

PHYS 326: Lecture 4

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Multiple D.O.F: The general case

Consider a system with a generalized coordinates $q_\alpha(t)$ for $\alpha = 1, 2, \dots, N$. Then

$$L = L(q_\alpha, \dot{q}_\alpha) \quad (1)$$

We want to find the equations of motion for small deviations from equilibrium. At equilibrium, $q_\alpha = c$.

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_\alpha} \right) - \frac{\partial L}{\partial q_\alpha} = 0 \quad (2)$$

At equilibrium, we have that nothing changes in time, therefore:

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_\alpha} \right) \Rightarrow \frac{\partial L}{\partial q_\alpha} = 0 \quad (3)$$

We study small deviations from equilibrium:

$$q_\alpha = \bar{q}_\alpha + x_\alpha(t) \quad (4)$$

For x_α small and \bar{q}_α is the equilibrium point. We plug this into the lagrangian, and we only retain up to the 2nd order:

$$L(x_1, x_2, \dots, \dot{x}_1, \dot{x}_2, \dots) = \frac{1}{2} \sum_{\alpha, \beta=1}^N M_{\alpha\beta} \dot{x}_\alpha \dot{x}_\beta + \frac{1}{2} \sum_{\alpha, \beta=1}^N \Gamma_{\alpha\beta} \dot{x}_\alpha x_\beta - \frac{1}{2} \sum_{\alpha, \beta=1}^N K_{\alpha\beta} x_\alpha x_\beta \quad (5)$$

In classical physics, Γ is usually 0. Then we plug into E.L. equations:

$$= \sum_{\beta} \left(\frac{M_{\gamma\beta} + M_{\beta\gamma}}{2} \ddot{x} \right) + \sum_{\beta} \left(\frac{\Gamma_{\gamma\beta} + \Gamma_{\beta\gamma}}{2} \dot{x} \right) + \sum_{\beta} \left(\frac{K_{\gamma\beta} + K_{\beta\gamma}}{2} x \right) \quad (6)$$

$$-\frac{\partial L}{\partial x_\gamma} = -\frac{\partial}{\partial x_\gamma} [L] \quad (7)$$

Note that $\frac{\partial x_\mu}{\partial x_\nu} = \delta_{\mu\nu}$. See lecture notes for derivation. But the key thing is that

$$M_{\alpha\beta} = \frac{\partial}{\partial x_\alpha \partial x_\beta} T \quad (8)$$

Since α and β are dummy variables, we have that M must be symmetric. In a similar fashion, K is also symmetric.

Cart + Spring + Pendulum