

PHYS 326: Lecture 6

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Symmetries

Definition 0.1. A symmetry is defined as a change of coordinates that leaves the lagrangian unchanged. \mathbf{S} is a symmetric matrix if

$$\mathbf{S}^T \mathbf{K} \mathbf{S} = \mathbf{K} \quad (1)$$

$$\mathbf{S}^T \mathbf{M} \mathbf{S} = \mathbf{M} \quad (2)$$

We define a new set of coordinates $\vec{q} = \mathbf{S}\vec{p}$, then

$$L(\vec{q}) = L(\vec{p}) \quad (3)$$

Then

$$PE = \frac{1}{2} \vec{q}^T \mathbf{K} \vec{q} \quad (4)$$

$$= \frac{1}{2} (\mathbf{S}\vec{p})^T \mathbf{K} \mathbf{S}\vec{p} \quad (5)$$

$$= \frac{1}{2} \vec{p}^T \mathbf{K} \vec{p} \quad (6)$$

We have various types of symmetries, namely cyclic and mirror symmetries.

Theorem 0.2. An eigenvector of the system is also an eigenvector of the symmetry matrix.

Then if

$$\mathbf{K}\vec{u} = \omega^2 \mathbf{M}\vec{u} \quad (7)$$

Then

$$\mathbf{S}\vec{u} = \lambda\vec{u} \quad (8)$$

Proof. Given

$$\mathbf{K}\vec{u} = \omega^2 \mathbf{M}\vec{u} \quad (9)$$

Multiply by \mathbf{S}^T and insert $\mathbf{S}\mathbf{S}^{-1}$:

$$\mathbf{S}^T \mathbf{K} \mathbf{S} \mathbf{S}^{-1} \vec{u} = \omega^2 \mathbf{S}^T \mathbf{M} \mathbf{S} \mathbf{S}^{-1} \vec{u} \quad (10)$$

$$\mathbf{K} \mathbf{S}^{-1} \vec{u} = \omega^2 \mathbf{M} \mathbf{S}^{-1} \vec{u} \quad (11)$$

Therefore,

$$\mathbf{S}^{-1} \vec{u} \quad (12)$$

obeys the E.O.M. with same eigenvalue as \vec{u} . In other words, the 2 vectors are same up to a constant. Therefore

$$\mathbf{S}\vec{u} = \lambda\vec{u} \quad (13)$$

□

We look into reflection, mirror, or parity symmetries. When applied twice, they recover the original state. That is $\mathbf{S}^2 = \mathbf{I}$. Then

$$\mathbf{S}\vec{x} = \lambda\vec{x} \quad (14)$$

$$\mathbf{S}^2\vec{x} = \lambda^2\vec{x} \quad (15)$$

Thus

$$\lambda = \pm 1 \quad (16)$$