

# An experimental investigation of the coupling between elastodynamic and hydraulic properties of naturally fractured rock at the laboratory scale

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**ABSTRACT:** We report on a series of laboratory experiments designed to simulate local effective stress field fluctuation and its influence on the evolution of permeability and dynamic stiffness in fractured samples of Westerly Granite. L-shaped samples are loaded with tri-axial stresses and fractured in situ. The fracture is subsequently sheared in two 4-mm steps. Oscillatory changes in the local effective stress field are imposed through application of normal stress or pore water pressure oscillations with varying amplitudes and frequencies. Active ultrasonic data (ultrasonic waves transmitted across the fracture) is used to monitor the evolution of wave velocity and attenuation before, during and after dynamic stressing. Throughout the experiment, the evolution of permeability is concurrently measured to determine the relationship between fracture permeability and nonlinear elastodynamic properties (stress-dependency of wave velocity and attenuation). Our results to date indicate that relative changes in wave velocity and permeability, due to both normal stress and pore pressure oscillations, are correlated, such that larger drops in wave velocity correspond to larger increases in permeability. Shearing of the fracture reduces the nonlinearity measured during normal stress oscillations for both rock samples. After shearing, the oscillations become generally less effective in enhancing the fracture permeability.

## 1. MOTIVATION AND BACKGROUND

In the context of induced seismicity due to fluid injection, of particular concern are dynamic stresses associated with injection, pumping and transport of supercritical H<sub>2</sub>O–CO<sub>2</sub> fluids. These stresses may pose significant risk associated with accelerated deformation, fault reactivation and possible damage to reservoir seals. Regional increases in seismic activity induced by such dynamic stresses have been reported in numerous studies (Healy et al., 1968; Raleigh et al., 1976; Simpson et al., 1988; Davis and Pennington, 1989; Sminchak and Gupta, 2003; Deichmann and Giardini, 2009; Frohlich, 2012; Horton, 2012; Zoback and Gorelick, 2012; Zoback, 2012; Brodsky and Lajoie, 2013; Ellsworth, 2013; Holland, 2013; van der Elst et al., 2013; McNamara et al., 2015; McGarr et al., 2015; Walsh and Zoback, 2015; Weingarten et al., 2015). The recent seismic activity in Ridgecrest California (Ross et al., 2019) and its impact on the community around Los Angeles and local geothermal energy production are stark reminders of the impact earthquakes can have on energy. However, in the context of energy recovery, these dynamic stresses may be beneficial via enhancing permeability. Transient permeability changes caused by dynamic stresses associated with passing seismic waves have been reported at the field scale and also demonstrated in laboratory

experiments (Elkhoury et al., 2006; Faoro et al., 2009; Elkhoury et al., 2011; Faoro et al., 2012; Candela et al., 2014, 2015; Carey et al., 2015; Frash et al., 2016; Im et al., 2018; Ishibashi et al., 2018; Zhang et al., 2018; Shi et al., 2018; Ye and Ghassemi, 2018; Shi et al., 2019). These studies show that fluid injection and dynamic stresses combine to produce significant changes in permeability, fault stability and poromechanical properties of rock. Such effects have significant implications for fossil fuel production and geothermal energy.

On the other hand, a number of studies have shown that earthquakes and the ensuing seismic waves cause transient changes in rock stiffness in the vicinity of faults (Li et al., 1998, 2006; Niu et al., 2008; Brenguier et al., 2008; Wu et al., 2009; Brenguier et al., 2014). Similar behavior has been recorded at laboratory scale (Shokouhi et al., 2017a). The parallels between the observed transient changes in permeability and stiffness due to dynamic stress perturbations suggest a physical link between the nonlinear stiffness and hydraulic characteristics of fractures (e.g., Witherspoon et al., 1980; Gale, 1982; Brown, 1987; Pyrak-Nolte et al., 1987; Durham & Bonner, 1995; Zimmerman & Bodvarsson, 1996; Berryman et al., 2002; Berkowitz, 2002; Guéguen & Schubnel, 2003; Hofmann et al., 2005; Batzle et al., 2006; Xu et al., 2006). Existing studies on coupling

between stiffness and poromechanical behavior of fractured samples focus primarily on quasi-static stress protocols (Mavko, 2013; Nara et al., 2011; Petrovitch et al., 2013; Pyrak-Nolte and Nolte, 2016). In this study, we carry out well-controlled laboratory experiments to simulate dynamic stressing and illuminate the coupling between permeability transients and nonlinear elastodynamic response. This ongoing effort addresses the critical need for improved understanding of how fluid injection and associated changes in stresses influence

## 2. EXPERIMENTAL SETUP

We conducted laboratory experiments designed to investigate the influence of dynamic effective stress oscillations on the evolution of permeability and elastodynamic stiffness of fractures in Granite subject to shearing. L-shaped samples sandwiched between steel forcing blocks (with embedded ultrasonic piezoelectric transducers) were jacketed and placed inside a pressure vessel, itself inserted into a biaxial load frame as shown in **Figure 1a**. The samples were first saturated with

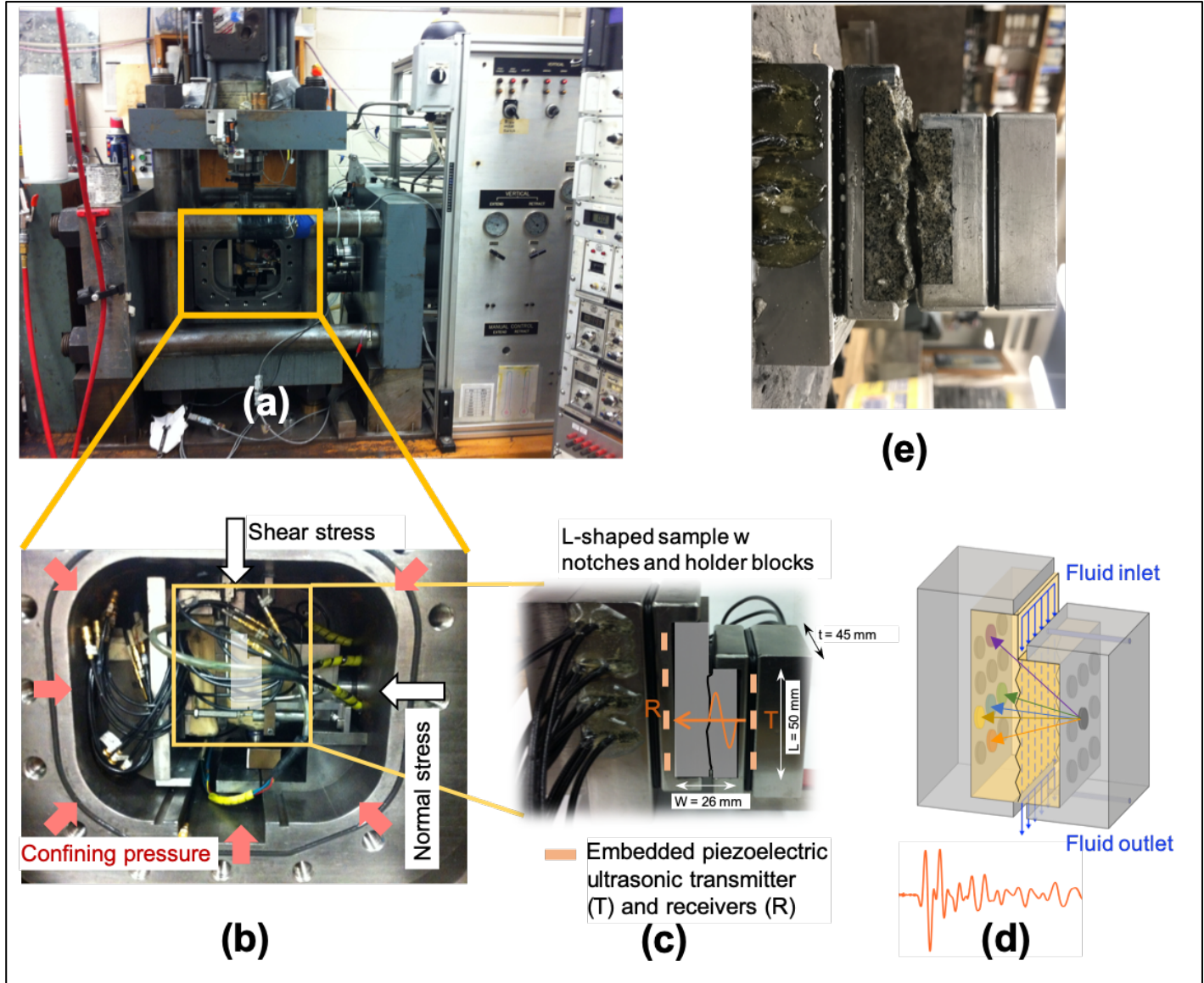


Figure 1 (a) The biaxial loading apparatus in Penn State's Rock Mechanics lab including (b) a pressure vessel to apply triaxial stress to the sample. (c) The L-shaped sample is sandwiched between two holder blocks with embedded piezoelectric transducers. (d) The experimental setup allows simultaneous measurements of flow and multi-channel ultrasonic data. (e) Picture of an in-situ fractured sample of Westerly Granite.

seismic and hydromechanical properties of rock.

This paper provides a brief description of the experimental approach and summarizes our main findings to date. A more in-depth discussion on some of the presented results can be found in Shokouhi et al. (2019).

deionized, de-aired water before being loaded to a normal stress and confining pressures of 20 MPa and 15 MPa, respectively. Subsequently, the samples were fractured in situ by applying a shear load in servo displacement-control. The fracture plane was guided by starter notches at the top and bottom of the sample in order to create reproducible quasi-planar fractures. Displacements of the

horizontal and vertical loading pistons were measured continuously using a pair of Direct Current Displacement Transducers (DCDTs). A Linear Variable Differential Transformer (LVDT) was placed inside the pressure vessel to allow for a redundant and more accurate measurement of displacement normal to the fracture plane. Deionized water was forced through the fracture via a differential pore pressure ( $P_{p-inlet} = 4 \text{ MPa}$  and  $P_{p-outlet} = 2 \text{ MPa}$ ) along the fracture length. The pore pressure intensifiers are independently servo-controlled and equipped with LVDTs to measure fluid flow rates into and out of the fracture plane (**Figure 1d**).

Changes in the applied effective stress field were imposed on the fractured samples via pore pressure ( $P_p$ ) and normal stress ( $\sigma_n$ ) oscillations. Pore pressure oscillations were applied by oscillating  $P_{p-inlet}$  while maintaining  $P_{p-outlet}$  constant. Conversely, normal stress oscillations were applied by oscillating the horizontal piston of the load frame at prescribed amplitudes and frequencies. Multiple sets of dynamic stress oscillations of varying amplitude ( $\sim 0.01$  up to about  $\pm 1 \text{ MPa}$ ) and frequency (0.1, 1, 10 and 40 Hz) were applied to investigate the repeatability as well as amplitude and frequency dependencies of the measured response. A multi-channel data acquisition system was used to record ultrasonic data in both active and passive modes. Ultrasonic waves (with a center frequency of  $\sim 500 \text{ kHz}$ ) transmitted across the fracture (active mode) were used to monitor the evolution of wave velocity and amplitude before, during and after dynamic stressing. The active ultrasonic data are used to determine changes in elastodynamic properties and fracture stiffness due to the imposed oscillations. After applying a series of normal stress and pore water pressure oscillations, the fracture was sheared in two 4-mm steps (a total of 8mm). The dynamic stress oscillation protocol was executed repeatedly after each shearing step. During fracturing and shearing, the seismicity along the fracture was monitored passively in acoustic emission (AE) mode. Throughout the experiment, the evolution of flow rates at inlet and outlet was used to determine the stress-induced changes in permeability and ultimately to investigate the relationship between fracture permeability, aperture and stiffness. Finally, high-resolution optical profilometry was used to measure the roughness across the two post-mortem fractured surfaces in order to reconstruct fracture aperture distribution (results not shown in this paper).

Here, we summarize the main findings from the experiments on two Westerly granite samples WG1 (experiment p4966) and WG2 (experiment p4975) obtained using multiple ultrasonic transmitter-receiver pairs shown in **Figure 1d**. Reporting the results from multiple transmitter-receiver pairs sets this paper apart from our recently-published closely related publication

(Shokouhi et al., 2019). Also in this paper, we compare the results before and after shearing the fracture.

### 3. NONLINEARITY OF THE ELASTODYNAMIC RESPONSE: TRANSIENT CHANGES IN STIFFNESS WITH DYNAMIC STRESSING AND SHEAR

We measure the nonlinearity of the elastodynamic response directly by quantifying the stress-dependency of wave velocity and amplitude. **Figure 2** depicts the typical response of our fractured samples to dynamic stress perturbations. Immediately after effective stress oscillations, the fractured rock exhibits an instantaneous softening i.e., the wave velocity  $c_0$  suddenly drops. However, if the dynamic perturbation is not too strong, this softening is transient; once the oscillation is removed, the velocity slowly recovers toward the pre-oscillation wave velocity  $c_0$ . Note that during the oscillations, the wave velocity oscillates in response to dynamic stressing at frequencies that are harmonics of the excitation frequency. The average amplitude of these oscillations (at the oscillation frequency) is denoted as  $dc$  in **Figure 2**. If the stress oscillations persist, the changes in wave velocity reach a non-equilibrium steady state. This characteristic response (transient softening, slow recovery and velocity fluctuations) is a signature of nonlinear mesoscopic elasticity (Geyer and Johnson, 2009) and rich

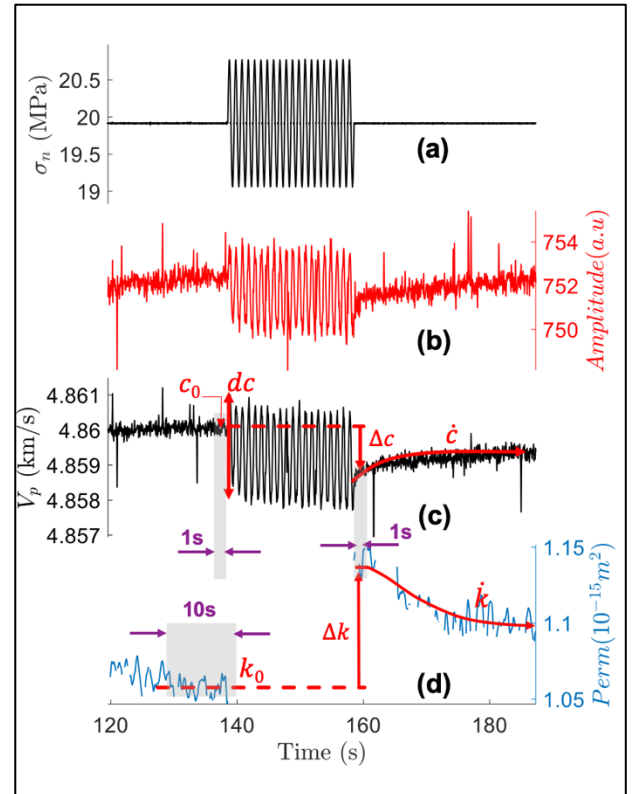


Figure 2. Typical response of fractured rock to stress perturbation: (a) imposed normal stress oscillations, the resulting evolutions of (b) ultrasonic wave amplitude, (c) velocity and (d) permeability with time.



in information on microstructure, fractures and contact mechanics. In comparison, the wave velocity in a linear elastic medium is stress-invariant i.e.,  $c = c_0$  before, during and after stress oscillations. Most intact rocks exhibit some degree of nonlinearity due to soft grain boundaries and microcracks in their matrix (Rivière et al., 2015). However, when fractured, their nonlinear electrodynamic responses are also influenced by the contact acoustic nonlinearity (CAN) at the rough fractured interface.

In order to quantify the nonlinearity of the fractures, we first obtain the evolution of ultrasonic wave amplitude (Figure 2b) and velocity (Figure 2c) from the continuously measured ultrasonic data (Shokouhi et al., 2019) and calculate three parameters: (1) the average or DC-change in wave velocity due to the imposed oscillations, (2) the amplitude of steady-state velocity fluctuation during the oscillations, and (3) the recovery rate of wave velocity post oscillation (Figure 2c). The relative velocity change ( $\Delta c/c_0$ ) is defined as the % change in velocity due to the imposed normal stress or pore pressure oscillation. We calculate  $\Delta c/c_0$  from the velocity before ( $c_0$ ) and after ( $c$ ) each oscillation averaged over 1-s time windows as depicted in Figure 2c. At a given oscillation amplitude, the more negative the  $\Delta c/c_0$  (larger absolute value), the more nonlinear is the fracture. The parameter  $dc/c_0$  is defined as the amplitude of velocity oscillations, averaged over the oscillation duration and normalized by pre-oscillation velocity  $c_0$  (Figure 2c). Similar to  $\Delta c/c_0$ , a larger magnitude of

$dc/c_0$  at a given oscillation amplitude is an indication of higher nonlinearity. The third parameter  $\dot{c}$  is defined as the (logarithmic) rate of recovery of wave velocity after the oscillation is removed (Figure 2c). Materials of higher nonlinearity are expected to have slower recoveries (Shokouhi et al., 2017b).

In the next section, we discuss the dependence of  $\Delta c/c_0$ ,  $dc/c_0$  and  $\dot{c}$  (Figure 3) on the imposed normal stress and pore pressure amplitudes as well as frequency. In addition, we compare the measured nonlinearities before and after shearing the fracture. This comparison provides insight into how changes in the aperture size distribution (due to shearing and wear) and the presence of fines alter the fracture stiffness and the stress-dependency of the elastodynamic response. Although not shown here, similar nonlinearity parameters may be extracted from the evolution of ultrasonic wave amplitude (Figure 2b).

### 3.1. Oscillation Amplitude- and Frequency-Dependencies

The details on amplitude- and frequency-dependencies of  $\Delta c/c_0$  and  $dc/c_0$  for a single transmitter-receiver pair are given in Shokouhi et al. (2019). Similar trends are observed for all transmitter-receiver pairs and therefore, not repeated here. In summary, the results given in Shokouhi et al. (2019) indicate that  $\Delta c/c_0$  and  $dc/c_0$  are modulated by dynamic stress amplitudes for both pore pressure oscillations and fracture normal stress oscillations (Figure 3 and Figure 4). In addition, the measurements exhibit frequency-dependence (not shown

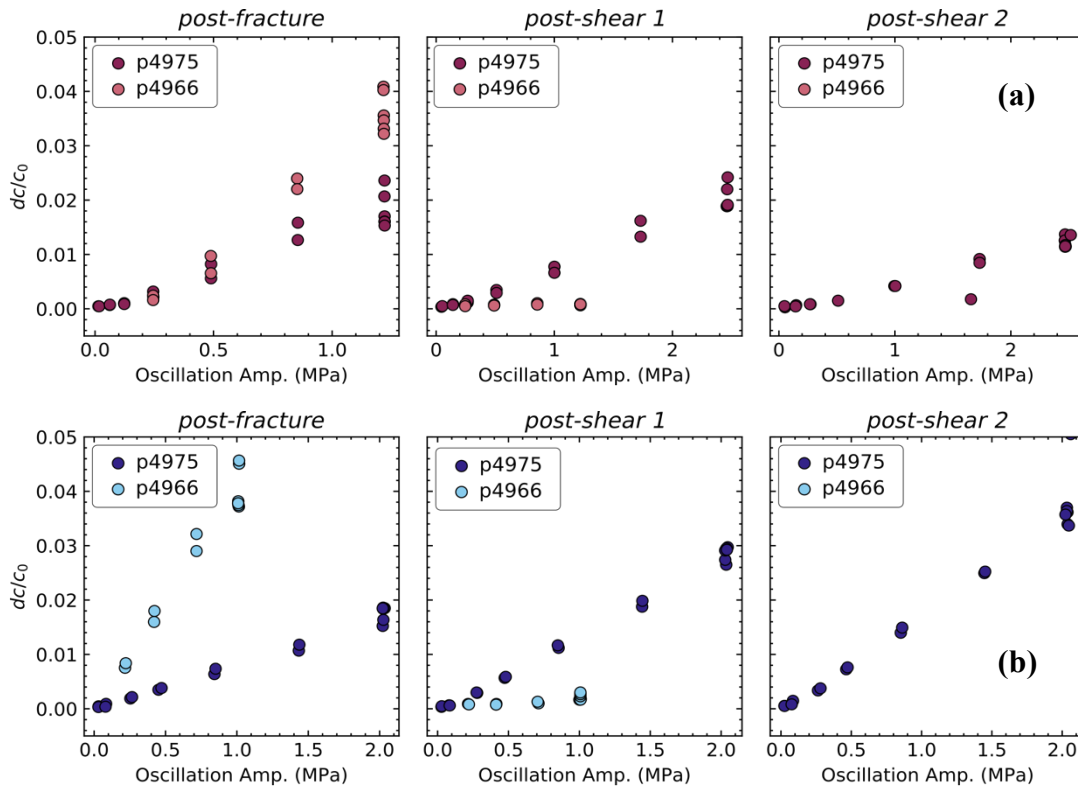


Figure 3 Amplitude-dependency of  $dc/c_0$  post fracture (left panels), after shear step 1 (middle panels) and after shear step 2 (right panels) calculated for (a) normal stress and (b) pore pressure oscillations.

here);  $\Delta c/c_0$  increases with the oscillation frequency, while  $dc/c_0$  mostly decreases. Observing different trends for the different nonlinearity parameters is not surprising. Although the origins of  $\Delta c/c_0$  and  $dc/c_0$  remain unclear, there is empirical evidence that they stem from different micro-mechanical mechanisms (Rivière et al., 2016, 2015). Although the  $\dot{c}$  values tend to generally increase with the oscillation amplitude, we do not observe a clear trend (not shown). This could be in part due to the variability in  $\dot{c}$  caused by the low signal-to-noise ratio. The recovery rate  $\dot{c}$  does not show a systematic dependency on frequency (not shown).

### 3.2. Influence of Shearing

Shearing of the fracture reduces  $dc/c_0$  during normal stress oscillations (**Figure 4a**). Interestingly, we have recorded different post-shear behavior for normal stress vs. pore pressure oscillations; the shearing increases the nonlinearity during pore pressure oscillations for one sample (p4975), while decreasing it for the other sample (p4966) (**Figure 4b**). This and other observations indicate possible differences in mechanisms activated during these

The stress-induced changes in permeability are captured by two parameters: (1) The transient change in permeability  $\Delta k/k_0$  defined as the % change due to the imposed normal stress or pore pressure oscillations, normalized by the pre-oscillation permeability  $k_0$  as illustrated in **Figure 2c** (Candela et al., 2014) and (2) The recovery rate of permeability after the transient increase  $\dot{k}$ .

### 4.1. Oscillation Amplitude- and Frequency-Dependencies

The dependencies of  $\Delta k/k_0$  on the amplitude (**Figure 5**) and frequency of normal stress and pore pressure oscillations are given in Shokouhi et al. (2019). One important observation is that although small-amplitude oscillations ( $< 0.5\text{MPa}$ ) may result in a decrease in permeability, large-amplitude oscillations generally increase permeability transiently. Furthermore, for the same oscillation amplitude and frequency, the pore pressure oscillations appear to be more effective in permeability enhancement ( $\Delta k/k_0 > 0$ ) than the normal stress oscillations by a factor of  $\sim 2.5$ , particularly for

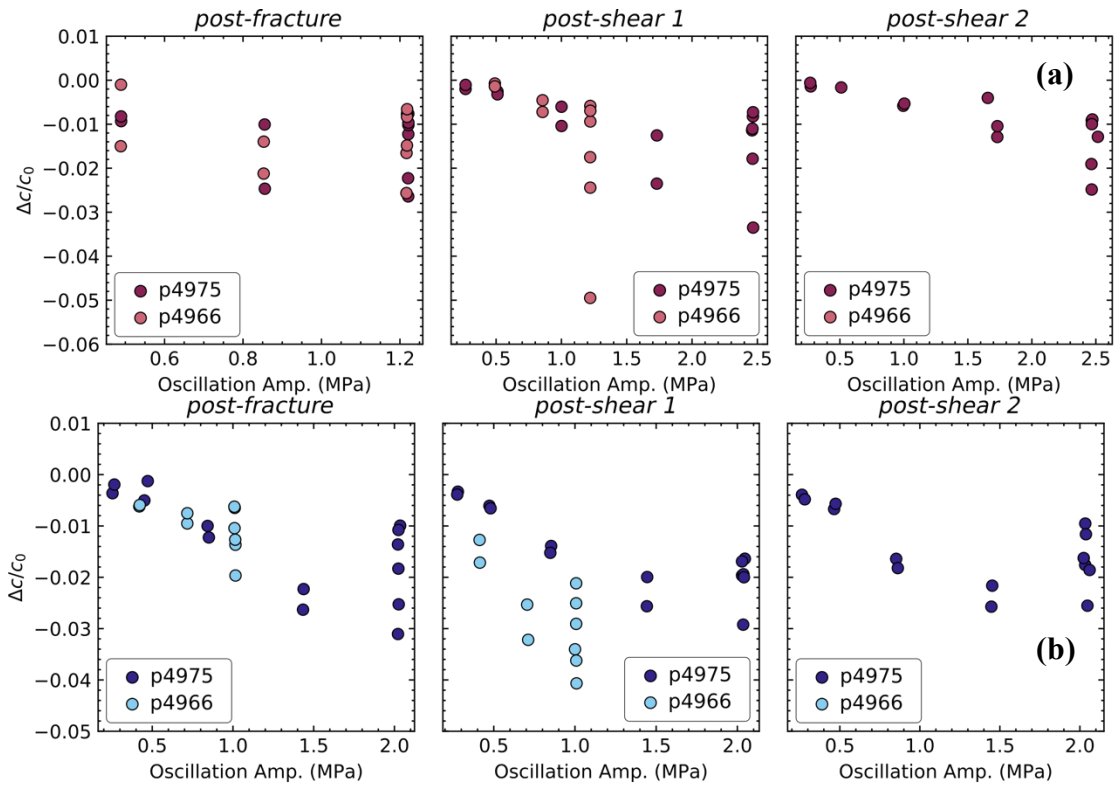


Figure 4 Amplitude-dependency of  $\Delta c/c_0$  post fracture (left panels), after shear step 1 (middle panels) and after shear step 2 (right panels) calculated for (a) normal stress and (b) pore pressure oscillations.

two modes of dynamic stressing. On the other hand, shearing does not have a significant influence on  $\Delta c/c_0$ . The second shearing step decreases the measured nonlinearity  $\Delta c/c_0$  but only slightly.

## 4. PERMEABILITY EVOLUTION: TRANSIENT CHANGES WITH DYNAMIC STRESSING AND SHEAR

sample WG1. In addition,  $\Delta k/k_0$  increases, albeit slightly, with the frequency of oscillations (not shown here). The only exception is the decrease in  $\Delta k/k_0$  for sample WG2, when we increase the fluid pressure oscillations from 1 Hz to 10 Hz. The details on amplitude- and frequency-dependencies of  $\Delta k/k_0$  are given in Shokouhi et al. (2019). The recovery rate  $\dot{k}$  does not show

clear amplitude- and frequency-dependencies (not shown here).

scatter. This observation further reinforces the presumed linkage between the evolution of fracture's elastodynamic

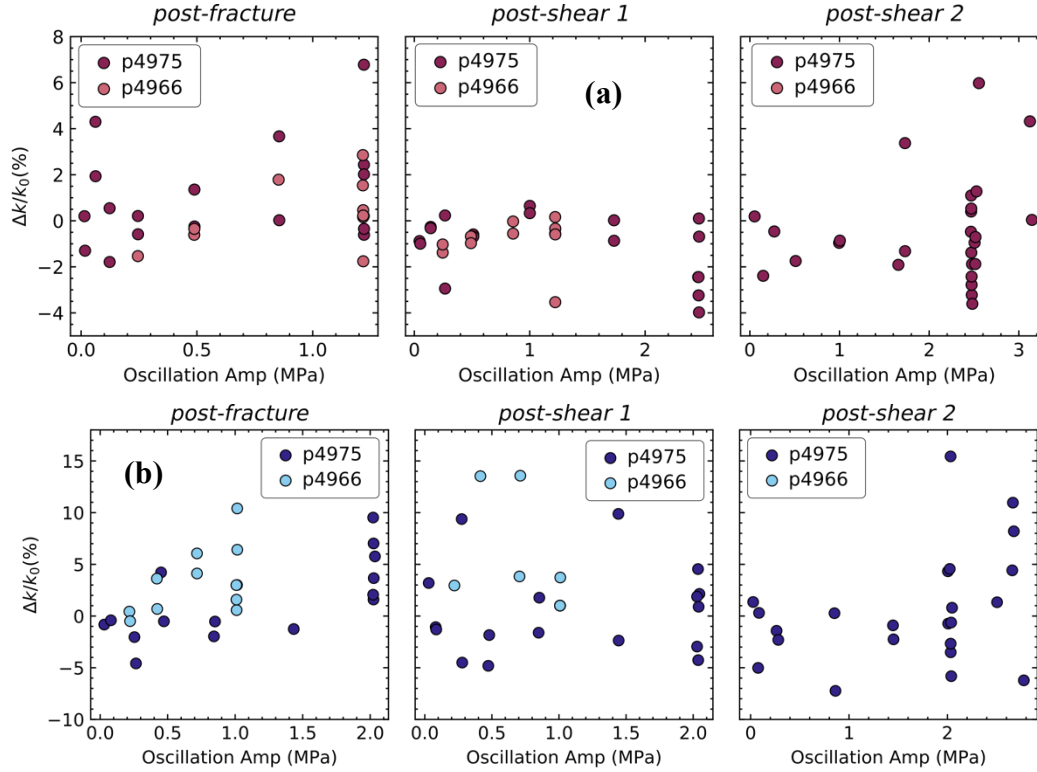


Figure 5. Amplitude-dependency of  $\Delta k/k_0$  post fracture (left panels), after shear step 1 (middle panels) and after shear step 2 (right panels) calculated for (a) normal stress and (b) pore pressure oscillations.

#### 4.2. Influence of Shearing

Generally speaking, after shearing the fracture, the normal stress oscillations become less effective in enhancing the fracture permeability (see **Figure 5a**). The trend is less obvious for pore water pressure oscillations (see **Figure 5b**).

### 5. CORRESPONDENCE BETWEEN STIFFNESS AND PERMEABILITY TRANSIENTS?

The main hypothesis driving this study was that the transient elastic softening is associated with a temporary increase in porosity and/or permeability, both of which are important for energy production and waste storage. Our results to date support this conjecture: the relative changes in wave velocity and permeability of fresh fractures (before shearing), due to both normal stress and pore pressure oscillations, are correlated, such that a larger drop in wave velocity (more negative  $\Delta c/c_0$ , equivalent to larger transient softening or higher nonlinearity) corresponds to a larger permeability enhancement  $\Delta k/k_0$ . As shown in **Figures 6a** and **6b (left panels)**, this overall correlation seems to hold for both normal stress and pore pressure oscillations. Furthermore, the recovery rates of wave velocity and permeability transients ( $\dot{c}$  and  $\dot{k}$ ) also seem to be correlated (**Figure 7**), although there is slightly more

and flow properties. In other words, the form of the measured permeability transients is similar to that for the nonlinear elastic effects and both recoveries follow the time-logarithmic trajectories observed in the field or in nature. We should note that the correlations shown in **Figure 6** correspond to  $\Delta c/c_0$  averaged for 5 transducer pairs shown in **Figure 1d**. Examining the results for individual transducer pairs (not shown here) reveals the strong dependency of the measurements on the wave path. This is not surprising as the fracture aperture is nonuniform and the ray paths sample different portions of that interface.

Shearing the fracture appears to alter these relations indicating the influence of contact stiffness and shear fabric on both sets of properties (**Figure 6a** and **6b**). Overall, shearing seems to weaken the correlation between  $\Delta c/c_0$  and  $\Delta k/k_0$  (**Figure 6a** and **6b, middle and right panels**). The question remains as to what underlying micro-mechanisms link the nonlinear elastodynamic behavior and hydraulic properties of fractures.

We considered two mechanisms for this observed correspondence: (1) the unclogging of fracture flow conduits via fluid pressure oscillations, and (2) the dependence of both properties on stress-induced changes in the fracture aperture. Previous work shows that particle clogging can alter permeability via dynamic stressing in

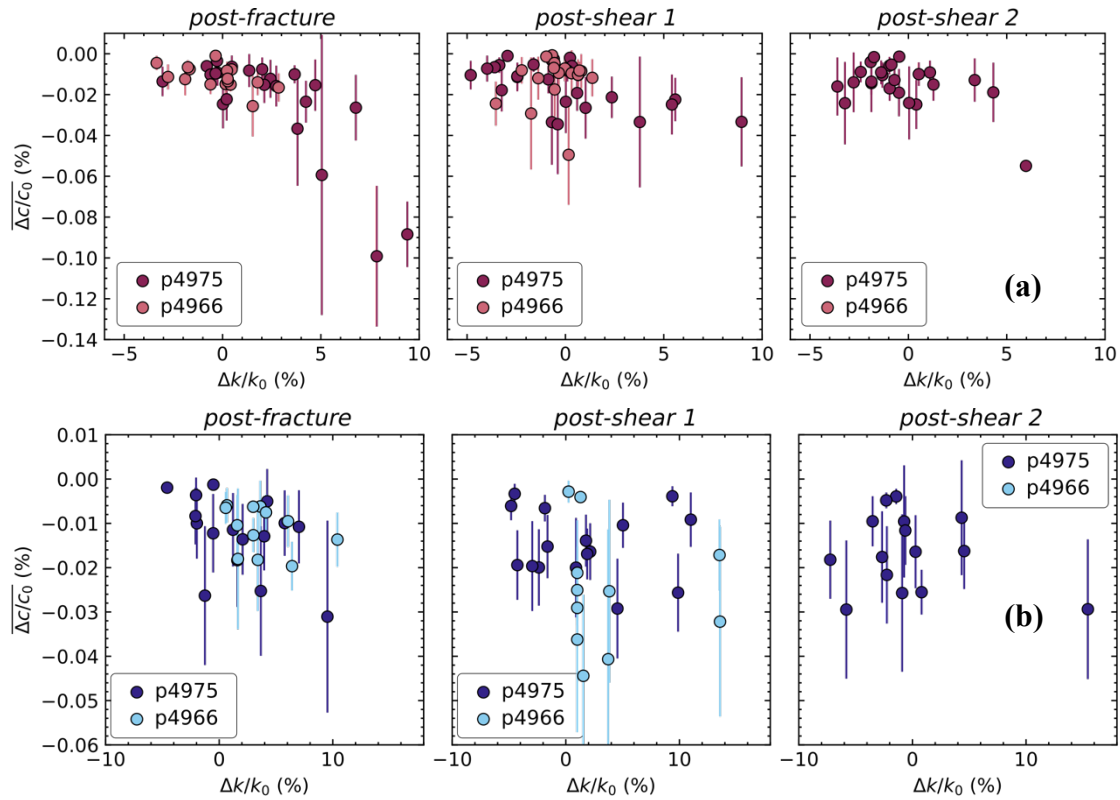


Figure 6. Correspondence between nonlinearity and permeability change due to (a) normal stress and (b) pore water pressure oscillations at three states: (left panel) post fracture, (middle panel) after shear step 1, and (right panel) after shear step 2. The nonlinearity is expressed in terms of average relative change in velocity due to imposed effective stress oscillations  $\Delta c/c_0$ . The  $\Delta c/c_0$  values are averaged over the five wave paths depicted in Figure 1d.

Berea sandstone (Elkhoury et al., 2011; Candela et al., 2014, 2015). However, for fractures in Westerly Granite we observe an increase in permeability for both pore pressure and normal stress oscillations of sufficiently large amplitude. With the new insights from the coupled ultrasonic data, one may conclude that unclogging is merely one of many potential mechanisms.

The other explanation for the linkage between stiffness and flow properties of fractured rocks could be the dependence of both properties on the fracture aperture change due to the imposed stress oscillations. To further examine this hypothesis, we investigated whether  $\Delta c/c_0$  or  $\Delta k/k_0$  correlate with the measured changes in sample thickness changes (proxy for aperture compaction/dilation) during imposed oscillations (Shokouhi et al., 2019). However, we did not observe a strong correlation. This lack of strong correlation suggests that changes in fracture aperture cannot be the sole driving mechanism for permeability change, especially when driven by pore pressure oscillations.

We also used high-resolution optical profilometry to measure the roughness across the post-mortem fractured surfaces in order to reconstruct the fracture aperture distribution. However, since the fractures were created in-situ and were sheared twice before the roughness measurements, we do not have information on the aperture or roughness of fracture interfaces immediately

post-fracture nor immediately post-shear (1<sup>st</sup> shearing step).

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