

PHYS 325: Lecture 5

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Lecture Span

- Non-linear Drag
- Time varying mass

Non-linear Drag

$$F(v) = -mg + cv^2 \iff -mg + \frac{c}{m}mv^2 \wedge \sigma = \frac{c}{m} \quad (1)$$

Thus, this simplifies to

$$F(v) = m(g - \sigma)v^2 \iff m\dot{v} = m(g - \sigma v^2) \iff \boxed{\dot{v} = g - \sigma v^2} \quad (2)$$

$$\frac{\dot{v}}{g - \sigma v^2} = 1 \quad (3)$$

$$\int_{v_0}^v \frac{\dot{v}}{g - \sigma v^2} dt = \int_{t_0}^t dt \quad (4)$$

$$\frac{-1}{\sqrt{g\sigma}} \tanh^{-1} \left(v \sqrt{\frac{\sigma}{g}} \right) = t \quad (5)$$

$$v(t) = -\sqrt{\frac{g}{\sigma}} \tanh(\sqrt{g\sigma}t) \quad (6)$$

$$t \rightarrow \infty \implies v(t) \rightarrow -\sqrt{\frac{g}{\sigma}} \quad (7)$$

In general, a special case of ODEs would be if

$$F(v) = f(v)g(t) \quad (8)$$

Then you can split the terms and achieve the following:

$$m \frac{dv}{dt} \frac{1}{f(v)} = g(t) \quad (9)$$

Then you're able to integrate both sides. Then what about

$$F = f(v)h(x) \quad (10)$$

Then we say

$$m \frac{dv}{dt} = m \frac{dv}{dx} \frac{dx}{dt} \frac{1}{f(v)} \iff m \frac{dv}{dx} v \frac{1}{f(v)} = g(x) \quad (11)$$

$$\int m \frac{v}{f(v)} dv = \int g(x) \quad (12)$$

Time Varying Mass, $M(t)$

Given this rocket with mass $= M(t)$ and velocity $= v(t)$ at time t , then say a mass dm is thrown out of the rocket at $t + \Delta t$ at some speed u relative to the rocket. We say that $v_f = v_i + dv$ and $M_f = M_i - dm$. Then

$$P_f = dm(v - u) + M_f v_f \iff dm(v - u) + (M_i + dm)(v_i + dv) \quad (13)$$

So expanding this out yields

$$P_f = dm(-u) + M_i v_i - dm dv + M dv = M_i v_i \quad (14)$$

We can assume that

$$-dm dv \sim 0 \text{ due to the 2nd order perturbation} \quad (15)$$

So we get

$$M dv = u dm \quad (16)$$

We know that

$$dm = -dM \text{ change of the mass of the ejected propellant} = \text{change of the mass of the rocket} \quad (17)$$

$$M dv = -u dM \quad (18)$$

$$dv = -u \frac{dM}{M} \quad (19)$$

$$v = v_0 - u \ln\left(\frac{M}{M_0}\right) \quad (20)$$

Where u is the speed of the propellant, we assume that this is constant. Then assume $M = 0.1M_0$, at some time t with $v_0 = 0$, then we get that

$$v_f = -u \ln(0.1) \sim 2.3u \quad (21)$$

Which means that about 90% of the rocket must be mass in order to achieve $2.3u$, which isn't a lot of velocity!

Time varying mass in gravity

Now we can't assume that

$$P_i = P_f \quad (22)$$

But rather

$$P_f - P_i = dp = -Mg dt \quad (23)$$

$$M dv + u dM = -Mg dt \quad (24)$$

$$\frac{dv}{dt} M + u \frac{dM}{dt} = -Mg \quad (25)$$

$$\frac{dv}{dt} = -\frac{u}{M} \frac{dM}{dt} - g \quad (26)$$

We have to specify $\frac{dM}{dt}$ because it is a degree of freedom, and can vary from rocket to rocket. For our cases, assume that

$$\frac{dM}{dt} = c \implies \dot{M} = \alpha \quad (27)$$

Then

$$M = M_0 - \alpha t \quad (28)$$

$$\frac{dv}{dt} = -g - u \frac{\alpha}{M_0 - \alpha t} \quad (29)$$

$$v(t) = v_0 - gt - u \ln\left(\frac{M_0 - \alpha t}{M_0}\right) \quad (30)$$

In general,

$$\vec{F}(\vec{r}, \dot{\vec{r}}, t) = m\vec{a} = m\ddot{\vec{r}} \quad (31)$$

In cartesian, we obtain a system of ODEs with respect to $m\ddot{x}$, $m\ddot{y}$, and $m\ddot{z}$. For example, assume that a particle lives in the B field, with a magnetic field that goes in the z direction. We have that

$$\vec{F} = q\vec{v} \times \vec{B} \quad (32)$$

The particle then spirals in a helix shape, with a radius dependent on its velocity.

$$= q \left[\vec{v} \times \vec{B} \right] \quad (33)$$

We take the determinant of the following matrix:

$$\begin{bmatrix} \vec{e}_x & \vec{e}_y & \vec{e}_z \\ v_x & v_y & v_z \\ B_x & B_y & B_z \end{bmatrix}$$

This simplifies down to

$$qB(v_y\vec{e}_x - v_x\vec{e}_y) = \vec{F} \quad (34)$$

$$F_x = qBv_y = m\ddot{x} \quad (35)$$

$$F_y = -qBv_x = m\ddot{y} \quad (36)$$

$$F_z = 0 = m\ddot{z} \implies \dot{z} = c \implies z = cx + b \quad (37)$$

$$qB\dot{y} = m\ddot{x} \quad (38)$$

$$-qB\dot{x} = m\ddot{y} \quad (39)$$

$$qB\ddot{y} = m\ddot{\dot{x}} \quad (40)$$

$$qB\left(-\frac{qB}{m}\right)\dot{x} = m\ddot{\dot{x}} \quad (41)$$

$$-\left(\frac{qB}{m}\right)^2 v_x = \ddot{v}_x \quad (42)$$

Then,

$$\ddot{v}_x = -\omega^2 v_x \wedge \omega^2 = \left(\frac{qB}{m}\right)^2 \quad (43)$$

Then the solution

$$v_x = A \sin(\omega t + \phi) \quad (44)$$

The velocity is going in a circle! Now let's try using complex numbers!

$$v_x + iv_y = \eta \quad (45)$$

Then

$$Re(\eta) = v_x \quad (46)$$

$$Im(\eta) = v_y \quad (47)$$

Then we have

$$\ddot{v}_x = -\omega^2 v_x \quad (48)$$

$$\ddot{v}_y = -\omega^2 v_y \quad (49)$$

$$\dot{\eta} = \dot{v}_x + i\dot{v}_y \quad (50)$$

$$= -i\omega(iv_y + v_x) \quad (51)$$

$$\dot{\eta} = -i\omega\eta \quad (52)$$

$$\eta = e^{-i\omega t} \quad (53)$$