Finite Volume Methods for Hyperbolic Problems

Linearization of Nonlinear Systems

- General form, Jacobian matrix
- Scalar Burgers' equation
- Compressible gas dynamics
- Linear acoustics equations

Linearization

General nonlinear conservation law: $q_t + f(q)_x = 0$

Suppose $q(x,t)=q_0+\tilde{q}(x,t)$ where $\|\tilde{q}(x,t)\|=\epsilon$ is small.

Linearization

General nonlinear conservation law: $q_t + f(q)_x = 0$

Suppose $q(x,t)=q_0+\tilde{q}(x,t)$ where $\|\tilde{q}(x,t)\|=\epsilon$ is small.

Then

$$\tilde{q}_t = q_t$$

$$= -f(q)_x$$

$$= -f'(q)q_x$$

$$= -f'(q_0 + \tilde{q})\tilde{q}_x$$

$$= -f'(q_0)\tilde{q}_x + \mathcal{O}(\epsilon^2).$$

Linearization

General nonlinear conservation law: $q_t + f(q)_x = 0$

Suppose $q(x,t) = q_0 + \tilde{q}(x,t)$ where $\|\tilde{q}(x,t)\| = \epsilon$ is small.

Then

$$\tilde{q}_t = q_t
= -f(q)_x
= -f'(q)q_x
= -f'(q_0 + \tilde{q})\tilde{q}_x
= -f'(q_0)\tilde{q}_x + \mathcal{O}(\epsilon^2).$$

Linearization: $\tilde{q}_t + A\tilde{q}_x = 0$ where $A = f'(q_0) = \text{Jacobian matrix}$

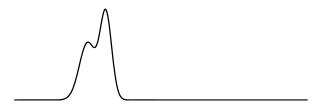
Scalar: Advection equation

Conservation form:
$$u_t + \left(\frac{1}{2}u^2\right)_x = 0$$
, $f(u) = \frac{1}{2}u^2$.

Quasi-linear form: $u_t + uu_x = 0$.

This looks like an advection equation with \boldsymbol{u} advected with speed \boldsymbol{u} .

True solution: u is constant along characteristic with speed u until the wave "breaks".



Conservation form:
$$u_t + \left(\frac{1}{2}u^2\right)_x = 0$$
, $f(u) = \frac{1}{2}u^2$.

Quasi-linear form:
$$u_t + uu_x = 0$$
.

This looks like an advection equation with \boldsymbol{u} advected with speed \boldsymbol{u} .

True solution: u is constant along characteristic with speed u until the wave "breaks".



Conservation form:
$$u_t + \left(\frac{1}{2}u^2\right)_x = 0$$
, $f(u) = \frac{1}{2}u^2$.

Quasi-linear form: $u_t + uu_x = 0$.

This looks like an advection equation with \boldsymbol{u} advected with speed \boldsymbol{u} .

True solution: u is constant along characteristic with speed u until the wave "breaks".



After breaking, the weak solution contains a shock wave.

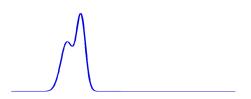
Conservation form: $u_t + \left(\frac{1}{2}u^2\right)_x = 0, \qquad f(u) = \frac{1}{2}u^2.$

Linearization about u_0 :

$$f(u) = \frac{1}{2}u^2 \implies f'(u_0) = u_0$$

So if $u(x,0)=u_0+\tilde{u}(x,0)$ with $\|\tilde{u}\|$ small, then $\tilde{u}(x,t)$ approximately satisfies advection equation

$$\tilde{u}_t + u_0 u_x = 0.$$



Conservation form: $u_t + \left(\frac{1}{2}u^2\right)_x = 0, \qquad f(u) = \frac{1}{2}u^2.$

Linearization about u_0 :

$$f(u) = \frac{1}{2}u^2 \implies f'(u_0) = u_0$$

So if $u(x,0)=u_0+\tilde{u}(x,0)$ with $\|\tilde{u}\|$ small, then $\tilde{u}(x,t)$ approximately satisfies advection equation

$$\tilde{u}_t + u_0 u_x = 0.$$



Conservation form: $u_t + \left(\frac{1}{2}u^2\right)_x = 0, \qquad f(u) = \frac{1}{2}u^2.$

Linearization about u_0 :

$$f(u) = \frac{1}{2}u^2 \implies f'(u_0) = u_0$$

So if $u(x,0)=u_0+\tilde{u}(x,0)$ with $\|\tilde{u}\|$ small, then $\tilde{u}(x,t)$ approximately satisfies advection equation

$$\tilde{u}_t + u_0 u_x = 0.$$



Compressible gas dynamics (simple case)

In one space dimension (e.g. in a pipe).

$$\begin{array}{ll} \rho(x,t)=\text{density,} & u(x,t)=\text{velocity,} \\ p(x,t)=\text{pressure,} & \rho(x,t)u(x,t)=\text{momentum.} \end{array}$$

Conservation of:

$$\begin{array}{lll} \text{mass:} & \rho & \text{flux:} & \rho u \\ \text{momentum:} & \rho u & \text{flux:} & (\rho u)u + p \end{array}$$

Conservation laws:

$$\rho_t + (\rho u)_x = 0$$
$$(\rho u)_t + (\rho u^2 + p)_x = 0$$

Equation of state:

$$p = P(\rho)$$
.

Suppose $\rho(x,t) \approx \rho_0$ and $u(x,t) \approx u_0$.

Model small perturbations to this steady state (sound waves).

$$\begin{bmatrix} \rho(x,t) \\ (\rho u)(x,t) \end{bmatrix} = \begin{bmatrix} \rho_0 \\ \rho_0 u_0 \end{bmatrix} + \begin{bmatrix} \widetilde{\rho}(x,t) \\ (\widetilde{\rho}u)(x,t) \end{bmatrix}$$

or $q(x,t)=q_0+\tilde{q}(x,t)$ where $\|\tilde{q}(x,t)\|=\epsilon$ is small.

Suppose $\rho(x,t) \approx \rho_0$ and $u(x,t) \approx u_0$.

Model small perturbations to this steady state (sound waves).

$$\left[\begin{array}{c} \rho(x,t) \\ (\rho u)(x,t) \end{array}\right] = \left[\begin{array}{c} \rho_0 \\ \rho_0 u_0 \end{array}\right] + \left[\begin{array}{c} \widetilde{\rho}(x,t) \\ (\widetilde{\rho u})(x,t) \end{array}\right]$$

or $q(x,t)=q_0+\tilde{q}(x,t)$ where $\|\tilde{q}(x,t)\|=\epsilon$ is small.

Then nonlinear equation $q_t + f(q)_x = 0$ leads to

$$\tilde{q}_t = q_t$$

$$= -f(q)_x$$

$$= -f'(q)q_x$$

$$= -f'(q_0 + \tilde{q})\tilde{q}_x$$

$$= -f'(q_0)\tilde{q}_x + \mathcal{O}(\epsilon^2).$$

Suppose $\rho(x,t) \approx \rho_0$ and $u(x,t) \approx u_0$.

Model small perturbations to this steady state (sound waves).

$$\left[\begin{array}{c} \rho(x,t) \\ (\rho u)(x,t) \end{array}\right] = \left[\begin{array}{c} \rho_0 \\ \rho_0 u_0 \end{array}\right] + \left[\begin{array}{c} \widetilde{\rho}(x,t) \\ (\widetilde{\rho u})(x,t) \end{array}\right]$$

or $q(x,t)=q_0+\tilde{q}(x,t)$ where $\|\tilde{q}(x,t)\|=\epsilon$ is small.

Then nonlinear equation $q_t + f(q)_x = 0$ leads to

$$\tilde{q}_t = q_t
= -f(q)_x
= -f'(q)q_x
= -f'(q_0 + \tilde{q})\tilde{q}_x
= -f'(q_0)\tilde{q}_x + \mathcal{O}(\epsilon^2).$$

Linearization: $\tilde{q}_t + A\tilde{q}_x = 0$ where $A = f'(q_0)$.

$$\begin{split} \rho_t + (\rho u)_x &= 0 \\ (\rho u)_t + (\rho u^2 + P(\rho))_x &= 0 \end{split}$$
 so
$$q = \left[\begin{array}{c} \rho \\ \rho u \end{array} \right] = \left[\begin{array}{c} q^1 \\ q^2 \end{array} \right],$$

$$f(q) = \left[\begin{array}{c} \rho u \\ \rho u^2 + P(\rho) \end{array} \right] = \left[\begin{array}{c} f^1(q) \\ f^2(q) \end{array} \right] = \left[\begin{array}{c} q^2 \\ (q^2)^2/q^1 + P(q^1) \end{array} \right].$$

$$\rho_t + (\rho u)_x = 0$$
$$(\rho u)_t + (\rho u^2 + P(\rho))_x = 0$$

so

$$q = \left[\begin{array}{c} \rho \\ \rho u \end{array} \right] = \left[\begin{array}{c} q^1 \\ q^2 \end{array} \right],$$

$$f(q) = \begin{bmatrix} \rho u \\ \rho u^2 + P(\rho) \end{bmatrix} = \begin{bmatrix} f^1(q) \\ f^2(q) \end{bmatrix} = \begin{bmatrix} q^2 \\ (q^2)^2/q^1 + P(q^1) \end{bmatrix}.$$

Jacobian:

$$f'(q) = \begin{bmatrix} \partial f^1/\partial q^1 & \partial f^1/\partial q^2 \\ \partial f^2/\partial q^1 & \partial f^2/\partial q^2 \end{bmatrix}.$$

$$f'(q_0) = \begin{bmatrix} 0 & 1 \\ -u_0^2 + P'(\rho_0) & 2u_0 \end{bmatrix}.$$

Linearization: $\tilde{q}_t + A\tilde{q}_x = 0$ where $A = f'(q_0)$.

$$A = f'(q_0) = \begin{bmatrix} 0 & 1 \\ -u_0^2 + P'(\rho_0) & 2u_0 \end{bmatrix}.$$

This can be written out as (2.47):

$$\widetilde{\rho}_t + (\widetilde{\rho u})_x = 0$$
$$(\widetilde{\rho u})_t + (-u_0^2 + P'(\rho_0))\widetilde{\rho}_x + 2u_0(\widetilde{\rho u})_x = 0.$$

Linearization: $\tilde{q}_t + A\tilde{q}_x = 0$ where $A = f'(q_0)$.

$$A = f'(q_0) = \begin{bmatrix} 0 & 1 \\ -u_0^2 + P'(\rho_0) & 2u_0 \end{bmatrix}.$$

This can be written out as (2.47):

$$\widetilde{\rho}_t + (\widetilde{\rho u})_x = 0$$
$$(\widetilde{\rho u})_t + (-u_0^2 + P'(\rho_0))\widetilde{\rho}_x + 2u_0(\widetilde{\rho u})_x = 0.$$

Rewrite in terms of p and u perturbations (Exer. 2.1):

$$\tilde{p}_t + u_0 \tilde{p}_x + K_0 \tilde{u}_x = 0,$$

$$\rho_0 \tilde{u}_t + \tilde{p}_x + \rho_0 u_0 \tilde{u}_x = 0,$$

where $K_0 = \rho_0 P'(\rho_0)$ is the bulk modulus.

$$\tilde{p}_t + u_0 \tilde{p}_x + K_0 \tilde{u}_x = 0,$$

$$\rho_0 \tilde{u}_t + \tilde{p}_x + \rho_0 u_0 \tilde{u}_x = 0,$$

gives the system $q_t + Aq_x = 0$ (Drop tildes)

$$q(x,t) = \begin{bmatrix} p(x,t) \\ u(x,t) \end{bmatrix}, \qquad A = \begin{bmatrix} u_0 & K_0 \\ 1/\rho_0 & u_0 \end{bmatrix}$$

$$\tilde{p}_t + u_0 \tilde{p}_x + K_0 \tilde{u}_x = 0,$$

$$\rho_0 \tilde{u}_t + \tilde{p}_x + \rho_0 u_0 \tilde{u}_x = 0,$$

gives the system $q_t + Aq_x = 0$ (Drop tildes)

$$q(x,t) = \begin{bmatrix} p(x,t) \\ u(x,t) \end{bmatrix}, \qquad A = \begin{bmatrix} u_0 & K_0 \\ 1/\rho_0 & u_0 \end{bmatrix}$$

Eigenvalues: $\lambda = u_0 \pm c_0$

where $c_0 = \sqrt{K_0/\rho_0} = \sqrt{P'(\rho_0)}$ is the linearized sound speed.

$$\tilde{p}_t + u_0 \tilde{p}_x + K_0 \tilde{u}_x = 0,$$

$$\rho_0 \tilde{u}_t + \tilde{p}_x + \rho_0 u_0 \tilde{u}_x = 0,$$

gives the system $q_t + Aq_x = 0$ (Drop tildes)

$$q(x,t) = \begin{bmatrix} p(x,t) \\ u(x,t) \end{bmatrix}, \qquad A = \begin{bmatrix} u_0 & K_0 \\ 1/\rho_0 & u_0 \end{bmatrix}$$

Eigenvalues: $\lambda = u_0 \pm c_0$

where $c_0 = \sqrt{K_0/\rho_0} = \sqrt{P'(\rho_0)}$ is the linearized sound speed.

Usually $u_0 = 0$ for linear acoustics. Then $\lambda^1 = -c_0$, $\lambda^2 = +c_0$.

Hyperbolicity

A system of m conservation laws $q_t+f(q)_x=0$ is called hyperbolic at some point \bar{q} is state space if

The $m \times m$ Jacobian matrix $f'(\bar{q})$ is diagonalizable with real eigenvalues $\lambda^1(q), \ldots, \lambda^m(q)$.

Then small disturbances about the steady state $q=\bar{q}$ satisfy a linear hyperbolic system and propagate as waves.

- Shallow water equations are hyperbolic for h > 0.
- Nonlinear elasticity hyperbolic if $\sigma'(\epsilon) > 0$.
- Gas dynamics hyperbolic if $P'(\rho) > 0$.

Quasi-linear form: $q_t + f'(q)q_x = 0$ Usually want to use conservation form!