

# PHYS 486: Lecture # 19

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## Recap: Harmonic Oscillator

The Hamiltonian is

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2} m \omega^2 x^2 \quad (1)$$

Introduced a couple of new oscillators:

$$a_{\pm} = \frac{1}{\sqrt{2\hbar m \omega}} (\mp i \hat{p} + m \omega \hat{x}) \quad (2)$$

This means that

$$\hat{H} = (a_+ a_- + \frac{1}{2}) \hbar \omega \quad (3)$$

Note that

$$|0\rangle, |1\rangle = a_+ |0\rangle, \dots \quad (4)$$

The energy difference is always  $\hbar \omega$ . Note that  $a_+$  (creation operator) and  $a_-$  (annihilation operator) are always called ladder operators. Solving

$$\hat{H} |0\rangle = \underbrace{\frac{\hbar \omega}{2}}_{E_0} |0\rangle$$

This  $E_0$  is interpreted as some sort of quantum noise. This is inevitable. Introduce

$$a_+ a_- = \frac{H}{\hbar \omega} - \frac{1}{2} = \hat{n} \text{ (number operator)} \quad (5)$$

This is the number of energy levels there are. Note that

$$\langle n | \hat{H} | n \rangle = \frac{1}{2} \hbar \omega + n \hbar \omega \quad (6)$$

Then

$$\langle a_+ a_- \rangle = n \quad (7)$$

Since this value is Hermitian, this means that it is measurable. We compute

$$\langle n | a_+ a_- | n \rangle \neq 0 \quad (8)$$

Since we are in an orthonormal eigenbasis, and this inner product is  $\neq 0$ , we can infer that

$$(|a_- n\rangle)^\dagger = (a_- |n\rangle)^\dagger = \langle n | (a_-)^\dagger = \langle n | a_+ \quad (9)$$

Therefore,

$$a_-^\dagger = a_+ \quad (10)$$

## Normalization

We know that  $\langle n|n \rangle = 1$ . But how do we normalize this? We show that:

$$\begin{aligned}
 a_+ |n\rangle &= c_{n+1} |n+1\rangle \\
 \langle n|a_- a_+ |n\rangle &= \langle n+1| c_{n+1}^* c_{n+1} |n+1\rangle \\
 &= \langle n| \frac{\hat{H}}{\hbar\omega} + \frac{1}{2} |n\rangle \\
 &= \langle n| \frac{n\hbar\omega}{\hbar\omega} + \frac{1}{2} |n\rangle \\
 &= n + 1
 \end{aligned}$$

This means that  $c_{n+1} = \sqrt{n+1}$ . This is nothing but a mathematical necessity. We repeat with  $a_-$ :

$$\begin{aligned}
 a_- |n\rangle &= c_{n-1} |n-1\rangle \\
 c_{n-1} &= \sqrt{n-1}
 \end{aligned}$$

### Brief Summary:

1.  $|0\rangle$  is the ground state
2.  $a_+$ :  $a_+ |n\rangle = \sqrt{n+1} |n+1\rangle$
3.  $a_-$ :  $a_- |n\rangle = \sqrt{n-1} |n-1\rangle$
4.  $a_+^\dagger = a_-$
5.  $\hat{n} = a_+ a_-$
6.  $\frac{\hat{H}}{\hbar\omega} = a_+ a_- + \underbrace{\frac{1}{2}}_{ZPE}$  (ZPE = zero point energy)
7.  $[a_-, a_+] = 1$

S.S in the position basis. The boundary conditions are

$$a_- \psi_0(x) = \frac{1}{\sqrt{2\hbar m\omega}} \left( \hbar \frac{\partial}{\partial x} + m\omega x \right) \psi_0(x) = 0 \quad (11)$$

This means that

$$\frac{\partial}{\partial x} \psi_0(x) = -\frac{m\omega x}{\hbar} \psi_0(x) \quad (12)$$

$$\implies \psi_0(x) = A \exp\left(-\frac{m\omega x^2}{2\hbar}\right) \quad (13)$$

The phase space of the position and momentum is aligned. We first normalize.

$$\langle \psi_0(x) | \psi_0(x) \rangle \implies A = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} \quad (14)$$

We find  $\psi_1(x)$ :

$$a_+ \psi_0 = \sqrt{0+1} \psi_1 \quad (15)$$

$$\psi_1(x) = \frac{1}{\sqrt{2\hbar m\omega}} \left( -\hbar \frac{\partial}{\partial x} + m\omega x \right) \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} \exp\left(-\frac{m\omega}{2\hbar} x^2\right) \quad (16)$$

Following an iterative procedure, we obtain

$$\psi_n = \left( \prod_{k=1}^n \frac{a_+}{\sqrt{k}} \right) \psi_0 = \frac{1}{\sqrt{n!}} a_+^n \psi_0 \quad (17)$$

$$\psi_n(x) = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n(x) e^{-x^2/2} \quad (18)$$