a custom looper and should be repackaged in case we want to use them more widely.
- Think more about the slightly iffy use of the P_T in the Atlas-like algorithm
- Talk to Hualin about one more attempt to include the P_T pdf in the standard FastMTT (ie: without ME) in a sensible way
- I wonder if the average of the the two algorithms leads to better resolution?
- There is some algebra involved in “FastMTT+ME”, needs to be double-checked
 - See backup

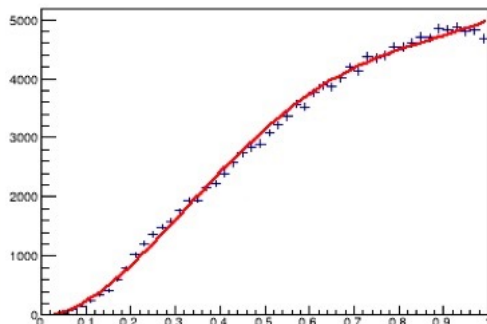
Backup

Algebra for FastMTT + ME

TAU DECAY PDFS

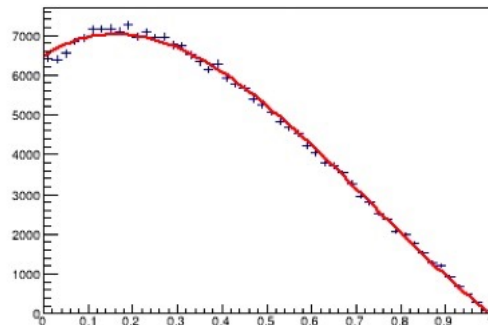
HADRONIC TAU

eFrac_vis_o_tau



LEPTONIC TAU

eFrac_lep_o_tau



Fit to "POL4"

```
*****
Minimizer is Linear / Migrad
Chi2      = 201.181
Ndf       = 44
p0        = -20.8584 +/- 5.60061
p1        = -103.46 +/- 181.624
p2        = 28912.2 +/- 1039.63
p3        = -40754 +/- 1916.33
p4        = 16937.9 +/- 1087.52
```

Fit to "POL3"

```
*****
Minimizer is Linear / Migrad
Chi2      = 118.649
Ndf       = 46
p0        = 6462.7 +/- 42.8792
p1        = 7119.87 +/- 319.185
p2        = -24764.4 +/- 642.921
p3        = 11189.4 +/- 371.598
```

$$f(x_i) = \sum_{j=0}^4 a_j x_i^j \quad \text{with } a_4 = 0 \text{ for leptonic}$$

We reuse some of the results from Appendix A of the FastMTT analysis note

$$m_{\tau\tau} \simeq \frac{m_{vis}}{\sqrt{x_1 x_2}}$$

$$\frac{2m_{vis}^2}{m_{test}^3} \frac{1}{x_2}, \quad x_1 x_2 = \left(\frac{m_{vis}}{m_{test}}\right)^2$$

$$x_{min} = \max\left(x_{2,min}, \left(\frac{m_{vis}}{m_{test}}\right)^2\right)$$

$$x_{max} = \min\left(1, \left(\frac{m_{vis}}{m_{test}}\right)^2 \frac{1}{x_{1,min}}\right)$$

$$V_{q\bar{q}}(m_{test}, m_{vis}) = \int_{x_{1,min}}^1 dx_1 \int_{x_{2,min}}^1 dx_2 \int_0^{2\pi} d\phi_1 \int_0^{2\pi} d\phi_2 \delta(m_{test} - \frac{m_{vis}}{\sqrt{x_1 x_2}}) =$$

$$4\pi^2 \int_{x_{1,min}}^1 dx_1 \int_{x_{2,min}}^1 dx_2 \delta(m_{test} - \frac{m_{vis}}{\sqrt{x_1 x_2}}) =$$

$$4\pi^2 \frac{2m_{vis}^2}{m_{test}^3} \int_{x_{min}}^{x_{max}} \frac{1}{x_2} dx_2 = 4\pi^2 \frac{2m_{vis}^2}{m_{test}^3} \log\left(\frac{x_{max}}{x_{min}}\right)$$

← This is for the case that $f(x_1)f(x_2)$ we ignore

$$x_{1,min} = \left(\frac{m_1}{m_\tau}\right)^2, \quad x_{2,min} = \left(\frac{m_2}{m_\tau}\right)^2,$$

m_1 and m_2 are the visible mass of the τ

The formula above works for hadronic taus, which have $m_{\nu\nu} = 0$, which is exactly what we care about

$$\int f_1(x_1) f_2(x_2) \delta(m_{test} - \frac{m_{vis}}{\sqrt{x_1 x_2}}) dx_1 dx_2$$

$$= \frac{2m_{vis}^2}{m_{test}^3} \int_{x_{min}}^{x_{max}} \frac{1}{x_2} f_1\left(\frac{m_{vis}^2}{m_{test}^2} \frac{1}{x_2}\right) f_2(x_2) dx_2$$

$$\text{Let } \alpha = \left(\frac{m_{vis}}{m_{test}} \right)^2$$

$$\text{Also } f_1(x) = \sum_{j=0}^4 a_{1j} X^j \quad f_2(x) = \sum_{k=0}^4 a_{2k} X^k$$

$$= \frac{2\alpha}{m_{test}} \int_{x_{min}}^{x_{max}} \sum_{j=0}^4 \sum_{k=0}^4 \frac{a_{1j}}{x_2} \alpha^j X_2^{-j} a_{2k} X_2^k dx_2$$

$$= \frac{2\alpha}{m_{test}} \int_{x_{min}}^{x_{max}} \sum_{j=0}^4 a_{1j} \alpha^j \sum_{k=0}^4 a_{2k} X_2^{k-j-1} dx_2$$

$$= \frac{2\alpha}{m_{test}} \sum_j a_{1j} \alpha^j \sum_k a_{2k} \left[\delta_{kj} \log \frac{x_{max}}{x_{min}} + \frac{1 - \delta_{kj}}{k-j} \left(x_{max}^{k-j} - x_{min}^{k-j} \right) \right]$$

(the factor of 2 is irrelevant for the purpose of maximizing the likelihood)