

PHYS 487: Lecture # 5

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Lecture Span

- Free electron gas and band structure

The free electron gas

Imagine some box of matter, and filling it with electrons. With dimensions L_x, L_y , & L_z . 3D infinite well. Independent, non-interacting electrons.

$$V(x, y, z) = \begin{cases} 0 & \text{inside} \\ \infty & \text{outside} \end{cases} \quad (1)$$

This is a good model for a metal. We first start with

$$-\frac{\hbar^2}{2m}\nabla^2\psi = E\psi \quad (2)$$

$$\psi(x, y, z) = X(x)Y(y)Z(z) \quad (3)$$

Then, we get

$$E = E_x + E_y + E_z \quad (4)$$

Define wavenumbers: $k_i = \sqrt{2mE_i}/\hbar$. Boundary conditions:

$$k_i L_i = n_i \pi \quad (5)$$

Therefore, the wavefunction is

$$\psi_{n_x, n_y, n_z} = \sqrt{\frac{8}{L_x L_y L_z}} \sin(n_x \pi / L_x x) \dots \quad (6)$$

With energies:

$$E_{n_x, n_y, n_z} = \frac{\hbar^2 \pi^2}{2m} \left(\frac{n_x^2}{L_x^2} + \frac{n_y^2}{L_y^2} + \frac{n_z^2}{L_z^2} \right) = \frac{\hbar^2 k^2}{2m} \quad (7)$$

Find the highest energy (fermi energy): Note that $|E| \sim |k|^2$. Consider a coordinate system with axis k_x, k_y, k_z . Then the distance between each k vector is π/L_i . Therefore, the volume per state is

$$V \sim \frac{\pi^3}{L_x L_y L_z} \quad (8)$$

Here, we find

$$\frac{1}{8} \left(\frac{4}{3} \pi k_f^3 \right) = \frac{N_e}{2} \frac{\pi^3}{V} \quad (9)$$

Then, we get

$$k_F = (3\rho\pi^2)^{1/3} \quad (10)$$

With

$$E_f = \frac{\hbar^2}{2m} (3\rho\pi^2)^{2/3} \quad (11)$$

Physically, the energy levels depend on the magnitude of k , which can be interpreted as a sphere in k space, where all the surface has the same k and therefore the same energy. Since we have a finite number of electrons, we find the maximal energy of the system. Here, using this volume approximation, we find that the maximal volume of the vector in k space of just one eighth of the vector is $N_e/2 \times \pi^3/V$. Doing some algebra, where $1/8$ represents just $1/8$ of the volume, we can find k .

Electron Shell

Consider a shell in k space. But how many states are in between $[k_i, k_i + dk]$. Consider a shell

$$V_{shell} = \frac{1}{8} (4\pi k^2) dk \quad (12)$$

Then, the number of energy states in this shell is (dividing this shell by the volume differential)

$$dN = \frac{\frac{1}{2}\pi k^2}{\pi^3/V} dk \quad (13)$$

Then, accounting for the pauli exclusion principle

$$dN = \frac{V k^2}{\pi^2} dk \quad (14)$$

Then

$$dE = \frac{\hbar^2 k^2}{2m} dN \quad (15)$$

It follows that

$$dE = \frac{\hbar^2 k^2}{2m} \frac{V}{\pi^2} k^2 dk \quad (16)$$

We can then integrate this in k space:

$$\int_{E_f}^0 dE = \int_0^{k_f} dk \frac{\hbar^2 k^2}{2m} \frac{V}{\pi^2} k^2 \quad (17)$$

$$= \frac{\hbar^2 (3\pi^2 N_e)^{5/3}}{10\pi^2 m} V^{-2/3} \quad (18)$$

Note, although the electrons are discrete with k , if $N_e \rightarrow \infty$, then we can approximate this with an integral. Consider

$$dE = -\frac{2}{3} \hbar^2 \frac{(2\pi^2 N_e)^{5/3}}{10\pi^2 m} V^{-5/3} dV = -\frac{2}{3} E_{tot} \frac{dV}{V} \quad (19)$$

Thermodynamics connection: $dW = PdV$, here $P = (2/3)E_{tot}/V$, this means that particles tend to stay escape (pressure pushes out).

Simple model of a crystal

Dirac comb

$$V(x) = \alpha \sum_{j=0}^N \delta(x - ja) \quad (20)$$

Then, we use Bloch's theorem which states that the wavefunction should be

$$\psi(x) = \exp(iqx)u(x) \quad (21)$$

and

$$\psi(x - a) = \exp(-iqa)\psi(x) \quad (22)$$

If we do a cyclical boundary condition,

$$\psi(x) = \exp(iNqa)\psi(x) \quad (23)$$

Therefore,

$$q = \frac{2\pi n}{Na} \quad (24)$$

In $(0, a)$,

$$\frac{\hbar^2}{2m}\psi'' = E\psi'' \iff \psi'' = -k^2\psi \quad (25)$$

Then

$$\psi(x) = A \sin(kx) + B \cos(kx) \quad (26)$$

And using Bloch's theorem, from $(-a, 0)$

$$\psi(x) = \exp(-iqa)(A \sin(kx) + B \cos(kx)) \quad (27)$$

Boundary condition at $x = 0$: ψ is continuous but not necessarily ψ' . Then, the boundary condition states that the constant B must be:

$$B = \exp(-iqa)(A \sin(kx) + B \cos(ka)) = \psi(0) \quad (28)$$

$$-\frac{\hbar^2}{2m} \int_{-\epsilon}^{\epsilon} \psi''(x) dx + \int_{-\epsilon}^{\epsilon} V(x)\psi(x) dx = E \int_{-\epsilon}^{\epsilon} \psi(x) dx \quad (29)$$

Take the limit as $\epsilon \rightarrow 0$

$$-\frac{\hbar^2}{2m} \Delta\psi' + \psi(0) = 0 \quad (30)$$

Then given $\psi'_R(0)$ and $\psi'_L(0)$, we get that $\psi'(0)$ is

$$KA - \exp(-iqa)k(A \cos(ka) - B \sin(ka)) = \frac{2m\alpha}{\hbar^2} B \quad (31)$$

Then, we can obtain

$$\cos(qa) = \cos(ka) + \frac{m\alpha}{\hbar^2 k} \sin(ka) \quad (32)$$

Note that because of the LHS, we have that

$$|\cos(ka) + \frac{m\alpha}{\hbar^2 k} \sin(ka)| \leq 1 \quad (33)$$

This forms bands on the ψ, ka space.