

PHY531 Problem Set 3. Due February 19, 2015

1. a. Using the method of images, calculate the solution for the Green's function in spherical coordinates for the interior of a sphere of radius  $a$  satisfying:

$$\begin{aligned}\nabla^2 G(\mathbf{r}, \mathbf{r}') &= -\delta^3(\mathbf{r} - \mathbf{r}') \\ G(a, \theta, \phi, \mathbf{r}') &= 0.\end{aligned}\tag{1}$$

- b. Use the Green's function of part a to calculate the potential inside a sphere of radius  $a$  if the potential on the sphere is given as  $V(a, \theta, \phi) = f(\theta, \phi)$ . Show your result can be manipulated into the form

$$\Phi(r, \theta, \phi) = r^{1/2} \frac{\partial}{\partial r} r^{1/2} \int_{-1}^1 d\cos\theta' \int_0^{2\pi} \frac{d\phi'}{2\pi} \frac{af(\theta', \phi')}{\sqrt{r^2 + a^2 - 2ra[\cos\theta\cos\theta' + \sin\theta\sin\theta'\cos(\phi - \phi')]}}\tag{2}$$

- c. Calculate the potential along the  $z$  axis inside the sphere for the case where the upper hemisphere is at potential  $V_0$  and the lower at  $-V_0$  by integrating Eq. 2.

Hint: Check your result for part b for a constant  $f$  where you should know the answer, and check that your result for part c gives the correct values at  $z = 0$  and  $z = \pm a$ .

2. Use the stress tensor to calculate the magnitude and direction of the forces below:
  - a. The force on the upper hemisphere of an overall neutral conducting sphere at the origin, with a charge  $q$ , a distance  $d$  along the  $\hat{\mathbf{z}}$  axis.
  - b. The force per area on the current of a long solenoid with  $n$  turns per unit length and current  $I$  in the wire. You can ignore the contribution to the field from the pitch of the winding that you calculated in Jackson problem 5.2.
3. Most undergraduate electromagnetism texts include a demonstration, using the Biot-Savart law and symmetry, that the static magnetic field of a uniformly wound toroid is circumferential. Here let's show this using the vector potential.

- a. Show that if the current density in cylindrical coordinates has the form

$$\mathbf{J}(r, \phi, z) = J_r(r, z)\hat{\mathbf{r}} + J_z(r, z)\hat{\mathbf{z}}\tag{3}$$

the vector potential has components

$$\begin{aligned}A_r(r, \phi, z) &= \frac{1}{4\pi c} \int_0^\infty r' dr' \int_0^{2\pi} d\phi' \int_{-\infty}^\infty dz' \frac{J_r(r', z') \cos(\phi')}{\sqrt{r^2 + r'^2 - 2rr' \cos(\phi') + (z - z')^2}} \\ A_\phi(r, \phi, z) &= \frac{1}{4\pi c} \int_0^\infty r' dr' \int_0^{2\pi} d\phi' \int_{-\infty}^\infty dz' \frac{J_r(r', z') \sin(\phi')}{\sqrt{r^2 + r'^2 - 2rr' \cos(\phi') + (z - z')^2}} = 0 \\ A_z(r, \phi, z) &= \frac{1}{4\pi c} \int_0^\infty r' dr' \int_0^{2\pi} d\phi' \int_{-\infty}^\infty dz' \frac{J_z(r', z')}{\sqrt{r^2 + r'^2 - 2rr' \cos(\phi') + (z - z')^2}}\end{aligned}\tag{4}$$

so that the form of  $\mathbf{A}$  is

$$\mathbf{A} = A_r(r, z)\hat{\mathbf{r}} + A_z(r, z)\hat{\mathbf{z}}.\tag{5}$$

- b. Explain why the uniformly wound toroid (many equally spaced, close wound single turns with equal current) has a current density of the form given in part a.

Show that the curl of  $\mathbf{A}$  of part a has only a  $\phi$  component which is independent of  $\phi$ , i.e. it is circumferential.

4. A toroidal core with unit relative permeability and permittivity (i.e. like vacuum) has an inner radius  $a$ , an outer radius  $b$ , and a height  $h$ . In cylindrical coordinates, the toroidal core is described by  $0 < z < h$ ,  $a < r < b$ , and  $\phi$  takes all values.

Each half is wound with a large number  $N$  of equally spaced turns of wire. Each turn is separately powered to give a static current  $\frac{I}{N}$ , all in the same sense. Since  $N$  is large, you can apply the result of problem 3.

- a. What is the easiest way to calculate the  $\mathbf{B}$  field? Calculate  $\mathbf{B}$ .  
b. The toroid is cut along the plane  $\phi = 0$  to  $\phi = \pi$ . Use the stress tensor to calculate the force holding the two halves together.  
c. Verify that the effective current density for this toroid is

$$\begin{aligned} \mathbf{J}(r, \phi, z) = & \frac{I}{\pi} \left\{ \hat{\mathbf{z}} [a^{-1}\delta(r-a) - b^{-1}\delta(r-b)] [\Theta(z) - \Theta(z-h)] \right. \\ & \left. + \hat{\mathbf{r}} r^{-1} [\delta(z-h) - \delta(z)] [\Theta(r-a) - \Theta(r-b)] \right\} \end{aligned} \quad (6)$$

and show that the force calculated from the Lorentz force law agrees with the result of part b.

5. The interior of a perfectly conducting circularly cylindrical cavity of radius  $R$  and height  $L$  has fields described by the Coulomb gauge potentials,  $\nabla \cdot \mathbf{A} = 0$ ,

$$\begin{aligned} \Phi(\mathbf{r}, t) &= 0 \\ \mathbf{A}(\mathbf{r}, t) &= \text{Re} [\mathbf{A}_c(\mathbf{r}) e^{-i\omega t}] \\ \mathbf{A}_c(r, \phi, z) &= A_0 J_0 \left( \frac{\gamma r}{R} \right) \hat{\mathbf{z}} \end{aligned} \quad (7)$$

where I use *cylindrical coordinates*  $x = r \cos \phi$ ,  $y = r \sin \phi$ , and  $\gamma$  is the first zero of the Bessel function  $J_0(\gamma) = 0$ , i.e.  $\gamma \sim 2.40482558$ . Note  $A_0$  is complex.

- a. Explain why the  $\mathbf{E}$ ,  $\mathbf{B}$  fields as well as the surface charge density and the surface current density can all be written as the real part of space dependent complex fields multiplied by  $e^{-i\omega t}$ , i.e. like  $\mathbf{A}$  in terms of  $\mathbf{A}_c$  above.

Calculate the complex fields  $\mathbf{E}_c(\mathbf{r})$  and  $\mathbf{B}_c(\mathbf{r})$  and solve for the value of  $\omega$  needed to satisfy Maxwell's equations.

- b. Verify that all the boundary conditions for a perfect conductor are satisfied.  
c. Explain why all quantities quadratic in the fields can be written as a part independent of time, and a part that oscillates at an angular frequency  $2\omega$ .  
Calculate the total electromagnetic energy inside the cavity.

- d. Use the stress tensor to calculate the time averaged outward force per unit area on the constant  $r$  wall, and then on the constant  $z$  walls.
- e. Calculate the complex  $\sigma_c$  needed to give the surface charge density on the cavity surface as a function of position and time. Calculate the complex  $\mathbf{K}_c$  needed to give the surface charge current density on the cavity surface as a function of position and time.

Hint: *Inside a perfect conductor, ohms law  $\mathbf{J} = \sigma \mathbf{E}$  with  $\sigma \rightarrow \infty$  tells you that  $\mathbf{E} = 0$ . Since  $\mathbf{E}$  is zero,  $\mathbf{B}$  must have no time derivative, so any nonzero  $\mathbf{B}$  must be static. The Bessel function recursion relations given in Jackson, and the Lommel integral notes on the class blackboard site may be helpful.*

6. You have studied black-body radiation in statistical mechanics and quantum mechanics. Obviously, since Planck invented quantum mechanics to describe black-body radiation, classical theory fails. However, it does give the correct relationship between the pressure, volume and total energy. If we look at radiation in a cavity surrounded by a perfect conductor (as in the previous problem) in thermal equilibrium and write the momentum conservation equation as in Jackson Eq. 6.121, we find the differential form

$$\frac{1}{c} \frac{\partial}{\partial t} \mathbf{E}(\mathbf{r}, t) \times \mathbf{B}(\mathbf{r}, t) + \mathbf{f}_{\text{mech}}(\mathbf{r}, t) = \nabla \cdot \overleftrightarrow{\mathbf{T}}(\mathbf{r}, t) \quad (8)$$

where  $\mathbf{f}_{\text{mech}}(\mathbf{r}, t)$  is the mechanical force per unit volume at  $\mathbf{r}$  and time  $t$ .

- a. Show that the time average of the time derivative of a periodic function is identically zero when averaged over a period. Show that the time average of the time derivative of a bounded function goes to zero when averaged over long times.
- b. In thermal equilibrium the time-averaged mechanical force on the walls from the electromagnetic field balances the pressure  $P$ . Therefore, the mechanical force is an outward force normal to the surface of the cavity with magnitude  $P \hat{\mathbf{n}} dS$  over an infinitesimal surface. Convince yourself that

$$\int_V d^3r \, \mathbf{r} \cdot \mathbf{f}_{\text{mech}}(\mathbf{r}, t) = P \int_S dS \, \mathbf{r} \cdot \hat{\mathbf{n}} \quad (9)$$

where  $V$  is the volume of the cavity extending infinitesimally into the metal walls so that  $\mathbf{f}_{\text{mech}}$  acts. Show using the divergence theorem that

$$\begin{aligned} P \int_S dS \, \mathbf{r} \cdot \hat{\mathbf{n}} &= 3PV \\ \int_V d^3r \, \mathbf{r} \cdot [\nabla \cdot \overleftrightarrow{\mathbf{T}}(\mathbf{r}, t)] &= -\text{tr} \int_V d^3r \, \overleftrightarrow{\mathbf{T}}(\mathbf{r}, t) \end{aligned} \quad (10)$$

where  $\text{tr}$  is the trace,  $\text{tr} \overleftrightarrow{\mathbf{T}} = \sum_{\alpha} T_{\alpha\alpha}$ .

- c. Put these pieces together to show that the equation of state is

$$PV = \frac{1}{3} W_{\text{em}}, \quad (11)$$

where  $W_{\text{em}}$  is the time averaged electromagnetic energy in the cavity.