Calculating the Trial Wave Function for AFDMC

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1 Trial Wave Function

The trial wave function for AFDMC must be simple to evaluate. In the past the simple Slater determinant with pair-wise correlations has been used as shown in [1],

$$\langle RS|\Psi_T\rangle = \langle RS|\left[\prod_{i<j} f_c(r_{ij})\right]\left[1+\sum_{i<j}\sum_p f_p(r_{ij})\mathcal{O}_{ij}^p\right]|\Phi\rangle\,, \tag{1}$$

where the $\mathcal{O}_{ij}^p$'s are $\boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j$, $\boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j$, and $t_{ij}\boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j$, where $t_{ij} = 3\boldsymbol{\sigma}_i \cdot \hat{r}_{ij}\boldsymbol{\sigma}_j \cdot \hat{r}_{ij} - \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j$. *Why weren't $\boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j$ and t_{ij} used in this paper?*

My goal is to add the additional independent pair correlations.

$$\langle RS|\Psi_T\rangle = \langle RS|\left[\prod_{i<j} f_c(r_{ij})\right]\left[1+\sum_{i<j}\sum_p f_p(r_{ij})\mathcal{O}_{ij}^p + \sum_{i<j}\sum_{k<l}\sum_p f_p(r_{ij})\mathcal{O}_{ij}^p f_p(r_{kl})\mathcal{O}_{kl}^p\right]|\Phi\rangle\,, \tag{2}$$

2 Evaluation the Trial Wave Function

To understand how to to this I'm going to just assume that $\mathcal{O}_{ij}^p$ only contains the term $\boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j$ and I'll start by looking at the trial wave function, equation 1, with only the linear term. So now

$$\langle RS|\Psi_T\rangle = \langle RS|\left[\prod_{i<j} f_c(r_{ij})\right]\left[1+\sum_{i<j} f_1(r_{ij})\boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j\right]|\Phi\rangle\,. \tag{3}$$

Also since the central correlations don't change the states by any more than a multiplicative factor I am going to ignore that term as well. I will also just look at one term in the sum (a particular i and j value). So we are just looking at

$$\langle RS|\left[1+f_1(r_{ij})\boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j\right]|\Phi\rangle\,. \tag{4}$$

Now we also know that the Slater determinant is defined as

$$\langle RS|\Phi\rangle = \det(S_{ij}) = \frac{1}{\sqrt{N!}}\begin{vmatrix} \phi_1(R_1S_1) & \phi_2(R_1S_1) & \cdots & \psi_N(R_1S_1) \\ \phi_1(R_2S_2) & \phi_2(R_2S_2) & \cdots & \phi_N(R_2S_2) \\ \vdots & \vdots & \ddots & \vdots \\ \phi_1(R_NS_N) & \phi_2(R_NS_N) & \cdots & \phi_N(R_NS_N) \end{vmatrix}\,, \tag{5}$$

where $\phi_i(R_jS_j) = \phi_i^r(R_j)\phi_i^s(S_j)$ and S_{ij} is called the Slated Matrix.

Now lets look at equation 4 again for an example.

$$\langle RS| [1 + f_1(r_{ij})\boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j] |\Phi\rangle \tag{6}$$
$$= \det(S_{ij}) + f_1(r_{ij}) \langle RS| \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j |\Phi\rangle \tag{7}$$
$$= \det(S_{ij}) + f_1(r_{ij})\det(S'_{ij}) \tag{8}$$

Here S'_{ij} is the updated matrix. It only has two columns different than S_{ij} and so we can get it's determinant of S'_{ij} easily once we have the determinant of S_{ij} by using the fact that

$$\det(S_{ij}^{-1}S'_{ij}) = \frac{\det(S'_{ij})}{\det(S_{ij})}. \tag{9}$$

When we solve for $\det(S_{ij})$ we finish solving for the inverse, S_{ij}^{-1} and the product $S_{ij}^{-1}S'_{ij}$ is 1 on the diagonal and 0 everywhere else except the two columns i and j. This makes the $\det(S_{ij}^{-1}S'_{ij})$ easy to solve for since it is simply the determinant of the submatrix. Thus once we have $\det(S_{ij})$ it is easier to solve for $\det(S'_{ij})$. All that is left is to do this over the pair loops and over each operator.

References

[1] S. Gandolfi, A. Lovato, J. Carlson, and Kevin E. Schmidt. From the lightest nuclei to the equation of state of asymmetric nuclear matter with realistic nuclear interactions. 2014. arXiv:1406.3388v1 [nucl-th].