# Study guide for qualifying exams

Cody L. Petrie

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### 1 Classical Mechanics

- 1. Newtonian Mechanics
  - (a) Newton's Laws/Kinematics
  - (b) Energy
  - (c) Momentum/Angular Momentum
- 2. Lagrangian Mechanics
  - (a) Calculous of Variations
  - (b) Principle of Least Action/Lagranges Equation
  - (c) Generalized Coordinates
  - (d) Holonomic/Non-Holonomic Constraints
  - (e) Noether's Theorem
  - (f) Rigid Body Motion
    - i. Inertia Tensor
    - ii. Euler's Equations
- 3. Hamiltonian Formalism
  - (a) Legendre Transformation/Hamilton's Equations
  - (b) Generalized Momenta/Cyclic Coordinates/Conserved Quantities
  - (c) Liouville's Theorem
  - (d) Poisson Brackets
  - (e) Canonical Transformations

### 2 Statistical Mechanics

- 1. Thermodynamics Review
  - (a) Laws of Thermodynamics
  - (b) Intensive vs Extensive Variables
  - (c) Thermodynamic Potentials and Ensembles
  - (d) Maxwell's Relations
  - (e) Various Definitions
    - i. Compressibility
    - ii. Heat Capacity etc.
- 2. Statistical Mechanics
  - (a) Statistical Review
  - (b) Partition Function/Trace
  - (c) Thermodynamic Limit
  - (d) Density Matrix
  - (e) Ideal Gas
  - (f) Ideal Bose Gas
  - (g) Ideal Fermi Gas
  - (h) Photons (BB)
  - (i) Phonons
  - (j) Bose-Einstein Condensates
  - (k) Cluster Expansion

## 3 Quantum Mechanics

- 1. Shankar Math Review
- 2. Postulates
- 3. Free Particle
- 4. Particle in a Box
- 5. Harmonic Oscillator
- 6. Angular Momentum
- 7. Hydrogen Atom
- 8. Spin

- 9. Angular Momentum Addition
- 10. Time-Independent Perturbation Theory
- 11. Time-Dependent Perturbation Theory
  - (a) Einstein A and B Coefficients
- 12. Scattering
- 13. WKB Formula
- 14. Dirac Equation

# 4 Electricity and Magnetism

- 1. Electrostatics
  - (a) Coulomb's Law
  - (b) Electrostatic Potentials
    - i. Poisson/Laplace's Equations
  - (c) Boundary Conditions
  - (d) Method of Images
  - (e) Multipole Expansion
  - (f) Work and Energy
  - (g) Electric Fields in Matter
- 2. Magnetostatics
  - (a) Lorentz Force Law
  - (b) Biot-Savart Law
  - (c) Vector Potential
  - (d) Magnetic Fields in Matter
- 3. Electrodynamics
  - (a) Faraday's Law
  - (b) Ampere's Law
  - (c) Ohm's Law
  - (d) Maxwell's Equations
  - (e) Boundary Conditions to Maxwell's Equations
  - (f) Continuity Equation

- (g) Poynting's Theorem
- (h) Maxwell Stress Tensor
- (i) Electromagnetic Waves
  - i. The Wave Equation from Maxwell's Eq.
  - ii. EM Waves in Matter
  - iii. Wave Guides
- 4. Scalar and Vector Potentials
- 5. Coulomb and Lorentz Gauge
- 6. Retarted Potentials
- 7. Lienard-Wiechert Potentials
- 8. Radiation
  - (a) Electric/Magnetic Dipole Radiation
- 9. Helmholtz Theorem
- 10. Special Relativity
  - (a) Einstein's Postulates
  - (b) Lorentz Transformation
  - (c) 4-Vectors
  - (d) Field Tensor and Transformation
  - (e) Relativistic Potentials

#### Classical Mechanics Equations 5

### Newtonian Mechanics

#### Newton's Laws:

- 1. An object will maintain it's current motion unless acted upon by an external force.
- $2. \vec{F} = m\vec{a}$
- 3. All forces occur in equal but directionally opposite pairs.

Second Law:  $\vec{F} = m\vec{a} = \dot{\vec{p}}$ 

Angular Position/Velocity/Acceleration:  $\theta = s/r$ ,  $\omega = v/r$ ,  $\alpha = a/r$ 

Angular Momentum:  $\vec{L} = \vec{r} \times \vec{p}$ 

Torque:  $\vec{\tau} = \vec{r} \times \vec{F} = \vec{L}$ 

Centripital Acceleration:  $a_c = v^2/r$ 

Centrifugal/Coriolis Forces:  $\vec{F}_{cent} = -m\vec{\omega} \times (\vec{\omega} \times \vec{r'}), \ \vec{F}_{cor} = -2m\vec{\omega} \times \dot{\vec{r'}}$ 

Work to go from positions  $\vec{a}$  to  $\vec{b}$ :  $W_{ab} = \int_{\vec{a}}^{b} \vec{F} \cdot d\vec{s}$ 

Conservative Force Field (2 eq):  $W_{ab}$  is the same regardless of path so  $\oint \vec{F} \cdot d\vec{s} = 0$ , and thus we can write the force as  $\vec{F} = -\nabla V(\vec{r})$ .

### Lagrangian Formalism

Functional Derivative:  $\frac{\delta F}{\delta u}[x_0] = \lim_{\epsilon \to 0} \frac{F[x_0 + \epsilon u] - F[x_0]}{\epsilon} \to \frac{\delta F}{\delta x(t)}[x(t')] = \lim_{\epsilon \to 0} \frac{F[x(t') + \epsilon \delta(t'-t)] - F[x(t')]}{\epsilon}$ 

**Principle of Least Action:**  $\delta S = 0$ , where  $S = \int_{t_i}^{t_f} L(\vec{q}, \dot{\vec{q}}, t) dt$ 

Lagranges Equation:  $\frac{d}{dt} \frac{\partial L}{\partial \dot{x}^A} - \frac{\partial L}{\partial x^A} = 0$ Holonomic Constraints:  $f_{\alpha}(x^A, t) = 0$ ,  $L' = L(x^A, \dot{x}^A) + \lambda_{\alpha} f_{\alpha}(x^A, t) \rightarrow \frac{d}{dt} \frac{\partial L}{\partial \dot{x}^A} - \frac{\partial L}{\partial x^A} = 0$ 

Noether's Theorem: A continuous symmetry in the Action (and thus Lagrangian) result in a conserved quantity.

Moment of Inertia Tensor:  $\vec{L} = \overleftrightarrow{T}\vec{\omega}, T = \frac{1}{2}\omega_a I_{ab}\omega_b, I_{ab} = \sum_i m_i ((\vec{r}_i \cdot \vec{r}_i)\delta_{ab} - (\vec{r}_i)_a (\vec{r}_i)_b$ Euler's Equations: Only look at rotation, not translation. Conservation of Angular Momentum gives  $I_i\dot{\omega}_i + \omega_i\omega_k(I_k - I_j) = 0$ , for i,j,k being cyclic permutations of 1,2,3.

5

### Hamiltonian Formalism

Generalized Momenta:  $p_i = \frac{\partial L}{\partial \dot{q}_i}, \, \dot{p}_i = \frac{\partial L}{\partial q_i}$ 

**Hamiltonian:**  $H(q_i, p_i, t) = \sum_{i=1}^n p_i \dot{q}_i - L(q_i, \dot{q}_i, t)$ 

Hamilton's Equations:

1. 
$$\dot{p}_i = -\frac{\partial H}{\partial q_i}$$

$$2. \ \dot{q}_i = \frac{\partial H}{\partial p_i}$$

3. 
$$-\frac{\partial L}{\partial t} = \frac{\partial H}{\partial t}$$

Cyclic/Ignorable Coordinates: q is ignorable if  $\frac{\partial L}{\partial q} = 0$ , i.e. if q does not appear in L. Thus  $p = \frac{\partial L}{\partial \dot{q}}$  is conserved. **Liousille's Theorem:** A volume of a region of phase space remains the same, even when

the refion changes.  $V = dq_1 \dots dq_n dp_1 \dots dp_n$ .

Poisson Bracket:  $\{f,g\} = \frac{\partial f}{\partial q_i} \frac{\partial g}{\partial p_i} - \frac{\partial f}{\partial p_i} \frac{\partial g}{\partial q_i}$ .

Constant of Motion from Poisson Bracket:  $\frac{df}{dt} = \{f, H\} + \frac{\partial f}{\partial t}$ . If I, H = 0, then I is a constant of motion.

Transformation  $(q_i \rightarrow Q_i(q, p), p_i \rightarrow P_i(q, p))$  that leaves Canonical Transformation: Hamilton's equations invariant.

#### Statistical Mechanics Equations 6

#### Thermodynamics 6.1

Laws of Thermodynamics:

- 1. Energy conservation. dE = dQ pdV. dQ just means that the heat is an inexact differential and the integral depends on the path.
- 2.  $\Delta S \geq \int \frac{dQ}{T}$ , where equality is for a process that is reversible (never leaves equilibrium).
- 3. Entropy at zero temperature is zero. In stat mech this means that the ground state is nondegenerate and  $S \propto \ln(W)$ , where W is the number of available states.

Intensive vs Extensive Variables: Intensive variables do NOT scale with system size  $(T, p, \mu)$ , while extensive do scale (E, S, V, N).

Thermodynamic Potentials:

- Internal Energy: U(S, V, N)
- Helmholtz Free Energy: F(T, V, N) = U TS
- Enthalpy: H(S, p, N) = U + pV
- Gibbs Free Energy: G(T, p, N) = U TS + pV
- Landau(Grand) Potential:  $\Omega(T, V, \mu) = U TS \mu_i N_i$

Thermodynamic Ensembles:

- 1. Microcanonical: Does not exchange energy or particles with environment. Fixed E, N
- 2. Canonical: Does not exchange particles, but can exchange energy (heat bath). Fixed N, T

3. Grand canonical: Can exchange energy and particles with environment. Fixed  $T, \mu$ .

Maxwell's Relations (4 main):

• 
$$\frac{\partial^2 U}{\partial S \partial V} = -\left(\frac{\partial p}{\partial S}\right)_V = \left(\frac{\partial T}{\partial V}\right)_S$$

• 
$$\frac{\partial^2 F}{\partial T \partial V} = \left(\frac{\partial p}{\partial T}\right)_V = \left(\frac{\partial S}{\partial V}\right)_T$$

• 
$$\frac{\partial^2 H}{\partial S \partial p} = \left(\frac{\partial V}{\partial S}\right)_p = \left(\frac{\partial T}{\partial p}\right)_S$$

• 
$$\frac{\partial^2 G}{\partial T \partial p} = \left(\frac{\partial V}{\partial T}\right)_p = -\left(\frac{\partial S}{\partial p}\right)_T$$

Engine Efficience:  $\eta = \frac{Q_{in} - Q_{out}}{Q_{in}} = 1 - \frac{T_{out}}{T_{in}}$ 

**Isobaric Thermal Expansion Coefficient:**  $\alpha = \frac{1}{V} \left( \frac{\partial V}{\partial T} \right)_P$ , How much the volume changes with a change in termperature.

**Isothermal Compressibility:**  $\kappa_T = -\frac{1}{V} \left( \frac{\partial V}{\partial P} \right)_T$ , How much the volume changes when the pressure changes.

Isentropic(Adiabatic) Compressibility:  $\kappa_S = -\frac{1}{V} \left( \frac{\partial V}{\partial P} \right)_S$ , Same as above. Specific Heat at Constant V:  $C_V = \left( \frac{\partial Q}{\partial T} \right)_V = \left( \frac{\partial U}{\partial T} \right)_V$ , Amount of heat per unit mass to raise the temp by 1 degree.

Specific Heat at Constant p:  $C_p = \left(\frac{\partial Q}{\partial T}\right)_p = \left(\frac{\partial H}{\partial T}\right)_p$ , Same as above.

Fermi Energy/Temperature: Chemical potential at T=0.  $\epsilon_F=\mu(T=0)$ 

#### 6.2 Statistical Mechanics

Number of microstates in a mactostate (ways to get n heads):  $\Omega = \frac{N!}{\prod_i n_i!}$ 

Stirling's Approximation:  $\ln n! = n \ln n - n$ 

How many order important ways to order n things: n!

How many order important waus to order n things r at a time:  $\frac{n!}{(n-r)!}$ 

How many NOT order important ways to order n things r at a time:  $\binom{n}{r} = \frac{n!}{r!(n-r)!}$ 

Microcanonical (Classical) Partition Function:  $Z_m = \sum_s g_s e^{-\beta E_s}$ 

Canonical Partition Function:  $Z_c = \operatorname{tr}\left(e^{-\beta \hat{H}}\right)$ 

Grand Canonical Partition Function:  $Z_{gc} = \operatorname{tr}\left(e^{-\beta(\hat{H}-\mu\hat{N})}\right)$ 

Geometric Series:  $\sum_{n=0}^{\infty} x^n = \frac{1}{1-x}$ 

Classical limit of the trace of an operator:  $\operatorname{tr}(\mathcal{O}) = \frac{1}{N!(2\pi\hbar)^{3N}} \int d^3r_1 \dots d^3r_N \int d^3p_1 \dots d^3p_N \mathcal{O}$ , N! is for identical particles.

Thermodynamic Limit:  $T \to \infty, V \to \infty, N/V = const$ 

Expectation value for pure/mixed:  $\langle \mathcal{O} \rangle_p = \langle \psi | \, \mathcal{O} \, | \psi \rangle \,, \\ \langle \mathcal{O} \rangle_m = \sum_i P_i \, \langle \psi_i | \, \mathcal{O} \, | \psi_i \rangle \,$ 

Density Matrix (ex. Canonical Ensemble):  $\rho = \sum_{n} P_n |\psi_n\rangle \langle \psi_n|, \rho_c = \frac{e^{-\beta \hat{H}}}{\operatorname{tr} e^{-\beta \hat{H}}}$ 

Expectation value with Density Matrix:  $\langle \mathcal{O} \rangle = \operatorname{tr}(\mathcal{O}\rho)$ 

Trace of Density matrix:  $tr(\rho) = 1$ 

Time evolution of density matrix:  $\frac{\partial}{\partial t} \rho = \frac{1}{i\hbar} \left| \hat{H}, \hat{\rho} \right|$ 

 $Z_{gc}$  for an ideal gas:  $Z_{gc} = \frac{V^N(2mT\pi)^{3N/2}}{N!(2\pi\hbar)^{3N}}e^{\beta\mu}$ 

 $Z_{gc}$  for ideal fermi gas:  $Z_{gc} = \prod_{k=1}^{N(2Rh)^{s}} \left(1 + e^{-\beta(\epsilon_k - \mu)}\right)$ 

 $Z_{gc}$  for ideal bose gas:  $Z_{gc} = \prod_{k}^{k} \frac{1}{\left(1 - e^{-\beta(\epsilon_k - \mu)}\right)}$ 

Stuff here for black-body and phonons and bose condensates.

Explain Bose-Condensates with Bose statistics:  $\lim_{T\to 0} n(p) = \lim_{\beta\to\infty} \frac{1}{1-e^{\beta(\epsilon-\mu)}} \to 0$  un-

less  $\epsilon \to \mu$ , which happens at the ground state. Is this true?

What is cluster expansion used for?: Systems of interacting particles.

# 7 Quantum Mechanics Equations

Properties of a vector space:

- Sum  $|V\rangle + |W\rangle$
- Scalar product with properties
  - 1. closure: results in another vector in the space.
  - 2. distributive:  $a(|V\rangle + |W\rangle = a|V\rangle + a|W\rangle$ ,  $(a+b)|V\rangle = a|V\rangle + b|V\rangle$
  - 3. associative:  $a(b|V\rangle) = ab|V\rangle, |V\rangle + (|W\rangle + |Z\rangle) = (|V\rangle + |W\rangle) + |Z\rangle$
  - 4. commutative:  $|V\rangle + |W\rangle = |W\rangle + |V\rangle$
  - 5. addative inverse:  $|V\rangle + |-V\rangle = |0\rangle$
  - 6. null vector:  $|V\rangle + |\rangle = |V\rangle$

Hilbert space: Vector space with defined inner product.

Expand in orthonormal basis:  $|V\rangle = \sum_{i} vi |i\rangle$ 

Hermitian operator:  $\mathcal{O}^{\dagger} = \mathcal{O}$ 

Anti-Hermitian operator:  $\mathcal{O}^{\dagger}=\mathcal{O}$ 

Unitary operator:  $UU^{\dagger} = \mathbb{I}$ Orthogonality:  $\langle i|j \rangle = \delta_{ij}$ Completeness:  $\sum_{i} i = \mathbb{I}$ 

Postulates of QM:

- 1. The state of a physical system, at some fixed time, is given by a normalized ray in a Hilbert space over the complex numbers. (ray is vector whose norm doesn't matter)
- 2. The ray evolves deterministically in time according to Schrödingers equation.
- 3. Observables correspond to self-adjoint (hermitian) operators.
- 4. If a particle is in the state  $|\psi\rangle$  then a measurement of  $\mathcal{O}$  will yield one of the eigenvalues of  $\mathcal{O}$ ,  $\omega$ . The state of the system changes to an eigenstate of  $\mathcal{O}$ ,  $|\omega\rangle$ .

8

Schrödinger equation:  $i\hbar \frac{\partial}{\partial t}\Psi = \hat{H}\Psi$ 

Free particle  $\psi_p$  and  $E_p$ :  $\psi_p = Ae^{ikx} + Be^{-ikx}$ ,  $k^2 = \frac{2mE_n}{\hbar^2}$ ,  $E_p = \frac{p^2}{2m}$ 

Particle in a box  $\psi_n$  and  $E_n$ :  $\psi_n = \sqrt{\frac{2}{L}} \sin(k_n x)$ ,  $k_n = \frac{n\pi}{L}$ ,  $E_n = \frac{\hbar^2 \pi^2 n^2}{2mL^2}$ 

**Harmonic** Oscillator  $\hat{H}$ ,  $\psi_n$  and  $E_n$ :  $\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2\hat{x}^2$ ,  $\psi_n = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \frac{1}{2^n n!} H_n(x) e^{-x^2/2}$ ,  $E_n = (n + \frac{1}{2})\hbar\omega$ 

Raising and lowering operators and how to affect  $|n\rangle$  (3-2):

• 
$$a = \left(\frac{m\omega}{2\hbar}\right)^{1/2} \left(\hat{x} + \frac{i}{m\omega}\hat{p}\right), \ a|n\rangle = \sqrt{n}|n-1\rangle, \ a|0\rangle = 0$$

• 
$$a^{\dagger} = \left(\frac{m\omega}{2\hbar}\right)^{1/2} \left(\hat{x} - \frac{i}{m\omega}\hat{p}\right), a^{\dagger} |n\rangle = \sqrt{n+1} |n+1\rangle$$

 $\hat{H}$  in terms of a and  $a^{\dagger}$ :  $\hat{H} = \hbar\omega(a^{\dagger}a + 1/2)$ Commutation relations for H, a,  $a^{\dagger}$ :

$$\bullet \ [\hat{H}, a] = -a$$

$$\bullet \ [\hat{H},a^{\dagger}]=a^{\dagger}$$

$$\bullet \ [a,a^{\dagger}] = 1$$

 $J^2$  and  $J_z$  on the angular momentum state  $|jm_j\rangle$ :

• 
$$\mathbf{J}^2 \mid = \rangle j(j+1)\hbar^2 \mid jm_j \rangle$$

• 
$$J_z |jm_i\rangle = m_i \hbar |jm_i\rangle$$

Commutation relations for  $J_i$  and  $J_j$  and for  $J^2$  and  $J_i$ :

• 
$$[J_i, J_j] = i\hbar J_k$$

$$\bullet \ [\mathbf{J}^2, J_i] = 0$$

 $J_z$  and  $J^2$  in position basis:

$$J_z = -i\hbar \frac{\partial}{\partial t}$$

• 
$$\mathbf{J}^2 = -\hbar^2 \left[ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]$$

Raising and Lowering Angular Momentum Operators on  $|i, m\rangle$ :

$$J_{\pm} |j, m\rangle = \hbar [j(j+1) - m(m \pm 1)]^{1}/2 |j, m \pm 1\rangle$$

 $J_x$  and  $J_y$  in terms of  $J_+$  and  $J_-$ :  $J_x = \frac{1}{2}(J_+ + J_-), \ J_y = \frac{1}{2i}(J_+ - J_i)$ Momentum eigenstate,  $\langle x|p\rangle$ :  $\langle x|p\rangle = \frac{1}{\sqrt{2\pi\hbar}}e^{ipx/\hbar}$ 

Hydrogen Atom V(r),  $\psi_n$ ,  $E_n(\mathbf{x4})$ :  $V(r) = -\frac{Ze^2}{4\pi\epsilon\sigma r}$ ,  $\psi_n = stuff*L_{n-l-1}^{2l+1}(\rho)Y_l^m(\theta,\phi)$  (Laguerre)

$$E_n = -\frac{1}{2n^2} \left( \frac{Ze^2}{4\pi\epsilon_0 \hbar} \right)^2 m_e = -\frac{1}{2n^2} \alpha^2 m_e c^2 = -\frac{1}{n^2} 13.6 eV = -\frac{1}{2n^2} \left( \frac{1}{4\pi\epsilon_0} \frac{e^2}{a_0} \right),$$

Pauli matricies and commutation relations:

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$[\sigma_a \sigma_b] = 2i\epsilon_{abc}\sigma_c$$

Non-Deg Time-Ind Perturbation,  $E_n^{(1)}$ ,  $|n^{(1)}\rangle$ ,  $E_n^{(2)}$ :

$$E_n^{(1)} = H'_{nn} = \langle n^{(0)} | H' | n^{(0)} \rangle$$

$$| n^{(1)} \rangle = \sum_{m \neq n} \frac{\langle n^{(0)} | H' | n^{(0)} \rangle}{\langle E_n^{(0)} - E_m^{(0)} \rangle} | m^{(0)} \rangle$$

$$E_n^{(2)} = \sum_{m \neq n} \frac{|\langle n^{(0)} | H' | n^{(0)} \rangle|^2}{\langle E_n^{(0)} - E_m^{(0)} \rangle}$$

**Deg Time-Ind Perturbation**,  $E_n^{(1)}$ : Diagonalize the perturbation hamiltonian in the degenerate subspace.

Time-Dep Perturbation,  $P_{i\to f}(t)$ :  $P_{i\to f}(t) = \frac{1}{\hbar^2} \left| \int_0^t dt' \left\langle f \right| H'(t') \left| i \right\rangle e^{i(E_f - E_i)t'/\hbar} \right|^2$ Fermi's golder rule, and  $g(E_f)as\delta$ :  $R_{i\to f}=\frac{2\pi}{\hbar}\left|\langle f|H'|i\rangle\right|^2g(E_f), g(E_f)\approx\delta(E_f^{(0)}-E_f^{(0$  $\hbar\omega$ )

Einstein's Stimulated/Spontaneous emission coefficients:

Stimulated: 
$$B_{if} = \frac{\pi e^2}{3\epsilon_0 \hbar^2} |\langle f|\mathbf{r}|i\rangle|^2$$

Stimulated: 
$$B_{if} = \frac{\pi e^2}{3\epsilon_0 \hbar^2} \left| \langle f | \mathbf{r} | i \rangle \right|^2$$
  
Spontaneous:  $A_{if} = \frac{e^2 \omega_{21}^3}{3\pi\epsilon_0 \hbar c^3} \left| \langle f | \mathbf{r} | i \rangle \right|^2$   
Total  $\psi(\mathbf{r})$  in scattering problem:

$$\psi(\mathbf{r}) = \psi_{inc}(\mathbf{r}) + \psi_s(\mathbf{r}) = \psi_{inc}(\mathbf{r}) + f(\theta, \phi) \frac{e^{i\mathbf{k}\cdot\mathbf{r}}}{r}$$
$$f(\theta, \phi) = -\frac{m}{2\pi\hbar^2} \int d^3r' e^{-i\mathbf{k}'\cdot\mathbf{r}'} V(\mathbf{r}') \psi(\mathbf{r}')$$

Differential Cross Section:  $\frac{d\sigma}{d\Omega} = |f(\theta, \phi)|^2$ 

Born Approximation: In the above integral for  $f(\theta, \phi)$  let  $\psi \to \psi_{inc}$ .

**Dirac Equation:** 

$$\begin{split} &(i\hbar\gamma^{\mu}\partial_{\mu}-mc)\psi=0\\ &\gamma^{0}=\beta,\,\gamma^{i}=\beta\alpha_{i}\\ &\beta=\begin{pmatrix}\mathbb{I}&0\\0&-\mathbb{I}\end{pmatrix},\,\alpha_{i}=\begin{pmatrix}0&\sigma_{i}\\-\sigma_{i}&0\end{pmatrix} \end{split}$$

#### Electricity and Magnetism Equations 8

Maxwell's Equations in Vacuum (SI):

$$\nabla \cdot \mathbf{E} = \rho / \epsilon_0, \, \nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}, \ \nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

 $\begin{array}{l} \nabla \cdot \mathbf{E} = \rho/\epsilon_0, \ \, \dot{\nabla} \cdot \mathbf{B} = 0 \\ \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}, \ \, \nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \\ \text{Maxwell's Equations in Matter (SI), and D and H:} \end{array}$ 

$$\nabla \cdot \mathbf{D} = \rho, \, \nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}, \nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial D}{\partial t}$$
$$\mathbf{D} = \epsilon_0 \mathbf{E}, \mathbf{B} = \mu_0 \mathbf{H}$$

$$\mathbf{D} = \epsilon_0 \mathbf{E}, \ \mathbf{B} = \mu_0 \mathbf{H}$$

Continuity Equation:  $\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{J} = 0$ 

Lorentz Force:  $\mathbf{F} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B})$ 

Coulomb's Law (x2):  $\mathbf{F} = \frac{1}{4\pi\epsilon_0} \frac{q_1q_2}{r^2}$ ,  $\mathbf{F} = Q\mathbf{E}$ 

Gauss' Law:  $\oint \mathbf{E} \cdot d\mathbf{A} = q/\epsilon_0$ 

Electrostatic Potential (x2):  $\mathbf{E} = -\nabla \Phi$ ,  $\Phi = \frac{1}{4\pi\epsilon_0} \int d^3r' \frac{\rho(\mathbf{r'})}{|\mathbf{r} - \mathbf{r'}|}$ 

Laplace's Equation & General Solution(Spherical Coordinates, no  $\phi$ ):  $\nabla^2 \Phi = 0$ 

$$\Phi(r,\theta) = \sum_{l=0}^{\infty} \left( A_l r^l + \frac{B_l}{r^{l+1}} \right) P_l(\cos \theta)$$

Poisson's Equation:  $abla^2\Phi = -\rho/\epsilon_0$ 

# 9 Miscellaneous Physics

Taylor Expansion:  $f(\vec{x} + \vec{a}) = f(\vec{x}) + a_i \partial_i f(\vec{x}) + \mathcal{O}(\vec{a}^2)$ 

Gaussian Integral:  $\int_{0}^{\infty} dx e^{-ax^2 + bx + c} = \sqrt{\frac{\pi}{a}} e^{b^2/(4a) + c}$ 

3 types of Boundary Conditions:

Dirichlet:  $\Phi(\mathbf{a}) = const$ 

Neumann:  $\frac{\partial \Phi(\mathbf{a})}{\partial \mathbf{n}} = \hat{\mathbf{n}} \cdot \nabla \Phi = const$ 

Robin: Linear combination of the first two

Value of fine structure constant:  $\alpha \approx \frac{1}{137}$ 

Mass of electron in eV:  $m_e c^2 = 0.511 eV$ 

Value of the Bohr radius:  $a_0 = 0.529 \text{ Å}$ 

**Explain the Method of Images:** Because of the uniqueness theorem you can add charges OUTSIDE of the computational area to meet the same boundary conditions. A solution to this new configuration is also a solution to the initial configuration.