

PHYS 486: Lecture # 19

Cliff Sun

October 9, 2025

Recap: Harmonic Oscillator

The Hamiltonian is

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2} m\omega^2 x^2 \quad (1)$$

Introduced a couple of new oscillators:

$$a_{\pm} = \frac{1}{\sqrt{2\hbar m\omega}} (\mp i\hat{p} + m\omega\hat{x}) \quad (2)$$

This means that

$$\hat{H} = (a_+ a_- + \frac{1}{2})\hbar\omega \quad (3)$$

Note that

$$|0\rangle, |1\rangle = a_+ |0\rangle, \dots \quad (4)$$

The energy difference is always $\hbar\omega$. Note that a_+ (creation operator) and a_- (annihilation operator) are always called ladder operators. Solving

$$\hat{H} |0\rangle = \underbrace{\frac{\hbar\omega}{2}}_{E_0} |0\rangle$$

This E_0 is interpreted as some sort of quantum noise. This is inevitable. Introduce

$$a_+ a_- = \frac{H}{\hbar\omega} - \frac{1}{2} = \hat{n} \text{ (number operator)} \quad (5)$$

This is the number of energy levels there are. Note that

$$\langle n | \hat{H} | n \rangle = \frac{1}{2} \hbar\omega + n\bar{\omega} \quad (6)$$

Then

$$\langle a_+ a_- \rangle = n \quad (7)$$

Since this value is Hermitian, this means that it is measurable. We compute

$$\langle n | a_+ a_- | n \rangle \neq 0 \quad (8)$$

Since we are in an orthonormal eigenbasis, and this inner product is $\neq 0$, we can infer that

$$(|a_- n\rangle)^\dagger = (a_- |n\rangle)^\dagger = \langle n | (a_-)^\dagger = \langle n | a_+ \quad (9)$$

Therefore,

$$a_-^\dagger = a_+ \quad (10)$$

Normalization

We know that $\langle n|n \rangle = 1$. But how do we normalize this? We show that:

$$\begin{aligned} a_+ |n\rangle &= c_{n+1} |n+1\rangle \\ \langle n|a_- a_+ |n\rangle &= \langle n+1| c_{n+1}^* c_{n+1} |n+1\rangle \\ &= \langle n| \frac{\hat{H}}{\hbar\omega} + \frac{1}{2}|n\rangle \\ &= \langle n| \frac{n\hbar\omega}{\hbar\omega} + \frac{1}{2}|n\rangle \\ &= n+1 \end{aligned}$$

This means that $c_{n+1} = \sqrt{n+1}$. This is nothing but a mathematical necessity. We repeat with a_- :

$$\begin{aligned} a_- |n\rangle &= c_{n-1} |n-1\rangle \\ c_{n-1} &= \sqrt{n-1} \end{aligned}$$

Brief Summary:

1. $|0\rangle$ is the ground state
2. a_+ : $a_+ |n\rangle = \sqrt{n+1} |n+1\rangle$
3. a_- : $a_- |n\rangle = \sqrt{n-1} |n-1\rangle$
4. $a_+^\dagger = a_-$
5. $\hat{n} = a_+ a_-$
6. $\frac{\hat{H}}{\hbar\omega} = a_+ a_- + \underbrace{\frac{1}{2}}_{ZPE}$ (ZPE = zero point energy)
7. $[a_-, a_+] = 1$

S.S in the position basis. The boundary conditions are

$$a_- \psi_0(x) = \frac{1}{\sqrt{2\hbar m\omega}} \left(\hbar \frac{\partial}{\partial x} + m\omega x \right) \psi_0(x) = 0 \quad (11)$$

This means that

$$\frac{\partial}{\partial x} \psi_0(x) = -\frac{m\omega x}{\hbar} \psi_0(x) \quad (12)$$

$$\implies \psi_0(x) = A \exp\left(-\frac{m\omega x^2}{2\hbar}\right) \quad (13)$$

The phase space of the position and momentum is aligned. We first normalize.

$$\langle \psi_0(x) | \psi_0(x) \rangle \implies A = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \quad (14)$$

We find $\psi_1(x)$:

$$a_+ \psi_0 = \sqrt{0+1} \psi_1 \quad (15)$$

$$\psi_1(x) = \frac{1}{\sqrt{2\hbar m\omega}} \left(-\hbar \frac{\partial}{\partial x} + m\omega x \right) \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} \exp\left(-\frac{m\omega x^2}{2\hbar}\right) \quad (16)$$

Following an iterative procedure, we obtain

$$\psi_n = \left(\prod_{k=1}^n \frac{a_+}{\sqrt{k}} \right) \psi_0 = \frac{1}{\sqrt{n!}} a_+^n \psi_0 \quad (17)$$

$$\psi_n(x) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n(x) e^{-x^2/2} \quad (18)$$