

PHYS 487 Lecture # 1

Cliff Sun

January 20, 2026

PHYS 487 Plan

We will use techniques and examples to analyze realistic simple systems.

Outline

1. Symmetries + conservation laws
 - (a) Central symmetric problems (rotational symmetries)
 - (b) Periodic potentials (crystal)
2. Time-independent Perturbation Theory
 - (a) Fine-structure of hydrogen
3. Dynamics
 - (a) Emission and absorption of light
4. Multi-particle quantum systems: approximations to place boundaries
 - (a) Ground state energies estimation
5. Perhaps some bonus topics on some cool stuff

QM Speedrun Recap

1. Physical system $|\psi\rangle \in \mathcal{H}$
2. Linearity/Superposition principle $|\psi\rangle = \sum_k c_k |\psi_k\rangle$ where $|\psi_k\rangle$ are the basis vectors of the same operator.
3. Inner product: $\langle\varphi|\psi\rangle = \langle\psi|\varphi\rangle^*$ such that $\langle\psi|\psi\rangle \leq 0$
4. Observables: \hat{O} operators; Linearity $\hat{O}|\psi\rangle = \hat{O}(\sum_k c_k |\psi_k\rangle) = \sum_k c_k (\hat{O}|\psi_k\rangle)$; Hermitian: $\hat{O}^\dagger = \hat{O}$, defined as $\langle\psi|\hat{O}^\dagger\psi\rangle = \langle\hat{O}\psi|\psi\rangle$
5. Eigenvalues/vectors via diagonalization: $\hat{O}|\alpha_n\rangle = a_n |\alpha_n\rangle$. Where a_n are the possible measurement outcomes and $|\alpha_n\rangle$ is a complete basis set for $|\psi\rangle$.
6. Born's Rule: $|c_k|^2$ is the probability for outcome c_k .
7. Wavefunction Collapse: $|\psi\rangle \rightarrow |\psi_k\rangle$ after measuring.
8. Incompatibility: Observables A, B . If $[A, B] \neq 0$, then there doesn't exist an eigenbasis where both A and B are simultaneously diagonal. This means that the measurement of A will disturb that of B and vice versa.
9. Uncertainty principle: recall $\langle A \rangle = \langle\psi|A\psi\rangle$ and $\sigma_A^2 = \langle A^2 \rangle - \langle A \rangle^2$, then

$$\sigma_A^2 \sigma_B^2 \leq \left(\frac{1}{2i} \langle [A, B] \rangle \right)^2 \quad (1)$$

Example Systems

1. Continuous systems: $\mathcal{H} = L^2$. $\langle \psi | \varphi \rangle = \int dx \varphi(x)^* \psi(x)$. In the 1-d position basis, $\psi(x) = \langle x | \psi \rangle$

When expressing an operator acting on a function in the position basis, consider (note that $\langle x | \hat{x} \psi \rangle$ can be thought of as (first act \hat{x} on ψ , then find this vector in the position basis))

$$\begin{aligned}\langle x | \hat{x} \psi \rangle &= \int dx' \langle x | \hat{x} | x' \rangle \langle x' | \psi \rangle \\ &= \int dx' x' \delta(x - x') \psi(x) \\ &= x' \psi(x)\end{aligned}$$

The expectation value is

$$\int dx \psi^*(x) x \psi(x) \quad (2)$$

Momentum is

$$\langle x | \hat{p} | \psi \rangle = -i\hbar \frac{\partial}{\partial x} \psi(x) \quad (3)$$

This is interpreted as acting the momentum operator on ψ (abstract), then making it more concrete by turning this abstract concept into the position basis. Then, the Hamiltonian is

$$\hat{H} = \frac{\hat{p}^2}{2m} + V(x) \quad (4)$$

Discrete Systems

$\mathcal{H} = \mathbb{C}^n$,

$$\vec{c} = \begin{pmatrix} c_1 \\ c_2 \\ \vdots \\ c_n \end{pmatrix}$$

And $\langle c | d \rangle = \sum_k c_k^* d_k$

Simplest Physical System - Spin 1/2

$$|\chi\rangle = c_1 \underbrace{| \uparrow \rangle}_{\begin{pmatrix} 1 \\ 0 \end{pmatrix}} + c_2 \underbrace{| \downarrow \rangle}_{\begin{pmatrix} 0 \\ 1 \end{pmatrix}} \quad (5)$$

Recall that the electron is $|\psi\rangle |\chi\rangle$. Operators for spin-1/2:

$$S_z = \frac{\hbar}{2} \sigma_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad (6)$$

S_x, S_y are defined similarly.

Angular momentum

$$L = r \times p \quad (7)$$

Where $r = \begin{pmatrix} \hat{x} \\ \hat{y} \\ \hat{z} \end{pmatrix}$ and $p = -i\hbar \begin{pmatrix} \partial_x \\ \partial_y \\ \partial_z \end{pmatrix}$. For angular momentum, $[J_i, J_j] = i\hbar \epsilon_{ijk} J_k$. Eigenstates $|j, m\rangle$ where j corresponds to the total angular momentum (J^2) and m is the z-projection of the angular momentum J_z . Then

$$J^2 |j, m\rangle = \hbar^2 j(j+1) |j, m\rangle \quad (8)$$

$$J_z |j, m\rangle = \hbar m |j, m\rangle \quad (9)$$

Note, $j = 0, 1/2, 1, 3/2, \dots$ and $m = -j, -j+1, \dots, j$. Ladder operators $J_{\pm} = J_x \pm iJ_y$ raise and lower m by ± 1 .

Dynamics

$$i\hbar\partial_t |\psi\rangle = \hat{H} |\psi\rangle \quad (10)$$

And the stationary states are the eigenstates of the Hamiltonian: $\hat{H} |\psi_n\rangle = E_n |\psi_n\rangle$ where $|\psi_n\rangle$ is the stationary state. Given $|\psi(0)\rangle$, then

1. Find $|\psi_n\rangle$
2. $|\psi(0)\rangle = \sum_n c_n |\psi_n\rangle$
3. $|\psi(t)\rangle = \sum_n c_n e^{-iE_n t/\hbar} |\psi_n\rangle$

Examples to recall

1. 1D ∞ square well
2. QHO
3. 3D coulomb potential
4. Spin 1/2 (qubit)