

From Micro to Macro

Example of Percolation

Chapter I. Lecture 3

Monday April 17

From micro to macro Example of Percolation

Chapter I — Lecture 3

Plan:

1) Introduction to Percolation

- a) Percolation: What is it?
- b) Continuous phase transition
- c) Applications

2) Percolation phase transition

- a) Control parameter and Order parameter
- b) 1d site percolation
- c) Bethe Lattice

3) Properties at the critical point, universality

- a) Different type of percolation problems
- b) P_c is not universal
- c) Critical exponents are universal

4) Universality

Expectations: Participate in the discussions, take notes, try to do the analytical derivations

References: Book “Complexity and Criticality”, K. Christensen, N. Moloney, **Chapter 1: Percolation**

“[Entropy, Order Parameters, and Complexity](#)”, by J. P. Sethna, **12.1 Universality**

Introduction to Percolation problems

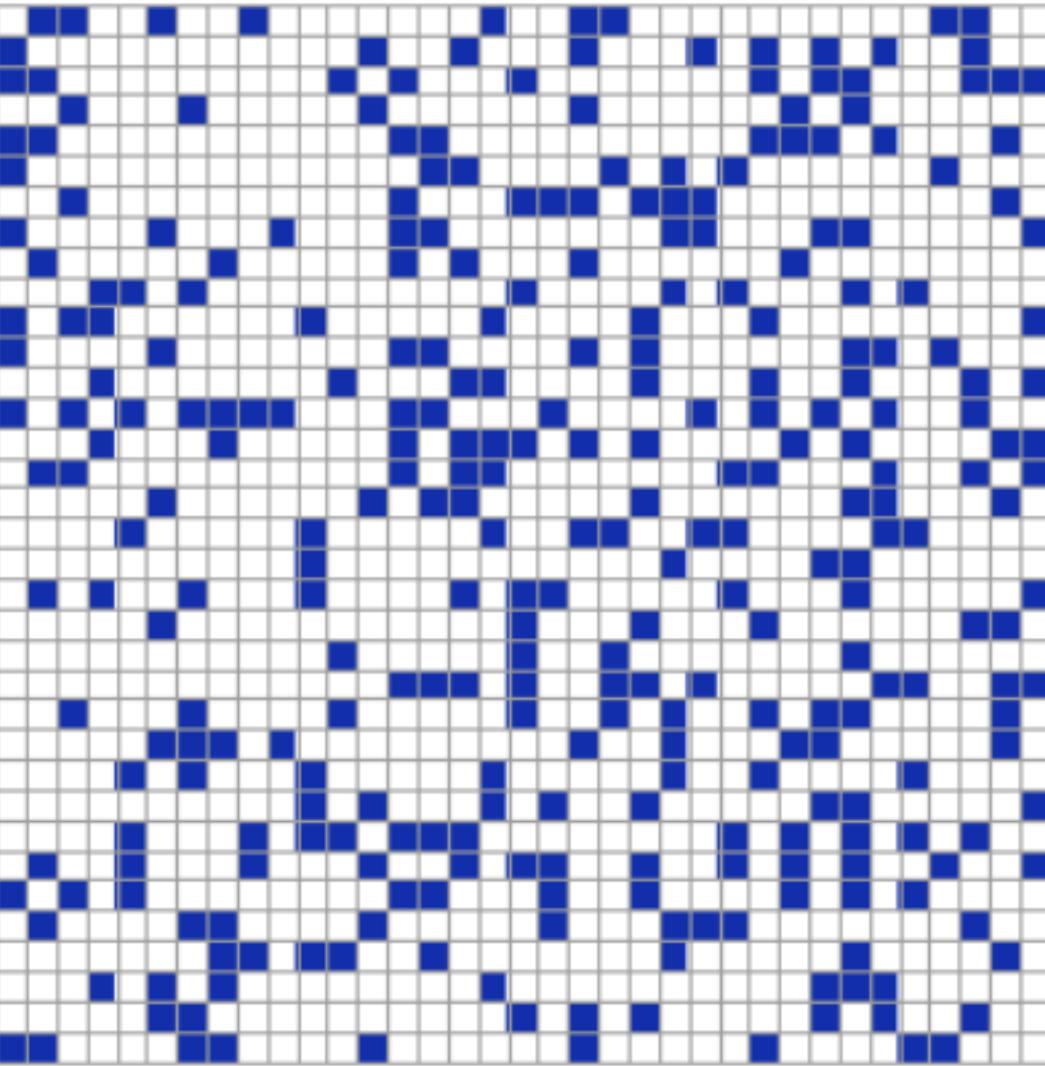
One of the **simplest models** undergoing a **continuous phase transition**.

Yet, relevant in **many applications!**

Percolation

Take a **squared paper** and **randomly black out** a **chosen fraction of the squares**.

Consider the **clusters** of adjacent black squares



A **small fraction** of the squares are **blacked out**,

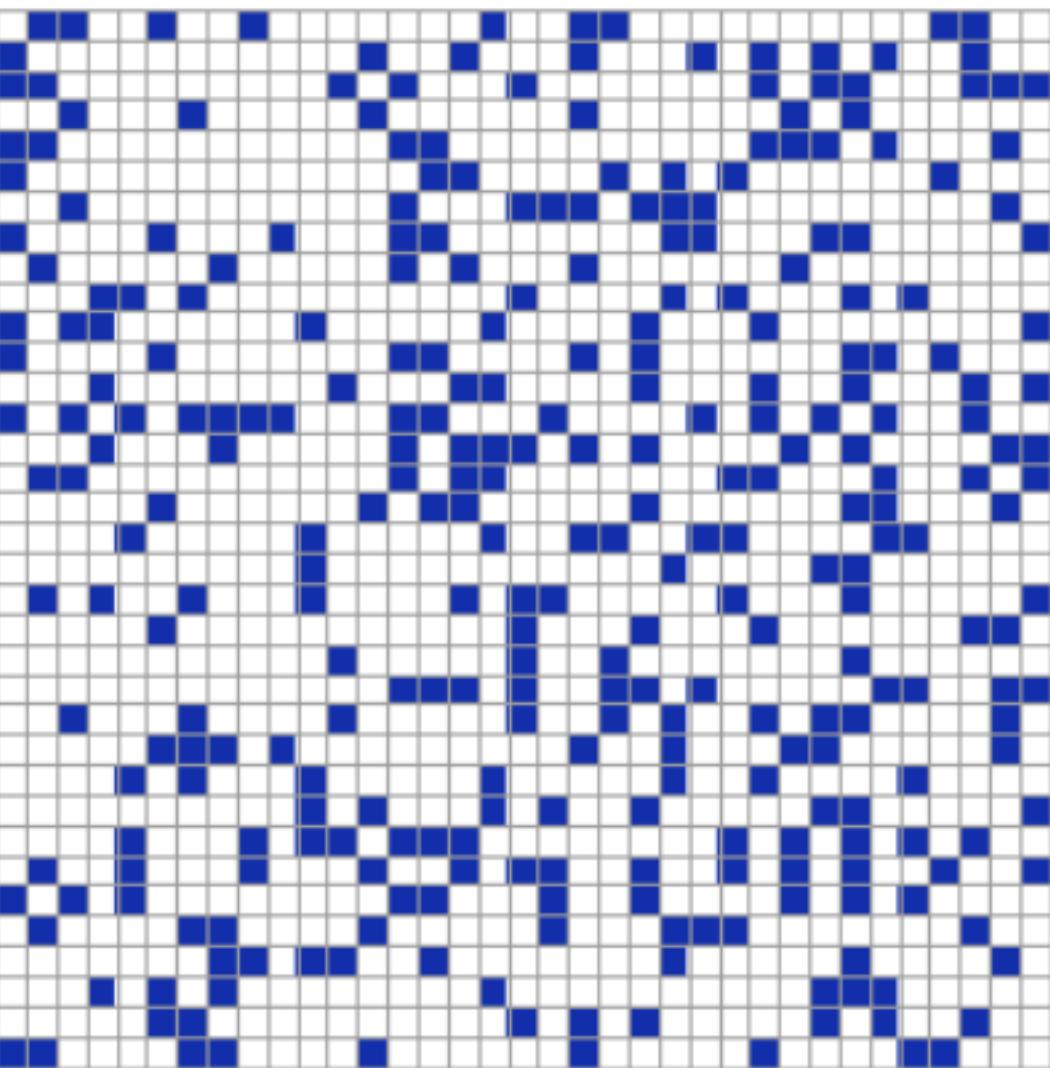


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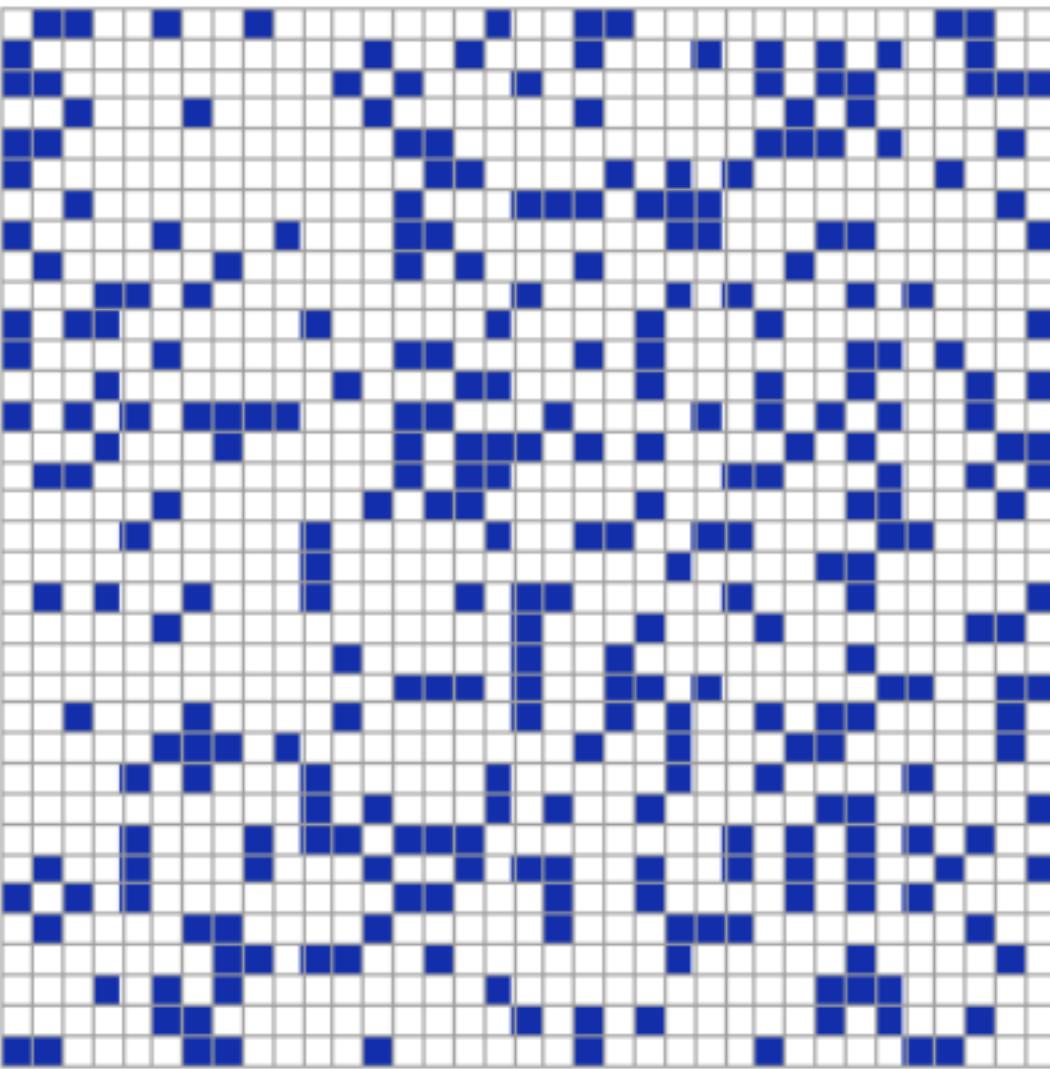


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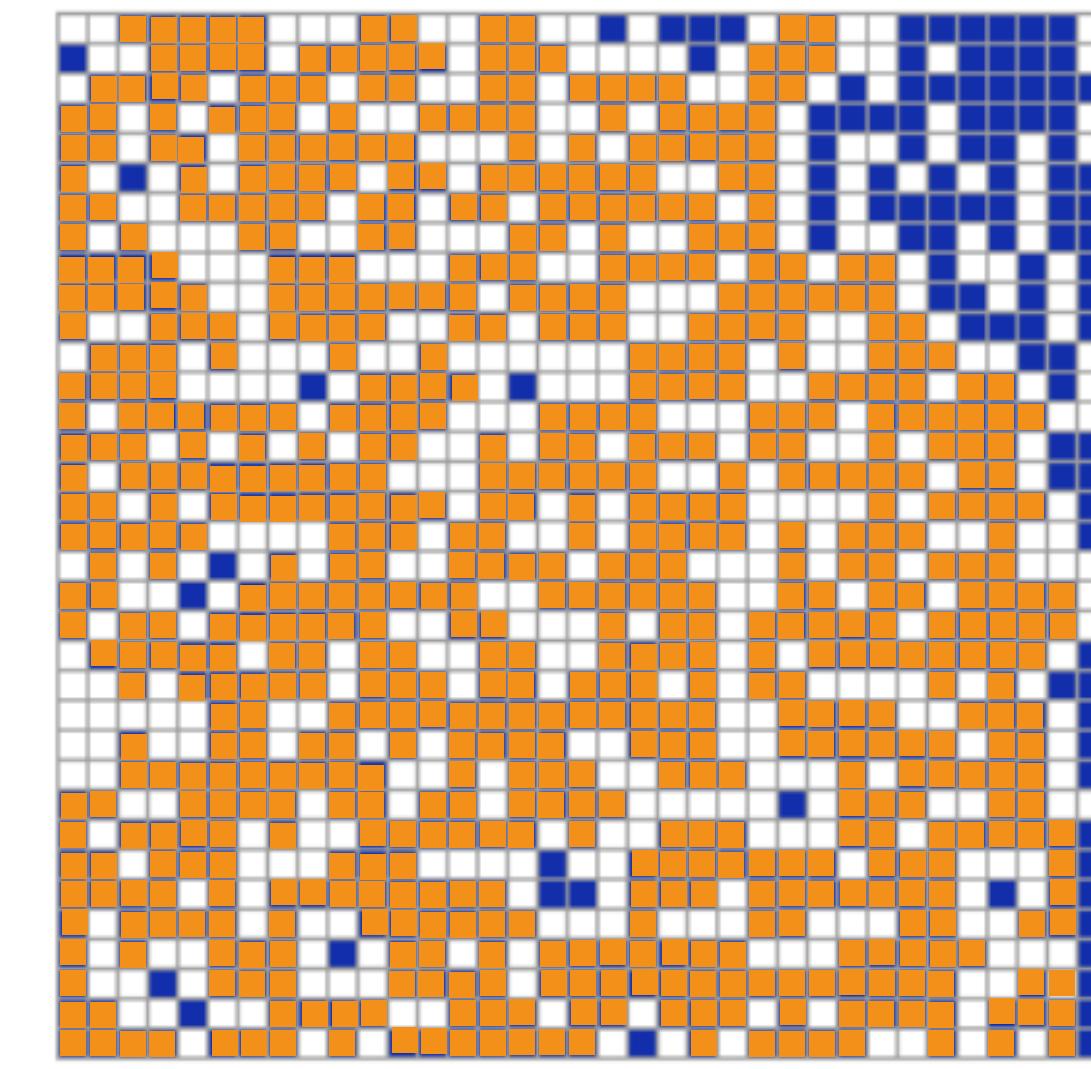
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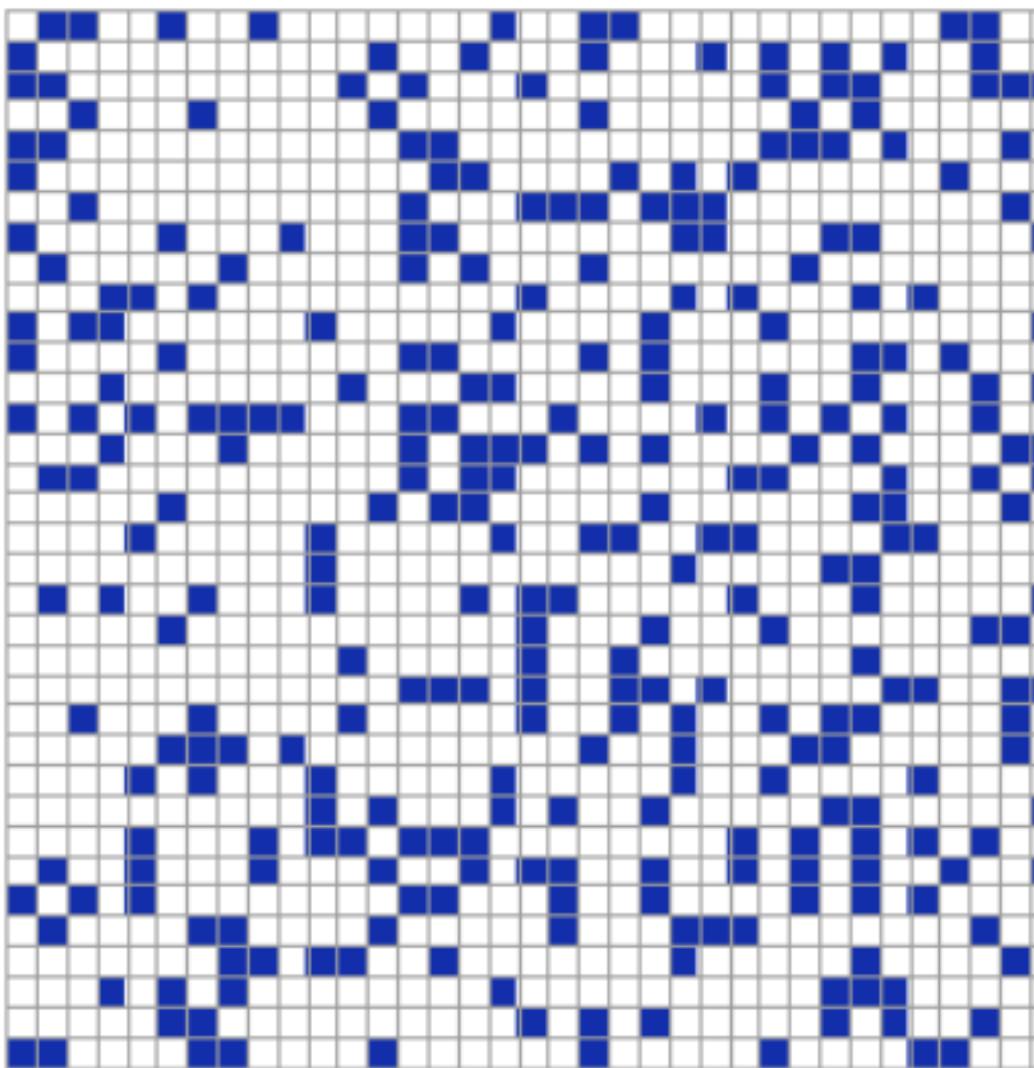
Like water **percolates** through the coffee machine

In physics/chemistry: **percolation** is the process of
filtering a liquid through a porous material.

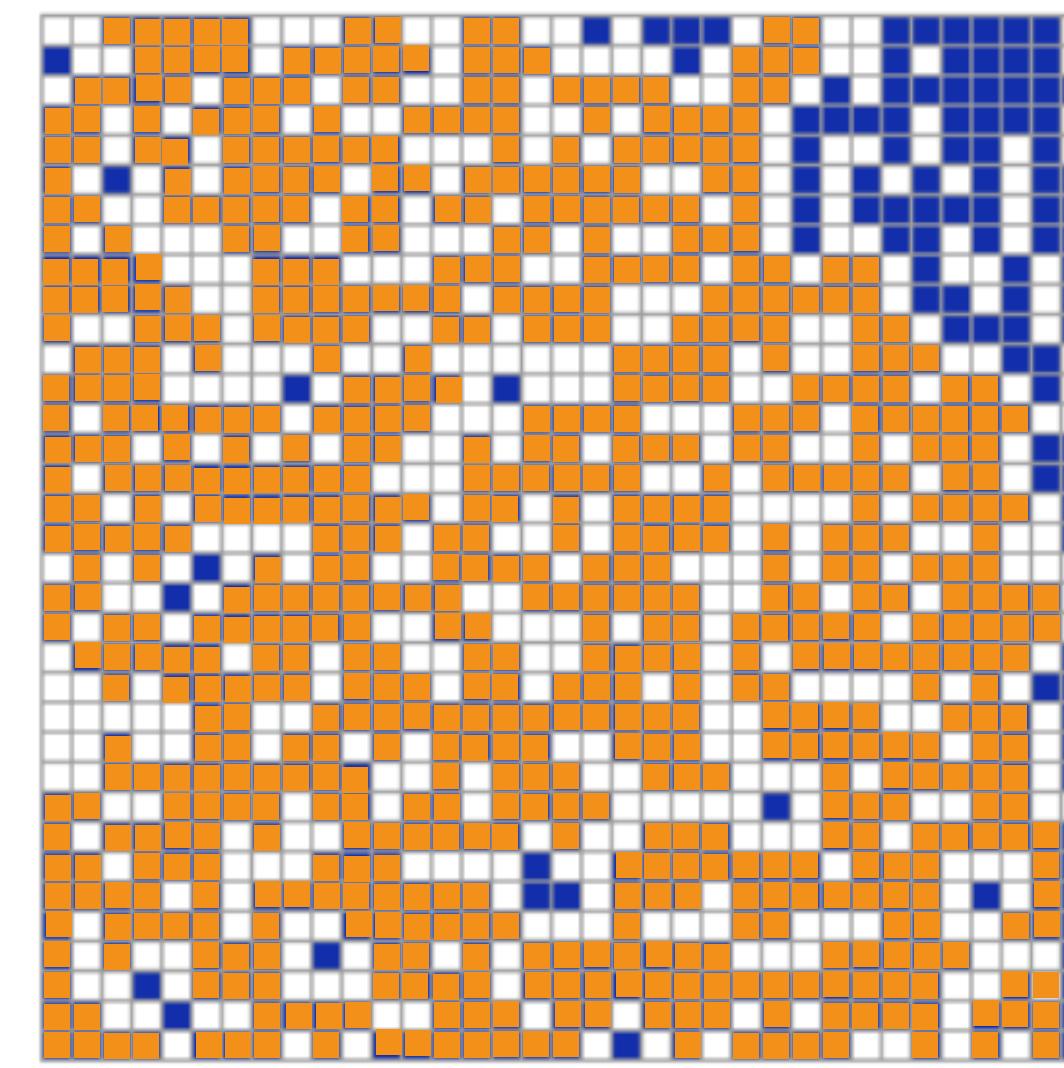
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phase transition
→
Control parameter
= probability to black out a square



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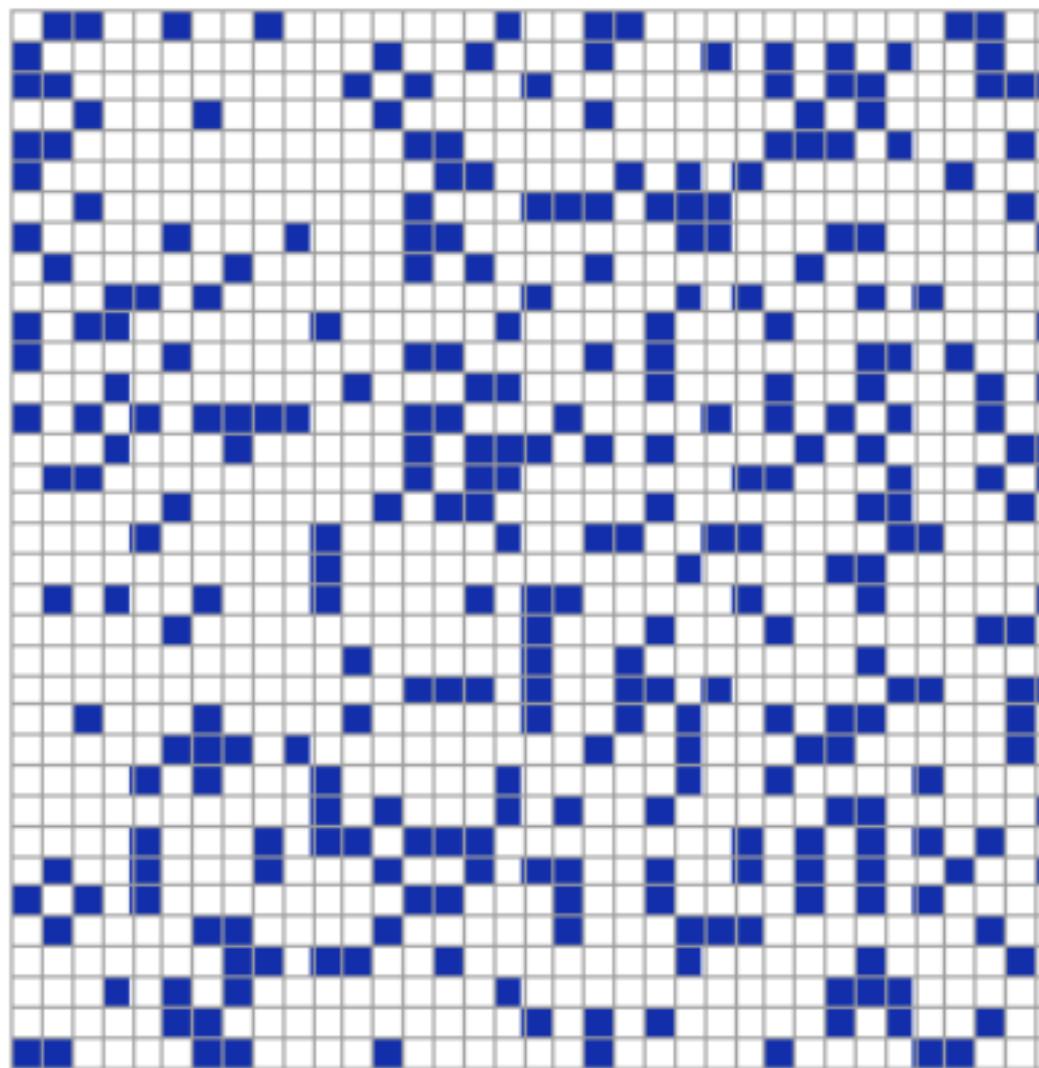
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There is a **fraction of randomly blacked out squares** that **marks** a **phase transition** from **non-percolating cluster** to **percolating cluster**.

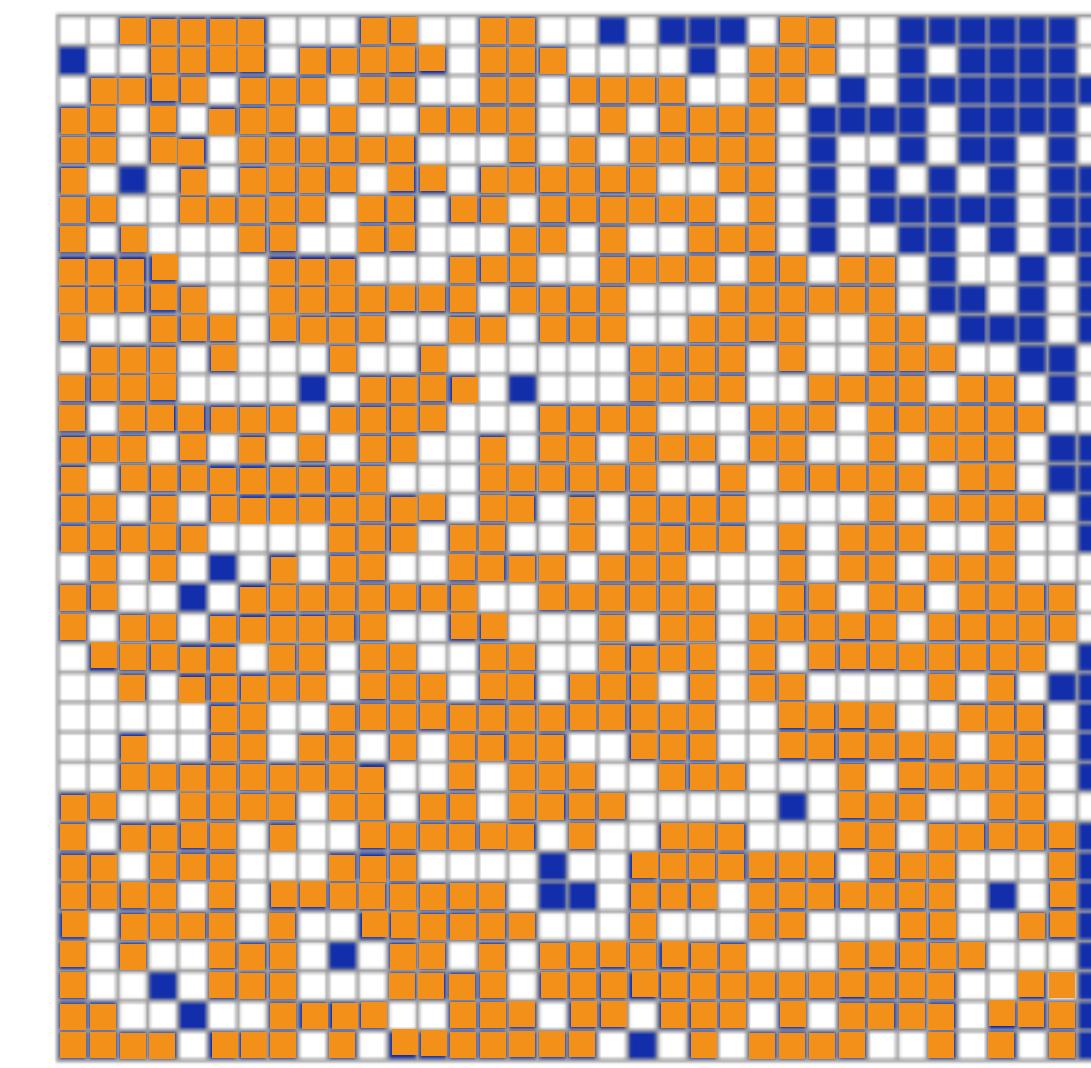
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One of the simplest model undergoing a **PHASE TRANSITION**

Which type
of phase transition?

Experiment: Percolation transition is a continuous phase transition

Experiment:

punching **holes at random** places **in a conducting sheet** of paper
and **measuring the conductance**.

(measures how easily electrical current can pass through a material)

Shows that it is a **continuous transition**:

- **conductance fell to a very small value** as the number of holes approached the **critical concentration**
- the conducting **paths** were **few** and **tortuous** just before the sheet fell apart —> singular properties

What happens at the transition?

One of the simplest model undergoing a **CONTINUOUS** phase transition.

Percolation Theory and Electrical Conductivity

B. J. Last and D. J. Thouless

Department of Mathematical Physics, University of Birmingham, Birmingham B15 2TT, England

(Received 4 October 1971)

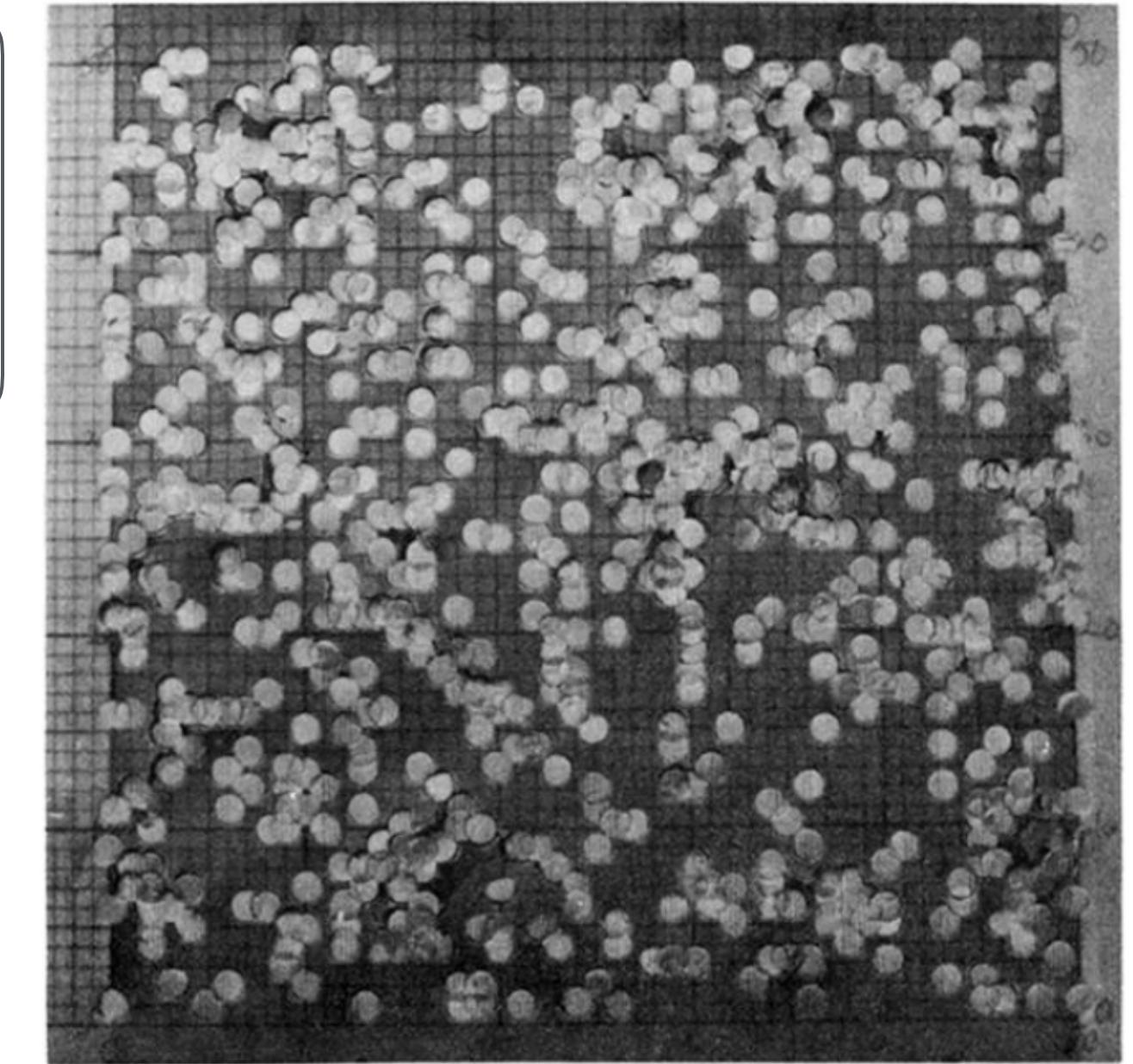
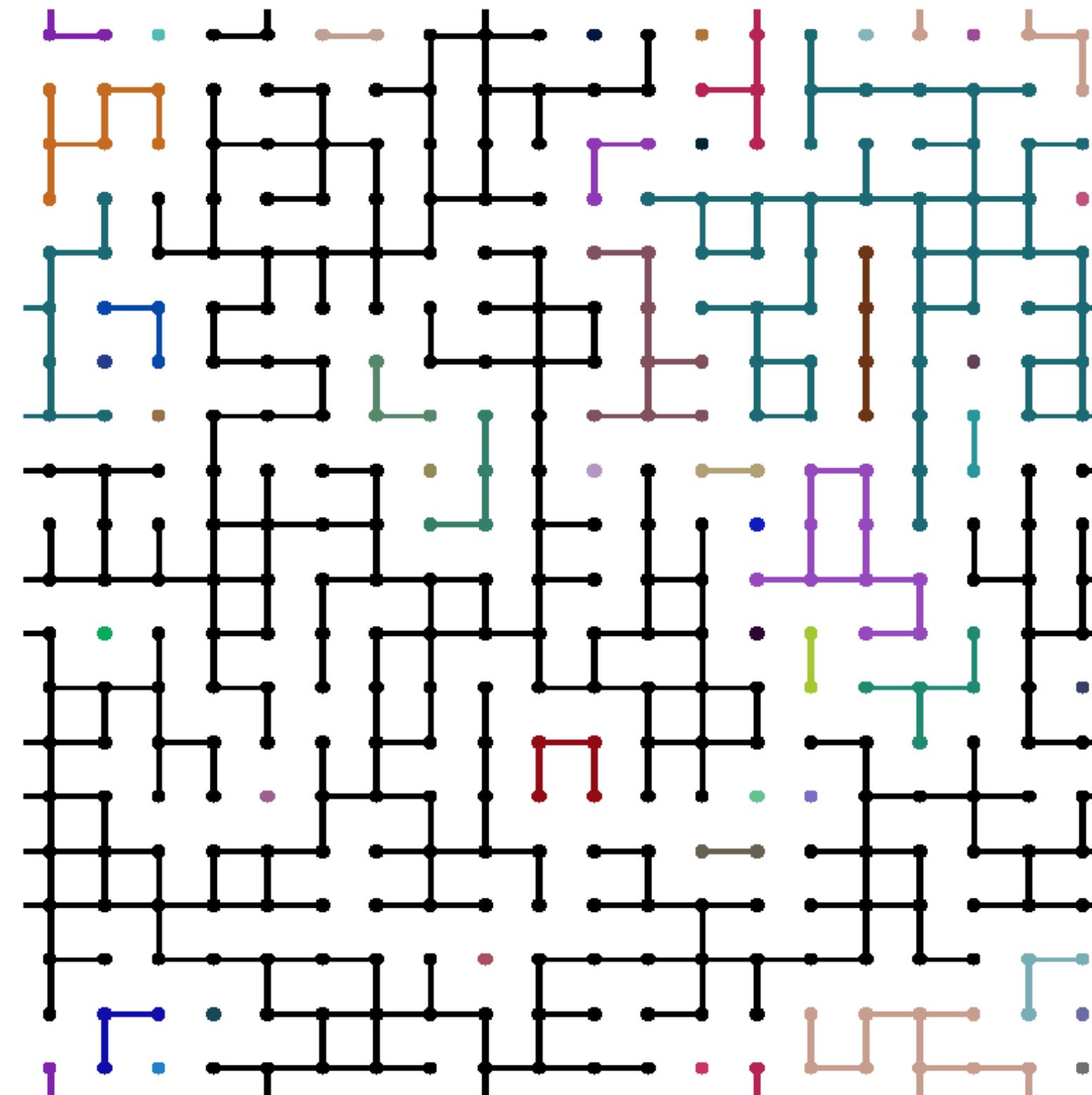


FIG. 1. Photograph of the sheet of conducting paper at the stage where the concentration of holes is 0.268.

What happens at the transition?

Ex. **Simulation of bond percolation:**

- Start from a grid and remove bonds (instead of making holes)
- Bonds between nodes are
 - kept with probability p
 - removed with probability $(1-p)$



In black: the largest cluster span the entire lattice

Other colors: highlight smaller clusters

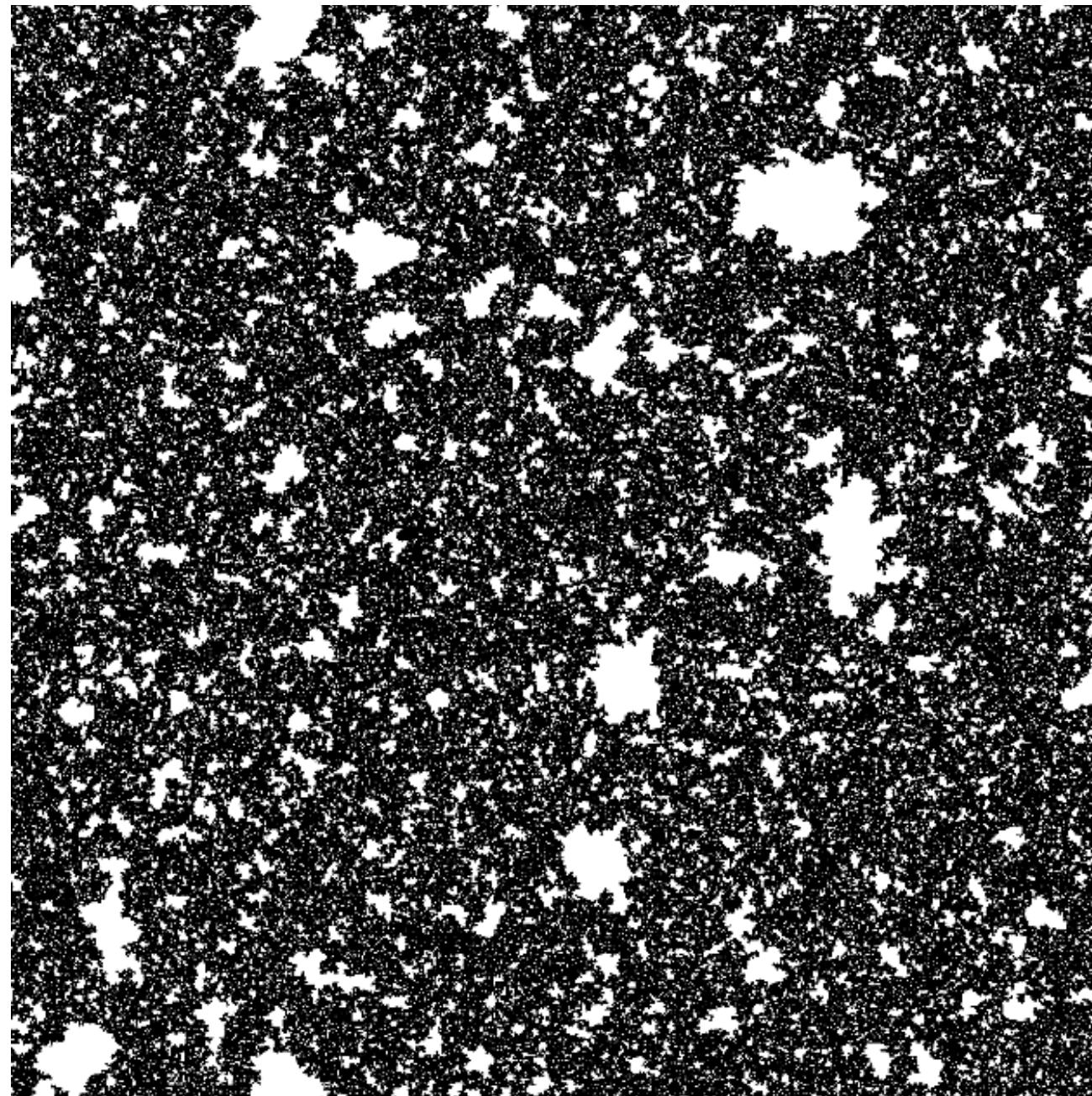
For p close to 1: large conductivity

Conductivity decreases as p decreases.

What happens at the transition?

Ex. **Simulation of bond percolation:**

- › Start from a grid and remove bonds (instead of making holes)
- › **Bonds** between nodes are
 - **kept** with probability p
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$$p = 0.51$$

The largest cluster is intact with only
small holes.

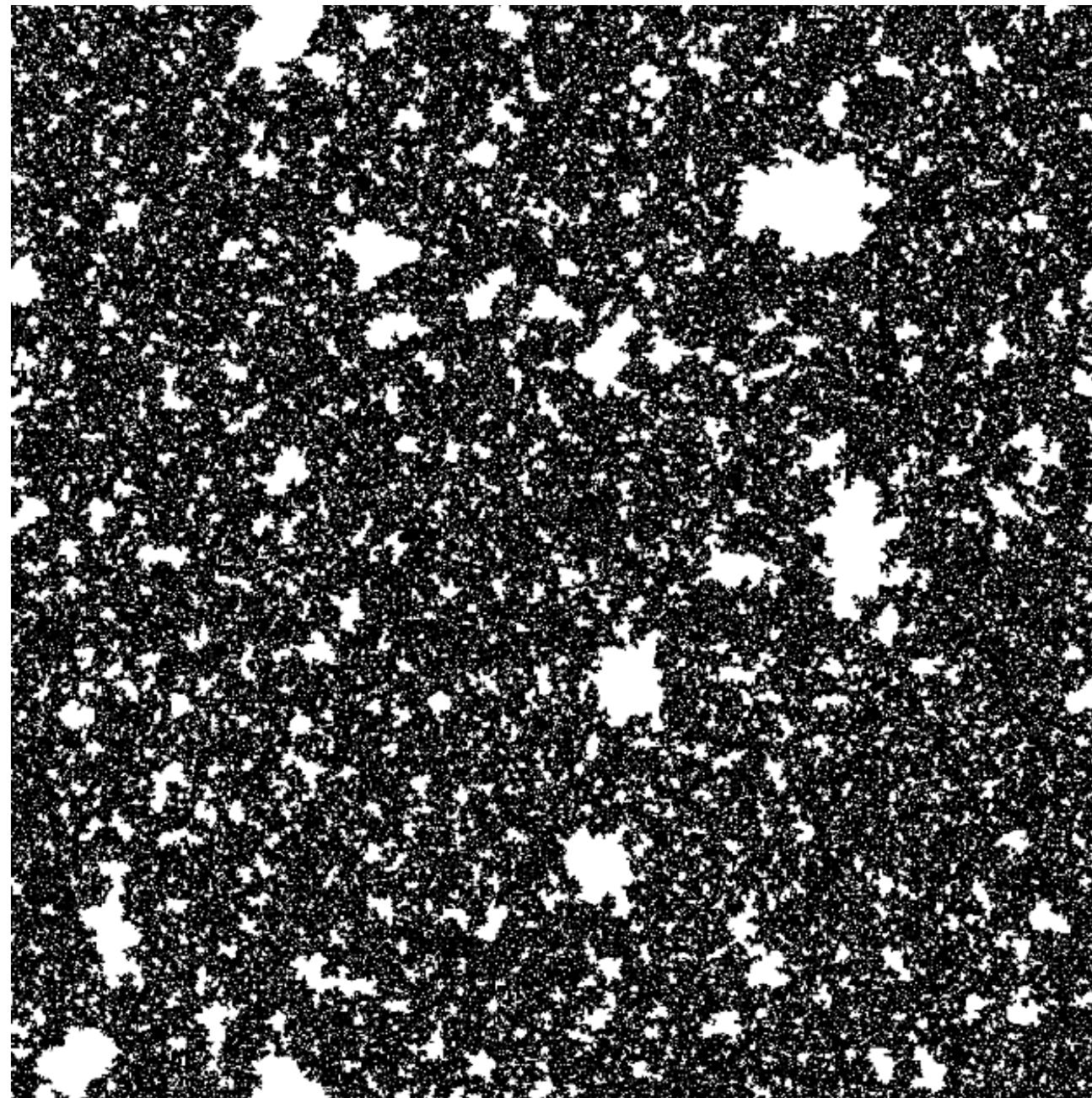
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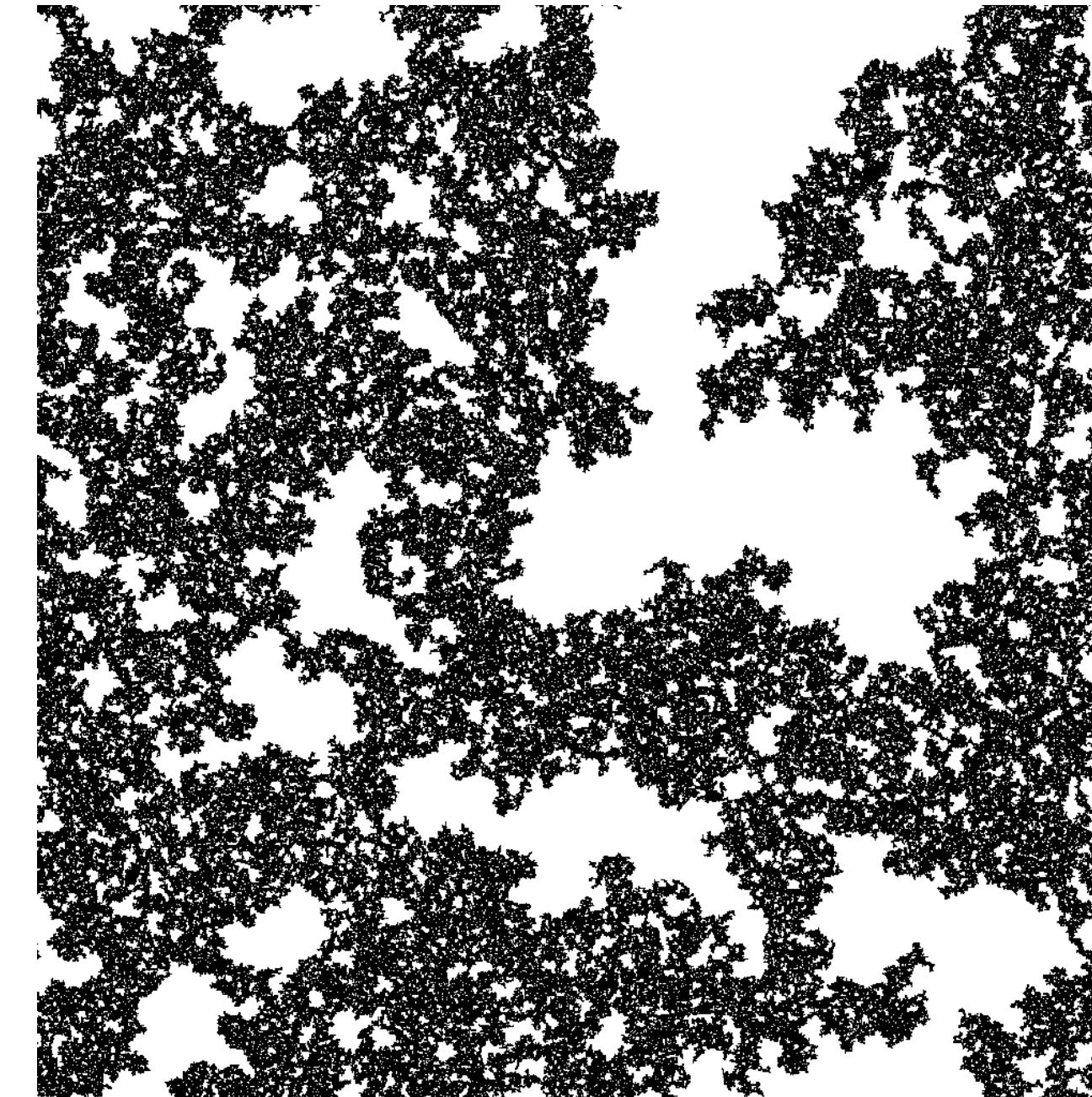
Bonds between nodes are **kept** with probability **p**

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The largest cluster is intact with only small holes.



$$p = p_c = 0.5$$

The biggest cluster barely hangs together, with holes on all length scales.

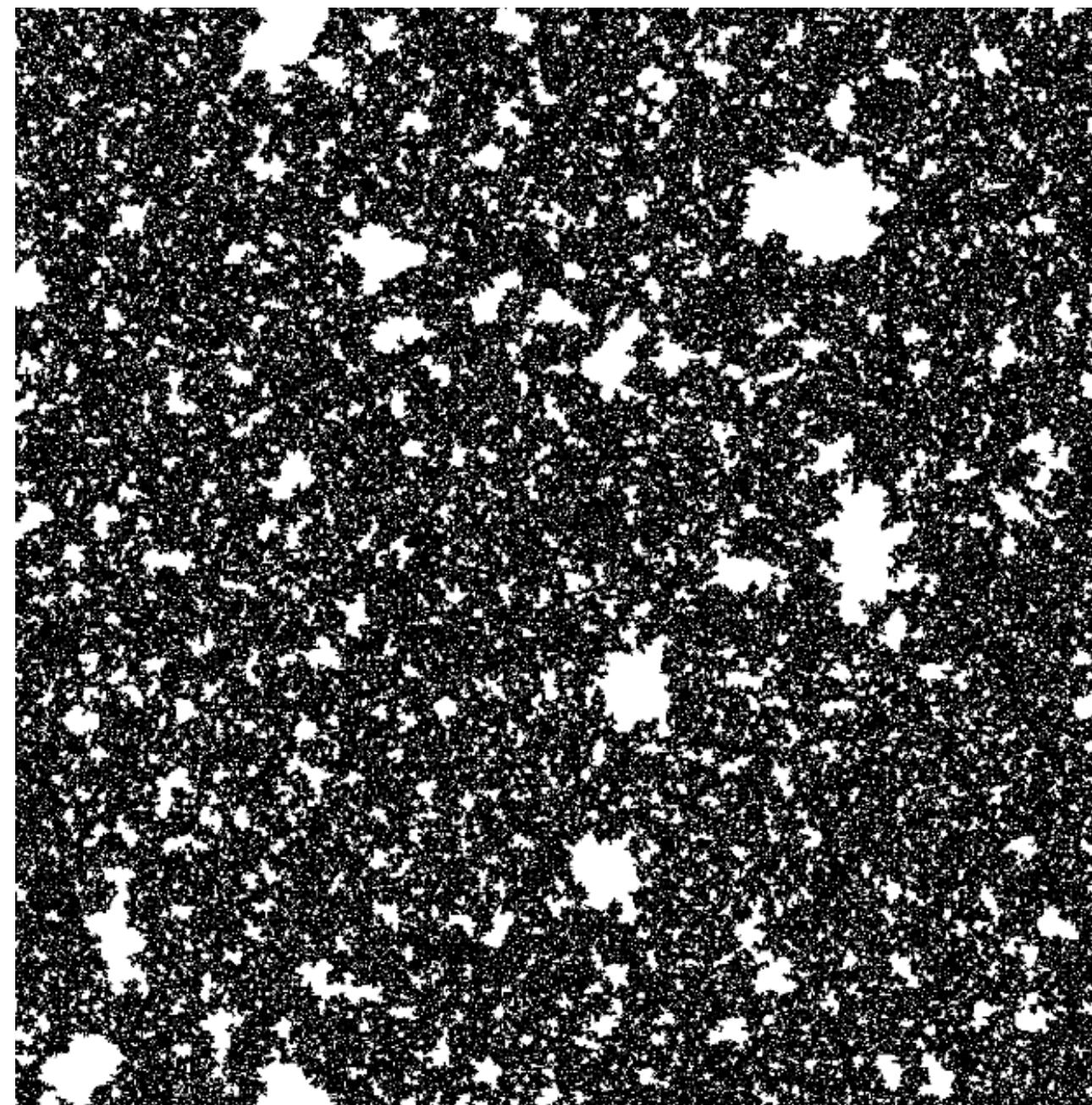
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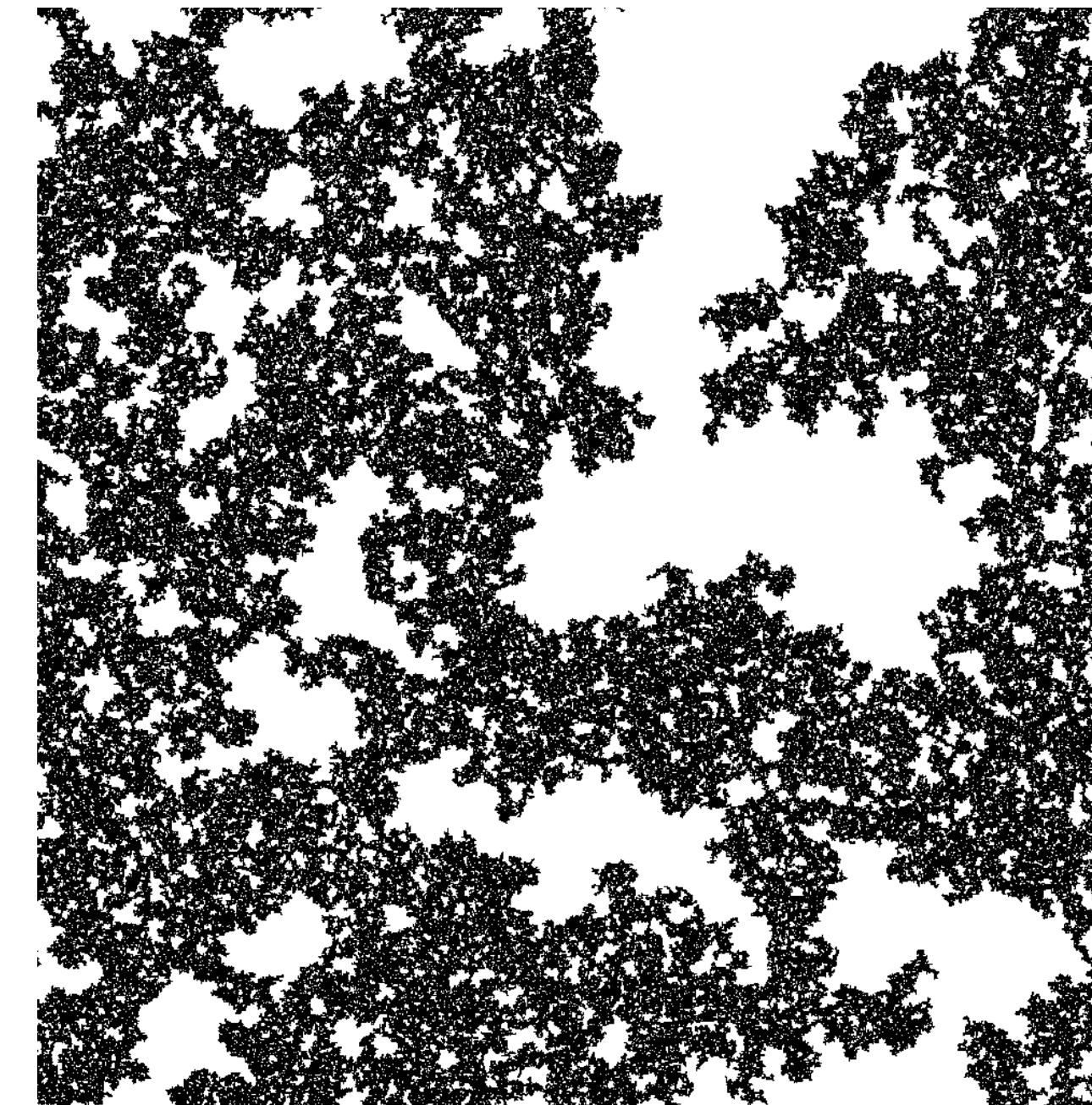
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The sheet falls into small fragments.
(colors indicate clusters)

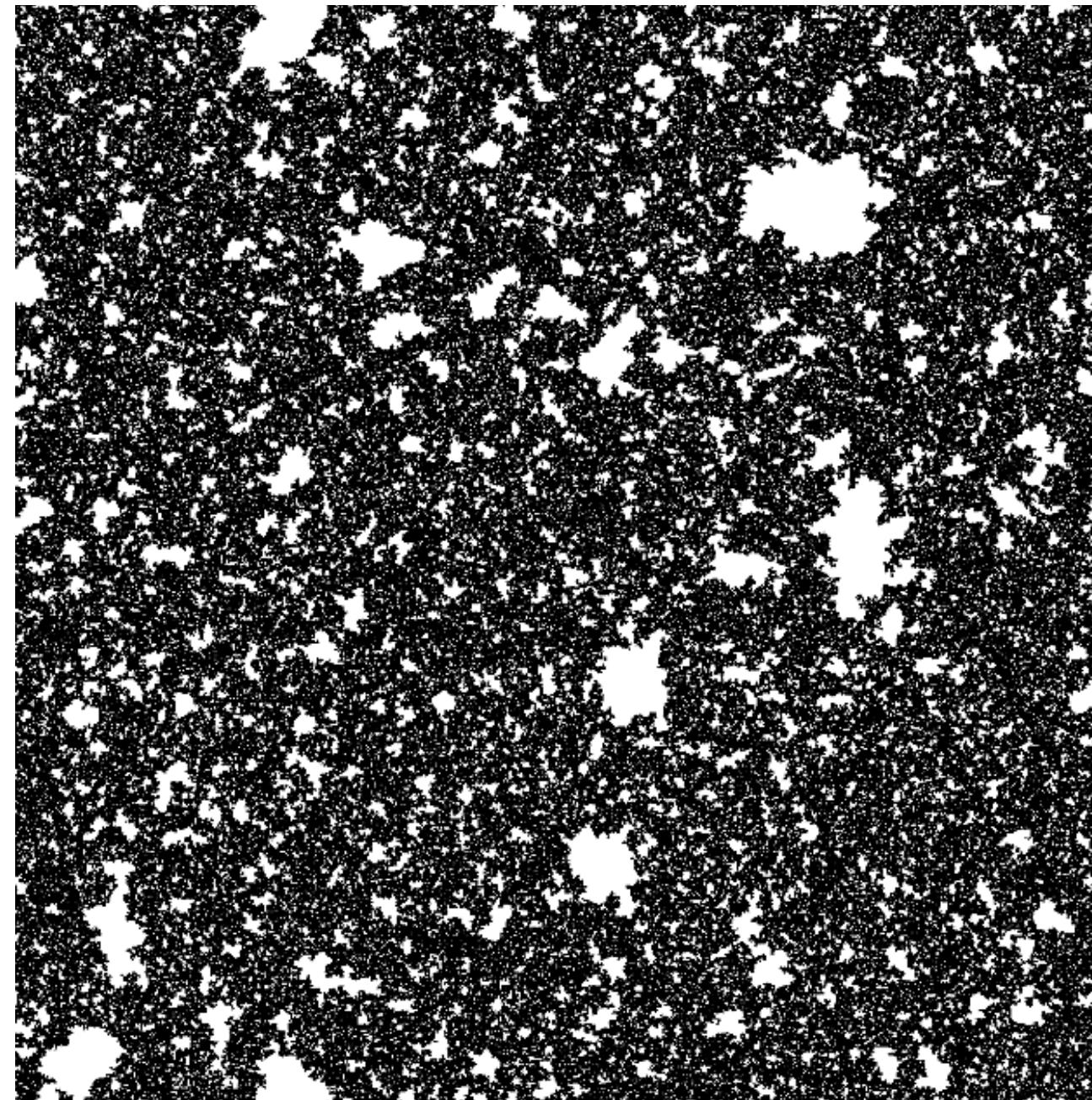
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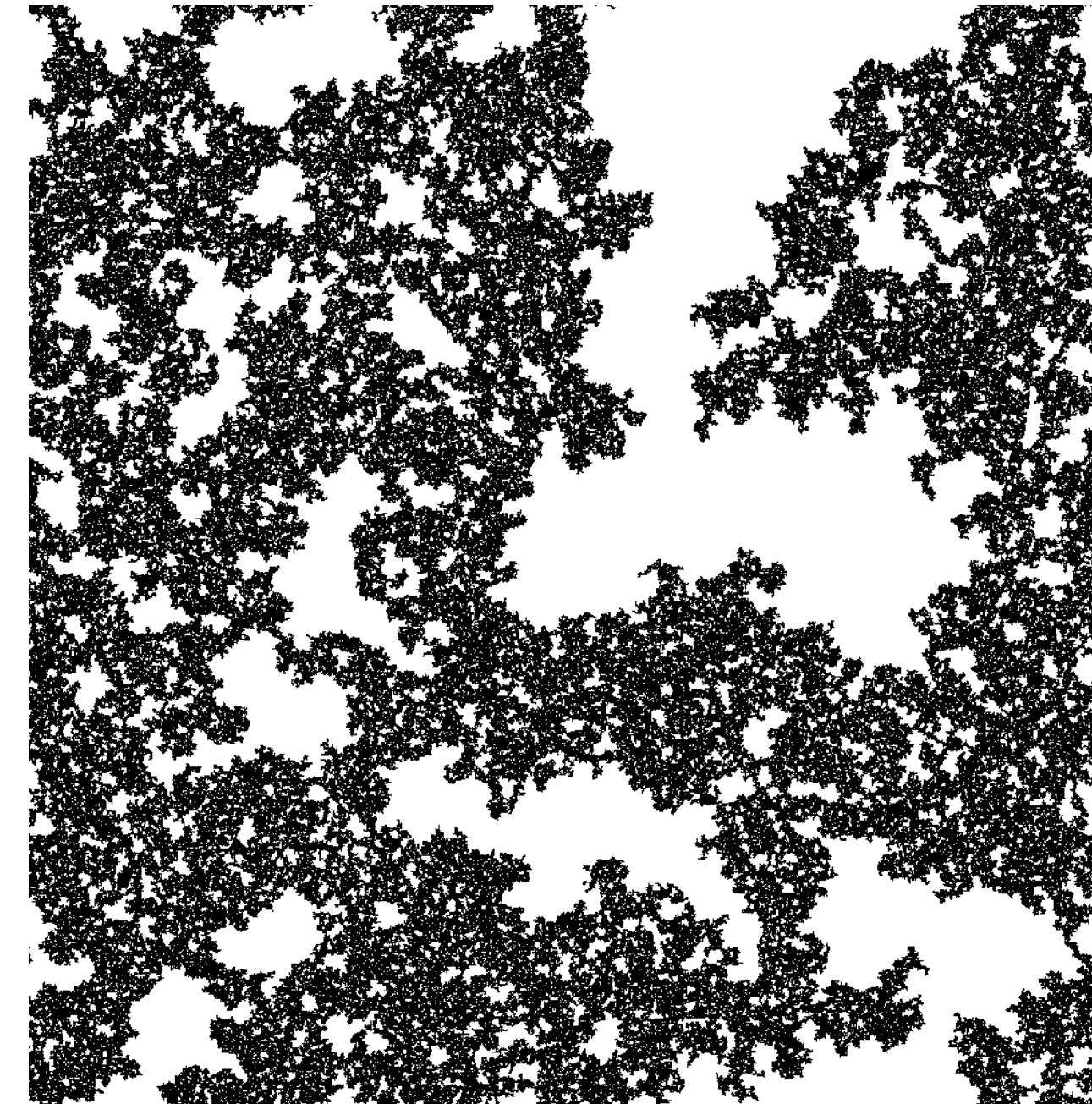
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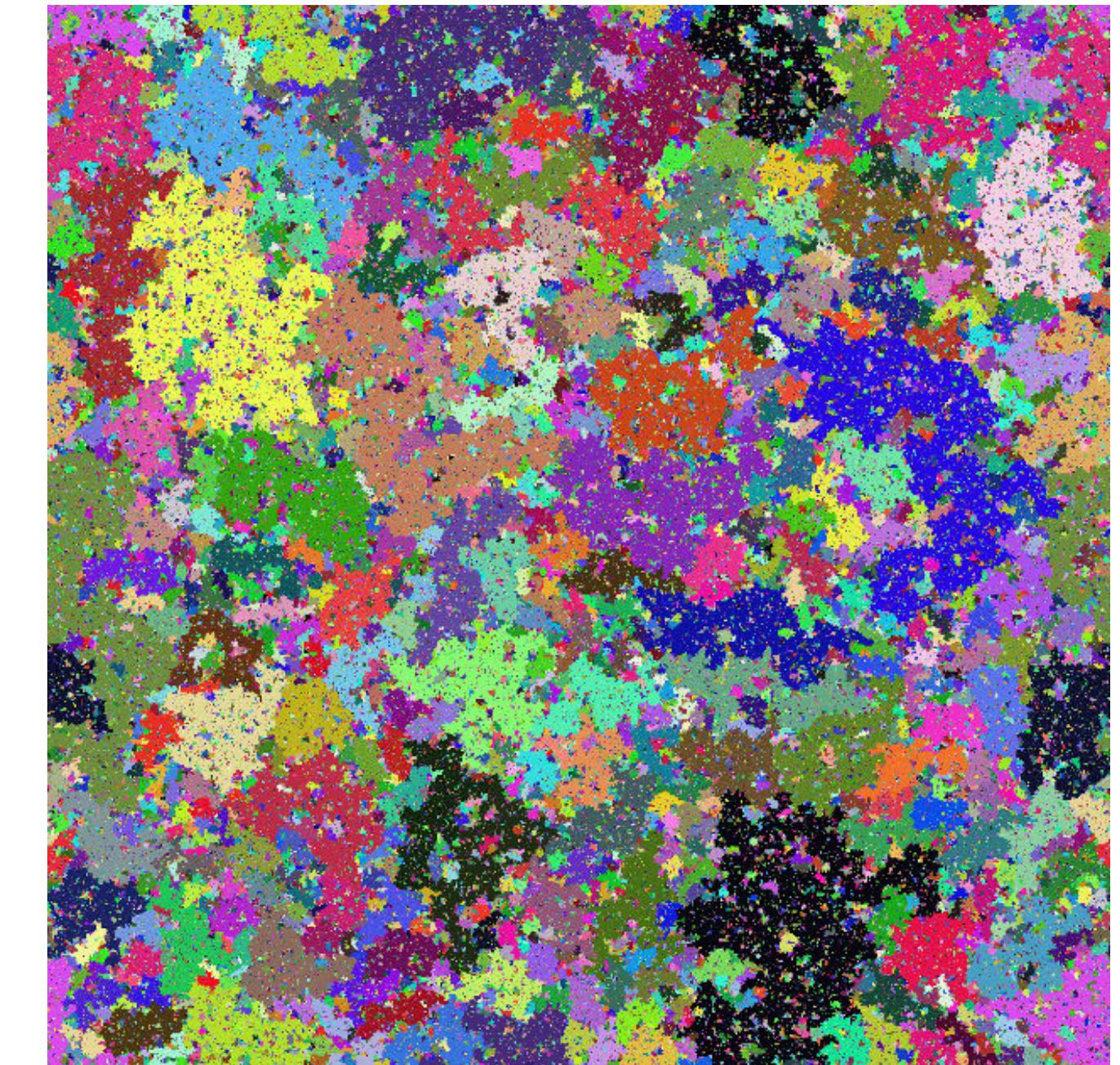
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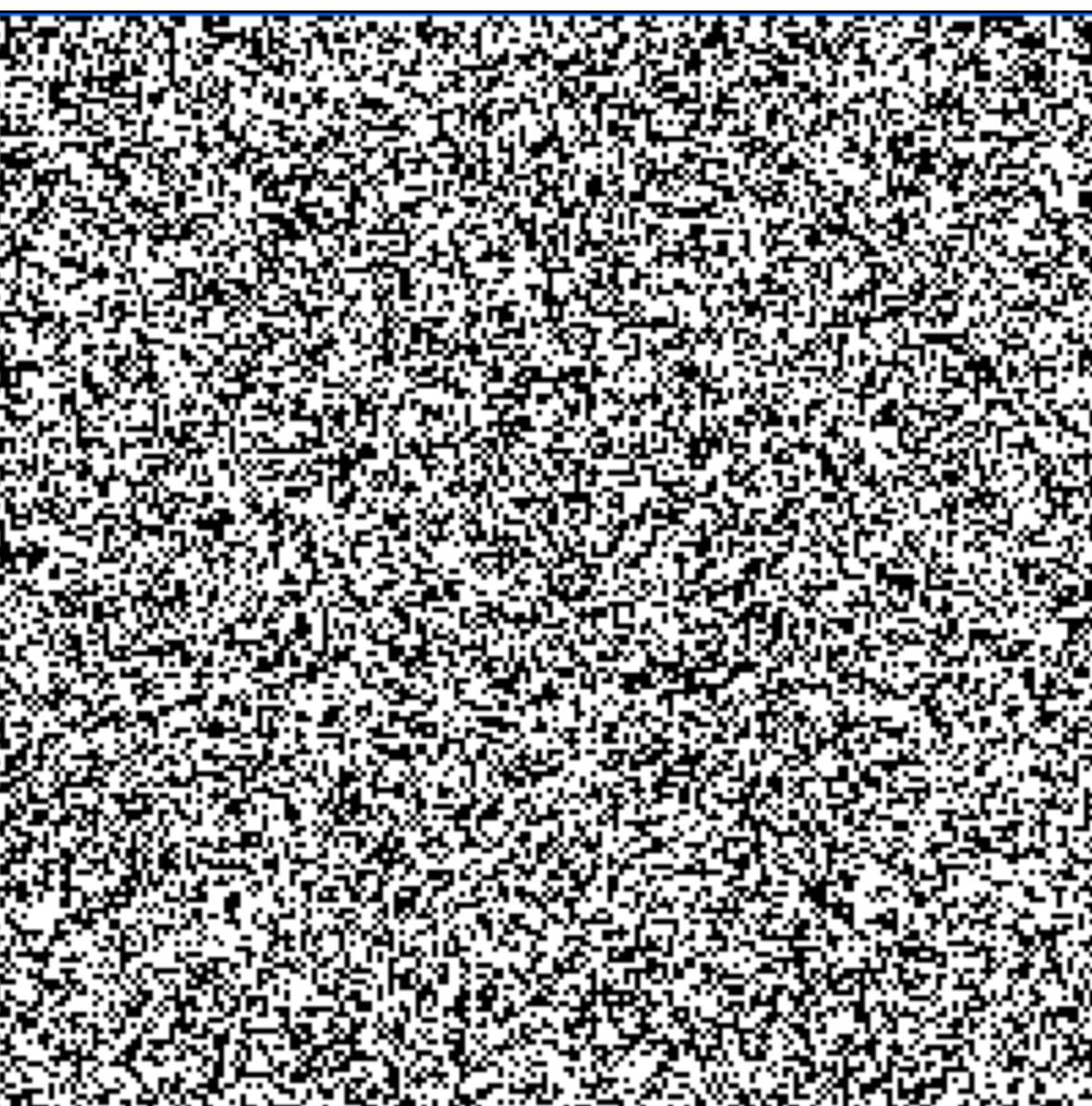
Percolation: transition between a **connected phase** and a **fragmented/disconnected phase**.

Multiple applications

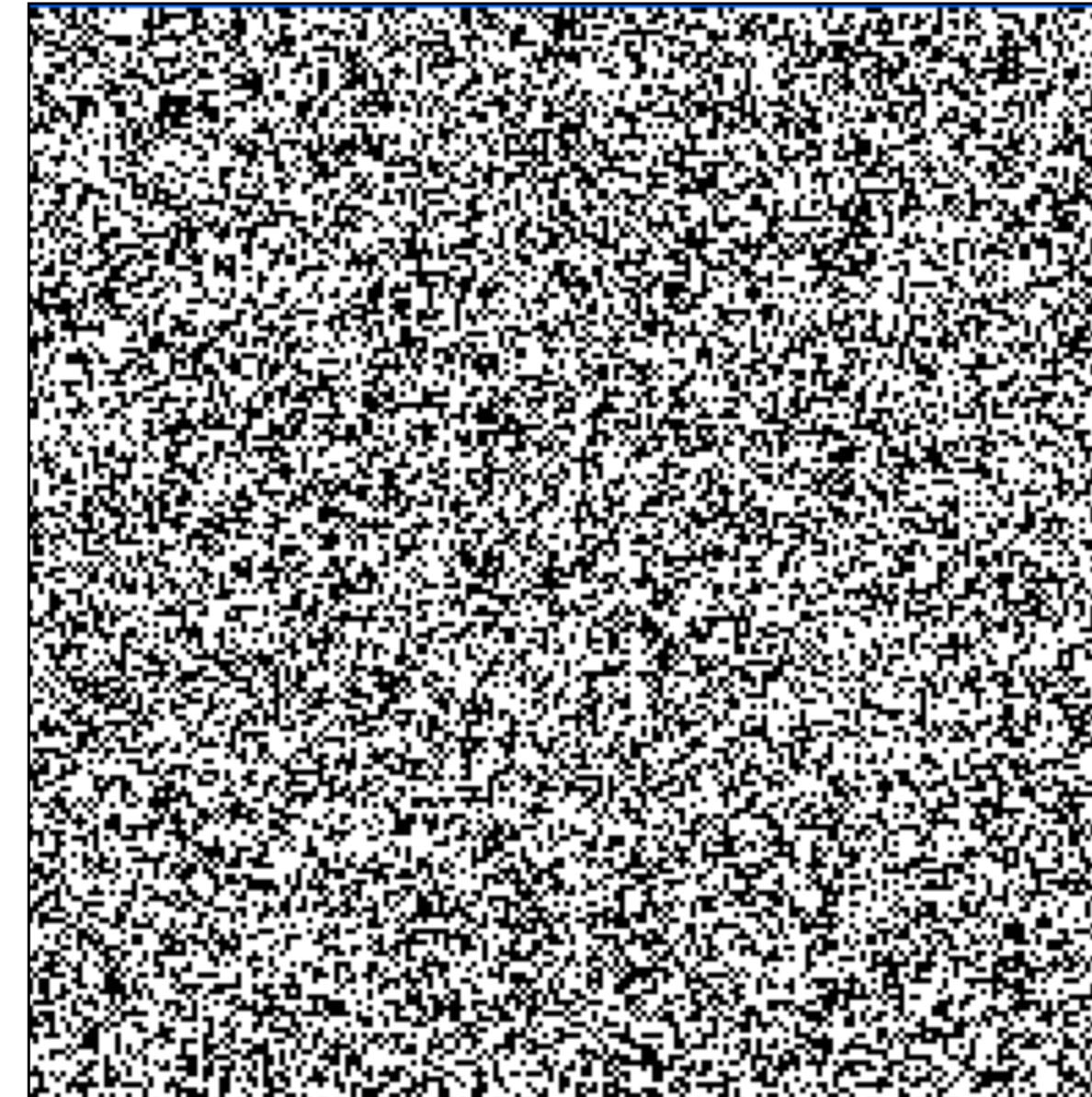
Ex. Oil recovery from a porous soil, making coffee (See simulations [here](#))

Take a **porous medium** (ground coffee, soil): **Can the liquid can make it through the medium?**

How does it depend on the porosity of the medium (i.e. the fraction of material that is empty)?



$$p = 0.59$$



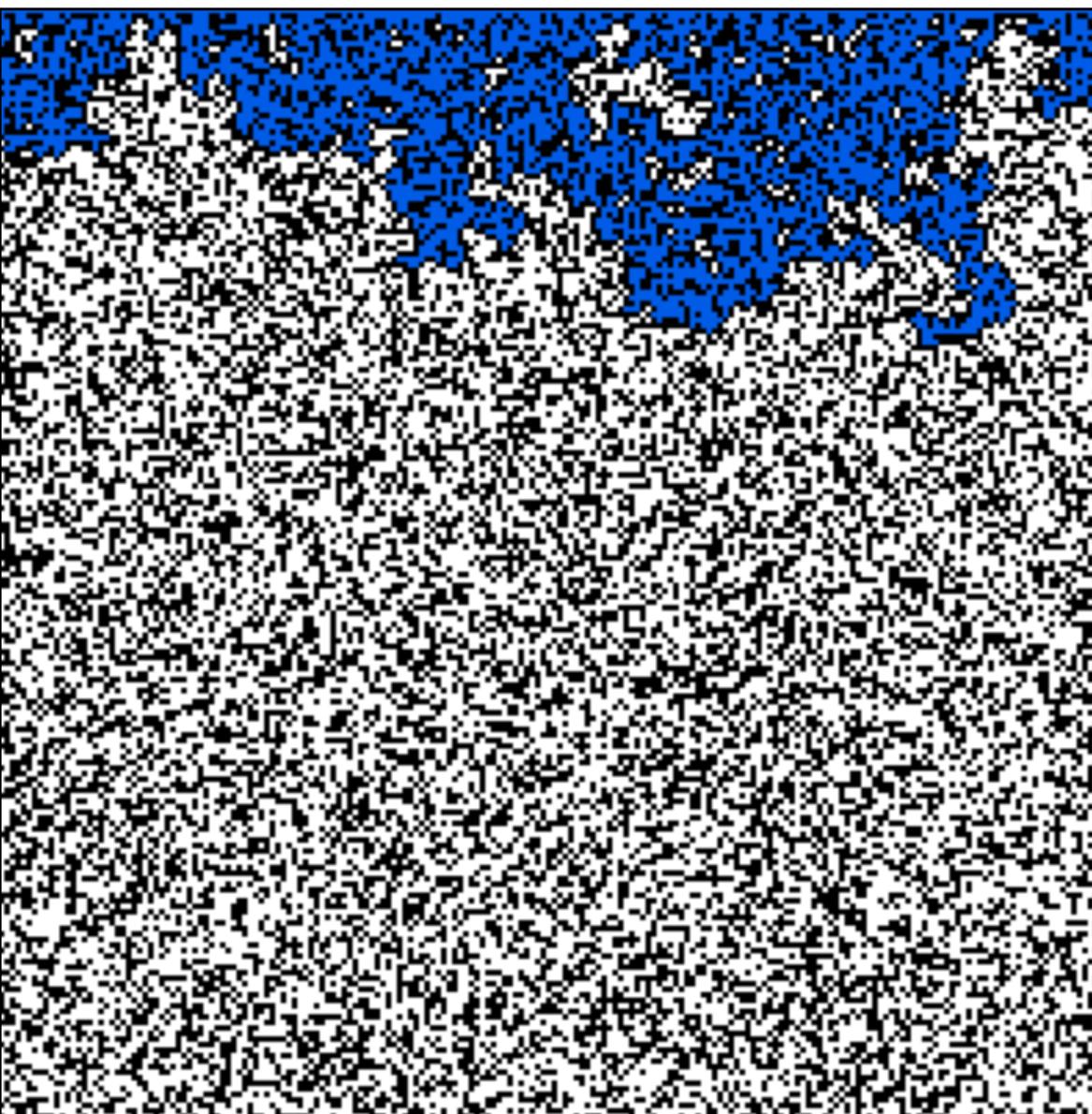
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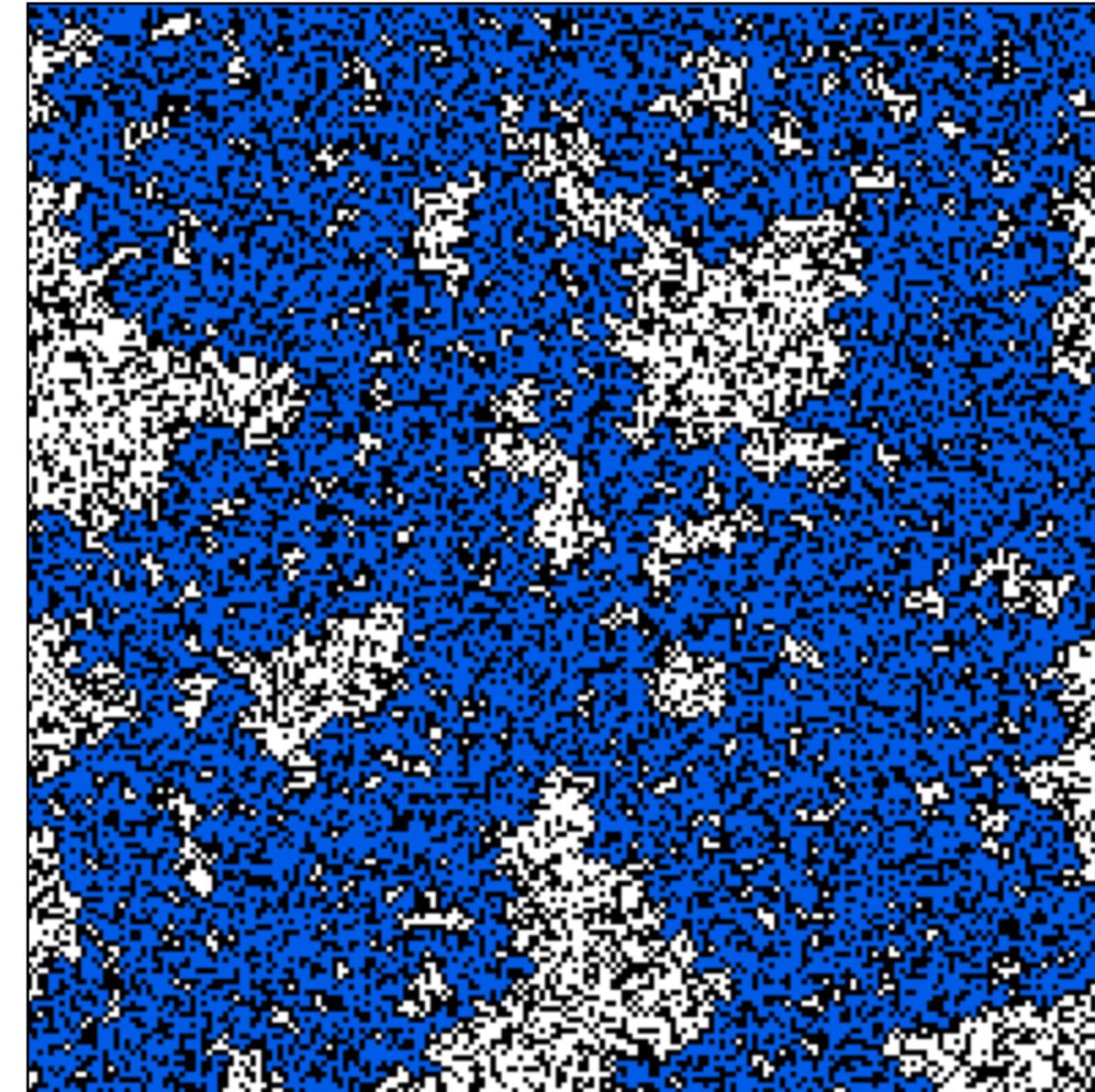
How does it depend on the porosity of the medium (i.e. the fraction of material that is empty)?



$$p = 0.59 < p_c$$

liquid almost never makes it through

(never make it through if the system has infinite size)



$$p = 0.60 > p_c$$

liquid almost always makes it through

Multiple applications

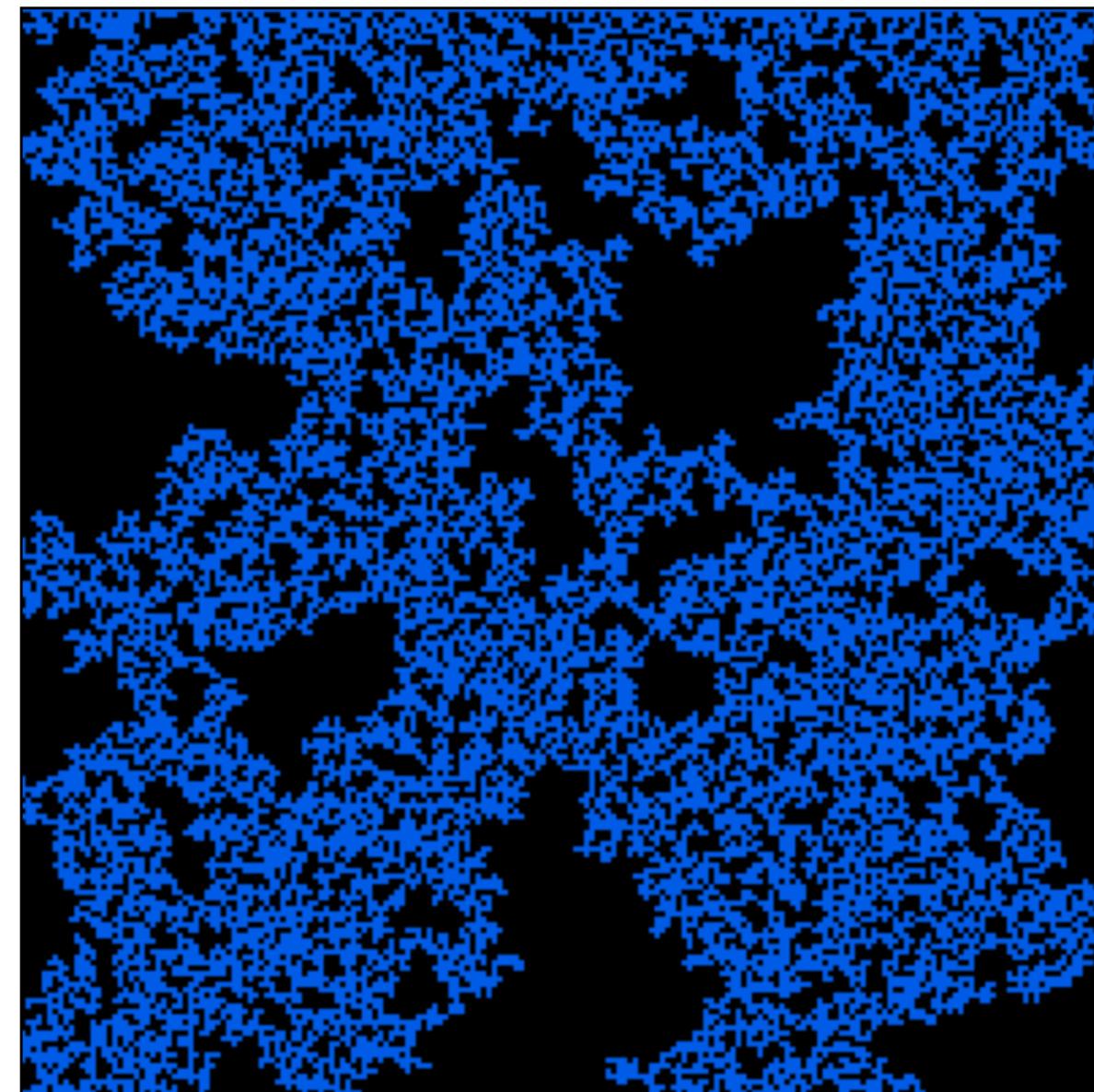
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$$p = 0.60 > p_c$$

Model of formation of gels (networks of chemical bonds spanning the whole system, ex. Boiling an egg become solid-like) —> **Flory-Stockmayer theory**

Forest Fire, Epidemic modeling, metal-insulator transition, networks fragmentation, fractures in rocks and earthquakes 'nucleation', ...

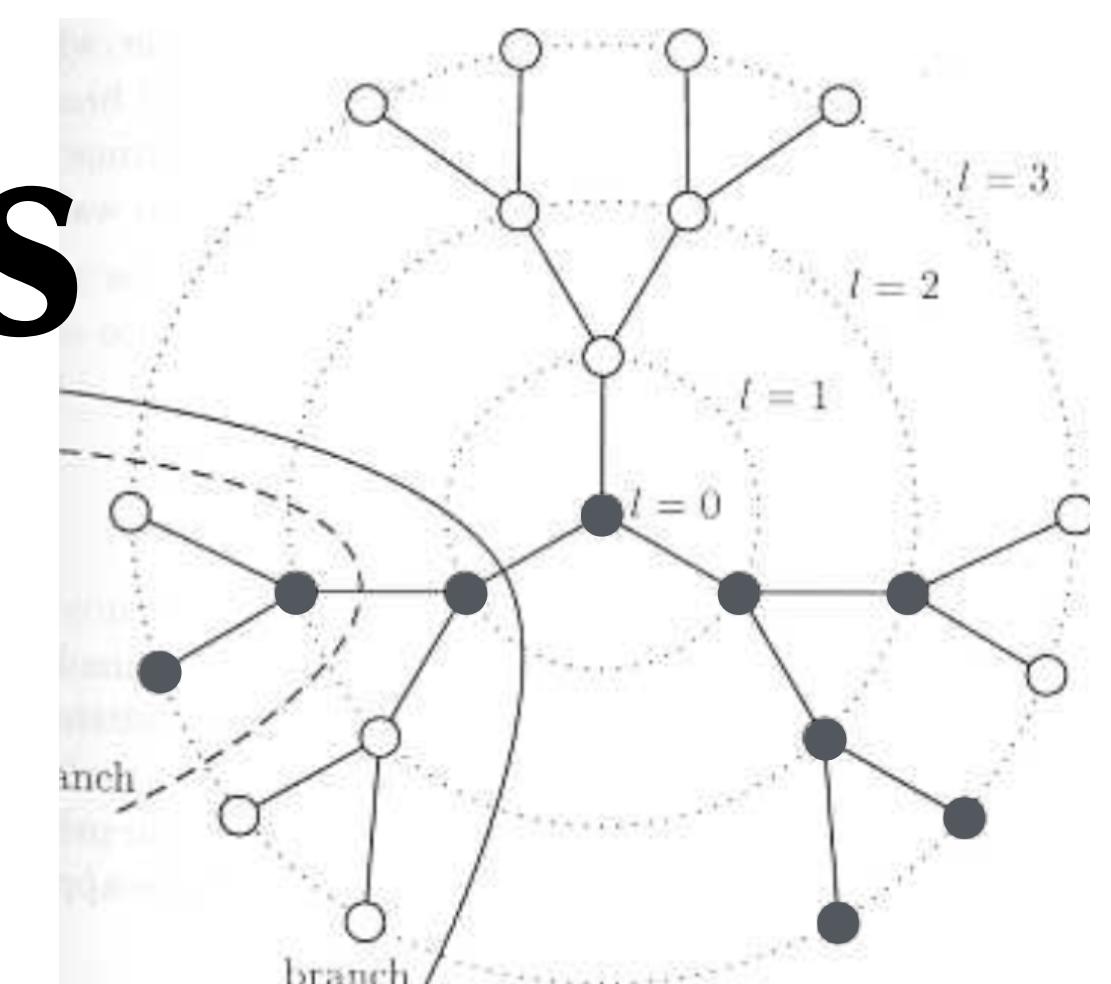
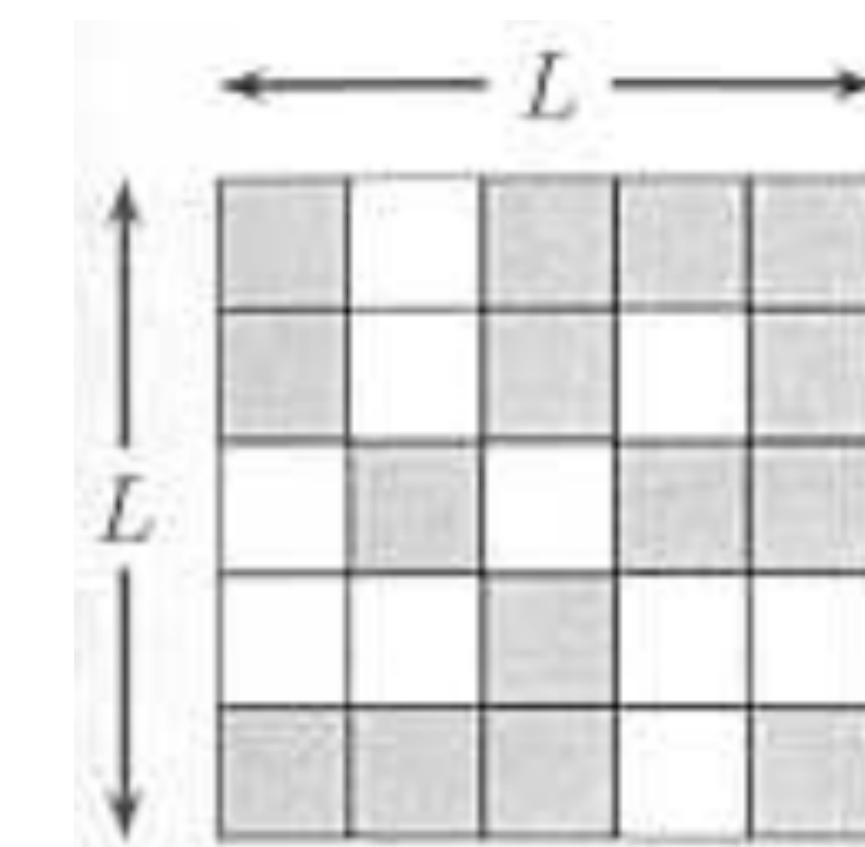
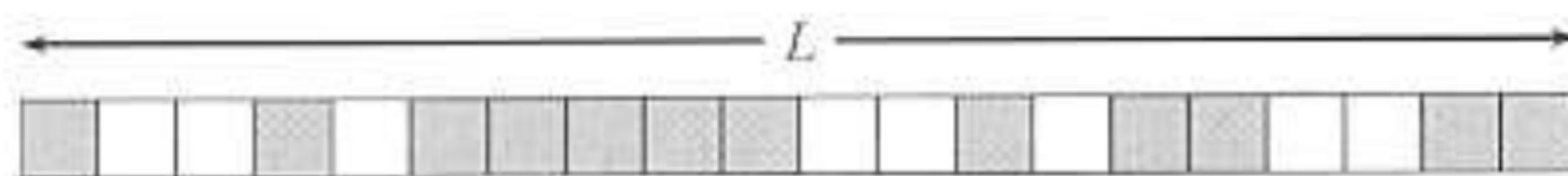
Cluster growth models (e.g., for growth of a cancer cell, a snowflake, or an iceberg)

2) Percolation phase transition

- a) Control parameter and critical value
- b) Order parameter and quantities of interests
- c) 1d site percolation
- d) Bethe Lattice

Theoretical solutions

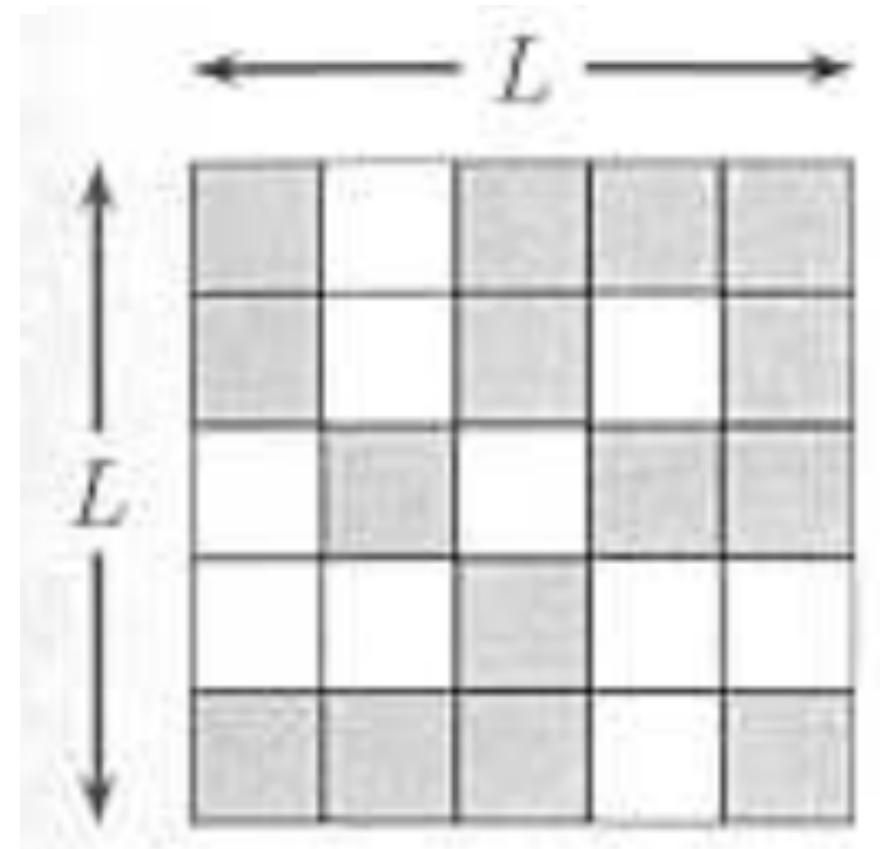
Results on Percolation problems



1) Control parameter and Critical value

Working model: Site percolation on a lattice

Ex. 2D lattice of $L \times L$ sites

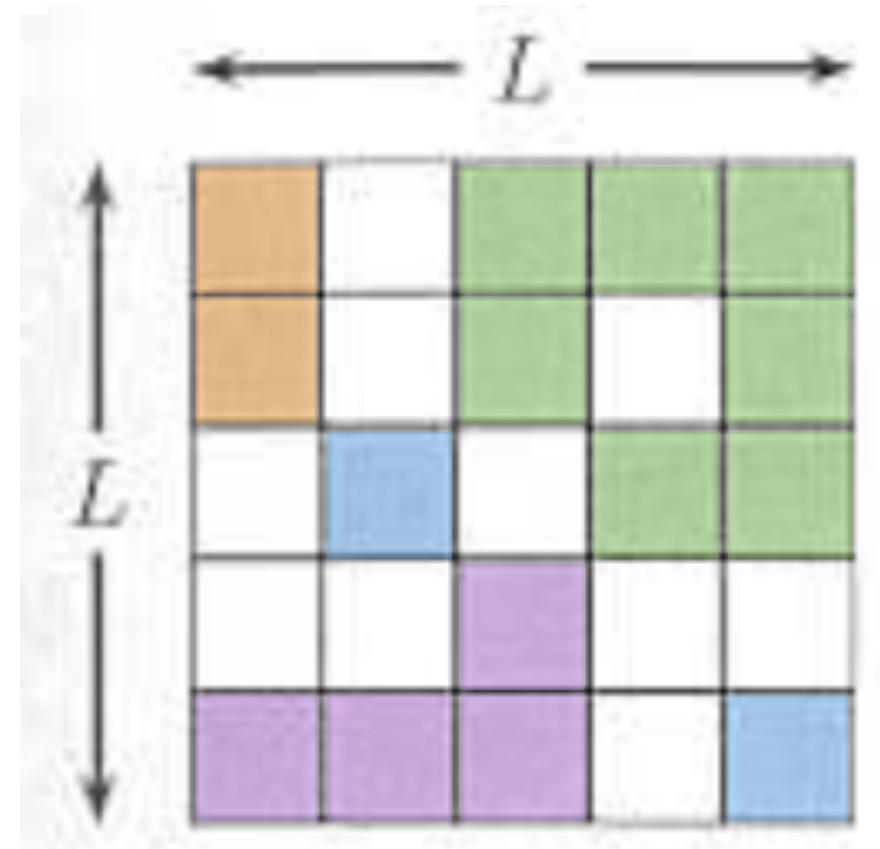


p = occupation probability = probability that a site is occupied

Average number of sites that are occupied = $L \times L \times p$

Working model: Site percolation on a lattice

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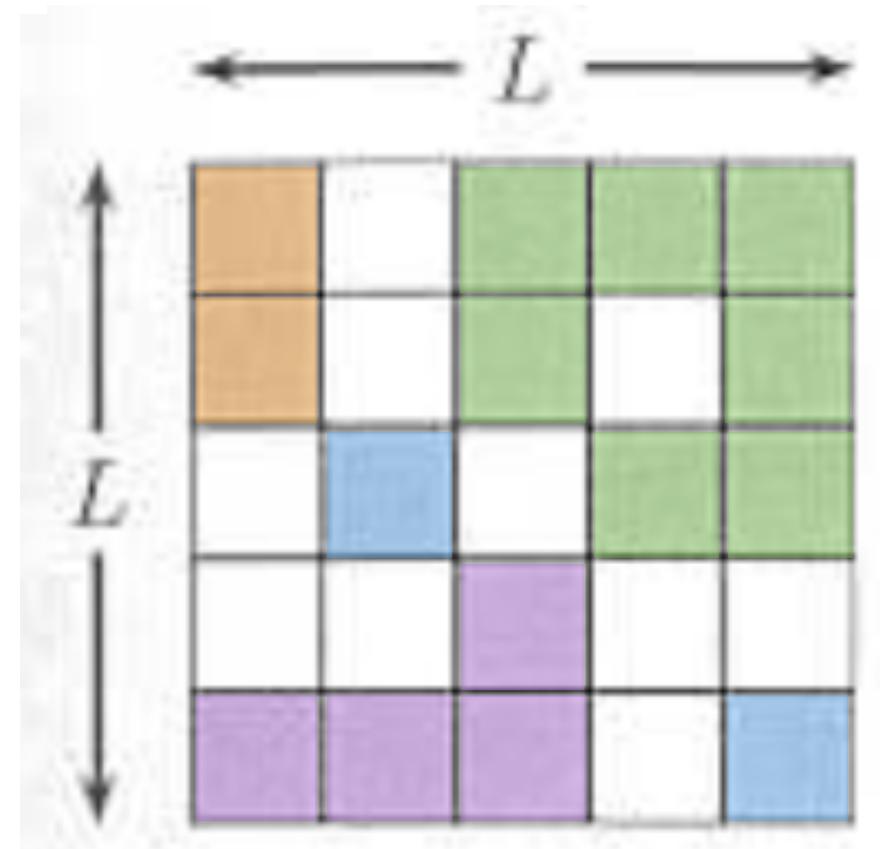
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Percolation theory: studies the **properties of the clusters** formed

cluster = group of nearest-neighboring occupied sites

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Cluster is percolating = **Cluster is infinite!**

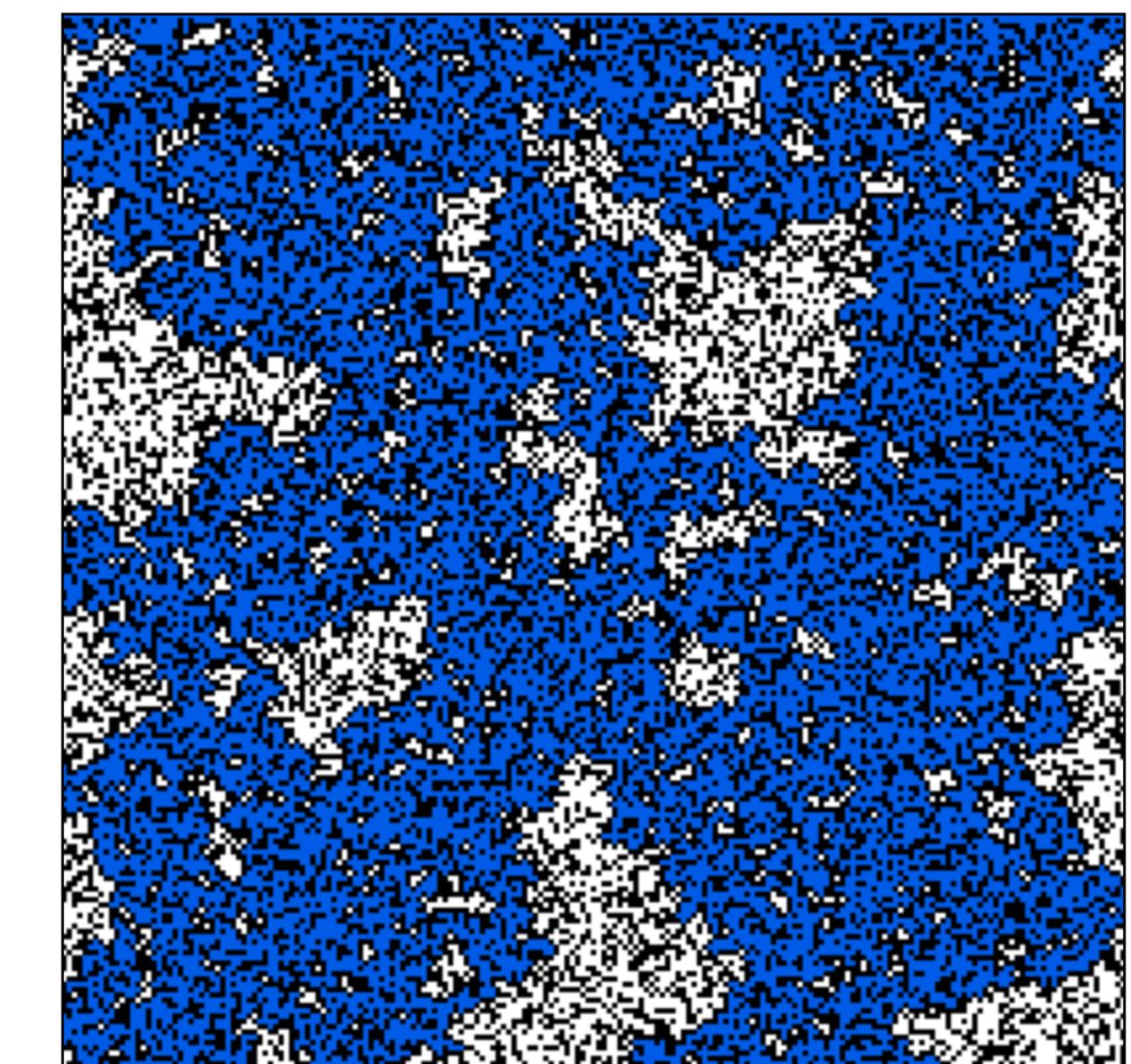


For finite size systems:

Clusters that span the entire lattice

from left to right boundaries
or from top to bottom boundaries

are good candidates of percolating clusters

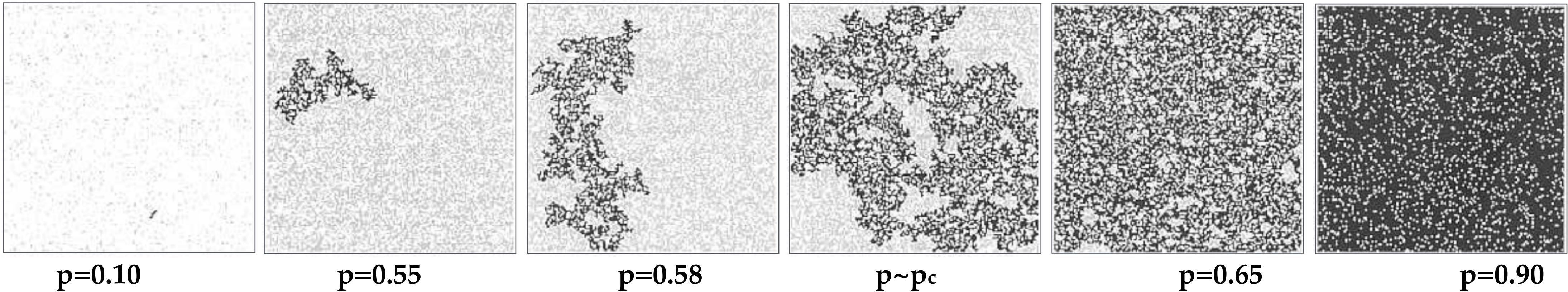


Control parameter

Expected behavior:

- For $p=0$: the lattice is empty
- For intermediate p : each realization will be different, but we expect the size of the largest cluster to increase with p .
- For $p=1$: there is only 1 cluster of size L^2

$L=150$

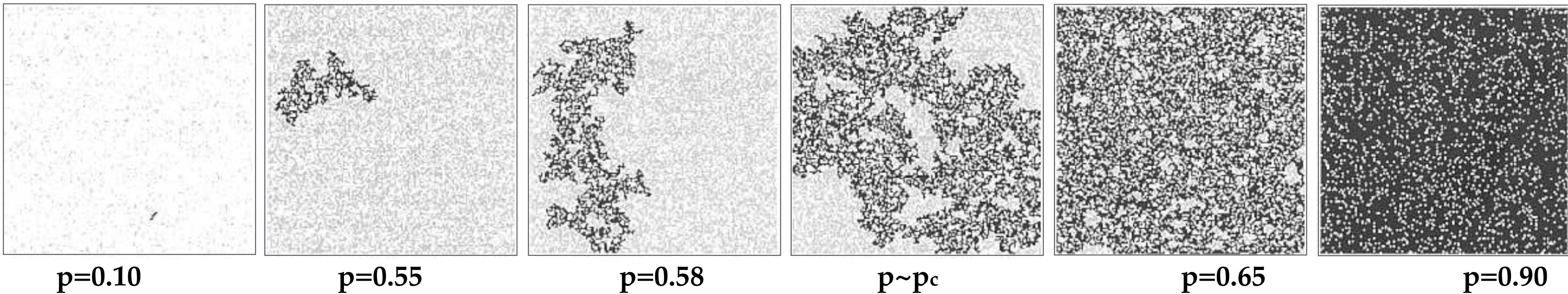


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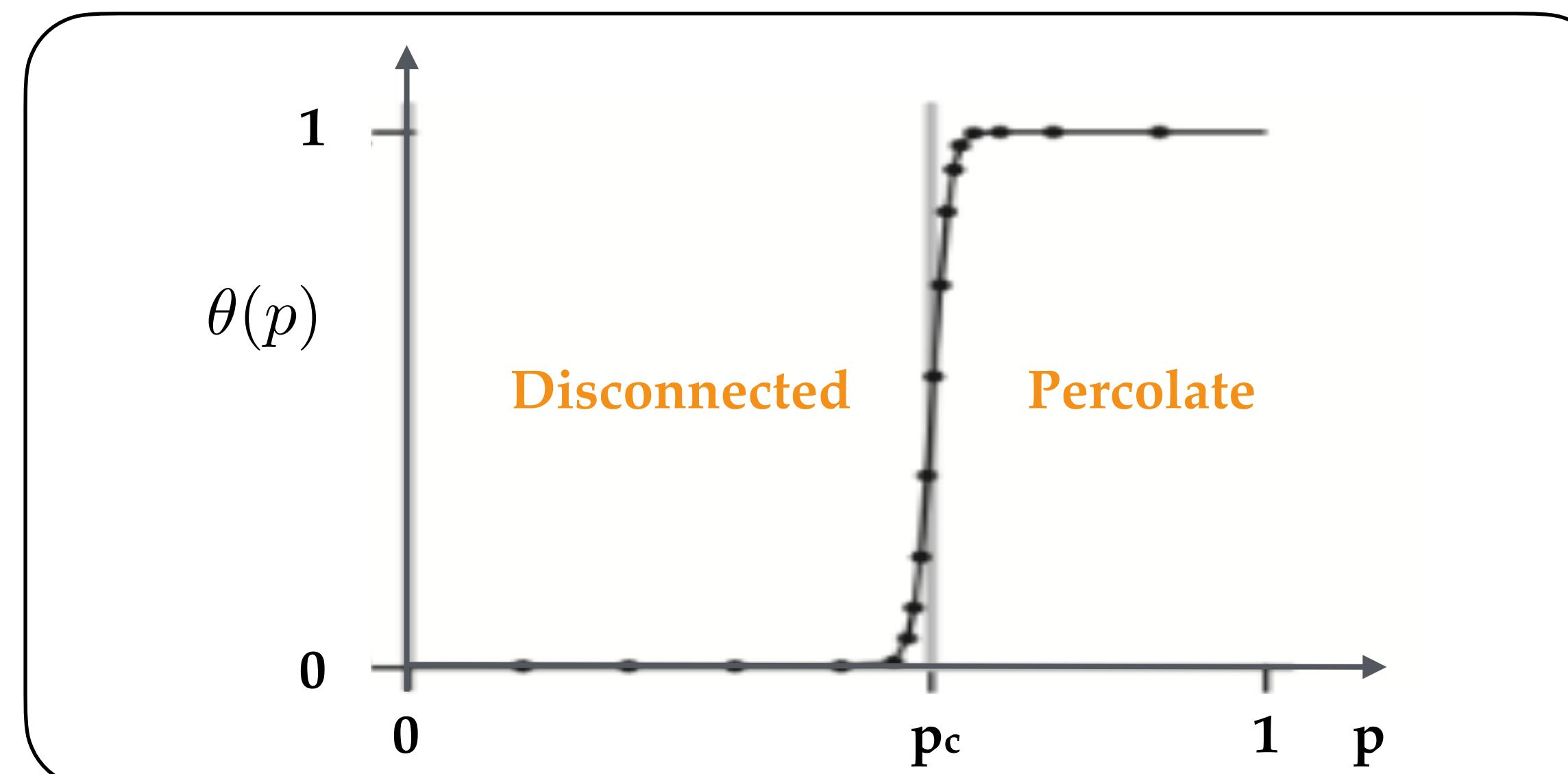
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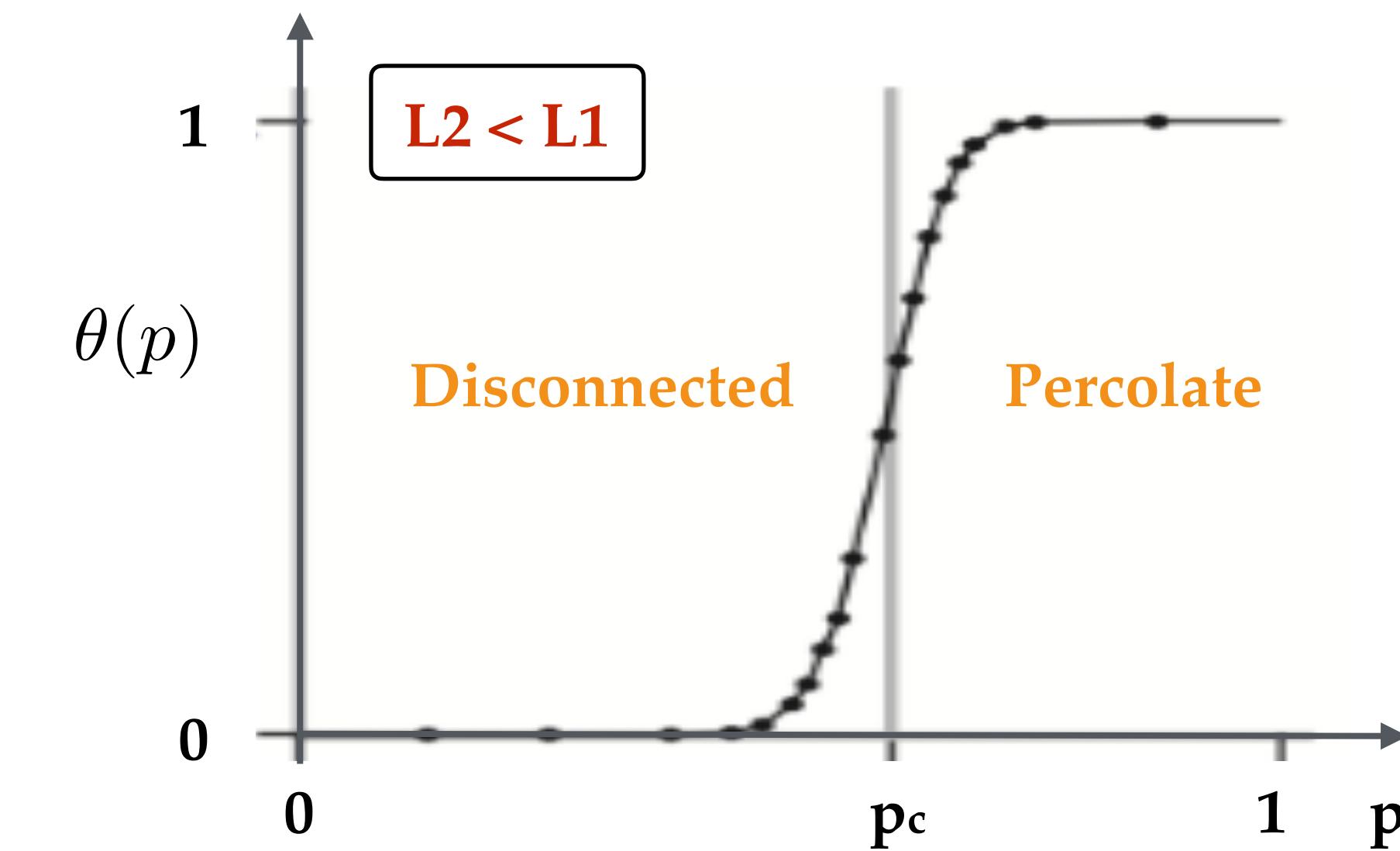
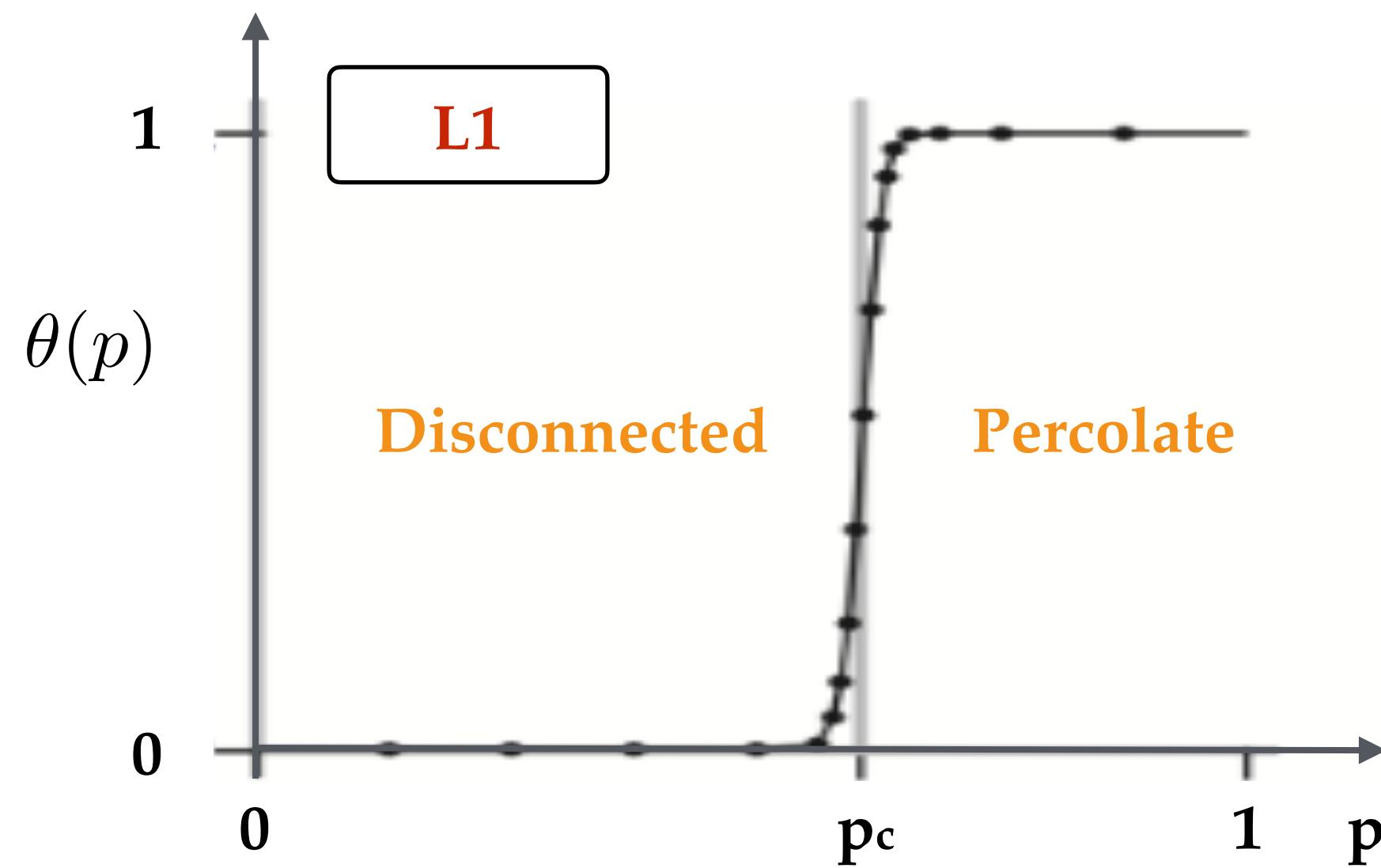


$\theta(p)$ = **percolation probability**
= probability that there is
a cluster of infinite size



Critical parameter

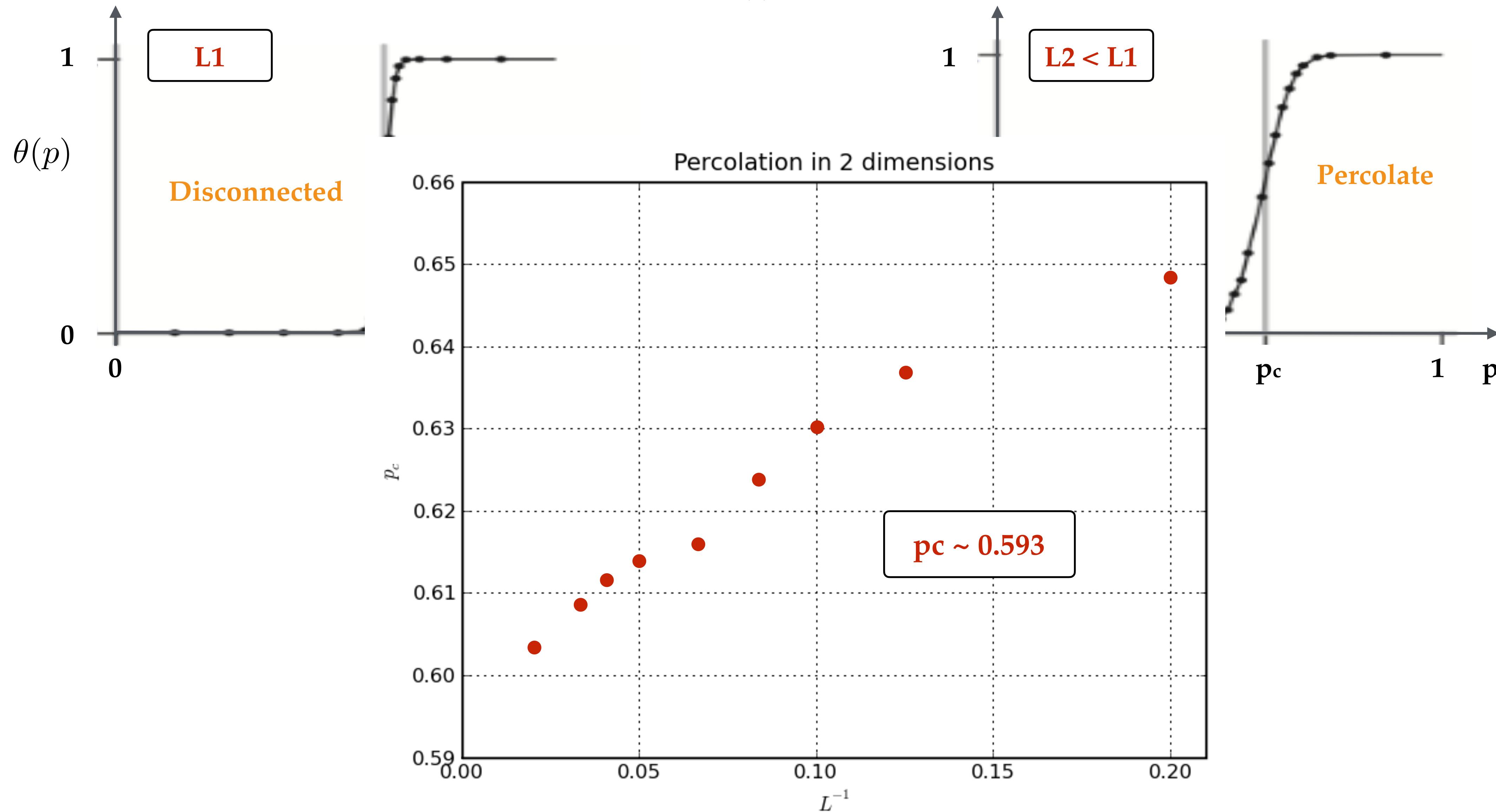
Critical parameter



Larger L: the curve increase at p_c becomes sharper

Critical parameter

Critical parameter



2) Parameters, and interesting quantities for the problem

Percolation Parameters

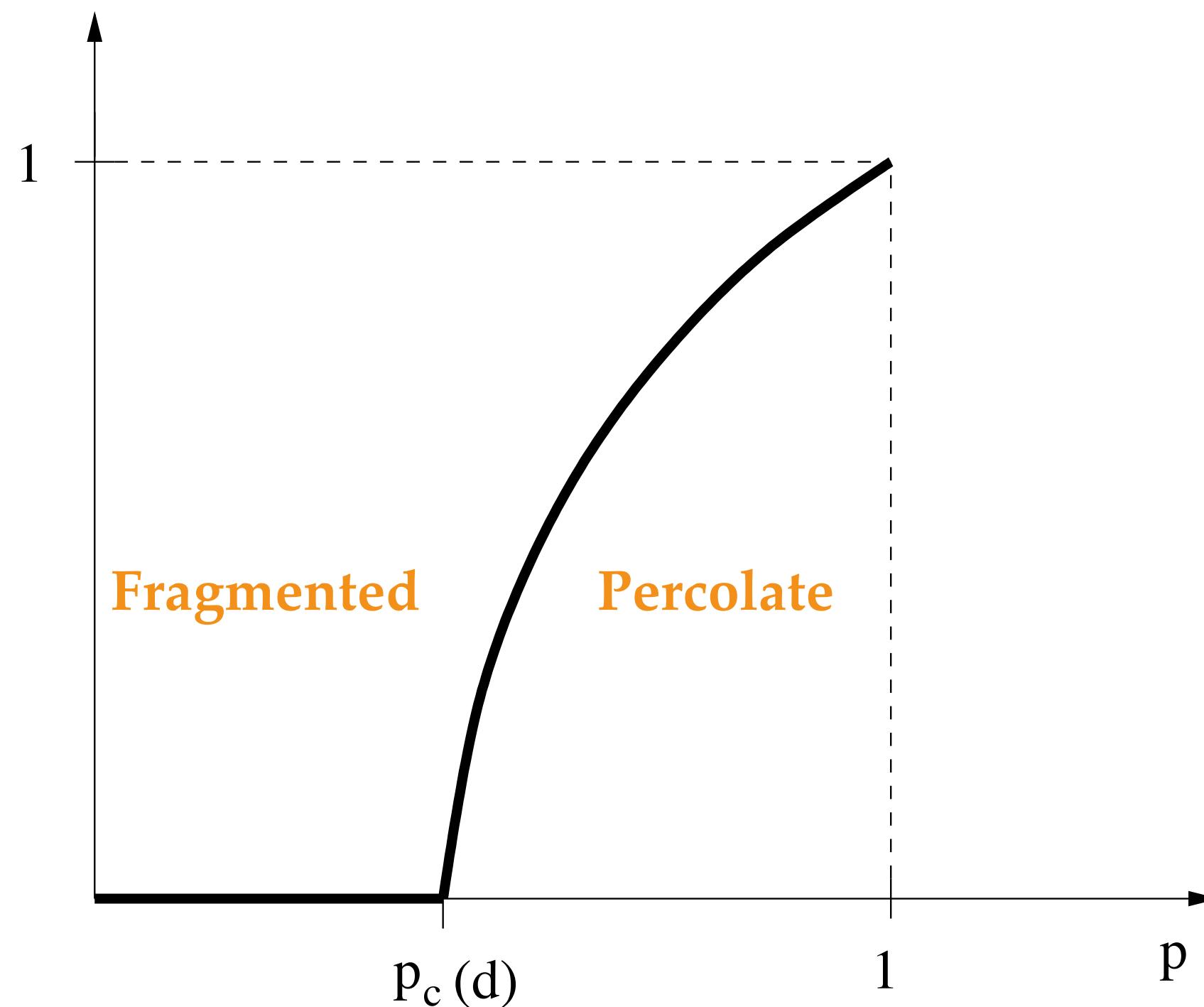
Control parameters:

Order parameter:

Percolation Parameters

Control parameters: p = **occupation probability** = probability that a site is occupied

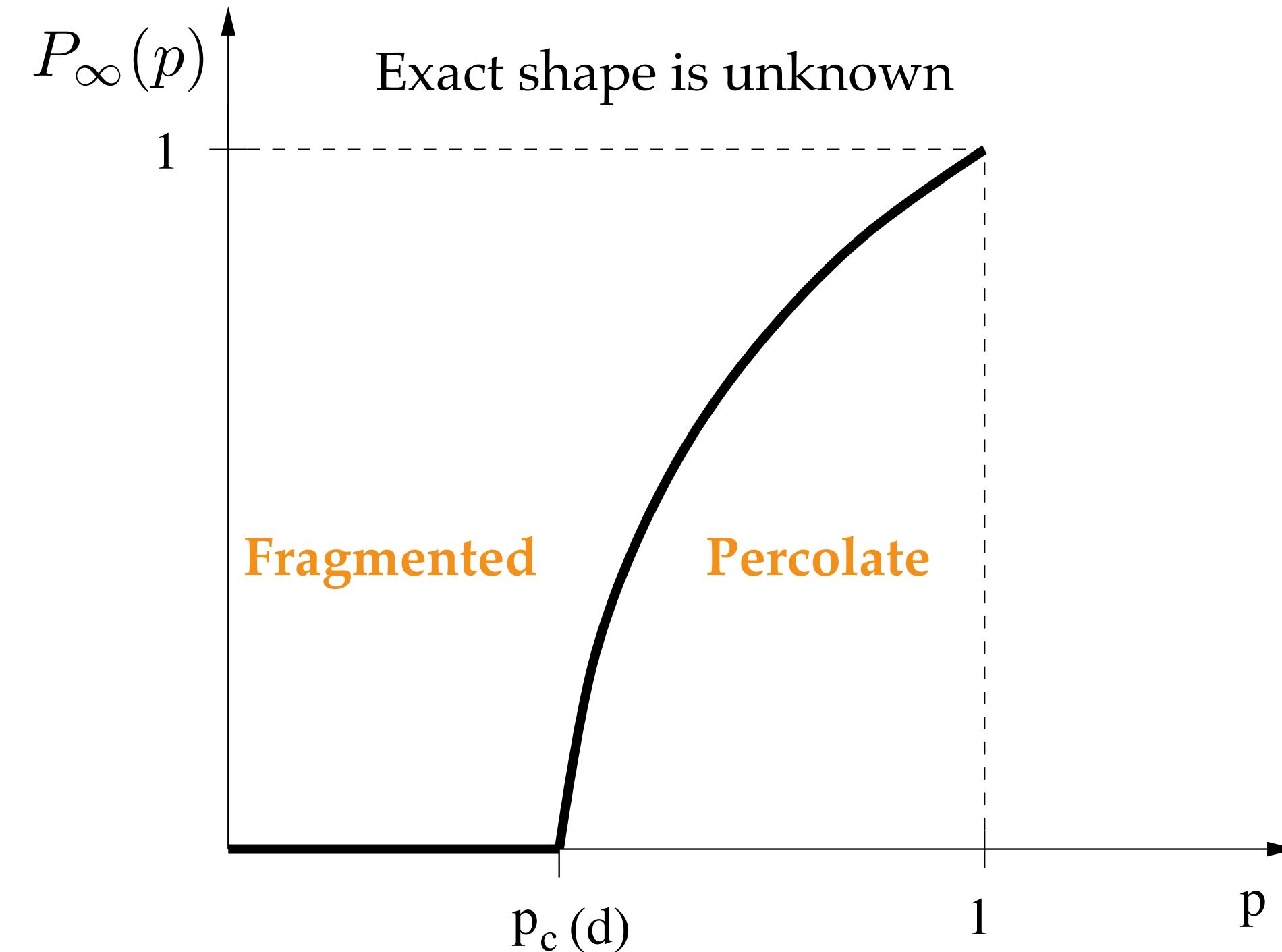
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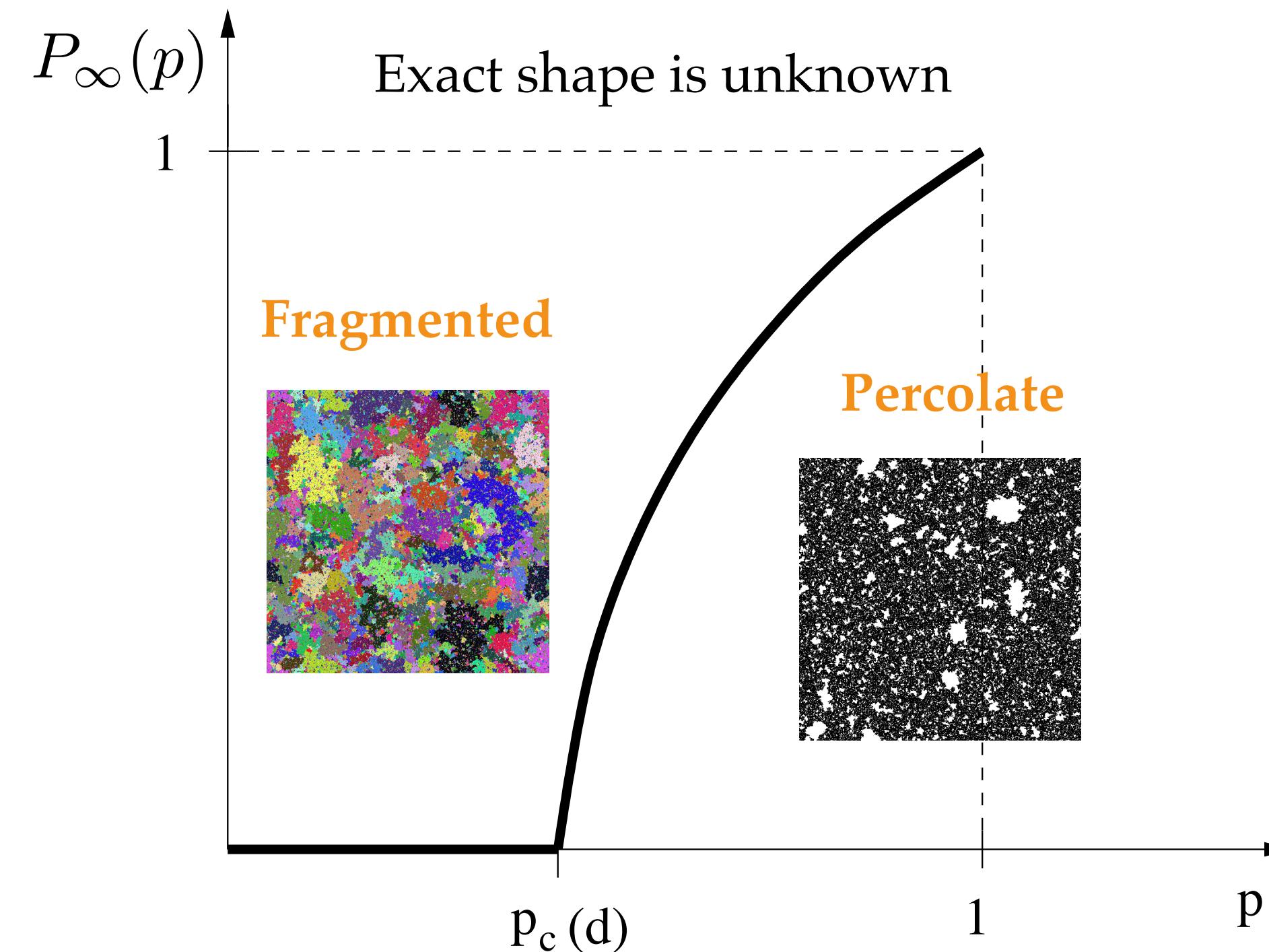
Order parameter: $P_\infty(p)$ = probability that any given site belongs to the percolating cluster



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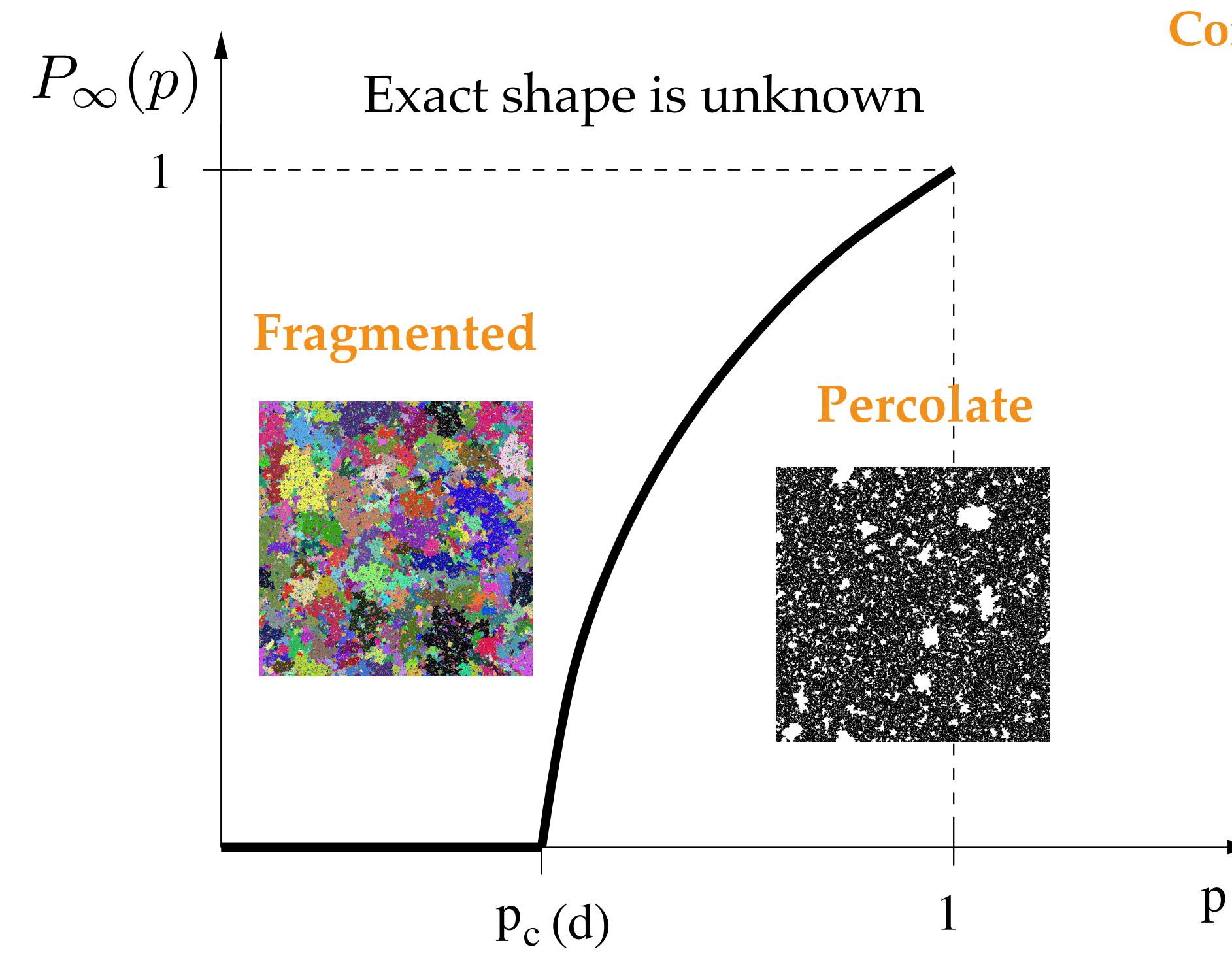
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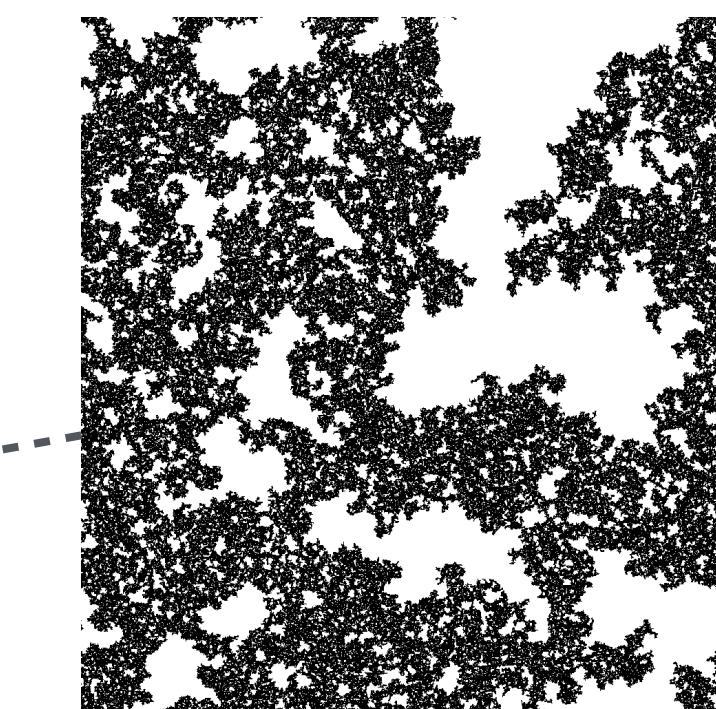


Continuous phase transition:

The fact that $P_\infty(p)$ is continuous at p_c is non-trivial

How to see this intuitively? **At the critical point:**

- the percolating cluster is “very fragile”, removing a few occupied site may break percolation
- the percolating cluster is very tortuous, very “thin”; the fraction of sites that belongs to the percolating cluster is very small



Interesting quantities

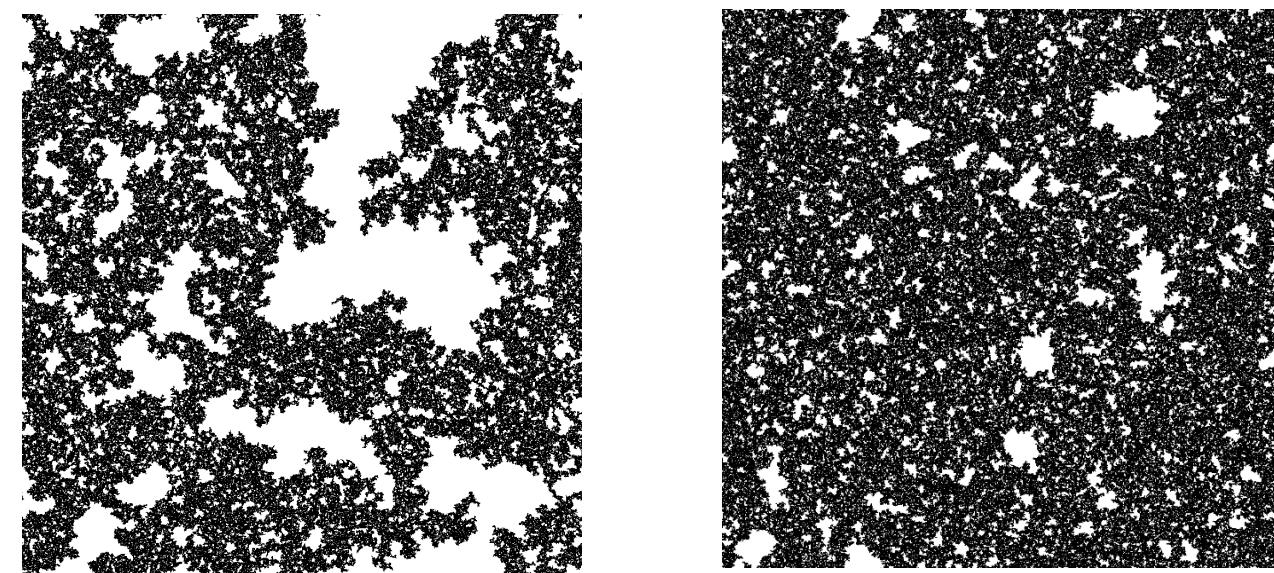
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General questions:

Size, Infinite Cluster: How much space does the percolating cluster occupy?

= What is the **probability that a site belong to the percolating cluster?** $P_\infty(p)$



Finite clusters:



Interesting quantities

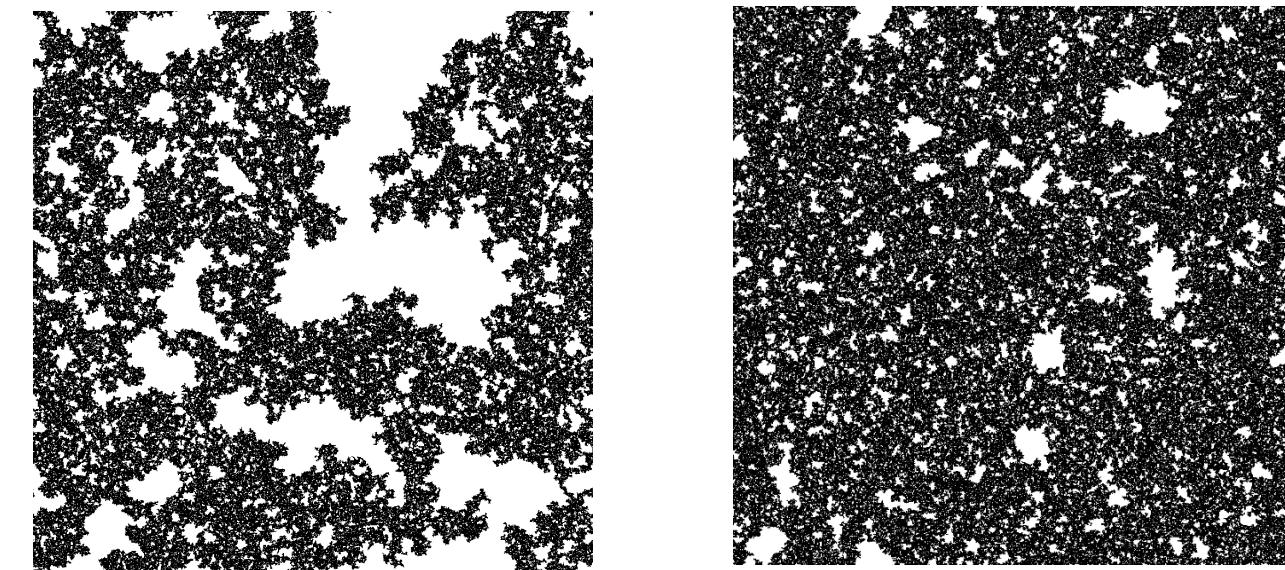
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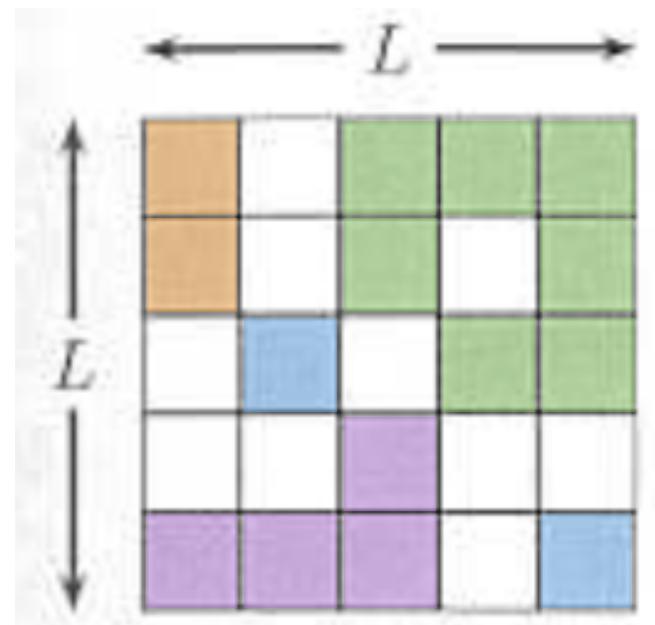
Finite clusters: What is the distribution of sizes of the finite clusters? $n(s, p)$

What is the **probability that a given site belongs to a cluster of size s?** $s n(s, p)$

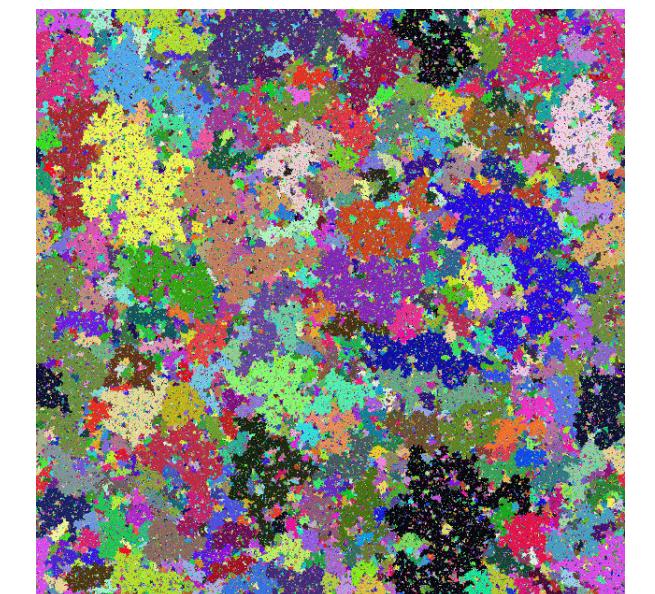
s = cluster size

= number of sites in a cluster

Ex.



s	$N(s, p)$
1	2
2	1
4	1
7	1



$N(s, p)$ = Number of clusters of size s

At 15h today:

Speaker: Jacopo Grilli

research scientist at the International Center for Theoretical Physics, Trieste, Italy.

Title: Competition without contingency and functional convergence

Abstract: Microbial communities are taxonomically diverse and variable: species presence and their abundances widely fluctuate over time and space, and even across biological replicates in experimental controlled conditions. On the other hand, environmental conditions exert a strong selection on the traits of community members and the function they perform, and similar environmental conditions are expected to correspond to functionally similar communities. The consequence of this environmental selection, together with taxonomic variability, lead to the influential concept of functional redundancy: the same function is performed by many species, so that one may assemble communities with different species but the same functional profile. The centrality of the concept of functional redundancy in microbial ecology does not parallel with a theoretical understanding of its origin. In this talk I will describe the eco-evolutionary dynamics of communities interacting through competition and cross-feeding. I will show that the eco-evolutionary trajectories rapidly converge to a "functional attractor", characterized by a functional composition uniquely determined by environmental conditions. The taxonomic composition instead follows non-reproducible dynamics, constrained by the conservation of the functional composition. This framework provides a deep theoretical foundation to the concepts of functional robustness and redundancy.

Location: Institute for Advanced Study (IAS) second floor library, Oude Turfmarkt 147

Zoom link: <https://uva-live.zoom.us/j/89587131703>

At 11h today: There will be the video, if some of you are interested

Speaker: Stefan Thurner

full professor Science of Complex Systems at the Medical University of Vienna,
president of the Complexity Science Hub Vienna, external faculty at the Santa Fe institute

Title: Towards a statistics of driven complex systems

Abstract: Most complex systems are not in equilibrium but driven. We argue that driven systems can be understood by so-called sample space reducing (SSR) processes. They provide an intuitive understanding of the origin and ubiquity of fat-tailed distributions in complex systems, power-laws in particular. SSR processes are mathematically simple and offer an alternative to Boltzmann-equation based approaches to non-equilibrium systems. We show that in many situations the statistics is determined by the details of the driving process and does not depend on the specific relaxation dynamics. Simple (homogeneous) driving strategies universally lead to Zipf's law and exact power laws. Other driving processes results in exponential, Gamma, normal, Weibull, Gompertz, and Pareto distributions. We discuss a number of examples of SRR processes, including fragmentation processes, language formation, cascading and search processes, as well as a classic in physics: inelastical collisions.

Location: Institute for Advanced Study (IAS) second floor library, Oude Turfmarkt 147

Zoom link: <https://uva-live.zoom.us/j/89587131703>

Interesting quantities

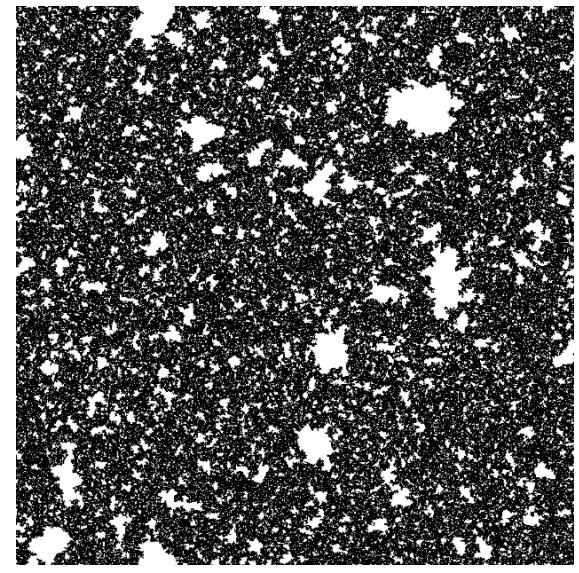
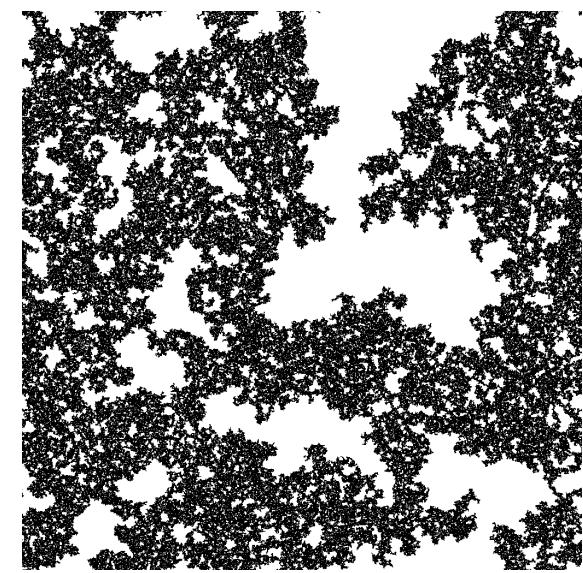
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Size, Infinite Cluster: How much space does the percolating cluster occupy?

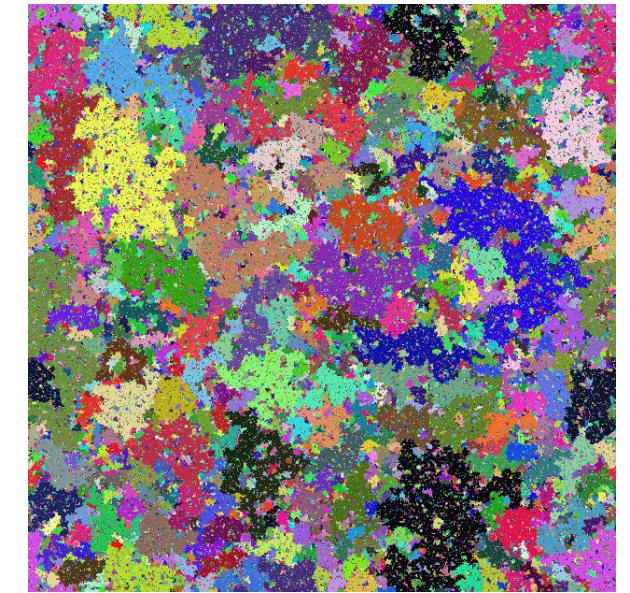
= What is the **probability that a site belong to the percolating cluster?** $P_\infty(p)$



Finite clusters: What is the distribution of sizes of the finite clusters? $n(s, p)$

What is the **probability that a given site belongs to a cluster of size s?** $s n(s, p)$

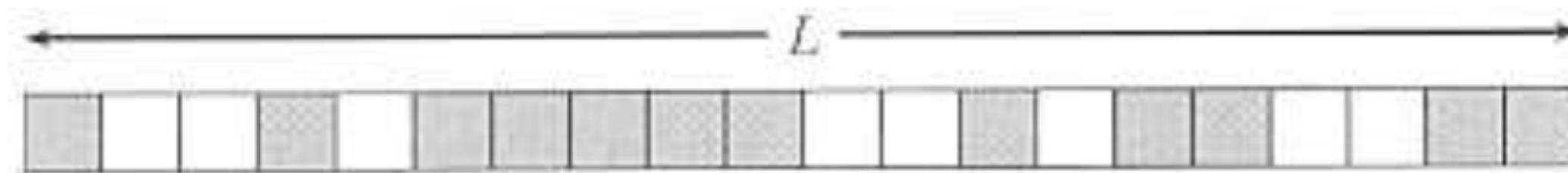
$n(s, p)$ → Will allow us to compute all the quantities of interest.



Average cluster size: Take an occupied site, what is the average size of the cluster it belongs to? $\chi(p)$

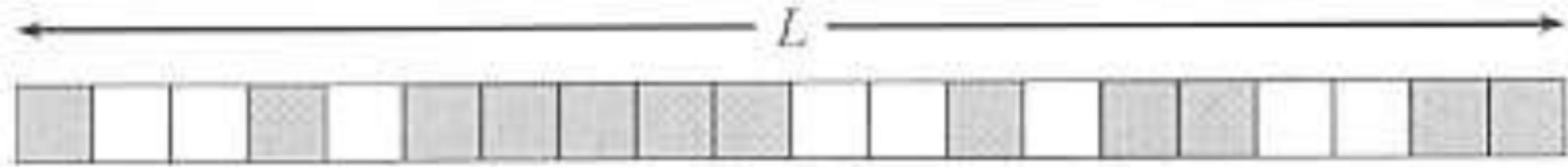
Correlation: Probability two sites at a distance r apart belong to the same cluster $g(r)$

3) Site Percolation in $d=1$



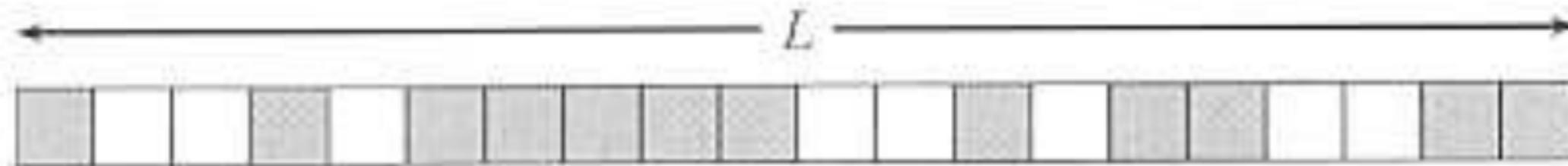
Ex. 1D site percolation

Critical probability



In 1d lattice: What is p_c ?

Critical probability



In 1d lattice: The percolating cluster must **include every sites in the lattice** to span from left to right

If $p < 1$: There is no percolating cluster

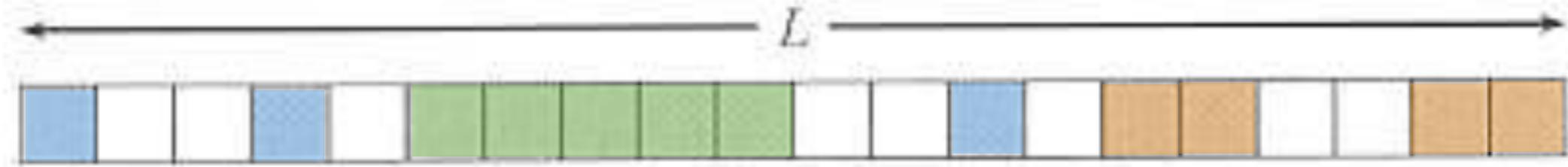
If $p = 1$: There is a percolating cluster



1D percolation: $p_c[1D] = 1$

$$P_\infty(p) \text{ = probability that any given site belongs to the percolating cluster} \quad P_\infty(p) = \begin{cases} 0 & \text{for } p < 1 \\ 1 & \text{for } p = 1 \end{cases}$$

Cluster number density

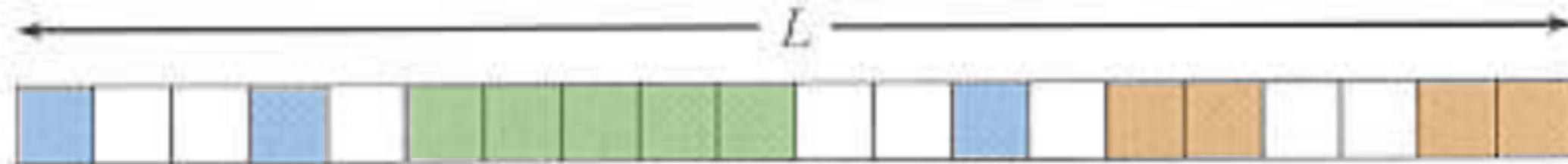


Note: boundary effects are ignored, as they become irrelevant at large L

s	$N(s, p)$
1	3
2	2
5	1

$N(s, p)$ = **cluster size frequency** = number of clusters of size s (in a system of size L)

Cluster number density



Note: boundary effects are ignored, as they become irrelevant at large L

s	$N(s, p)$
1	3
2	2
5	1

$N(s, p)$ = **cluster size frequency** = number of clusters of size s (in a system of size L)

$$n(s, p) = \frac{N(s, p)}{L}$$

= **cluster size density** = density of clusters of size s

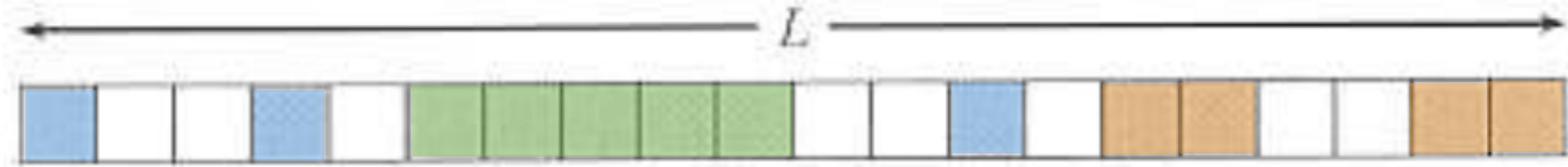
= **number of clusters of size s per unit length**

Extensive quantity

→ We want to compute $n(s, p)$

It is better to consider $n(s, p)$ than $N(s, p)$,
as we want to study what happens as we take L very large
And N depends on L

Cluster number density



Note: boundary effects are ignored, as they become irrelevant at large L

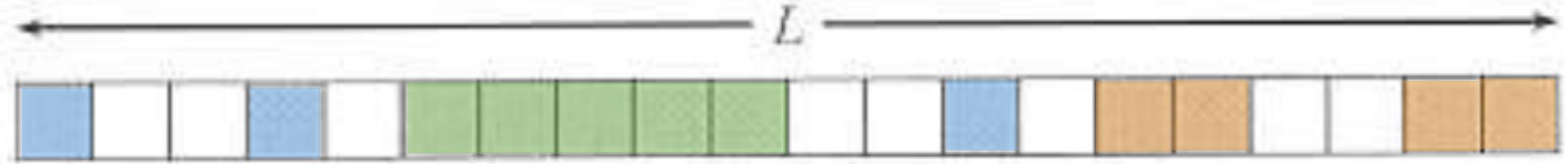
→ We want to compute $n(s, p)$ = **number of clusters of size s , per unit length**

Observe that:

$s n(s, p)$ = probability that a given site belong to a cluster of size s

$n(s, p)$ = probability that a given site belongs to the LHS of a cluster of size s

Cluster number density



Note: boundary effects are ignored, as they become irrelevant at large L

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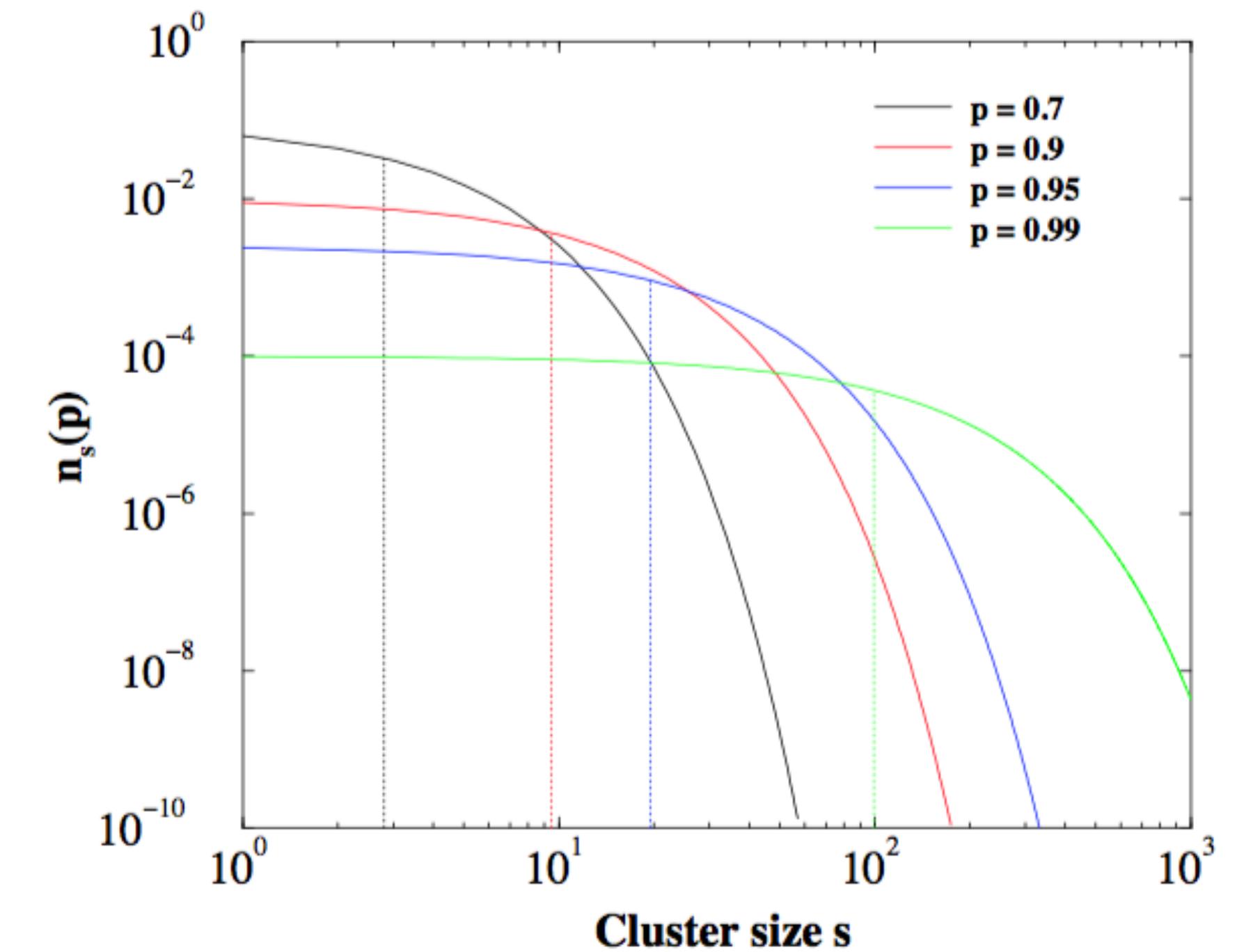
$n(s, p)$ = probability that a given site belongs to the LHS of a cluster of size s

Now $n(s, p)$ is simple to compute: $n(s, p) = p^s (1 - p)^2$

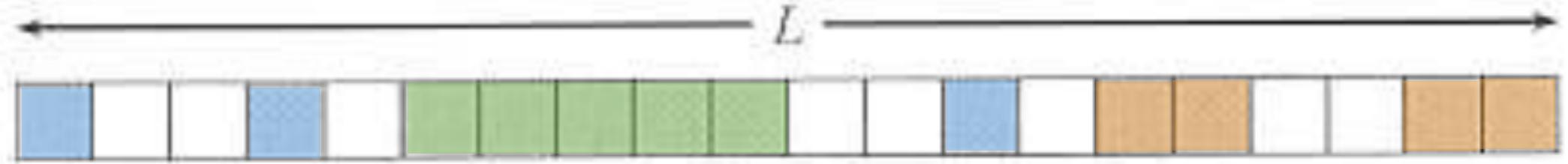
Behavior of $n(s, p)$ for large s : **there is a cutoff!**

very unlikely to find a cluster of size much larger than the cutoff value

Cutoff depends on p : the larger p , the larger the cutoff value



Cluster number density



Note: boundary effects are ignored, as they become irrelevant at large L

→ We want to compute $n(s, p)$ = **number of clusters of size s, per unit length**

Observe that:

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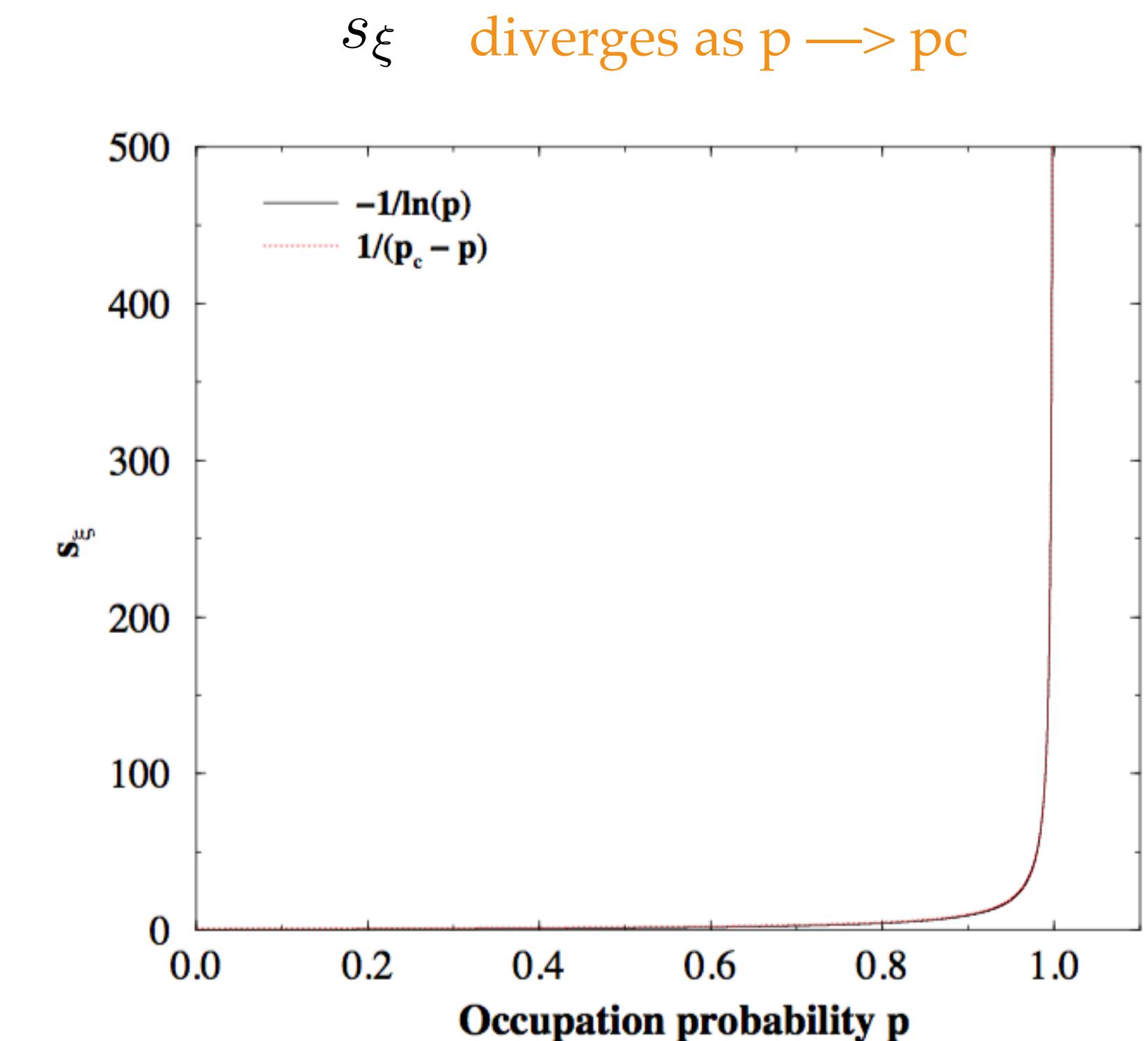
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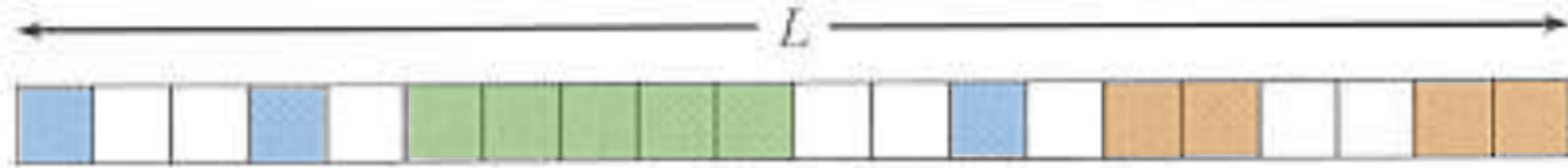
$$= (p_c - p)^2 \exp\left(-\frac{s}{s_\xi}\right)$$

s_ξ = **cutoff cluster size** = characteristic cluster size

$$s_\xi = \frac{-1}{\ln(p)}$$



Cluster number density



Note: boundary effects are ignored, as they become irrelevant at large L

→ We want to compute $n(s, p)$ = **number of clusters of size s , per unit length**

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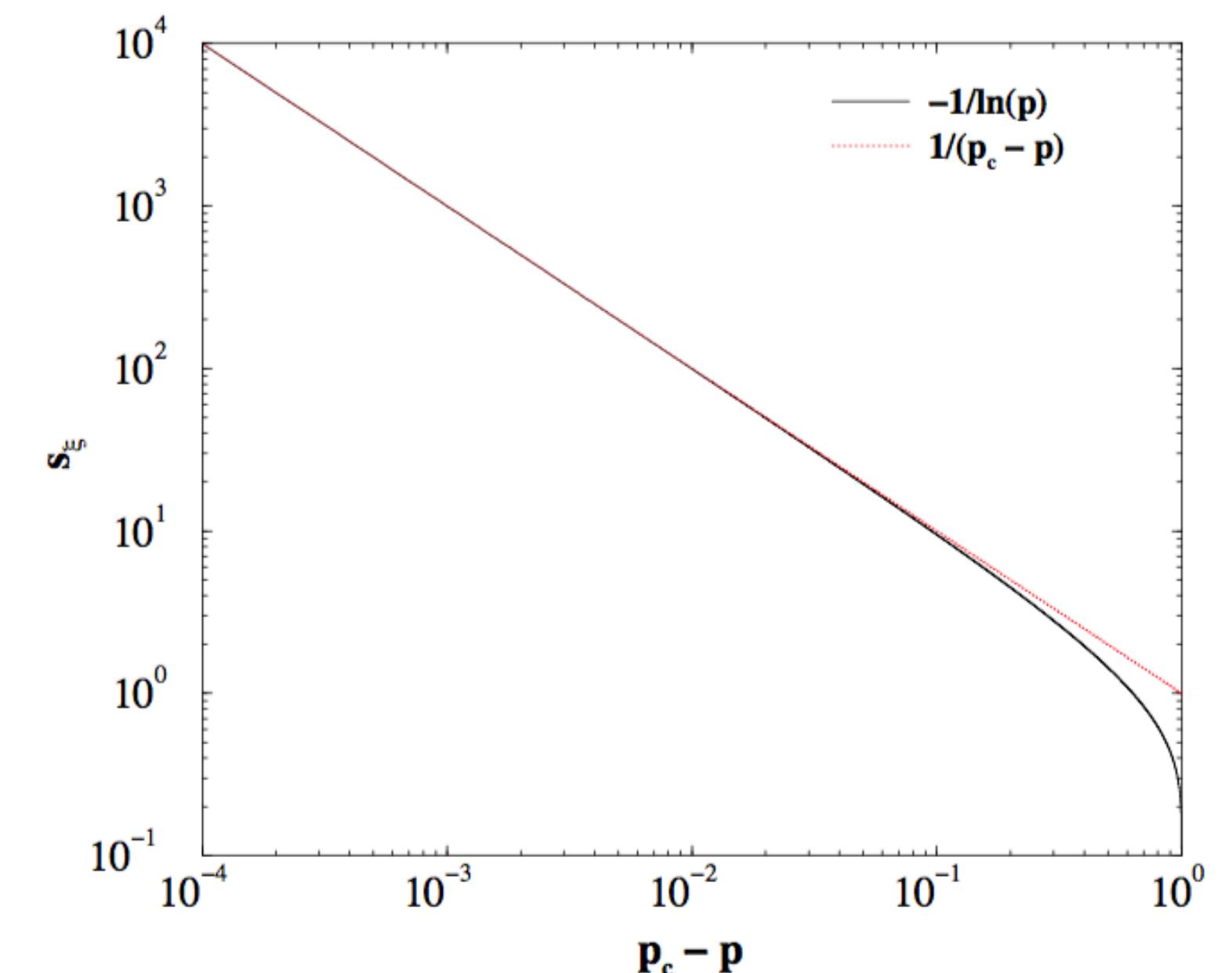
$$= (p_c - p)^2 \exp\left(-\frac{s}{s_\xi}\right)$$

s_ξ = **cutoff cluster size** = characteristic cluster size

$$\ln(1 - x) = -x + o(x) \quad s_\xi = \frac{-1}{\ln(p)} \xrightarrow{p \rightarrow p_c} \frac{1}{p_c - p} = (p_c - p)^{-1}$$

Cutoff cluster size s_ξ diverges for $p \rightarrow p_c$ as a power law!

$$s_\xi \xrightarrow{p \rightarrow p_c} |p_c - p|^{-1/\sigma}$$



Critical exponent: $\sigma = 1$

Useful equality

$s n(s, p)$ = probability that a given site belong to a cluster of size s

Probability that a site belongs to any cluster is: $\sum_s s n(s, p) = p$ for $p < p_c$

Valid at any dimension, and for any lattice size

Mean cluster size

Take a site that is occupied: **What is the average size of the cluster it belongs to?**

We are taking an occupied site, because if not,
the mean will be biased by the value of 0 for all the
sites that are not occupied

Probability that a site is occupied: $p = \sum s n(s, p)$

Probability that a site belongs to a cluster of size s : $s n(s, p)$

Mean cluster size

Take a site that is occupied: What is the average size of the cluster it belongs to?

Probability that a site is occupied: $p = \sum s n(s, p)$

Probability that a site belongs to a cluster of size s : $s n(s, p)$

Probability that an occupied site belongs to a cluster of size s : $w(s, p) = \frac{s n(s, p)}{p} = \frac{s n(s, p)}{\sum_s s n(s, p)}$

Mean cluster size

Take a site that is occupied: how large is the cluster it belongs to?

Probability that a site is occupied: $p = \sum s n(s, p)$

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Probability that an occupied site belongs to a cluster of size s : $w(s, p) = \frac{s n(s, p)}{p} = \frac{s n(s, p)}{\sum_s s n(s, p)}$

Mean cluster size: $\chi(p) = \sum_{s=1}^{\infty} s w(s, p)$

Mean cluster size

Take a site that is occupied: how large is the cluster it belongs to?

Probability that a site is occupied: $p = \sum s n(s, p)$

Probability that a site belongs to a cluster of size s : $s n(s, p)$

Probability that an occupied site belongs to a cluster of size s : $w(s, p) = \frac{s n(s, p)}{p} = \frac{s n(s, p)}{\sum_s s n(s, p)}$

$$\begin{aligned} \text{Mean cluster size: } \chi(p) &= \sum_{s=1}^{\infty} s w(s, p) = \sum_{s=1}^{\infty} \frac{s^2 n(s, p)}{p} = \frac{(1-p)^2}{p} \sum_{s=1}^{\infty} s^2 p^s \\ &= \frac{(1-p)^2}{p} \left(p \frac{d}{dp} \right) \left(p \frac{d}{dp} \right) \left(\sum_{s=1}^{\infty} p^s \right) \end{aligned}$$

$$n(s, p) = p^s (1-p)^2$$

$$\chi(p) = \frac{1+p}{1-p} = \frac{p_c + p}{p_c - p} \xrightarrow{p \rightarrow p_c} \frac{2p_c}{p_c - p} \propto (p_c - p)^{-1}$$

Mean cluster size

Take a site that is occupied: how large is the cluster it belongs to?

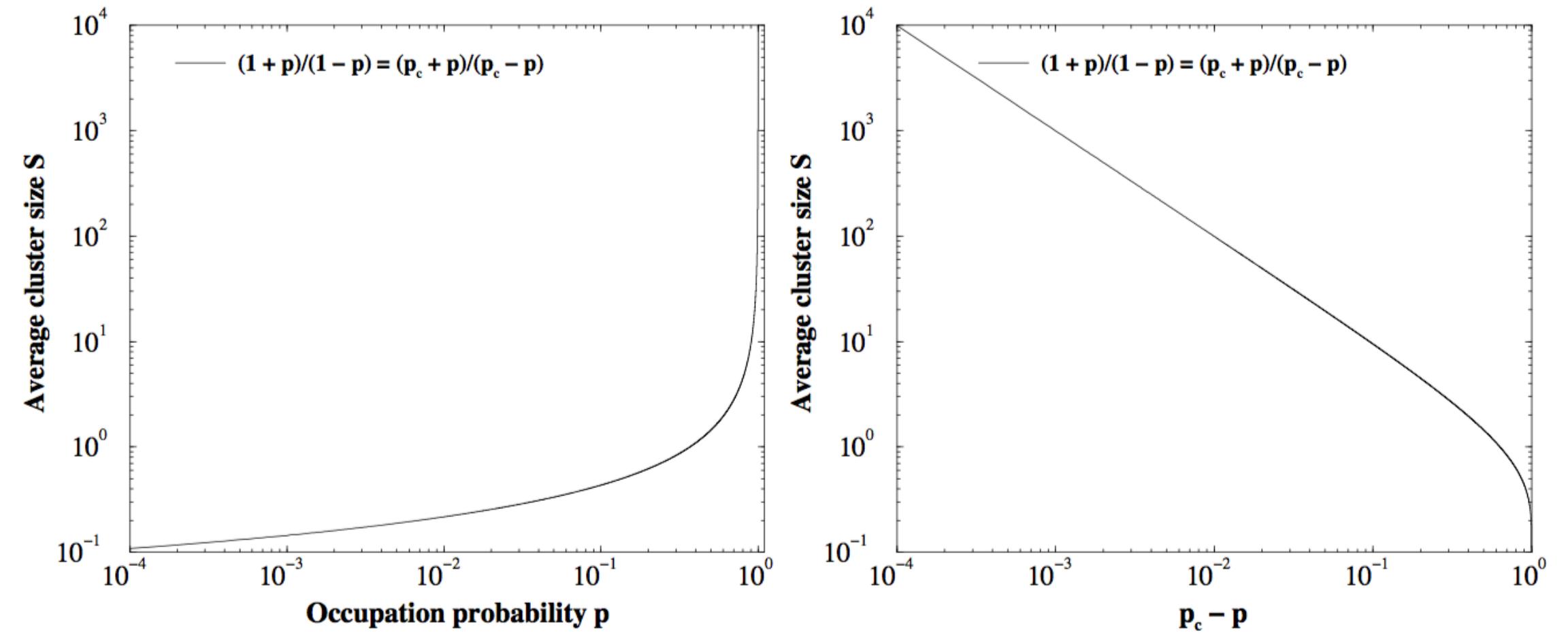
Probability that a site is occupied: $p = \sum s n(s, p)$

Probability that a site belongs to a cluster of size s : $s n(s, p)$

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Mean cluster size:

$$\chi(p) = \frac{1+p}{1-p} = \frac{p_c + p}{p_c - p} \xrightarrow{p \rightarrow p_c} \frac{2p_c}{p_c - p} \propto (p_c - p)^{-1}$$



In 1d, the mean cluster size diverges like a power law in the quantity $(p_c - p)$ when $p \rightarrow p_c$ $\chi(p) \propto (p_c - p)^{-\gamma}$ Critical exponent: $\gamma = 1$

Correlation function

Correlation function: probability that a site at a distance r apart from an occupied site belongs to the same cluster.

$$g(0) =$$

In 1d, for a site at position r to be occupied and belong to the same (finite) cluster, this site and the $(r - 1)$ intermediate sites must be occupied:

$$g(r) =$$

Correlation function

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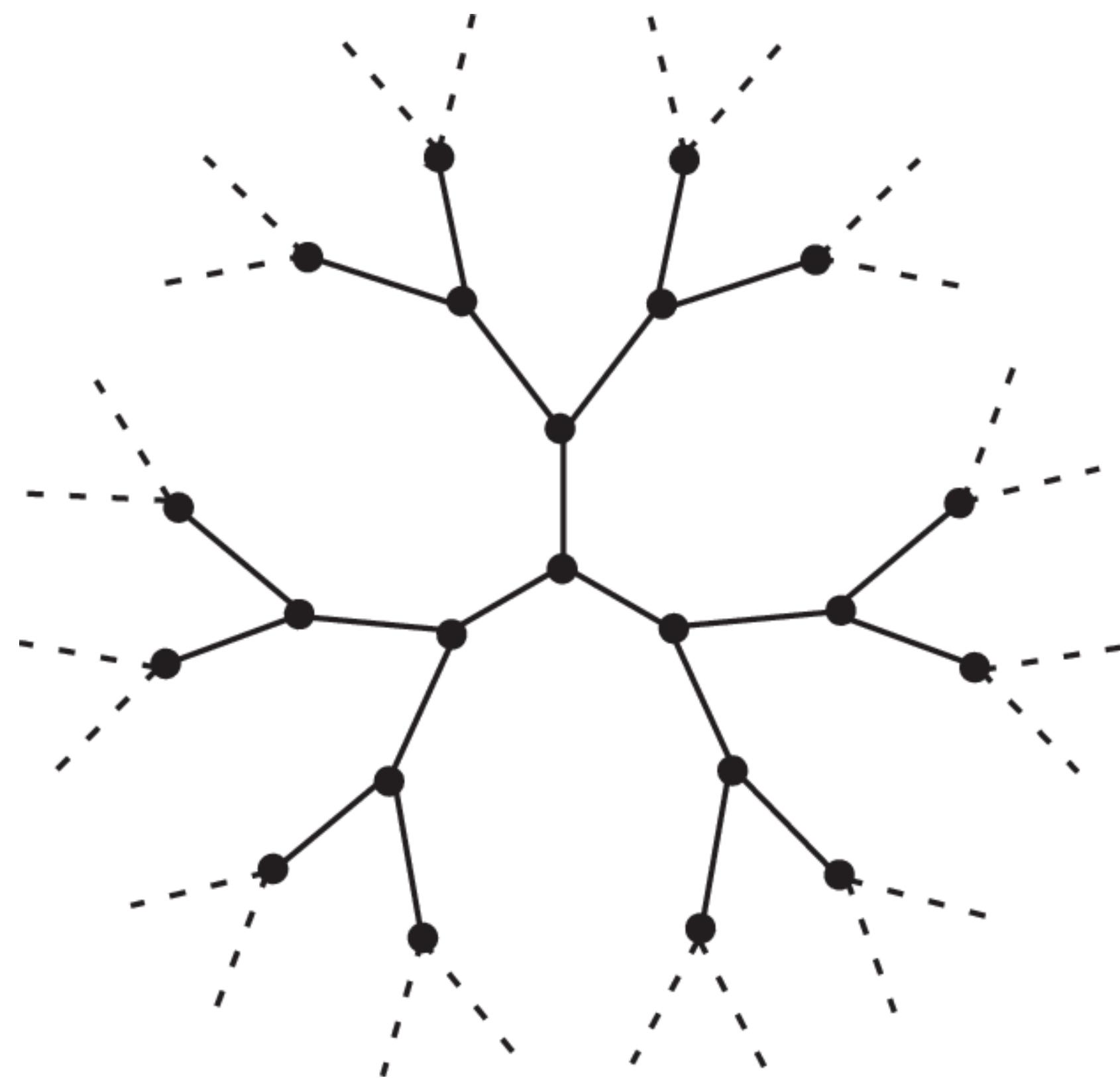
$$g(r) = \exp(r \ln(p)) = \exp(-r/\xi)$$

Correlation length: $\xi = -\frac{1}{\ln(p)} = -\frac{1}{\ln(p_c - (p_c - p))} \rightarrow \frac{1}{(p_c - p)} = (p_c - p)^{-1}$

Correlation length diverges for $p \rightarrow p_c$:

$$\xi = |p_c - p|^{-\nu} \quad \text{critical exponent} \quad \nu = 1$$

4) Percolation on Bethe Lattice

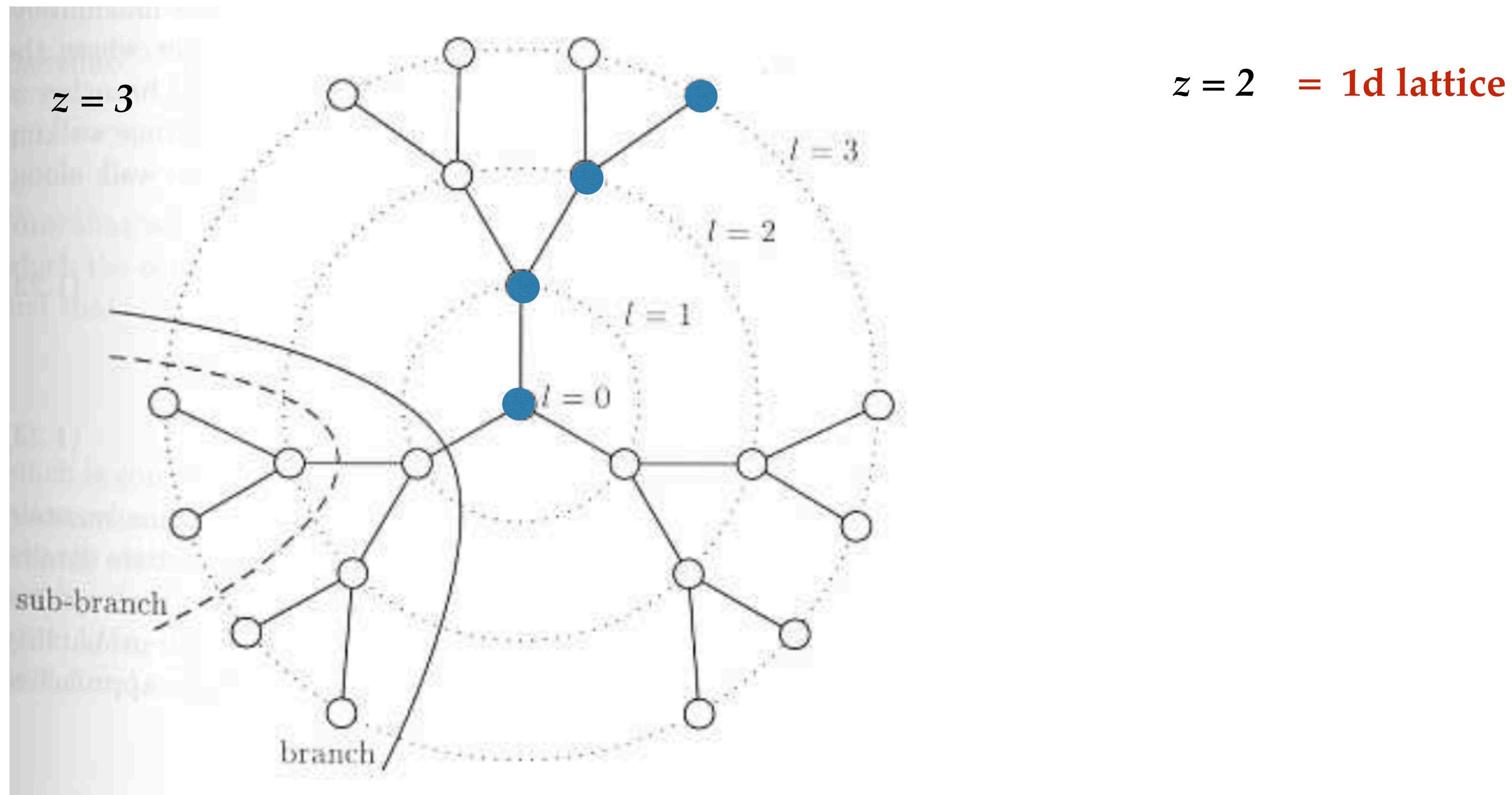


Bethe Lattice

Bethe Lattice: Infinite tree in which each site has exactly z neighbours

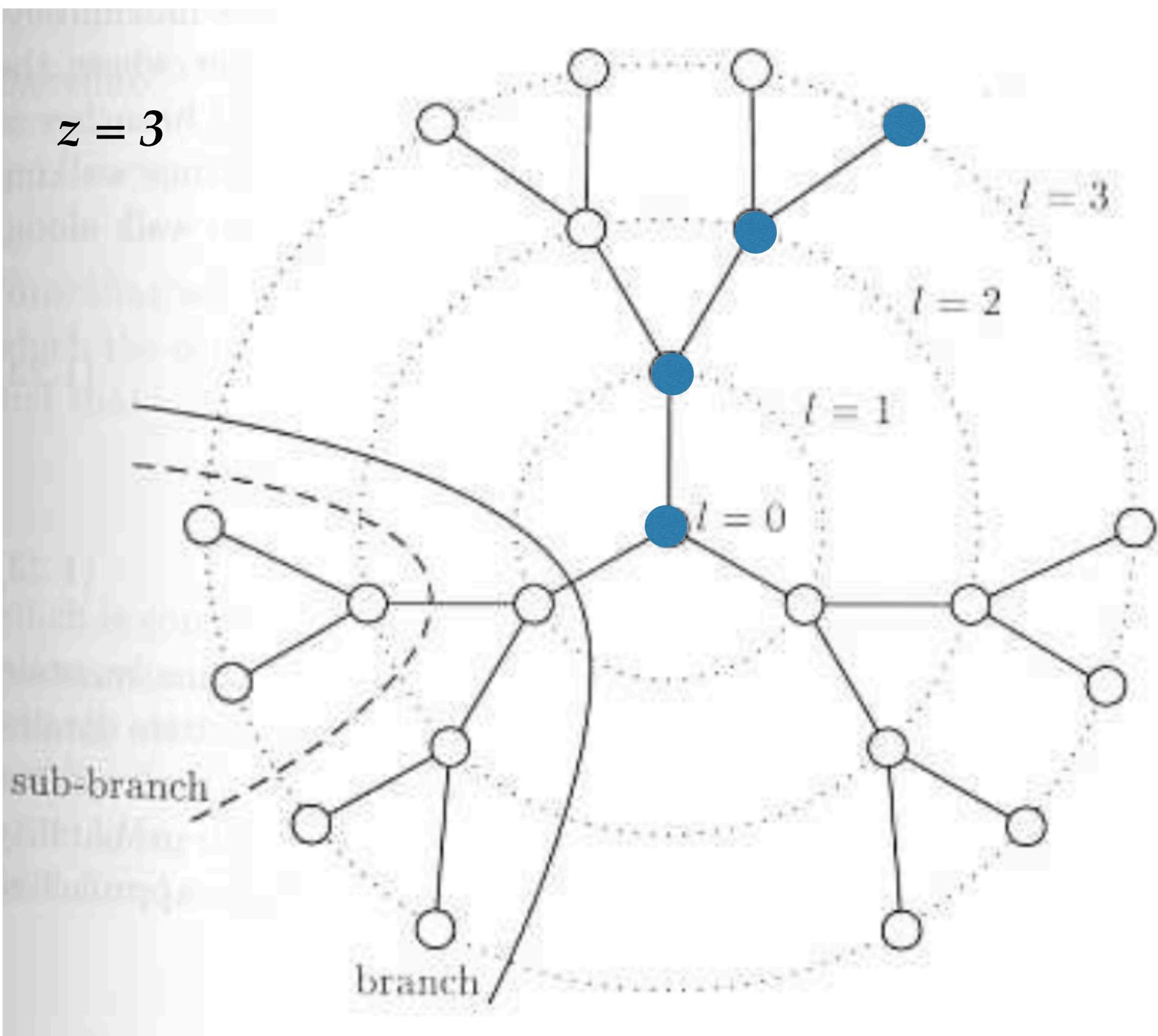
Observation:

- it is a cycle free
- all nodes are the same, there is no special root



Bethe Lattice

Bethe Lattice (Cayley tree): Tree in which each site has exactly z neighbours



$z = 2$ = 1d lattice

Observation:

- it is a cycle free
- all nodes are the same, there is no special root

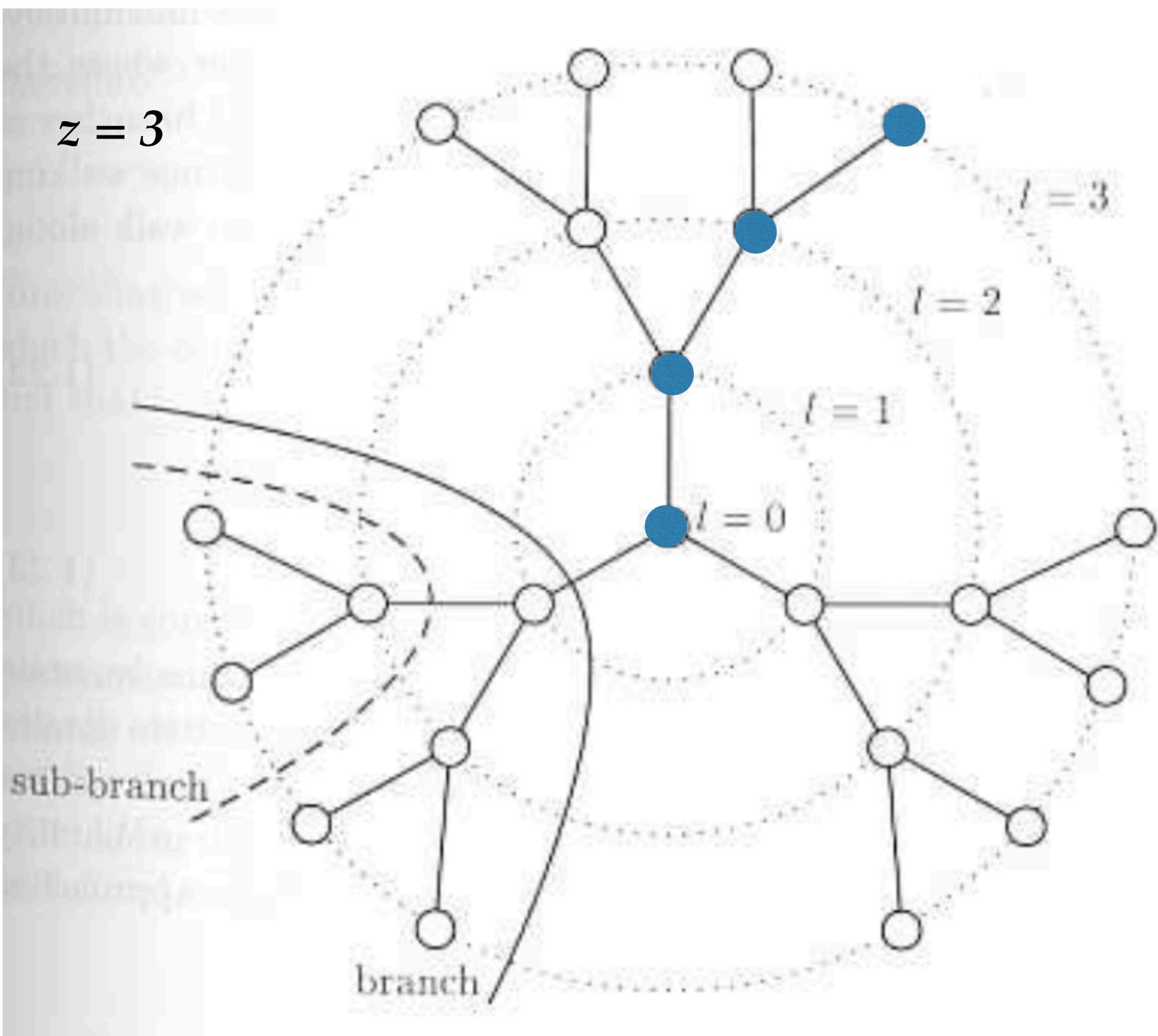
Critical value of p:

Take an occupied site on a path, place it “at the center”

Average number of nodes that can be in the next step of the path = ?

Bethe Lattice

Bethe Lattice (Cayley tree): Tree in which each site has exactly z neighbours



$z = 2$ = 1d lattice

Observation:

- it is a cycle free
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Critical value of p :

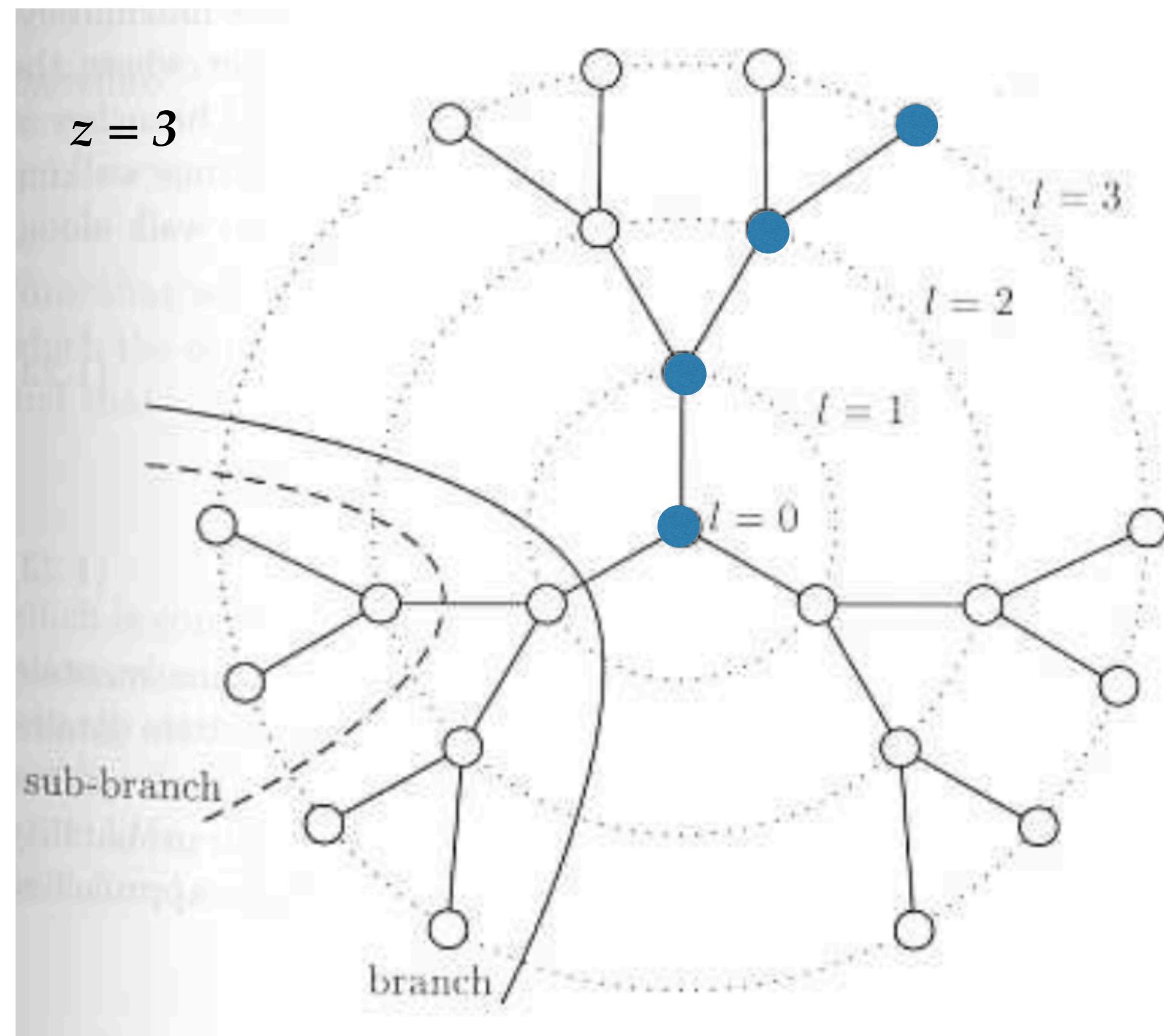
Take an occupied site on a path, place it “at the center”

Average number of nodes that can be in the next step of the path = ?

$$= p(z-1)$$

Bethe Lattice

Bethe Lattice (Cayley tree): Tree in which each site has exactly z neighbours



$z = 2$ = 1d lattice

Observation:

- it is a cycle free
- all nodes are the same, there is no special root

Critical value of p :

Take an occupied site on a path, place it “at the center”

Average number of nodes that can be in the next step of the path = ?

$$= p(z-1)$$

For the path to percolate we need at least $p(z-1) \geq 1$

$$p_c = 1 / (z - 1)$$

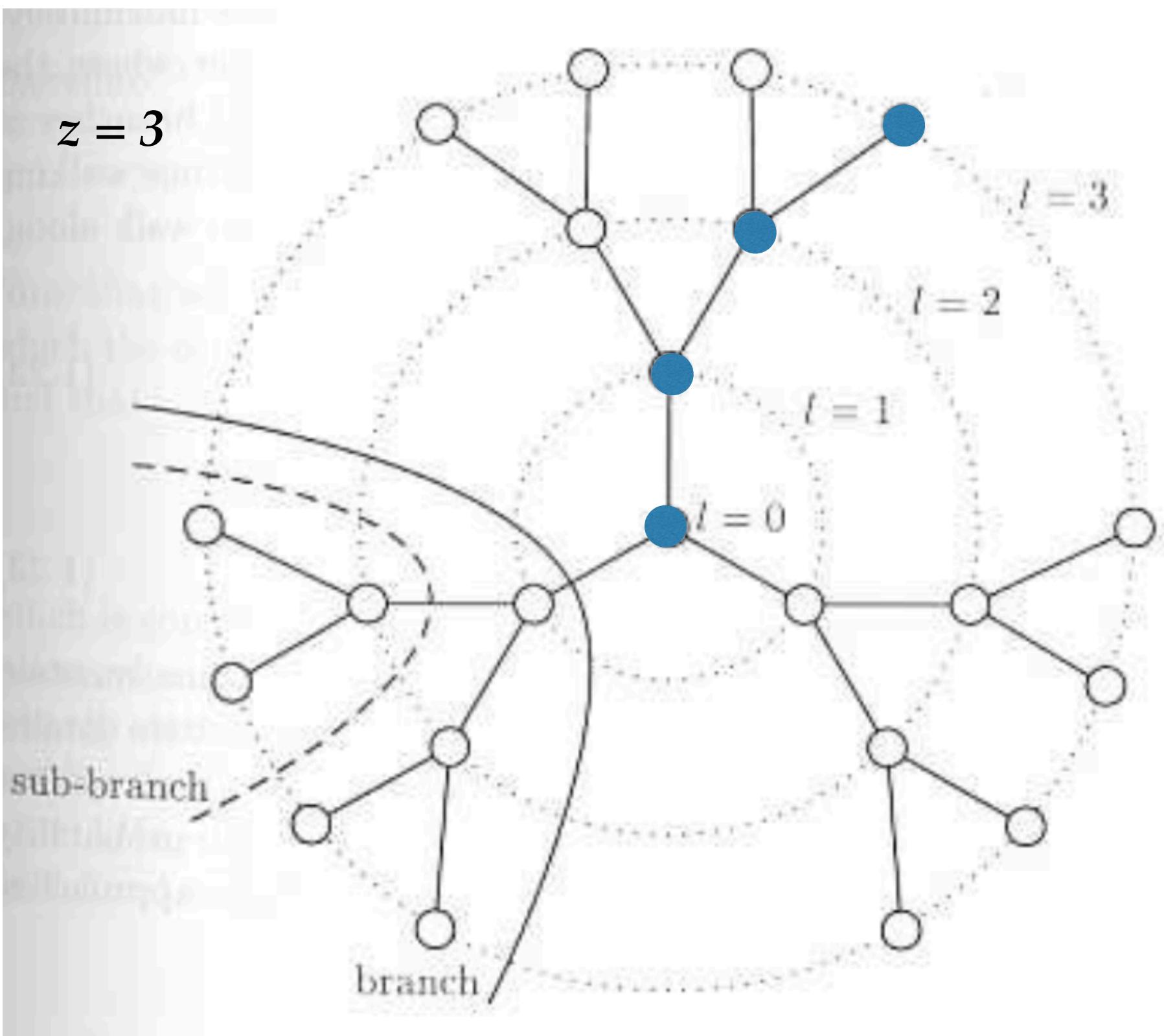
For $z = 2$: $p_c = 1$ \rightarrow 1D lattice

For $z > 2$: $p_c < 1$

The value of p_c depends on the lattice structure

Bethe Lattice

Bethe Lattice (Cayley tree): Tree in which each site has exactly z neighbours



$z = 2$

= 1d lattice

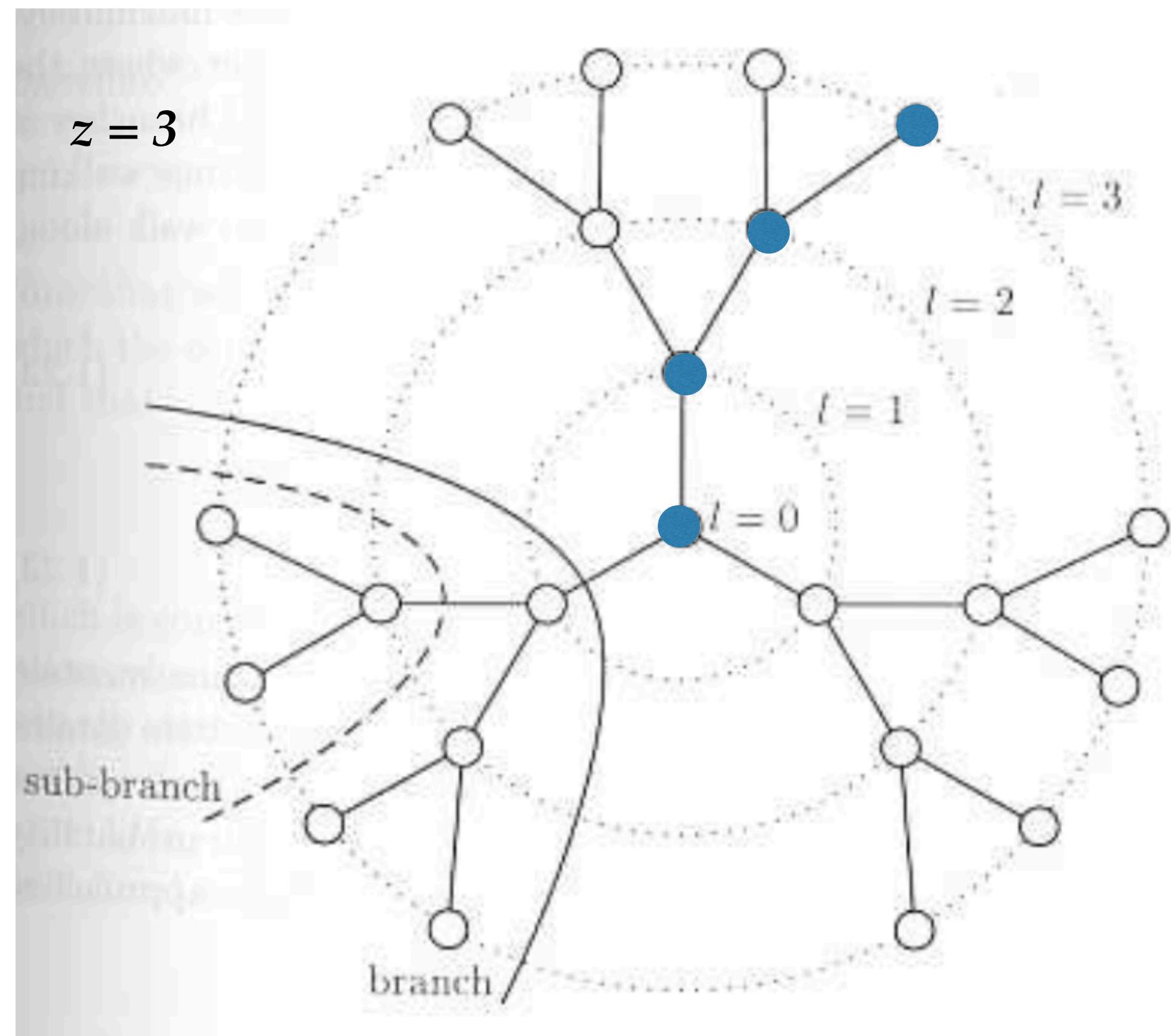
Observation:

- it is a cycle free
- all nodes are the same, there is no special root

Mean Cluster size: ?

Bethe Lattice

Bethe Lattice (Cayley tree): Tree in which each site has exactly z neighbours



$z = 2$

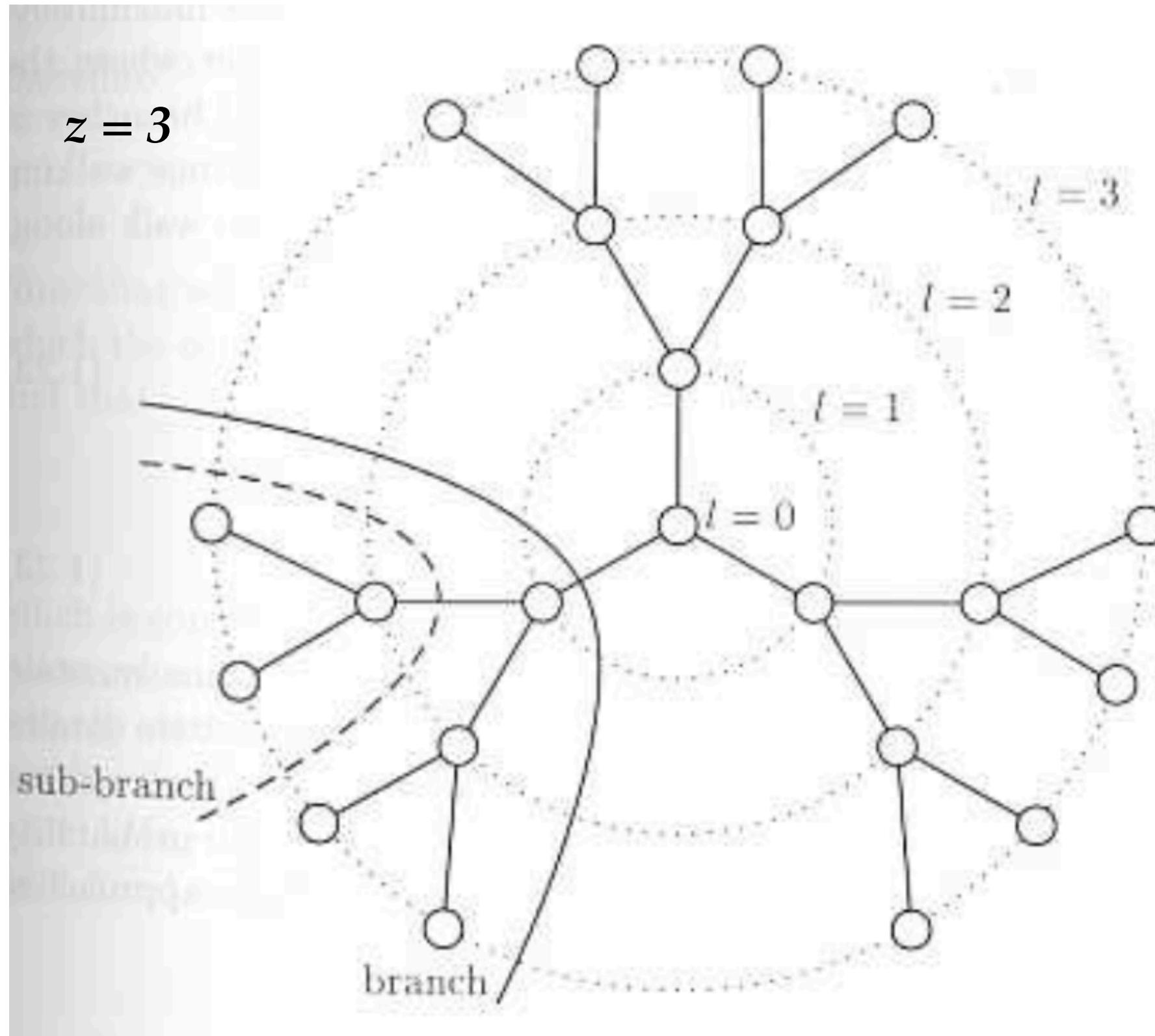
= 1d lattice

Mean Cluster size: ?

$\gamma = 1$

$$\chi(p) \propto |p - p_c|^{-\gamma}$$

Bethe Lattice



The value of p_c depends on the lattice structure

$$p_c = 1 / (z - 1)$$

Interesting quantities are also power laws;

which defines critical exponents:

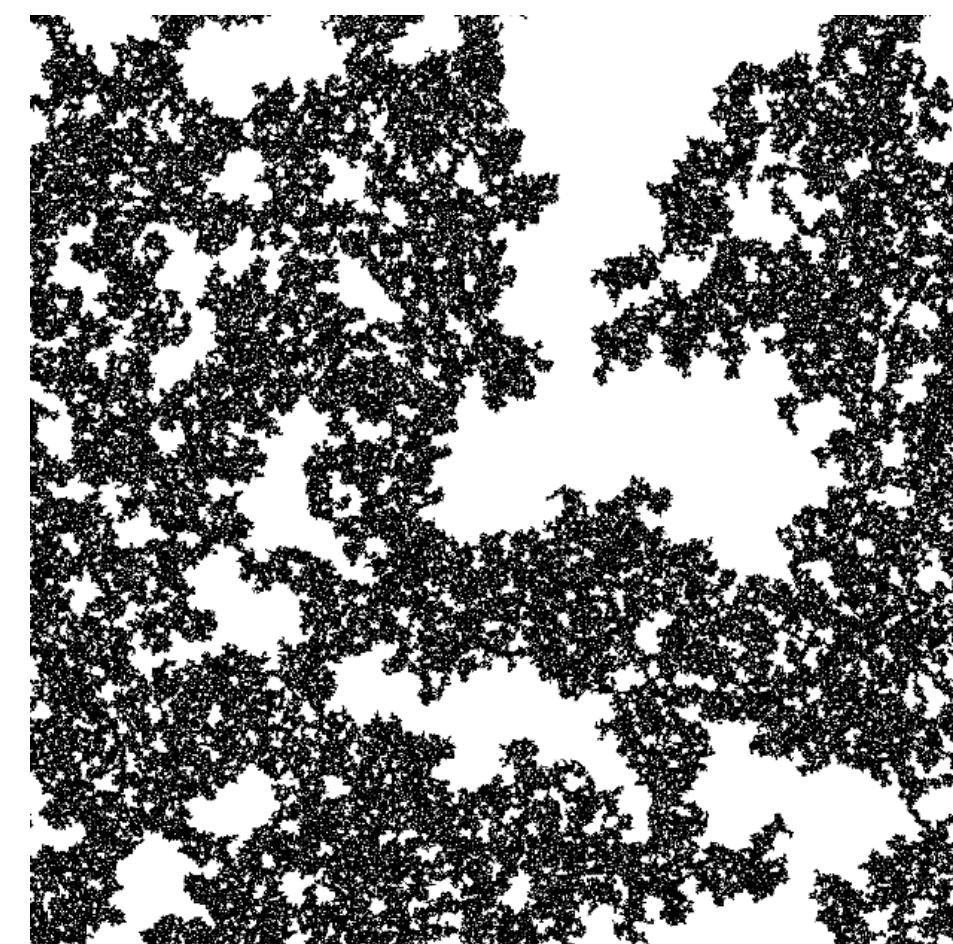
Exponent	Quantity	Bethe	
β	$P_\infty(p) \propto (p - p_c)^\beta$	1	Probability that a site belongs to percolating cluster
γ	$\chi(p) \propto p - p_c ^{-\gamma}$	1	Mean cluster size
ν	$\xi(p) \propto p - p_c ^{-\nu}$	1/2	Correlation length
σ	$\sigma_\xi(p) \propto p - p_c ^{-1/\sigma}$	1/2	Cutoff cluster size
D	$\sigma_\xi \propto \xi^D$	4	

3) Properties at the critical point; Universality

- a) Different type of percolation problems
- b) P_c is not universal
- c) Critical exponents are universal

Properties at the Critical Point

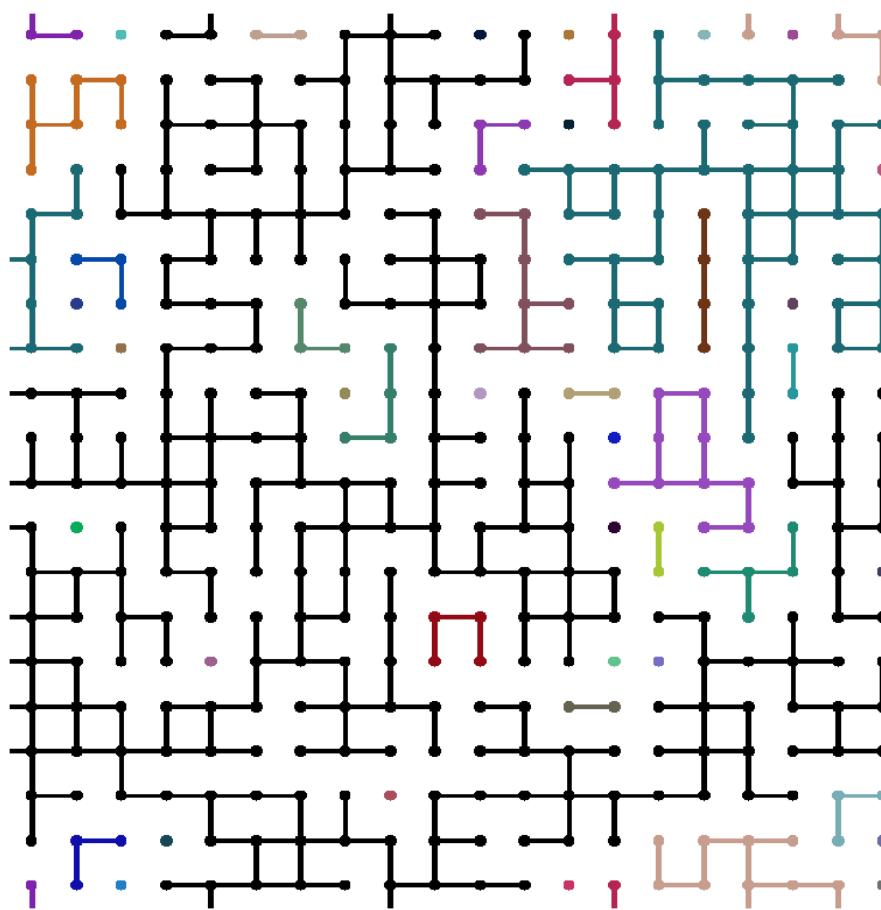
Universality



Percolation and Universality

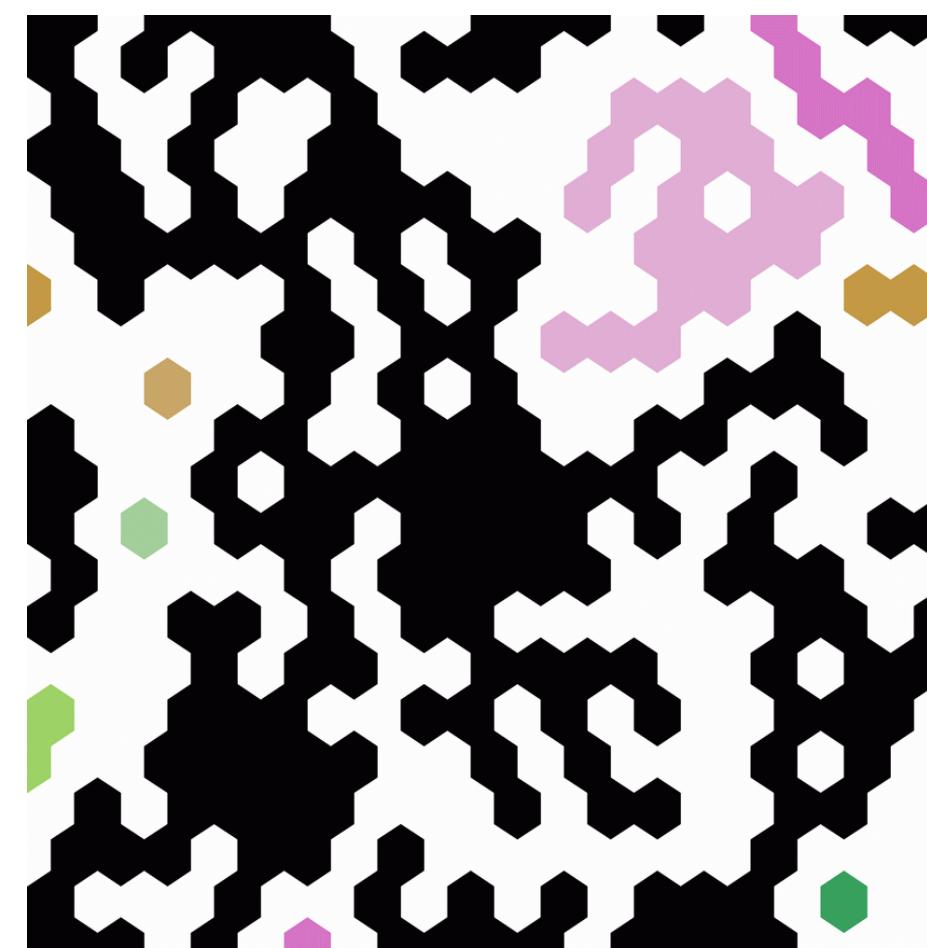
Bond percolation on a lattice:

bonds are randomly **occupied** with probability p



Site percolation on a triangular lattice:

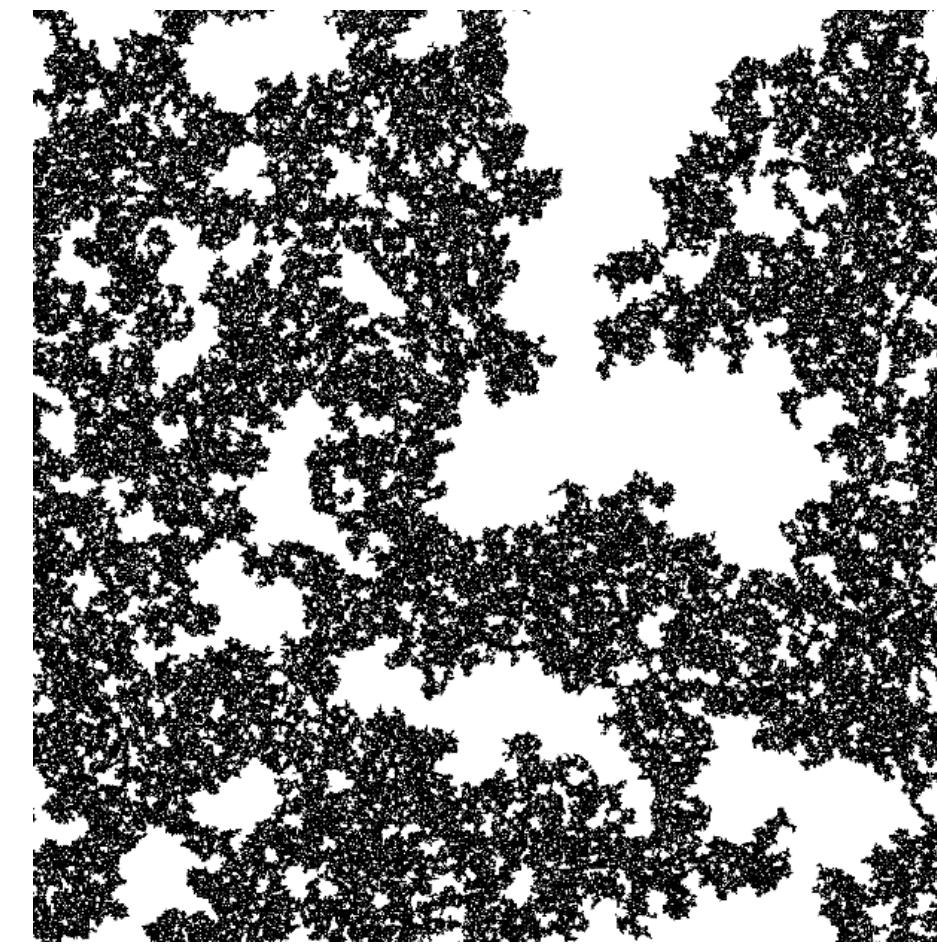
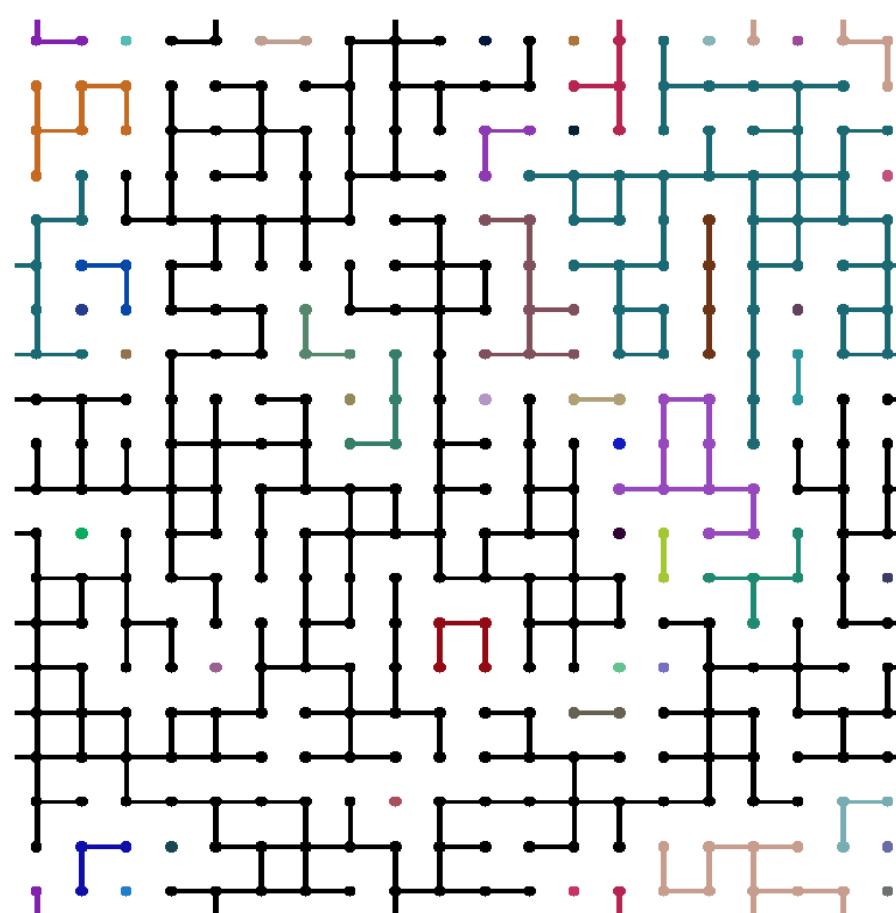
hexagonal sites that are **occupied** with probability p



Percolation and Universality

Bond percolation on a lattice:

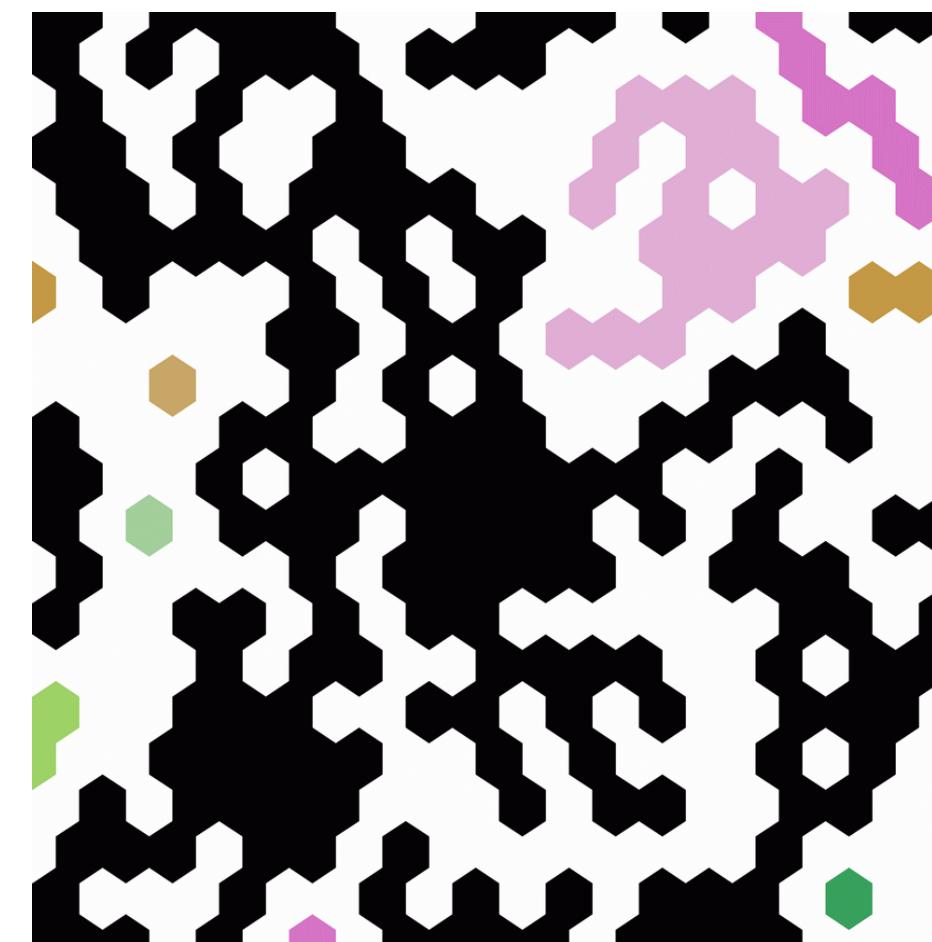
bonds are randomly **occupied** with probability **p**



Infinite cluster
at $p_c = 0.5$

Site percolation on a triangular lattice:

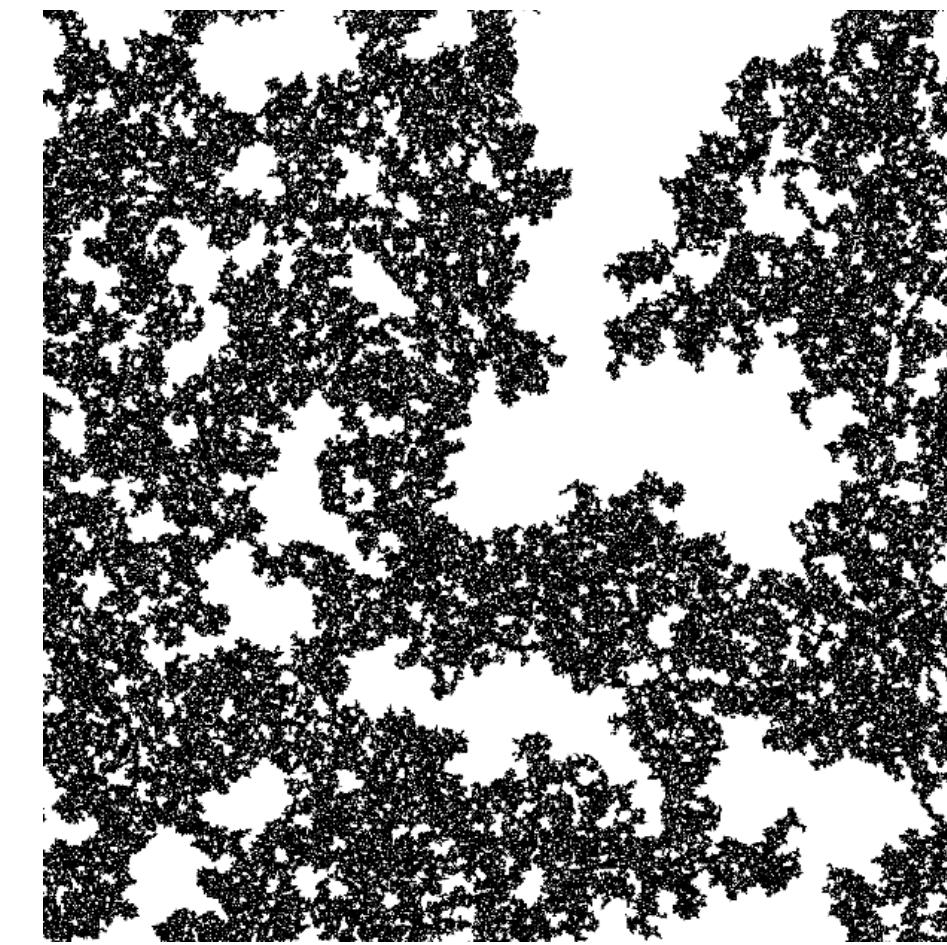
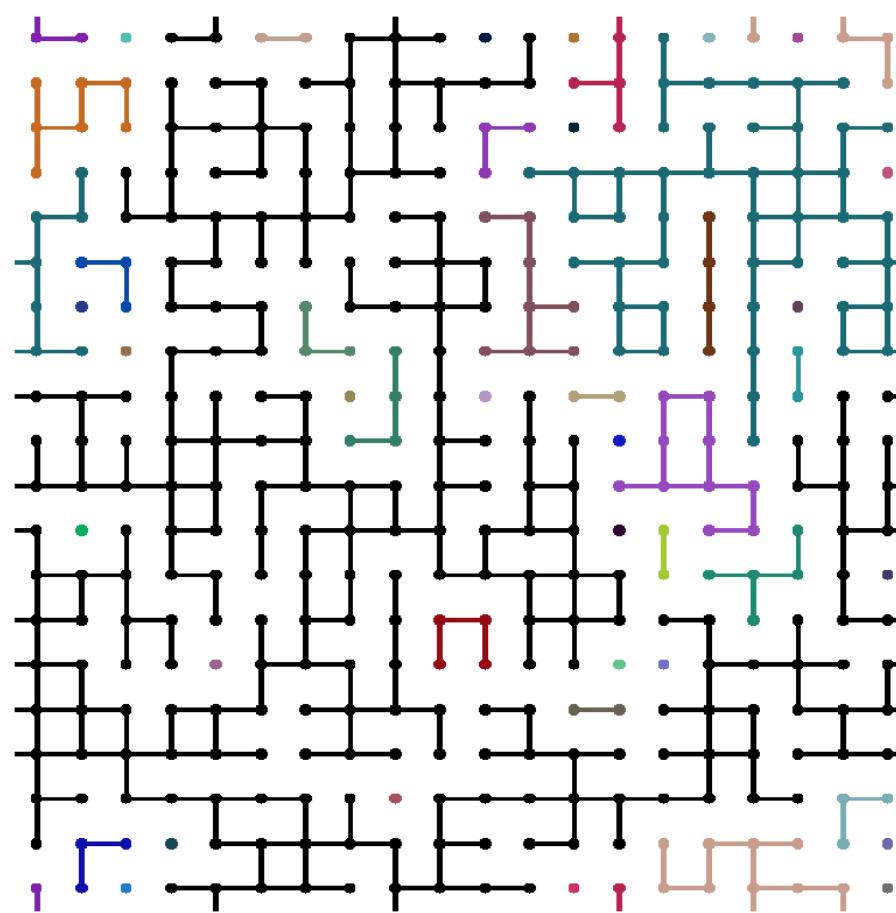
hexagonal sites that are **occupied** with probability **p**



Percolation and Universality

Bond percolation on a lattice:

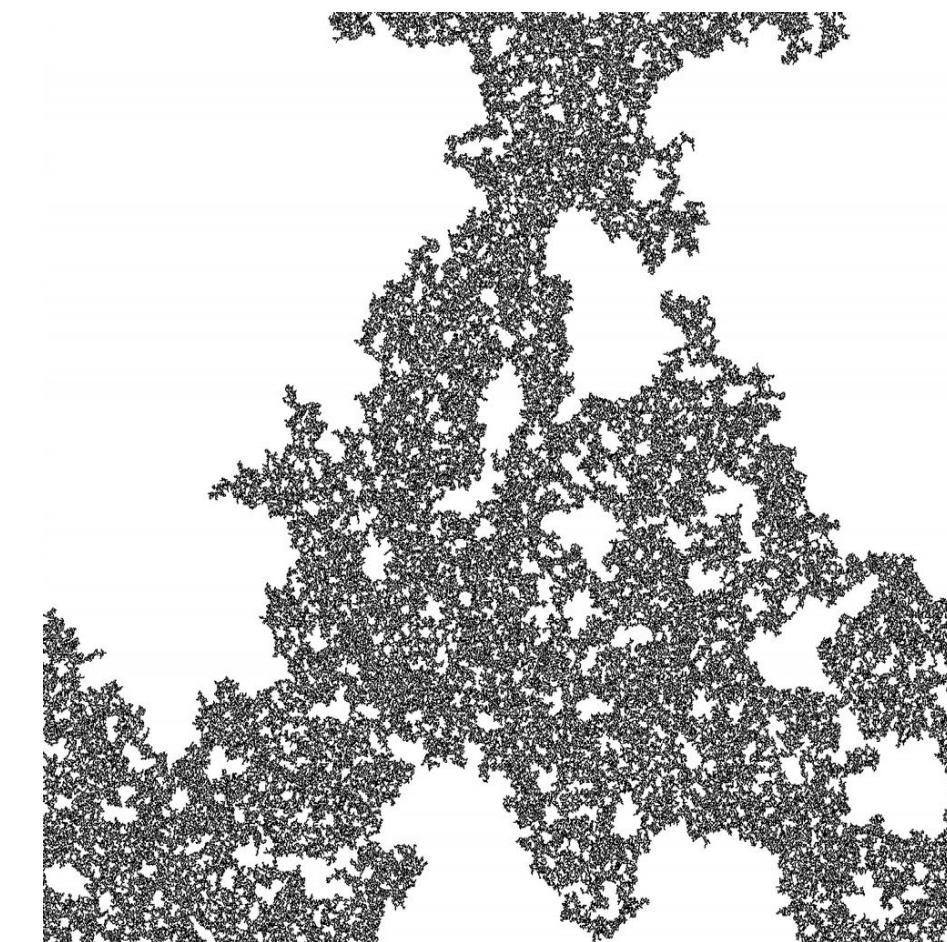
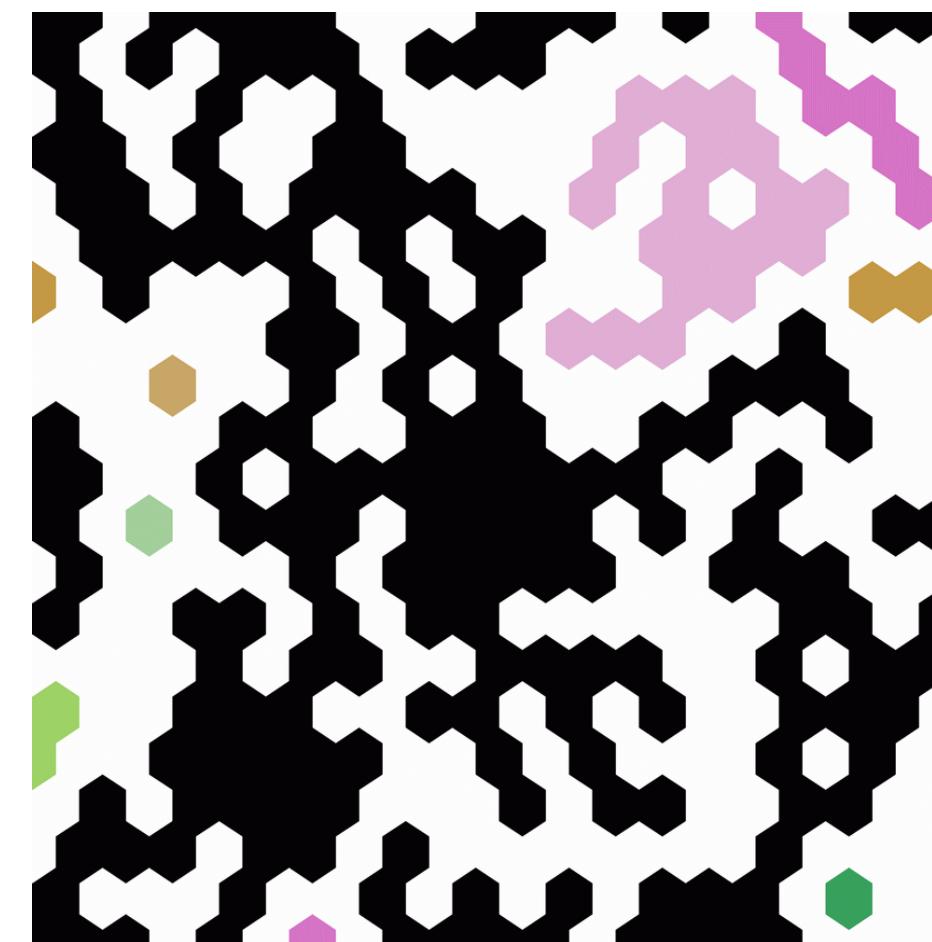
bonds are randomly **occupied** with probability p



Infinite cluster
at p_c

Site percolation on a triangular lattice:

hexagonal sites that are **occupied** with probability p

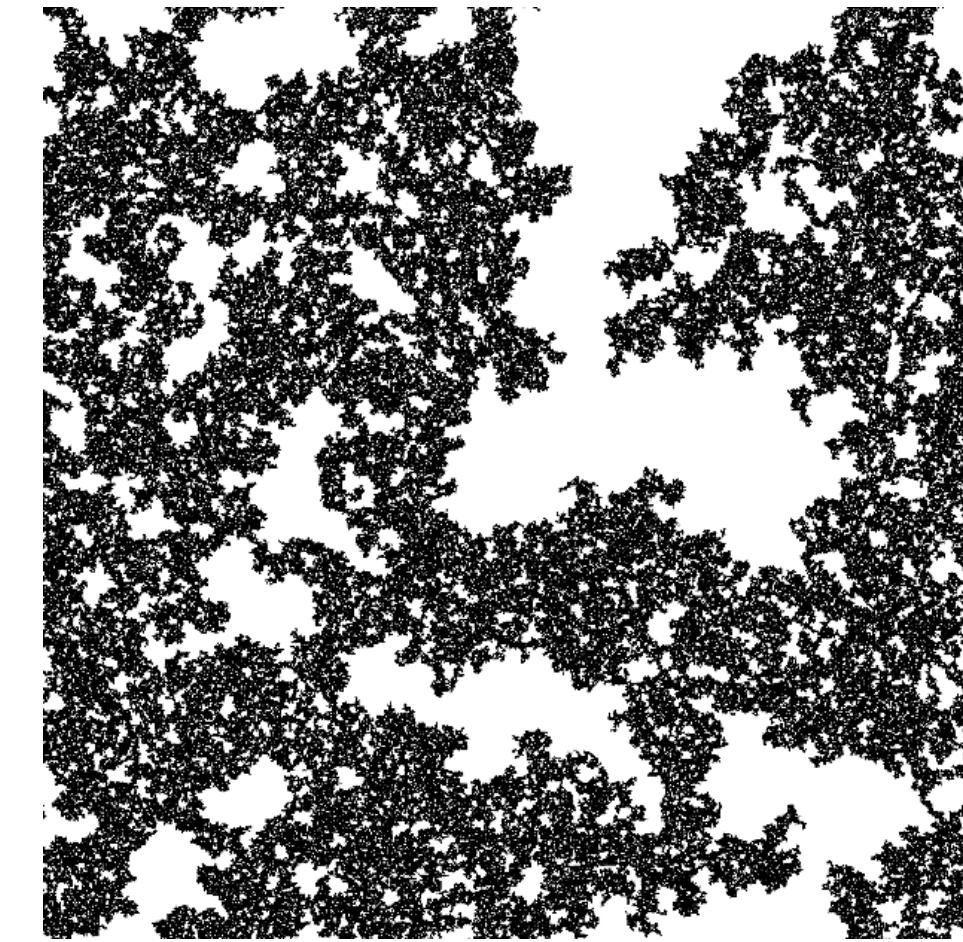
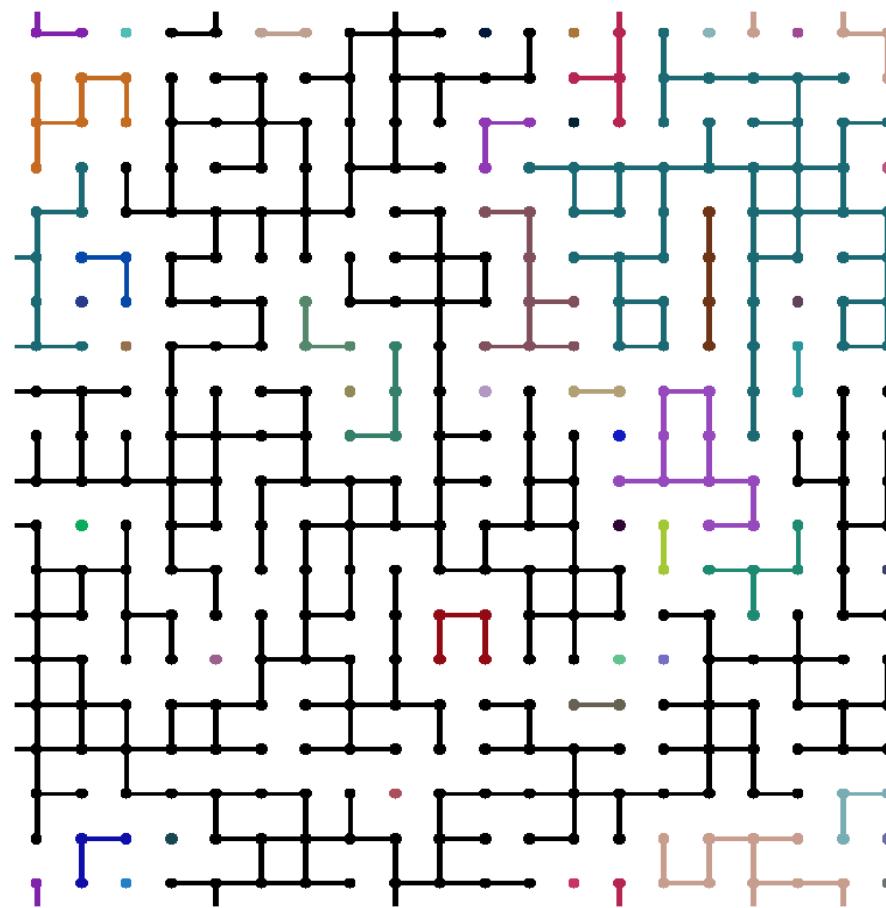


Infinite cluster
at p_c

Percolation and Universality

Bond percolation on a lattice:

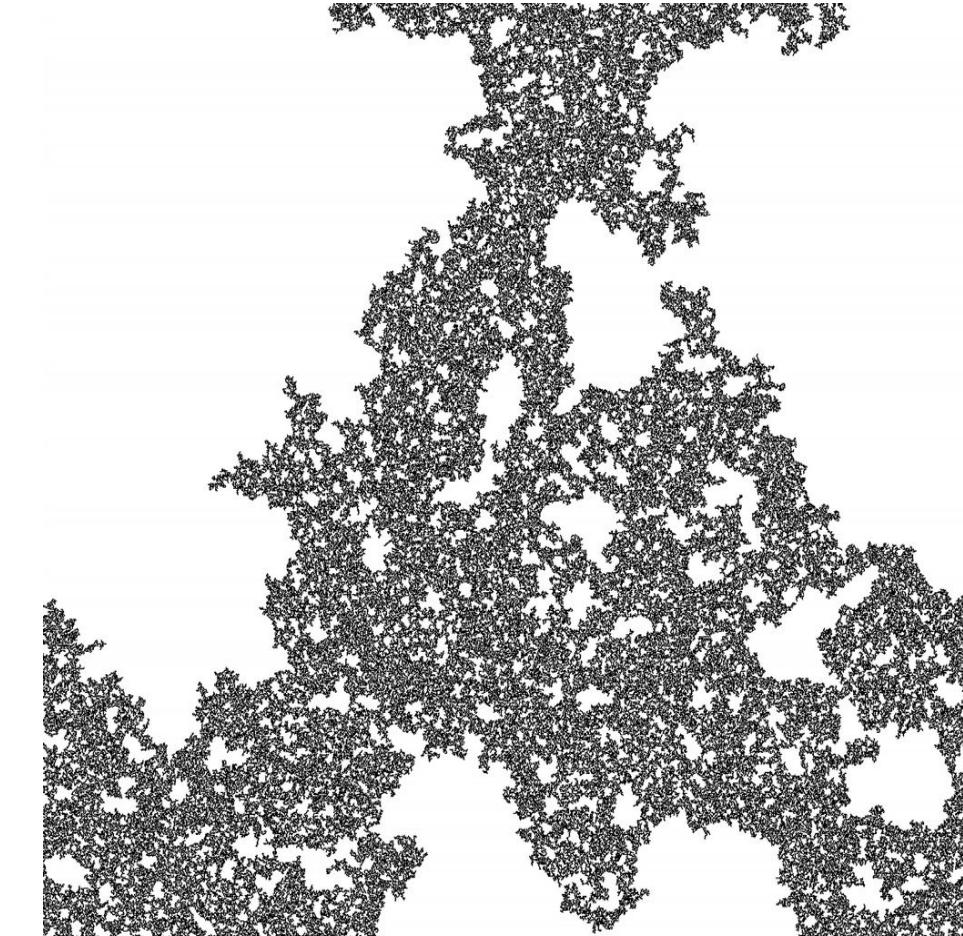
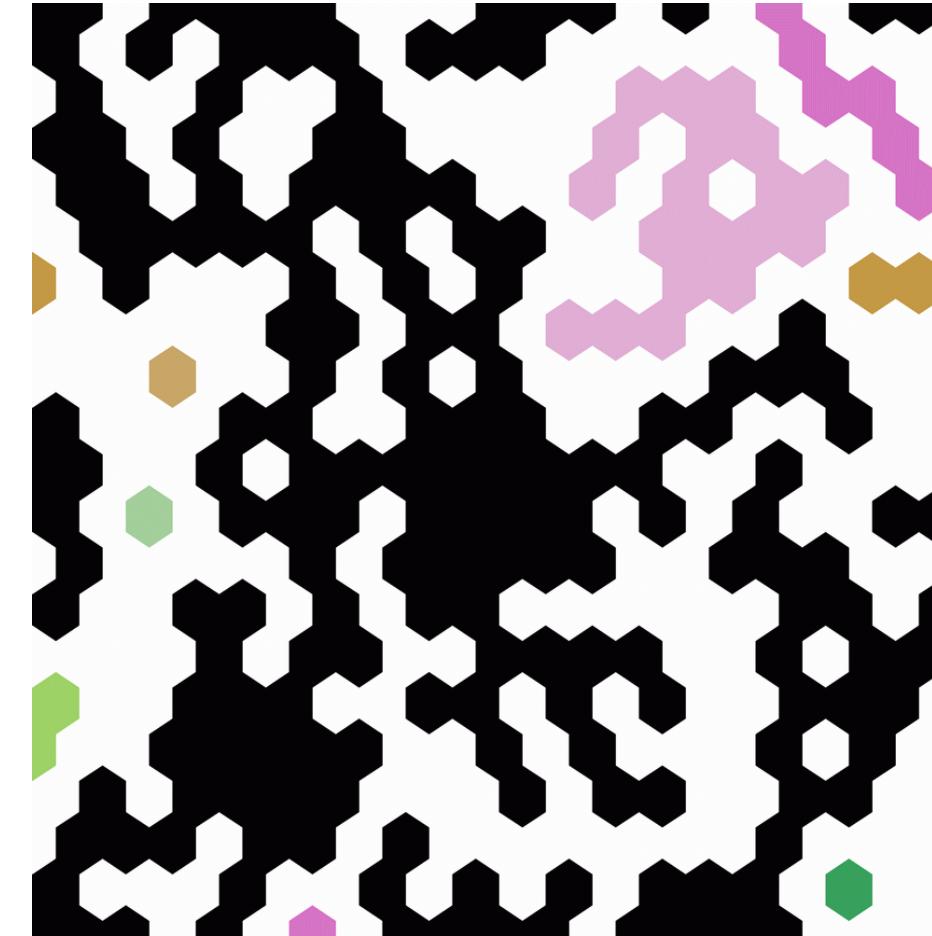
bonds are randomly **occupied** with probability p



Infinite cluster
at p_c

Site percolation on a triangular lattice:

hexagonal sites that are **occupied** with probability p



Infinite cluster
at p_c

Even though the microscopic **lattices** and **occupation rules** are **different**, the resulting **clusters at p_c look identical**.

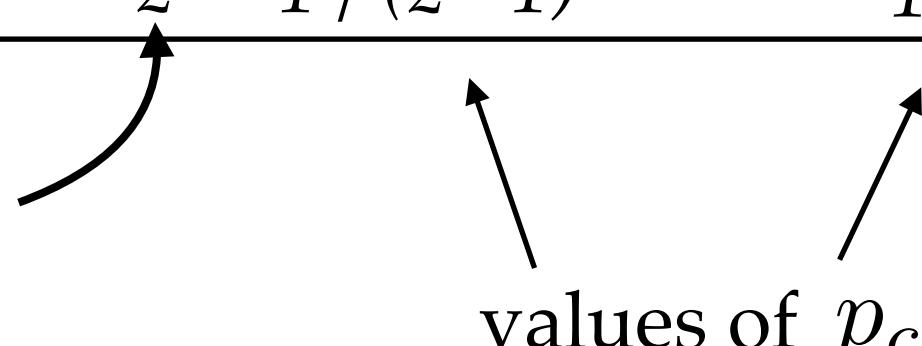
Universality: the morphology of the percolation cluster at p_c is **independent of microscopic details**.

Non-universal critical occupation probability

p_c is not universal

Lattice	z	Site percolation	Bond percolation
$d = 1$ Line	2	1	1
$d = 2$ Hexagonal	3	0.6971	$1 - 2 \sin(\pi/18)$
Square	4	0.5927	$1/2$
Triangular	6	$1/2$	$2 \sin(\pi/18)$
$d = 3$ Diamond	4	0.4301	0.3893
Simple cubique	6	0.3116	0.2488
$d = 4$ Hypercubic	8	0.1968	0.1601
$d = 5$ Hypercubic	10	0.1408	0.1181
$d = 6$ Hypercubic	12	0.1090	0.0942
$d = 7$ Hypercubic	14	0.0889	0.0786
Bethe	z	$1 / (z - 1)$	$1 / (z - 1)$

z = number of nearest neighbors



Depend on the lattice structure and on the occupation rule

Universal critical exponents

Critical exponents are universal

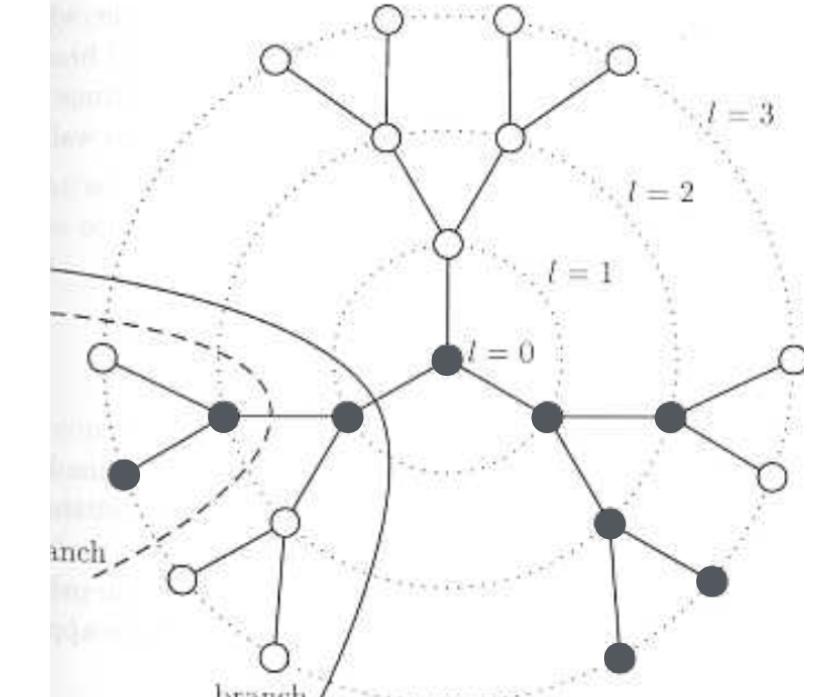
They are **insensitive to the details** of the underlying lattice (such as bound or site, value of z)

They only depend on dimensionality

Exponent	Quantity	$d=1$	$d=2$	$d=3$	$d=4$	$d=5$	$d \geq 6$	Bethe
Pba site belongs to percolating cluster:	$P_\infty(p) \propto (p - p_c)^\beta$	0	5/36	0.4181(8)	0.657(9)	0.830(10)	1	1
Mean Cluster Size:	$\chi(p) \propto p - p_c ^{-\gamma}$	1	43/18	1.793(3)	1.442(16)	1.185(5)	1	1
Correlation length:	$\xi(p) \propto p - p_c ^{-\nu}$	1	4/3	0.8765(16)	0.689(10)	0.569(5)	1/2	1/2
Cutoff cluster size: has exponent:	$s_\xi(p) \propto p - p_c ^{-1/\sigma}$	1	36/91	0.4522(8)	0.476(5)	0.496(4)	1/2	1/2
D	$\sigma_\xi \propto \xi^D$	1	91/48	2.523(6)	3.05(5)	3.54(4)	4	4

Values are independent of
the geometry of the lattice

Upper critical dimension $d_u = 6$
above d_u critical exponent remain unchanged



Bethe lattice:
Bethe lattice: correspond to
an infinite-dimensional lattice,
as there are no loops of occupied sites

The behavior near the critical threshold, p_c , is characterized by **universal critical exponents**

Introduction

Critical exponents and Universality

[“Entropy, Order Parameters, and Complexity”](#), by J. P. Sethna, Section 12.1

Last week: 2D Ising model

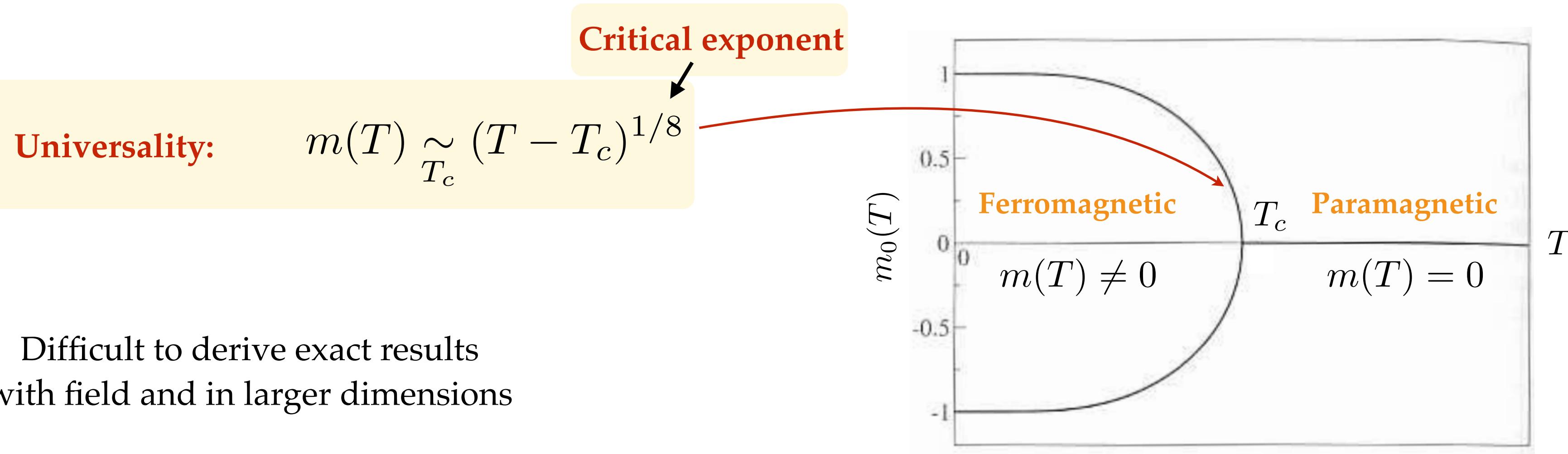
Exact solution for H=0 by Onsager (1944):

proves the existence of a critical phase transition at

$$k_B T_c = \frac{2J}{\log(1 + \sqrt{2})} \simeq 2.269 J$$

with a magnetization:

$$m(T) = \left(1 - \sinh^{-4} \left(\frac{2J}{k_B T} \right) \right)^{1/8}$$



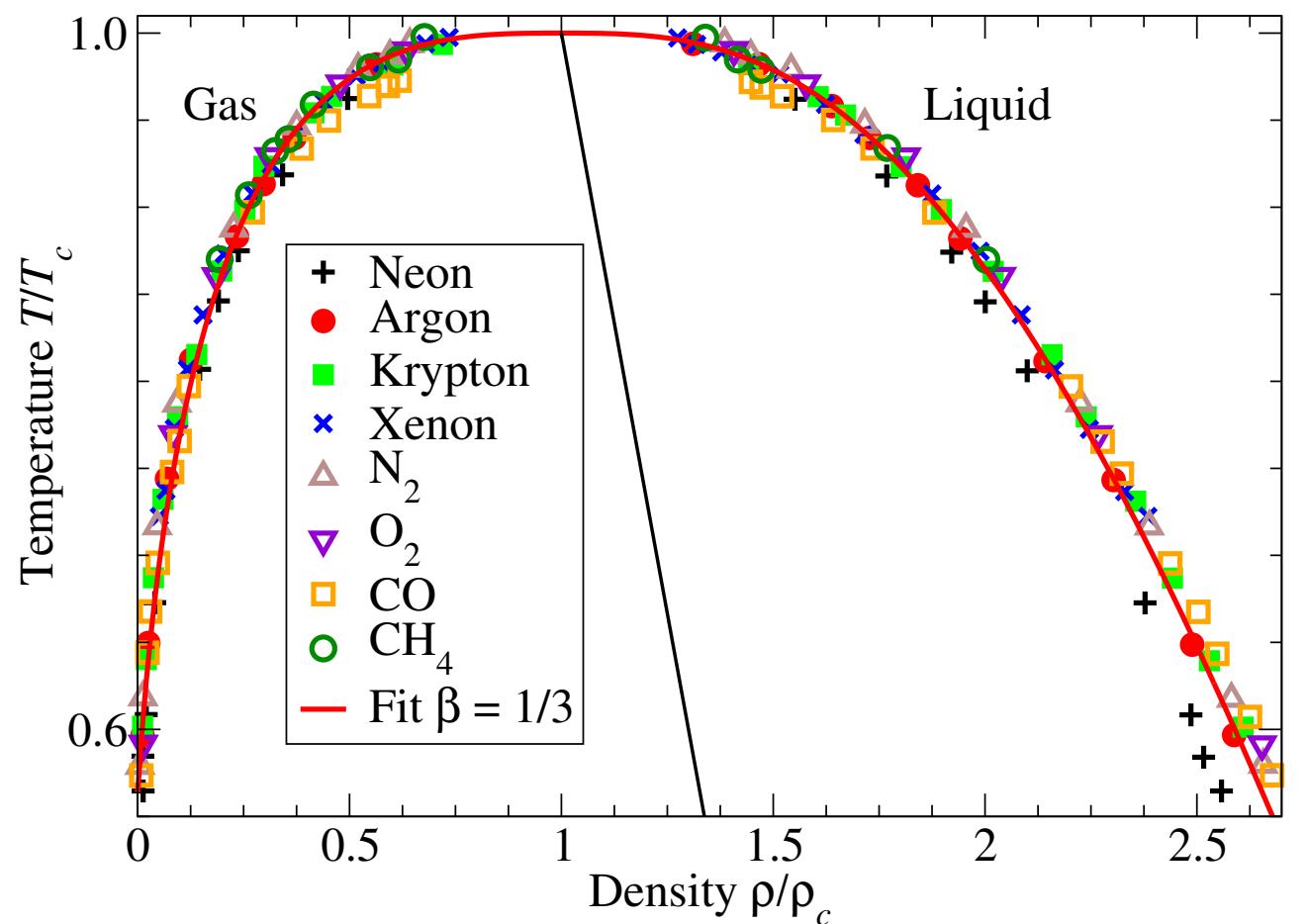
Difficult to derive exact results
with field and in larger dimensions

Universality

Universality: two systems, **microscopically very different**, can exhibit the **same behavior at criticality**

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rescaled by their critical density

rescaled by their critical temperature

Ex. Universality at the liquid-gas critical point:

Liquid-gas coexistence lines for a variety of atoms and small molecules near their critical point.

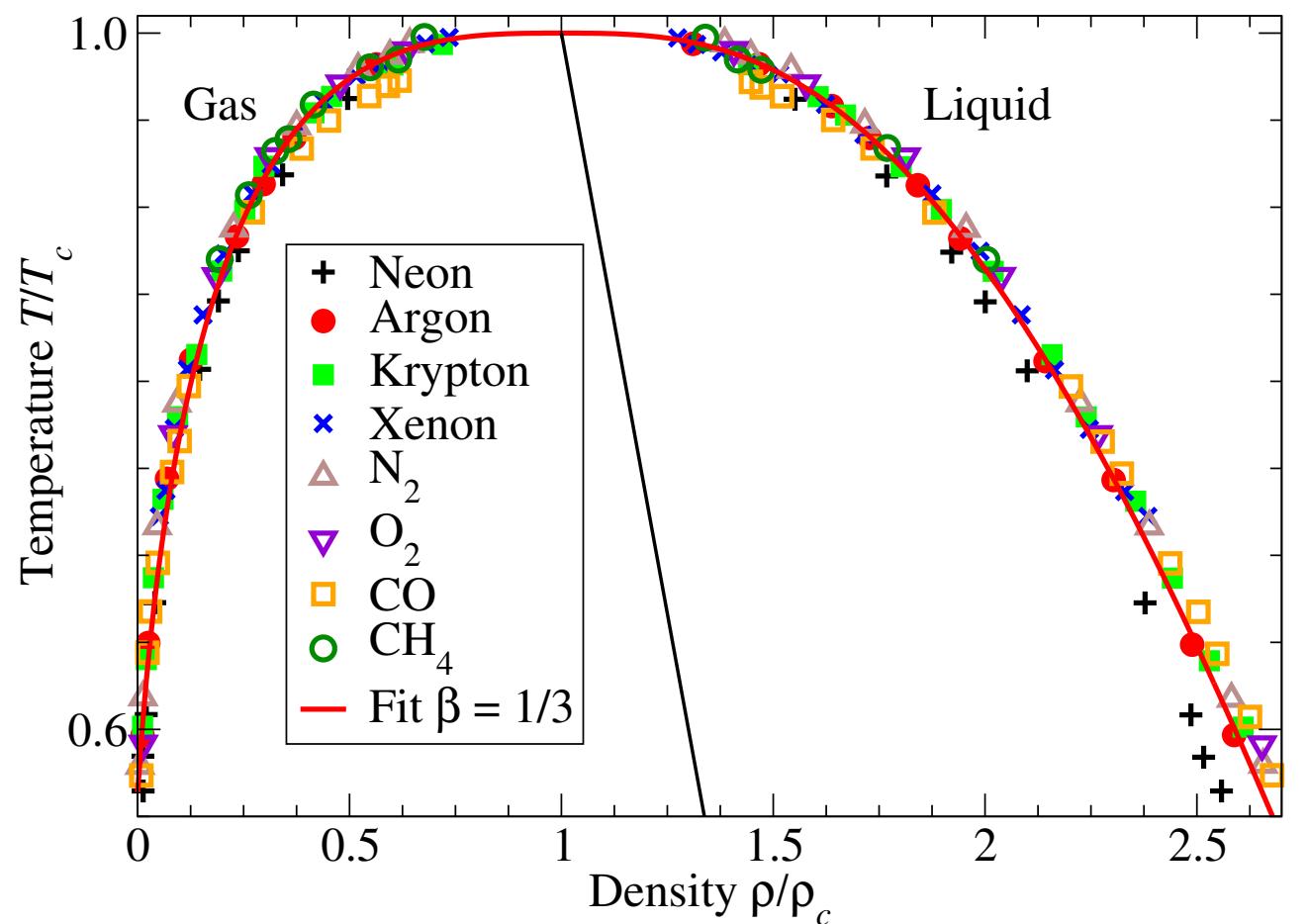
Fit to Argon data:

$$\frac{\rho}{\rho_c} = 1 + s \left(1 - \frac{T}{T_c}\right) \pm \rho_0 \left(1 - \frac{T}{T_c}\right)^\beta$$

$$\begin{cases} s = 0.75 \\ \rho_0 = 1.75 \\ \beta = 1/3 \end{cases}$$

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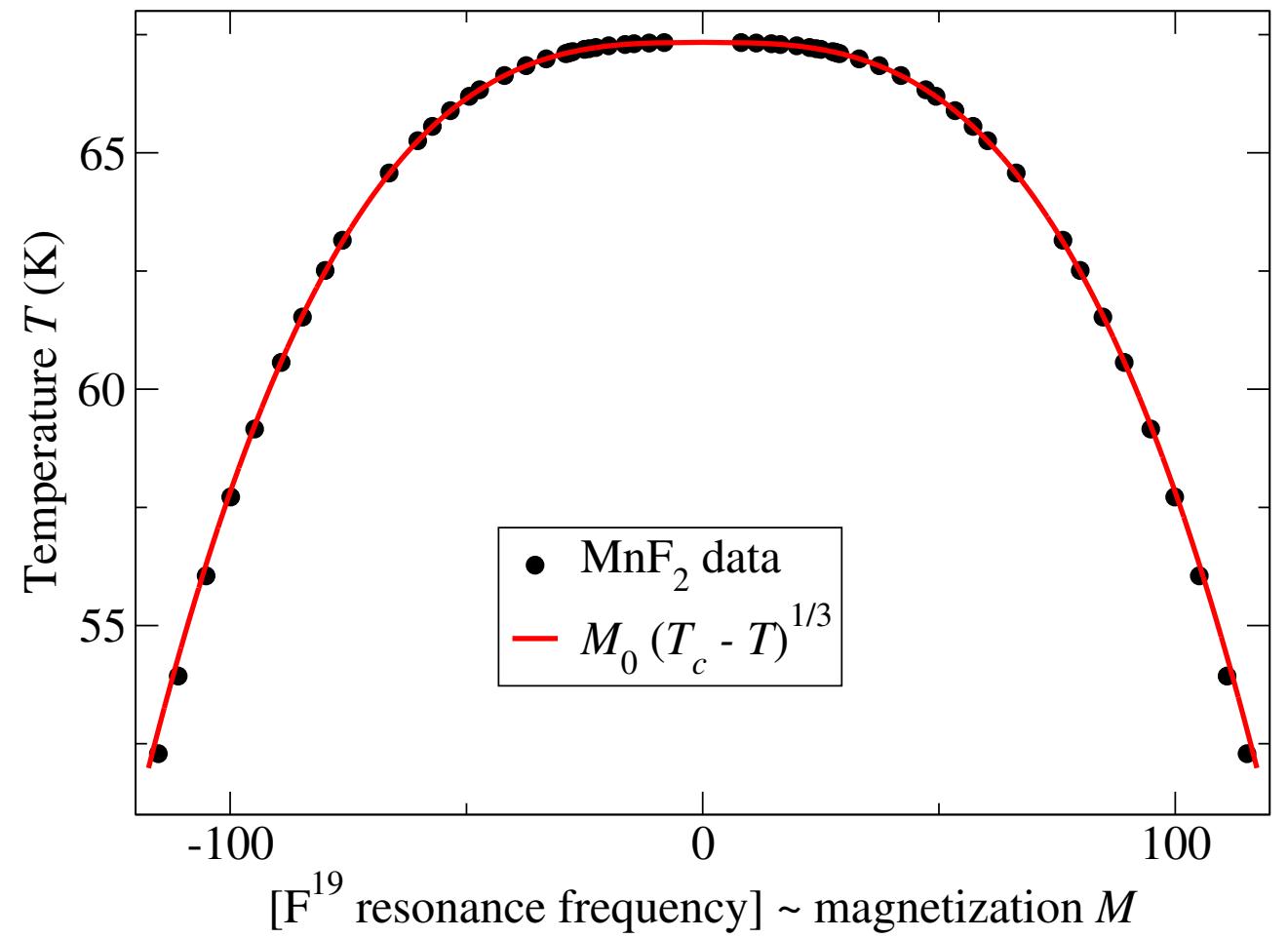
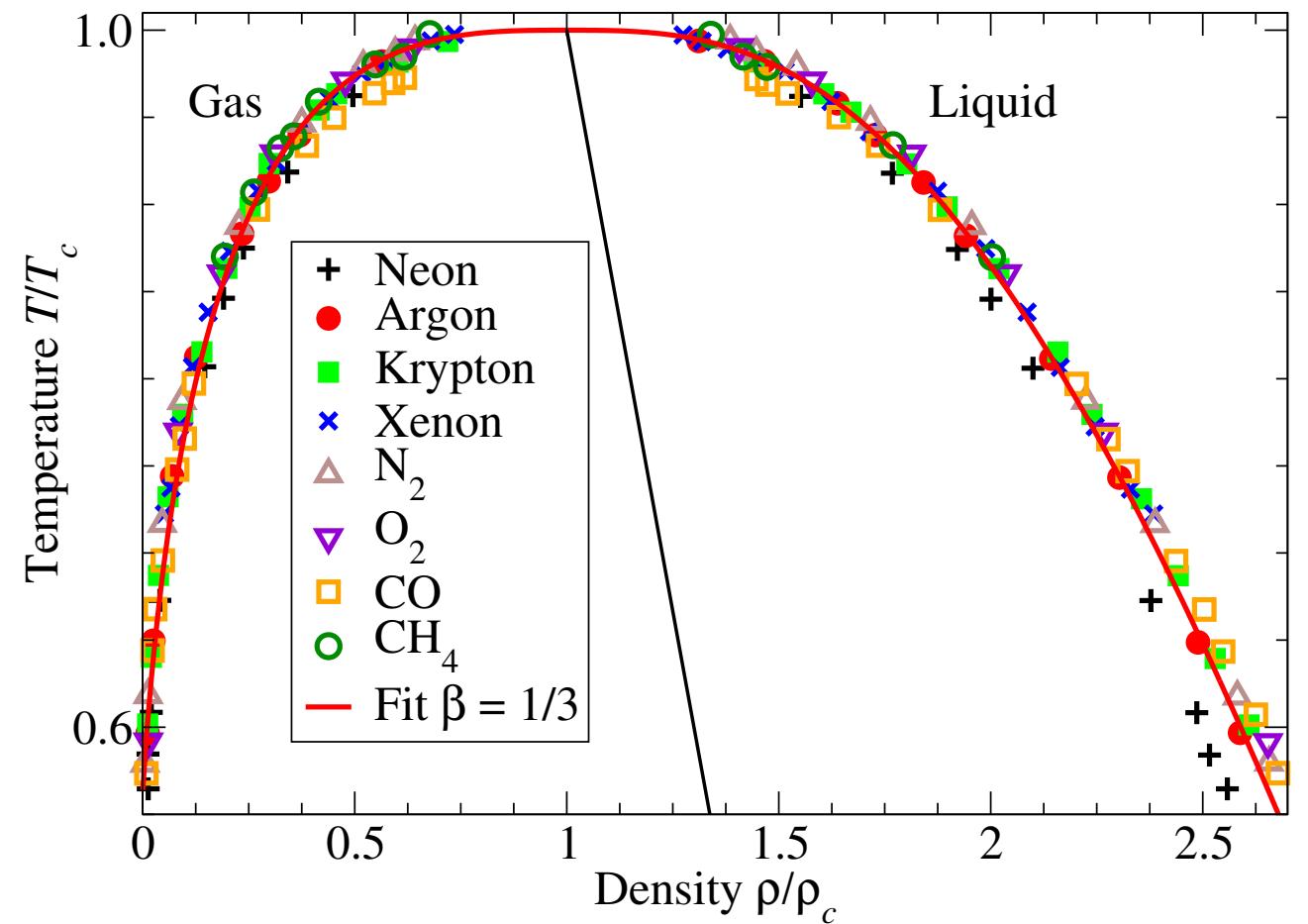
Fit to Argon data:

$$\frac{\Delta\rho}{\rho_c} = \frac{\rho_L - \rho_G}{\rho_c} = 2\rho_0 \left(1 - \frac{T}{T_c}\right)^\beta$$

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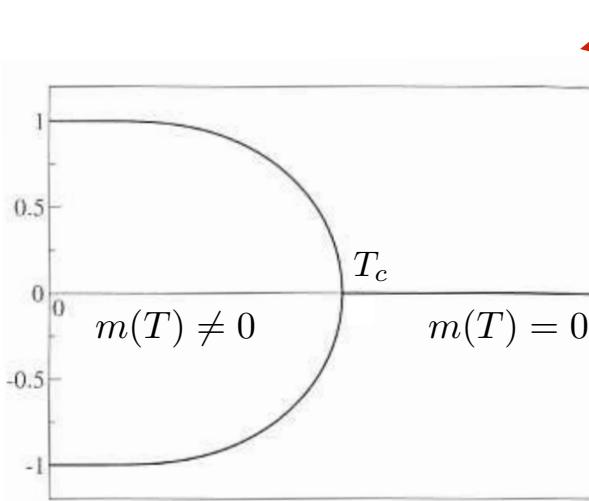
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Ex. Universality: ferromagnetic-paramagnetic critical point

Uniaxial antiferromagnet MnF₂

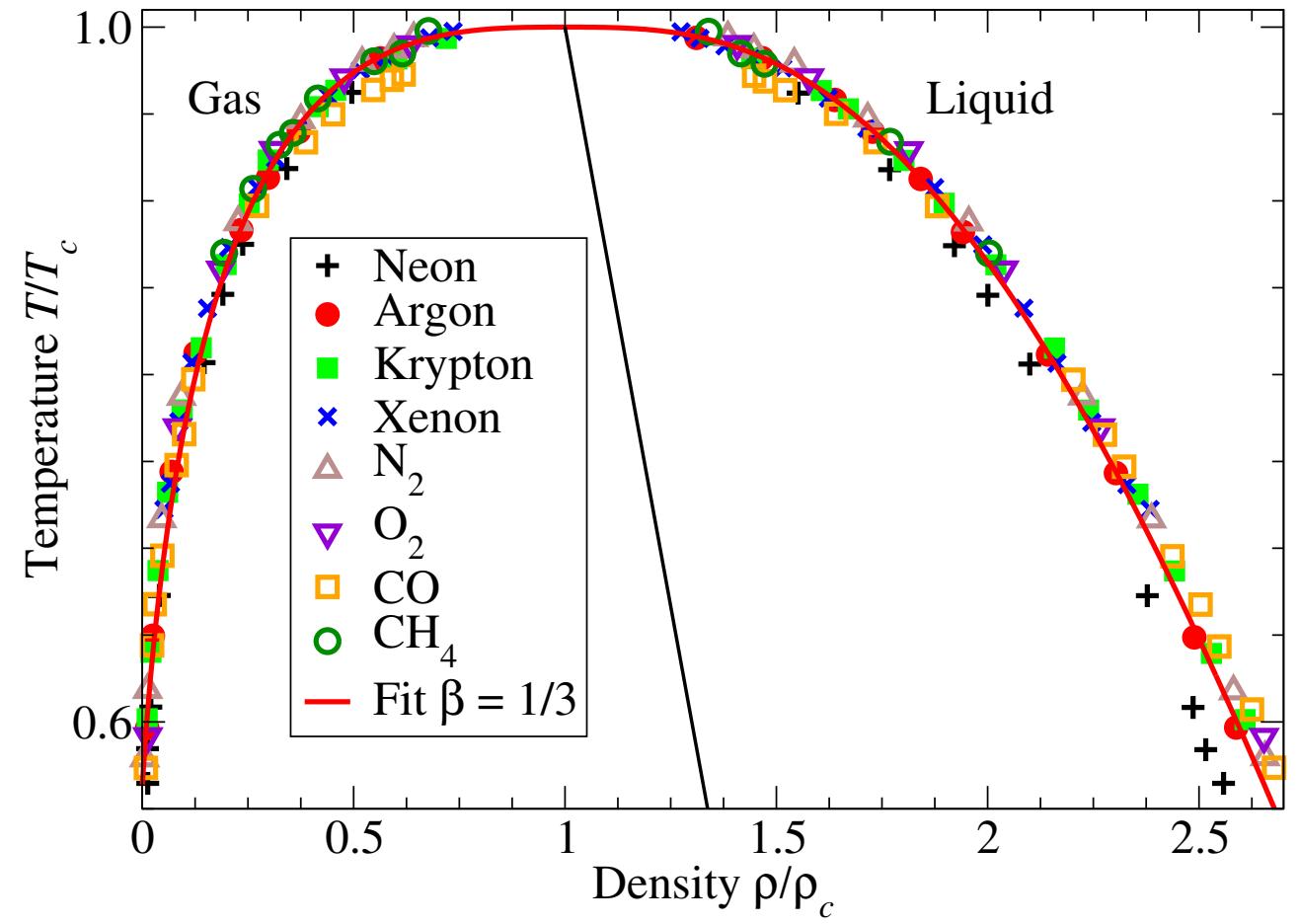
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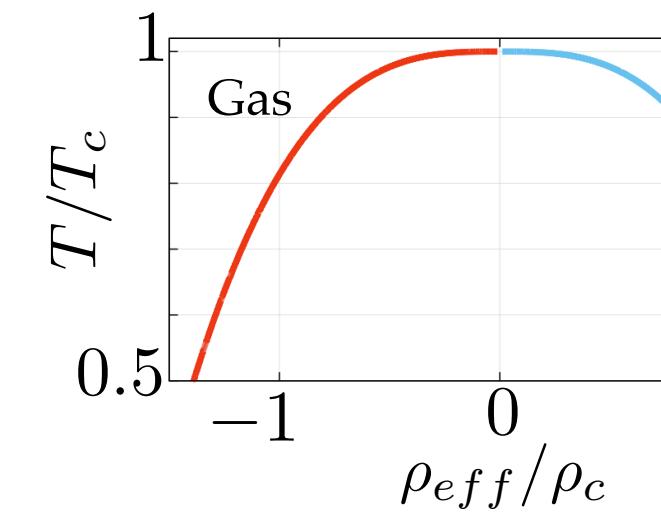
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Tilted →

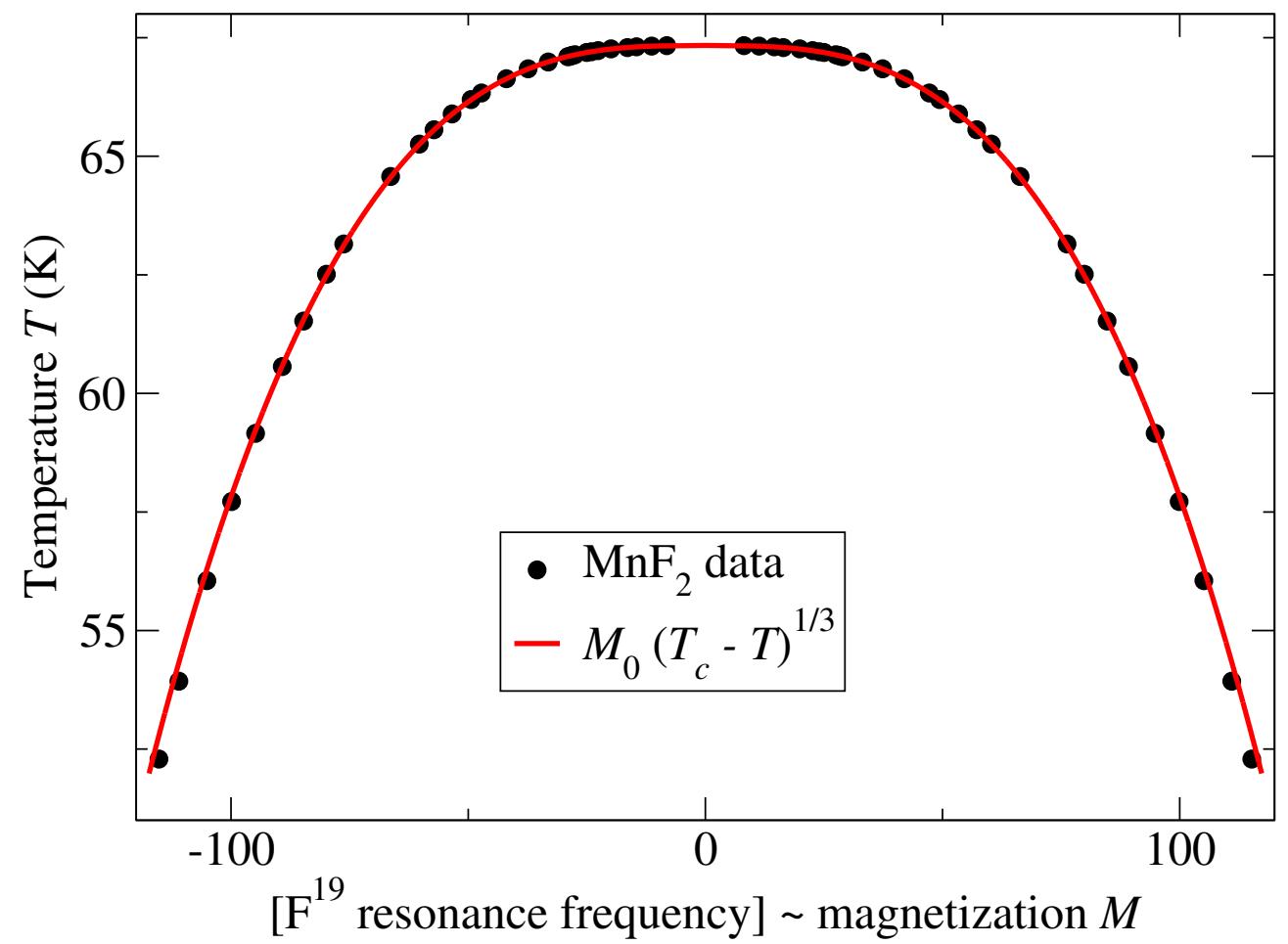
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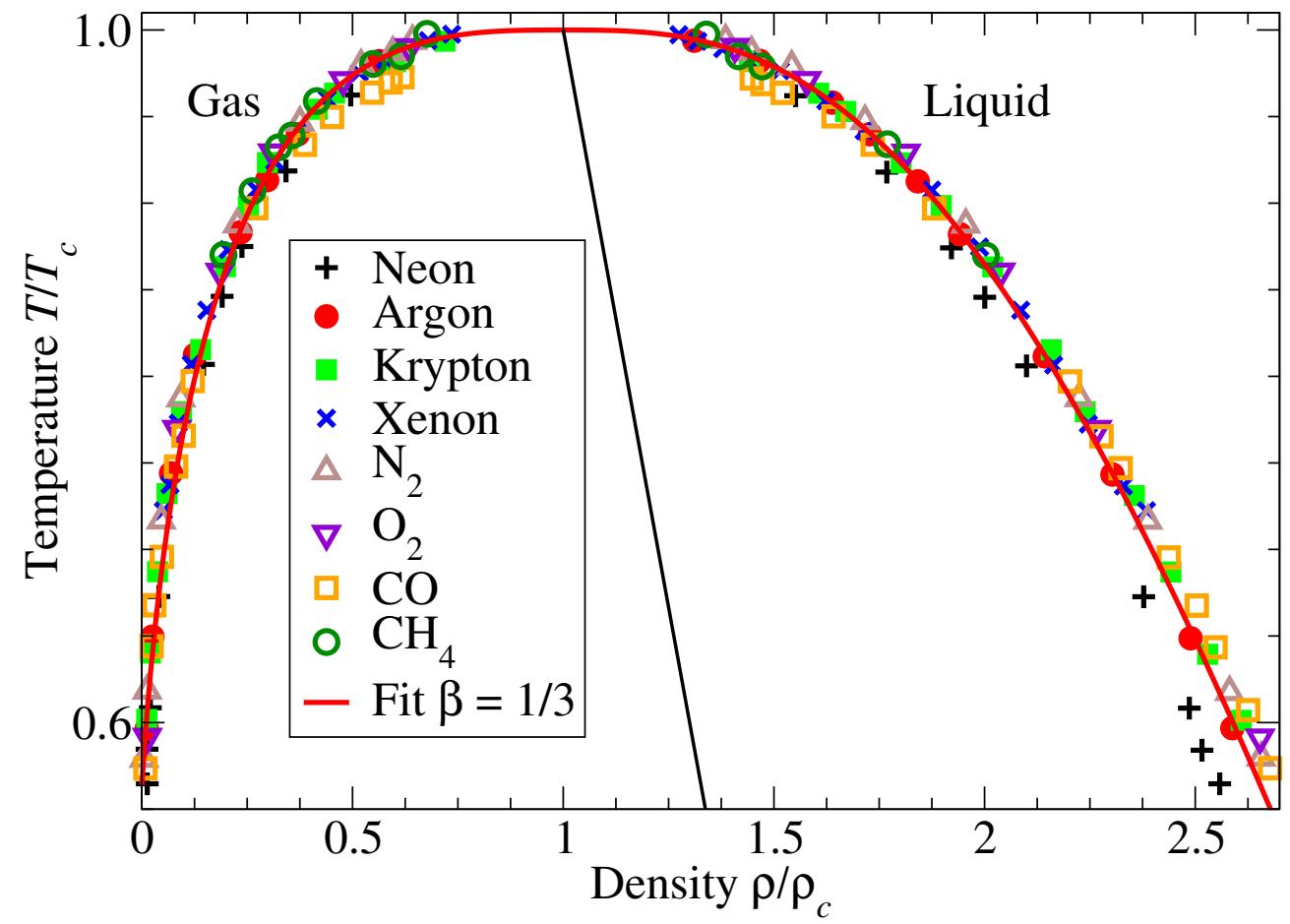
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Current estimate: $\beta = 0.3264\dots$

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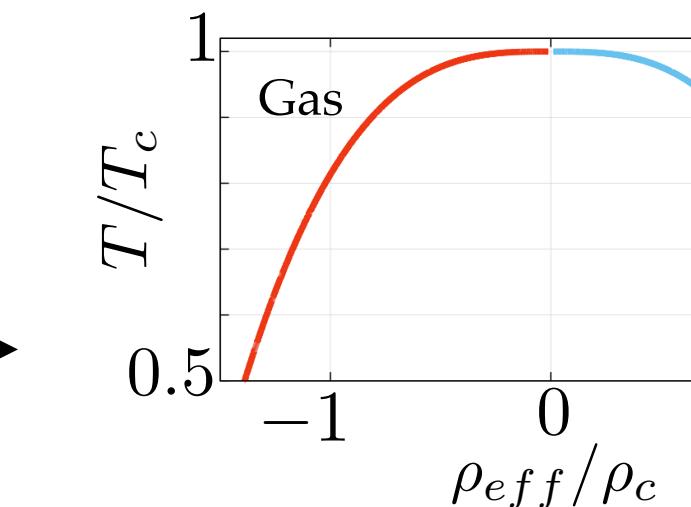
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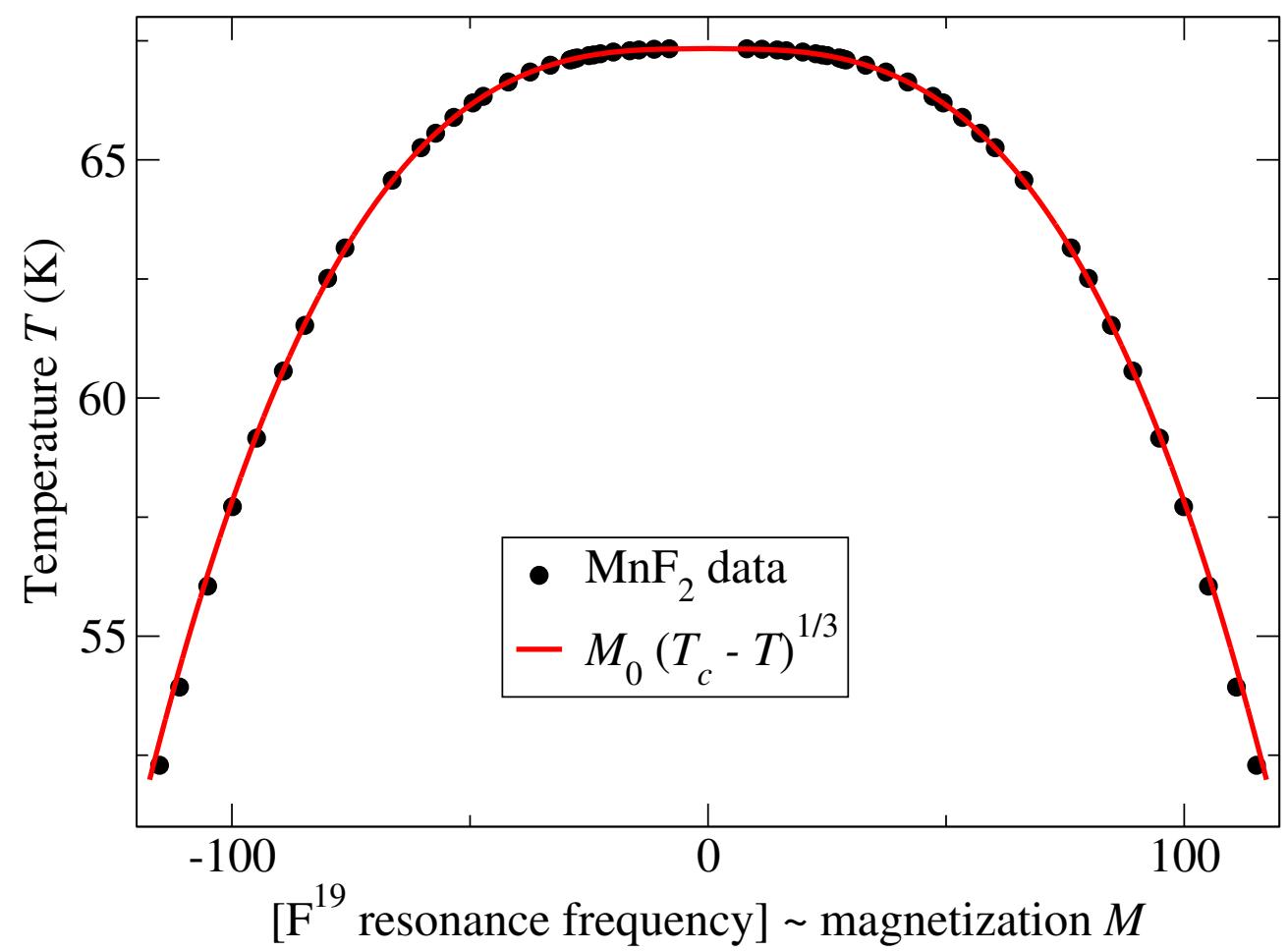
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Tilted



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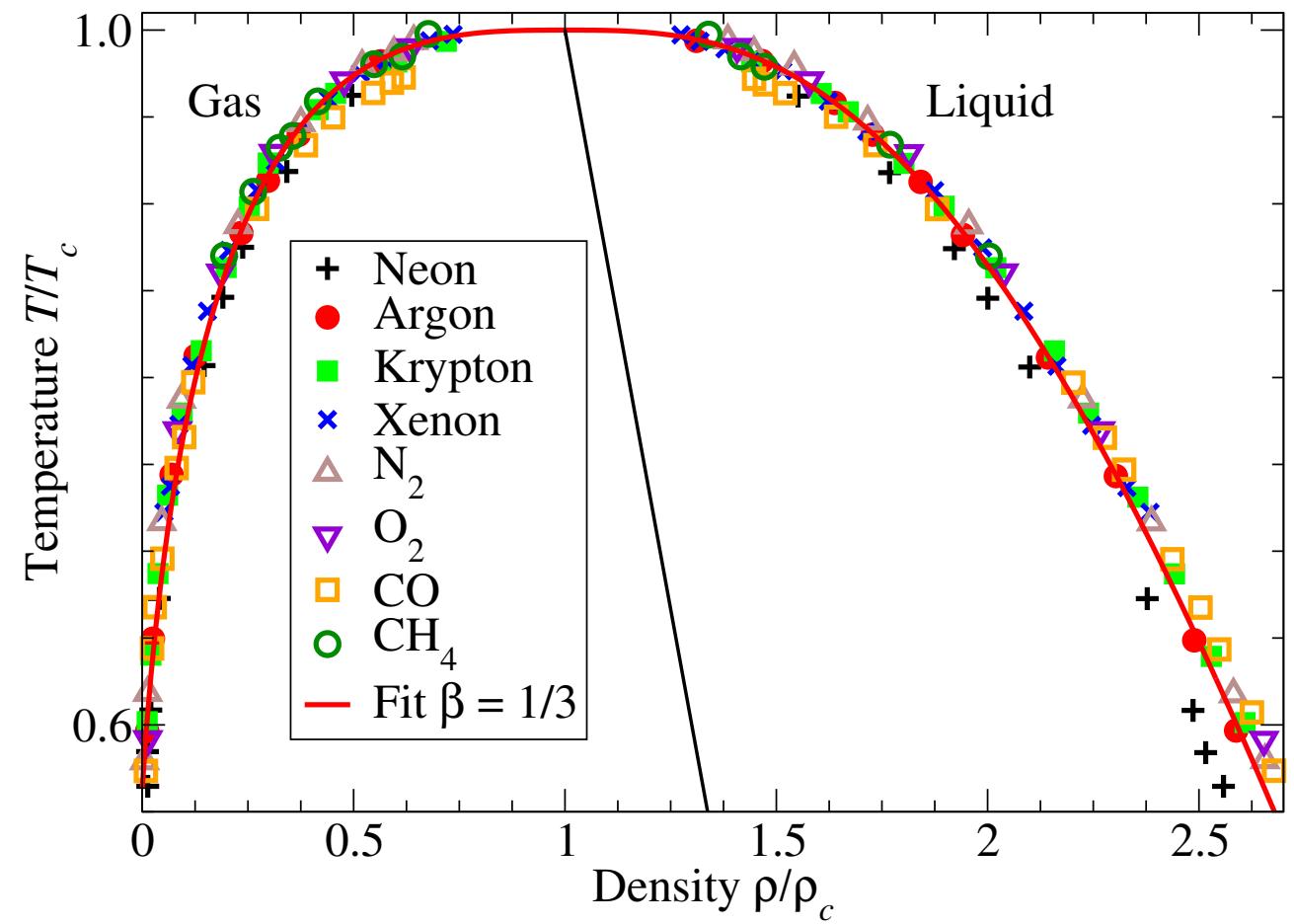
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Critical exponents all identical: at the critical point, correlation length, susceptibility, specific heat have power-law singularities with the same exponents.

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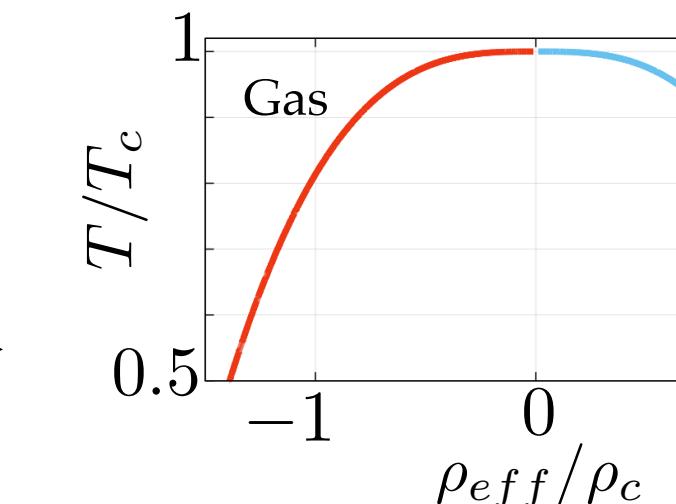
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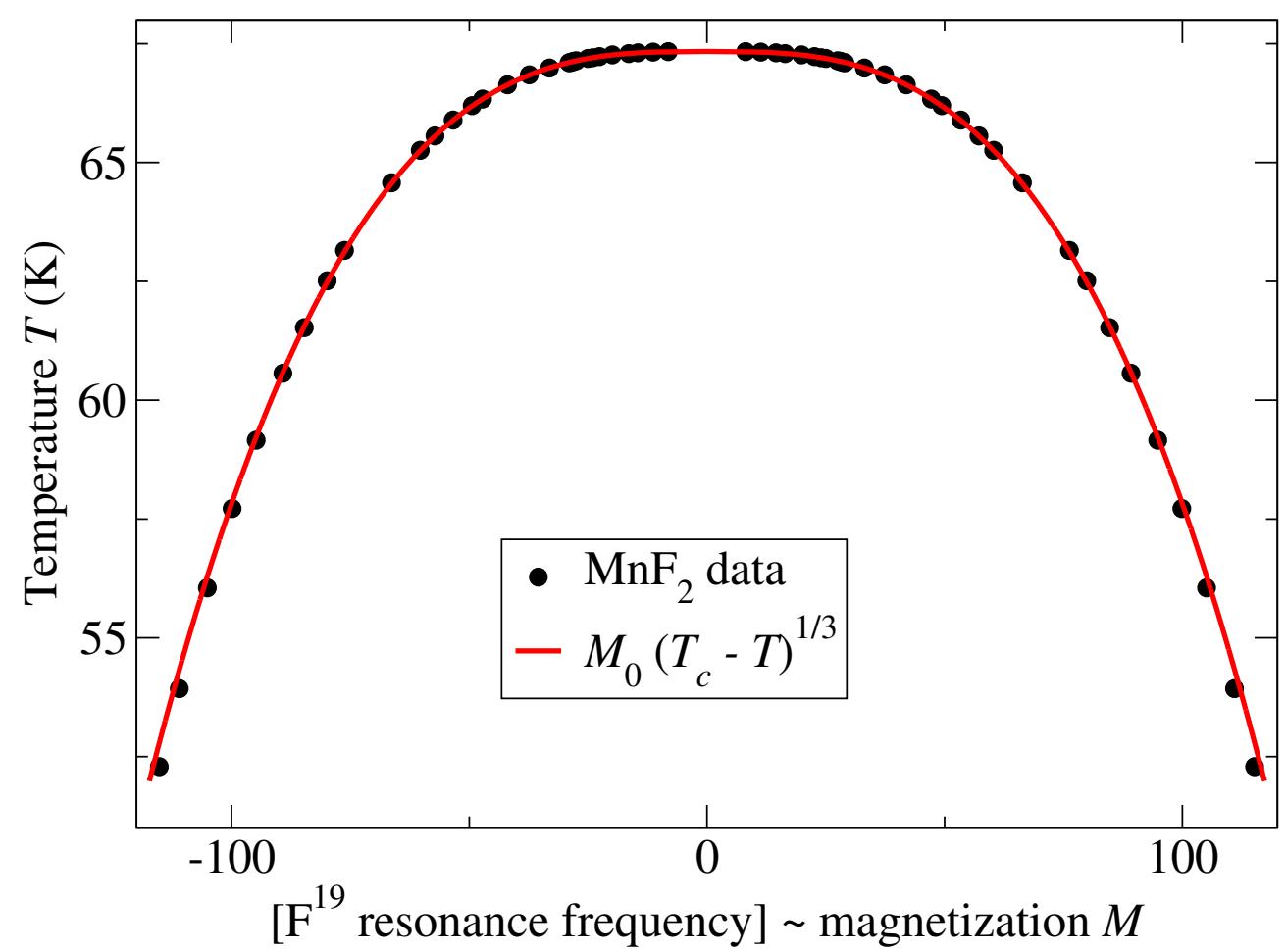
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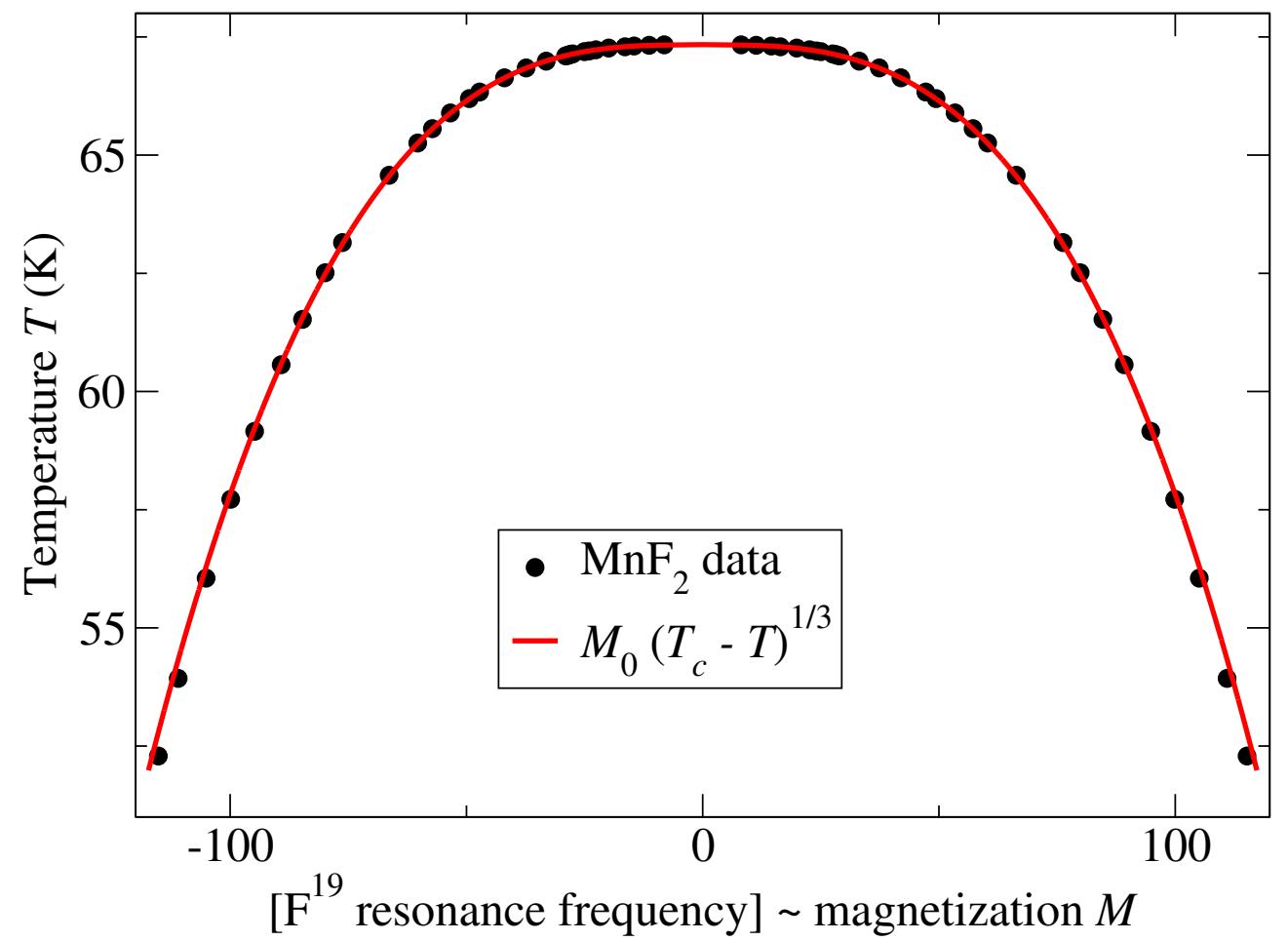
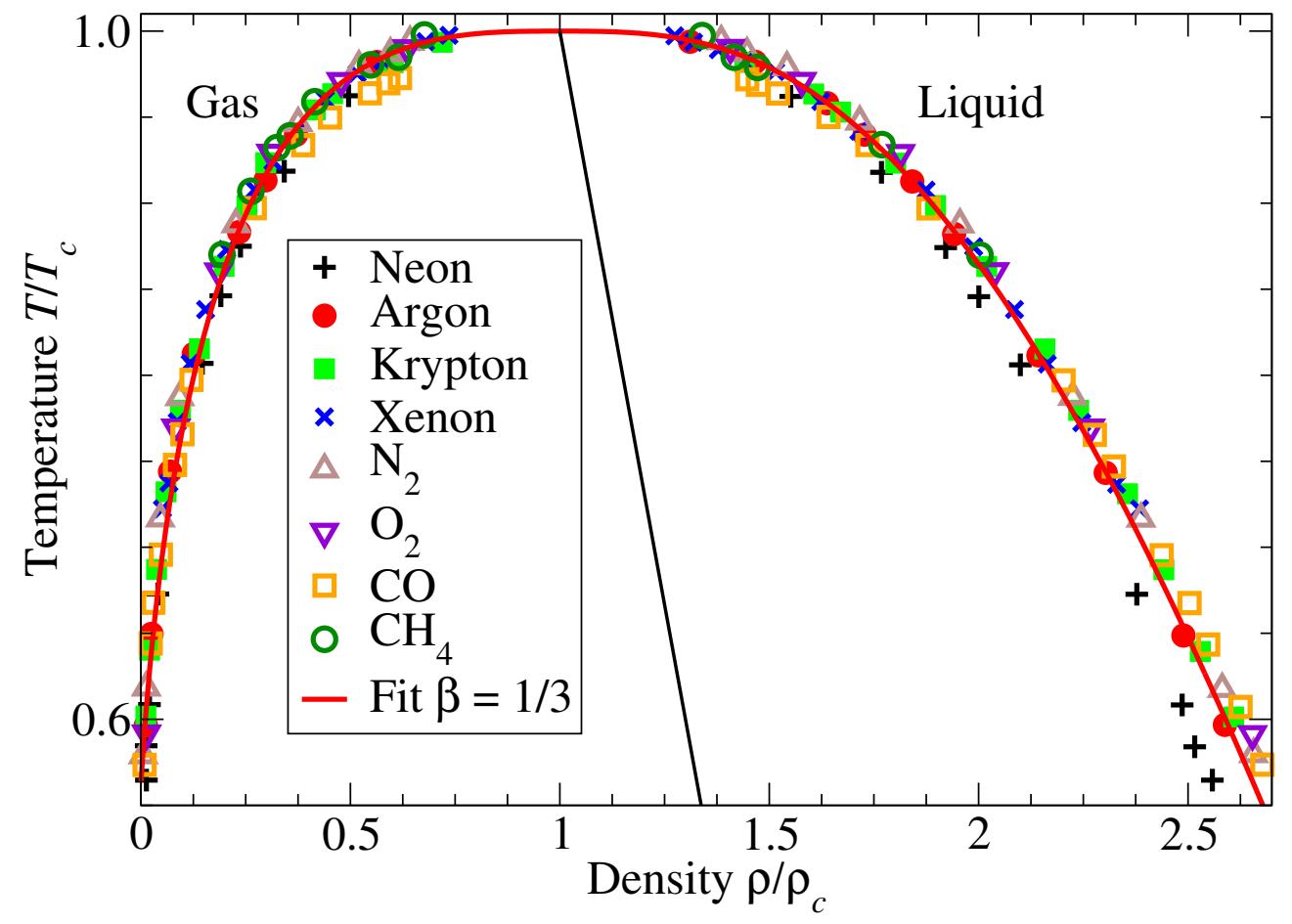
Critical exponents all identical: at the critical point, correlation length, susceptibility, specific heat have power-law singularities with the same exponents.

They are in the **same universality class**, along with the **3D Ising model!**

(despite drastic simplifications!)

Coming up

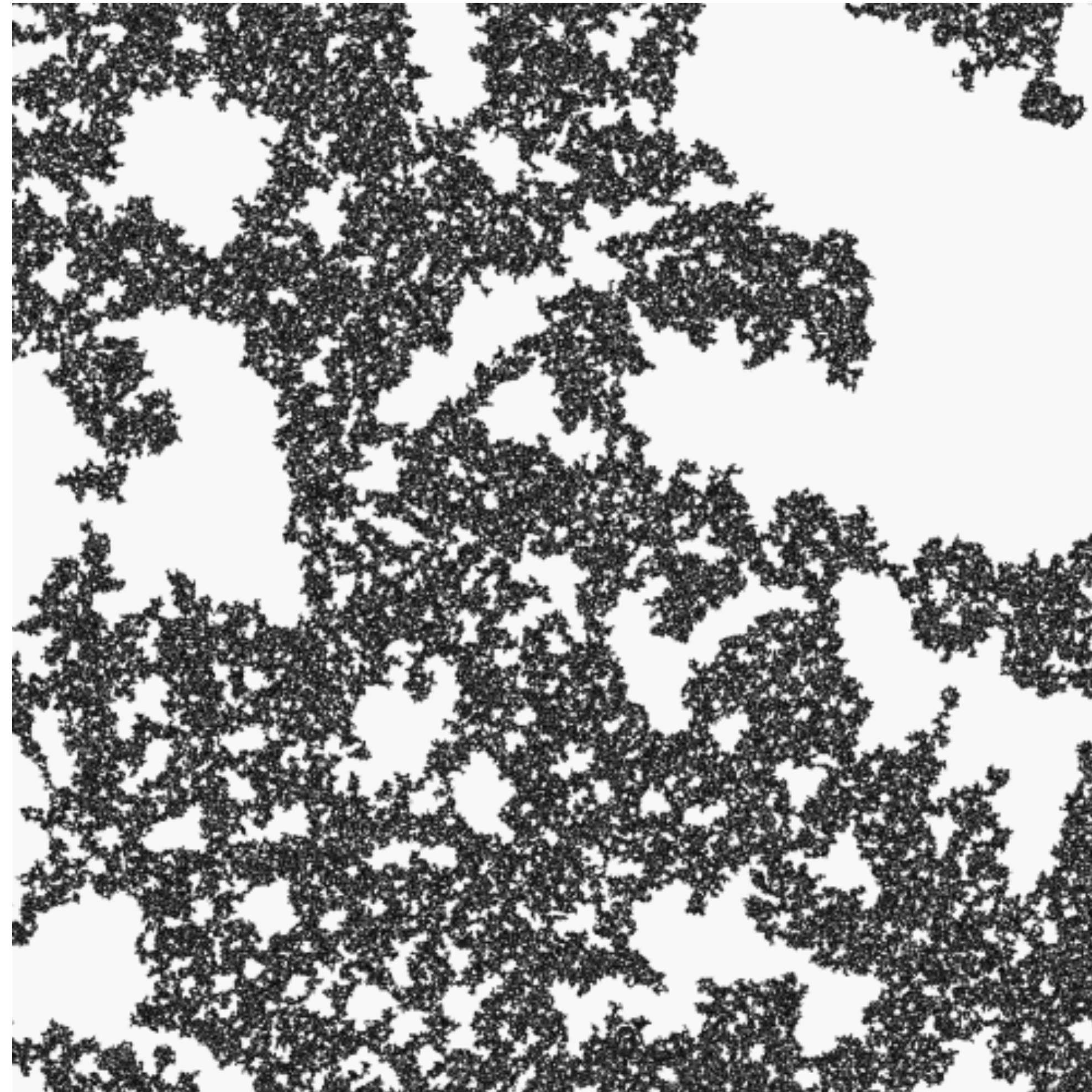
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In the coming chapters: We will discuss

- How to obtain critical exponent for $d > d_u$ using a **mean-field approximation**
- systems belonging to **universality class** of percolation problems: epidemic models

Coming up



CPCS (and others interested) please stay a few more minutes: **Question date for master project bazaar**