

# Brief Article

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## 1 Quasispin operators

We begin by defining the quasispin operators

$$J_+ = \sum_p a_{p+}^\dagger a_{p-} \quad (1)$$

$$J_- = \sum_p a_{p-}^\dagger a_{p+} \quad (2)$$

$$J_z = \sum_{p\sigma} \sigma a_{p\sigma}^\dagger a_{p\sigma} \quad (3)$$

$$J^2 = J_+ J_- + J_z^2 - J_z \quad (4)$$

We now set to show that these operators follow the canonical angular momenta commutator relations, which are as follows(setting  $\hbar$  to 1):

### 1.1 Angular Momenta Commutator Relations

$$[J_+, J_-] = 2J_z \quad (5)$$

$$[J_+, J_z] = -J_+ \quad (6)$$

$$[J_-, J_z] = J_- \quad (7)$$

$$[J^2, J_z] = 0 \quad (8)$$

## 1.2 Proof of commutator relations

### 1.2.1 $[J_+, J_-] = 2J_z$

First we consider the commutator relation (5)

$$[J_+, J_-] = \sum_p a_{p+}^\dagger a_{p-} - \sum_k a_{k-}^\dagger a_{k+} - \sum_k a_{k-}^\dagger a_{k+} + \sum_p a_{p+}^\dagger a_{p-} \quad (9)$$

We then immediately note that since for  $k \neq p$

$$[a_{p\sigma}, a_{k\sigma'}] = 0 \quad (10)$$

It follows that all the terms where  $p \neq k$  fall out of the commutator. Hence:

$$[J_+, J_-] = \sum_p a_{p+}^\dagger a_{p-} - a_{p-}^\dagger a_{p+} - \sum_p a_{p-}^\dagger a_{p+} + a_{p+}^\dagger a_{p-} \quad (11)$$

Now we note that because of the difference in spins, some of these terms commute and the sums can be merged and we can rewrite this as

$$[J_+, J_-] = \sum_p a_{p+}^\dagger a_{p+} a_{p-} - a_{p-}^\dagger a_{p-} a_{p+} \quad (12)$$

Finally we note that since

$$[a, a^\dagger] = 1 \quad (13)$$

This thus means that with further rearrangement

$$[J_+, J_-] = \sum_p a_{p+}^\dagger a_{p+} (a_{p-}^\dagger a_{p-} + 1) - (a_{p-}^\dagger a_{p-} + 1) a_{p-}^\dagger a_{p-} \quad (14)$$

And finally simplifying this is

$$[J_+, J_-] = \sum_p a_{p+}^\dagger a_{p+} - a_{p-}^\dagger a_{p-} \quad (15)$$

Which we note is just

$$[J_+, J_-] = \sum_{p\sigma} a_{p\sigma}^\dagger a_{p\sigma} = 2J_z \quad (16)$$

And this is as expected in agreement with the relationship (5) above.

### 1.2.2 $[J_+, J_z] = -J_+$

Again we write this out in full using the creation and annihilation operators.

$$[J_+, J_z] = \frac{1}{2} \left( \sum_k a_{k+}^\dagger a_{k-} - \sum_{p\sigma} \sigma a_{p\sigma}^\dagger a_{p\sigma} - \sum_{p\sigma} \sigma a_{p\sigma}^\dagger a_{p\sigma} \sum_k a_{k+}^\dagger a_{k-} \right) \quad (17)$$

From this it follows that if  $k \neq p$  it follows that those terms cancel and similarly for the sigmas.

$$[J_+, J_z] = \frac{1}{2} \sum_{p\sigma} \sigma \left( a_{p+}^\dagger a_{p-} a_{p\sigma}^\dagger a_{p\sigma} - a_{p\sigma}^\dagger a_{p\sigma} a_{p+}^\dagger a_{p-} \right) \quad (18)$$

We then expand this sum over sigma as

$$[J_+, J_z] = \frac{1}{2} \sum_p \left( a_{p+}^\dagger a_{p-} a_{p+}^\dagger a_{p+} - a_{p+}^\dagger a_{p+} a_{p+}^\dagger a_{p-} - a_{p+}^\dagger a_{p-} a_{p-}^\dagger a_{p-} + a_{p-}^\dagger a_{p-} a_{p+}^\dagger a_{p-} \right) \quad (19)$$

Again taking advantage of the creation/annihilation operators' commutation relations it follows that

$$[J_+, J_z] = \frac{1}{2} \sum_p \left( a_{p+}^\dagger a_{p+}^\dagger a_{p+} a_{p-} - a_{p+}^\dagger (a_{p+}^\dagger a_{p+} + 1) a_{p-} - a_{p+}^\dagger (a_{p-}^\dagger a_{p-} + 1) a_{p-} + a_{p+}^\dagger a_{p-}^\dagger a_{p-} a_{p-} \right) \quad (20)$$

And finally cancelling like terms

$$[J_+, J_z] = - \sum_p a_{p+}^\dagger a_{p-} = -J_+ \quad (21)$$

### 1.2.3 $[J_-, J_z] = J_-$

This is fundamentally the same as the previous argument and so I skip most of the explanations:

$$[J_-, J_z] = \frac{1}{2} \left( \sum_k a_{k-}^\dagger a_{k+} - \sum_{p\sigma} \sigma a_{p\sigma}^\dagger a_{p\sigma} - \sum_{p\sigma} \sigma a_{p\sigma}^\dagger a_{p\sigma} \sum_k a_{k-}^\dagger a_{k+} \right) \quad (22)$$

$$[J_-, J_z] = \frac{1}{2} \sum_{p\sigma} \sigma \left( a_{p-}^\dagger a_{p+} a_{p\sigma}^\dagger a_{p\sigma} - a_{p\sigma}^\dagger a_{p\sigma} a_{p-}^\dagger a_{p+} \right) \quad (23)$$

$$[J_-, J_z] = \frac{1}{2} \sum_p \left( a_{p-}^\dagger a_{p+} a_{p+}^\dagger a_{p+} - a_{p+}^\dagger a_{p+} a_{p-}^\dagger a_{p+} - a_{p-}^\dagger a_{p+} a_{p-}^\dagger a_{p-} + a_{p-}^\dagger a_{p-} a_{p-}^\dagger a_{p+} \right) \quad (24)$$

Again taking advantage of the creation/annihilation operators' commutation relations it follows that

$$[J_-, J_z] = \frac{1}{2} \sum_p \left( a_{p+}^\dagger a_{p+}^\dagger a_{p-} a_{p+} - a_{p+}^\dagger (a_{p+}^\dagger a_{p+} - 1) a_{p-} - a_{p-}^\dagger (a_{p-}^\dagger a_{p-} - 1) a_{p+} + a_{p-}^\dagger a_{p-}^\dagger a_{p-} a_{p+} \right) \quad (25)$$

And finally cancelling like terms

$$[J_-, J_z] = \sum_p a_{p-}^\dagger a_p = J_- \quad (26)$$

#### 1.2.4 $[J^2, J_z] = 0$

Here we take advantage of commutator algebra. Noting that

$$J^2 = J_+ J_- + J_z^2 - J_z \quad (27)$$

It is immediately clear that  $J_z$  commutes with the second and third terms.

Then

$$[J^2, J_z] = [J_+ J_-, J_z] = J_+ [J_-, J_z] + [J_+, J_z] J_- = J_+ J_- - J_+ J_- = 0 \quad (28)$$