

# Path Integral Monte Carlo Studies of Hard Sphere Systems

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July 29, 2016

# Overview of Path Integral Monte Carlo

- Used to calculate thermodynamic properties of systems
- Systems are at finite temperature
- Use Feynman path integral formalism to describe partition function as integrals over coordinate space
- Use Monte Carlo methods to calculate these integrals

# Path Integral Form of Partition Function

- PIMC formalism begins with the thermal density matrix of the system:

$$\hat{\rho} = e^{-\beta\hat{H}} = \sum_i e^{-\beta E_i} |\psi_i\rangle\langle\psi_i|$$

- The partition function is the trace of the thermal density matrix:

$$Z = \text{Tr}(\hat{\rho})$$

- The trace can be performed in the position basis:

$$Z = \int dR \langle R | \hat{\rho} | R \rangle$$

- Note:  $R = \{r_1, r_2, \dots, r_i\}$

# Path Integral Form of Partition Function

- Take advantage of these two properties:

$$e^{-(\beta_1+\beta_2)\hat{H}} = e^{-\beta_1\hat{H}}e^{-\beta_2\hat{H}} \quad \hat{1} = \int |R\rangle\langle R|dr$$

- With these it is possible to expand the matrix elements:

$$Z = \int dR \langle R | e^{\beta\hat{H}} | R \rangle$$

$$Z = \int dR dR_2 \langle R | e^{-(\beta/2)\hat{H}} | R_2 \rangle \langle R_2 | e^{-(\beta/2)\hat{H}} | R \rangle$$

$$Z = \int \dots \int dR dR_1 \dots dR_{M-1} \rho(R, R_1; \beta/M) \rho(R_1, R_2; \beta/M) \dots \rho(R_{M-1}, R; \beta/M)$$

# Evaluation of Intermediate Matrix Elements

- The intermediate matrix elements are referred to as propagators

$$\rho(R_N, R_{N+1}; \tau) = \langle R_N | e^{-\tau \hat{H}} | R_{N+1} \rangle, \quad \beta/M = \tau$$

- Use Trotter decomposition:

$$\rho(R_N, R_{N+1}; \tau) = \langle R_N | e^{-\tau \hat{K}} e^{-\tau \hat{V}} | R_{N+1} \rangle$$

$$\rho(R_N, R_{N+1}; \tau) = \frac{1}{(4\pi\lambda\tau)^{3N/2}} \exp\left(-\frac{(R_N - R_{N+1})^2}{4\lambda\tau}\right) \langle R_N | e^{-\tau \hat{V}} | R_{N+1} \rangle$$

- The first term is denoted as  $\rho_{free}$

# Evaluation of Intermediate Matrix Elements

- In many body simulations, the potential component of the propagator takes the form:

$$\rho(R_N, R_{N+1}; \tau) = \rho_{free} \prod_{i,j} \rho_{ij}(r_{ij}, r'_{ij})$$

- Where  $r_{ij}$  and  $r'_{ij}$  are the separation distances between particles at consecutive configurations
- The two body density matrix is often times difficult to evaluate exactly, often times approximations are used
- Results converge with increasing M, referred to as 'time slices'

# Hard Sphere Systems

- The hard sphere potential is infinite when the radii of the particles overlap, but zero otherwise
- Can add a hard cavity, infinite potential when the separation distance is greater than a specified distance
- Hard Sphere systems can model atoms or spherical molecules with minimal computational effort
- Previous work has been done deriving two body density matrices for the hard sphere problem

# Hard Sphere Systems

- For the regular hard sphere problem:

$$\rho_{IA}(r_{ij}, r'_{ij}) = (1 - e^{m/\hbar^2 \tau (r-\sigma)(r'-\sigma)})$$

$$\rho_{CB}(r_{ij}, r'_{ij}) = (1 - \frac{\sigma(r+r')-\sigma^2}{rr'}) e^{m/2\hbar^2 \tau (r-\sigma)(r'-\sigma)(1-\cos\chi)})$$

- Current work in this project is to study comparative effectiveness of both propagators



# Calculating Observables

- The full functional form of the partition function is now known

$$Z = \int d\mathcal{R} \prod_{N=1}^M \rho_{free}(R_N, R_{N+1}) \prod_{i,j} \rho_{ij}(r_{ij}, r'_{ij}) \quad , \quad \mathcal{R} = \{R_N\}$$

$$Z = \int d\mathcal{R} e^{-U_{eff}(\mathcal{R})}$$

- Thermodynamic observables can be written as a function of  $\mathcal{R}$  by taking the appropriate derivatives of the partition function

# Calculating Observables

- The expectation values of the operators can be calculated as:

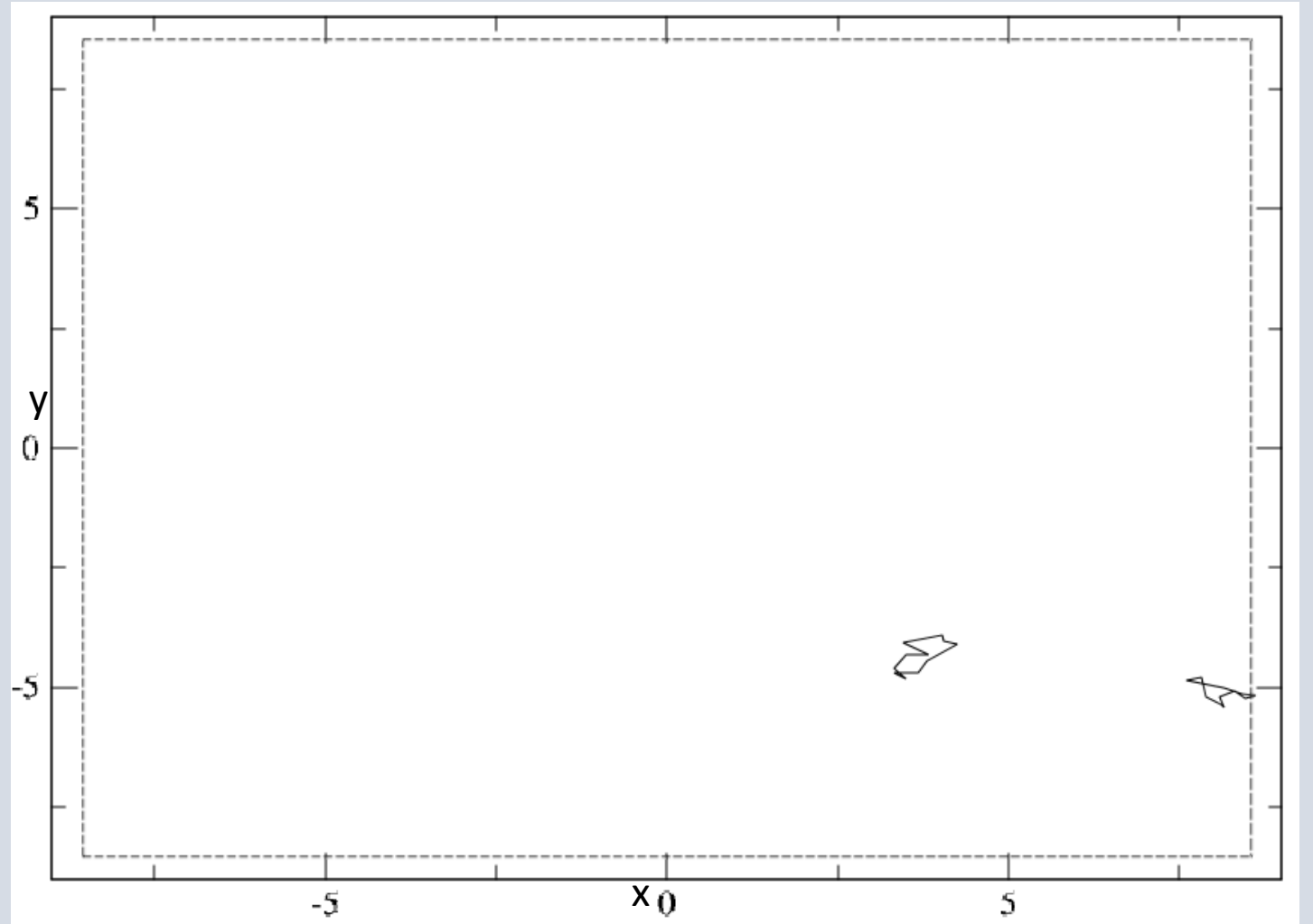
$$\langle O \rangle = Z^{-1} \int d\mathcal{R} O(\mathcal{R}) e^{-U_{eff}(\mathcal{R})}$$

- Integrate using Metropolis algorithm to sample configurations  $\mathcal{R}$  from the probability distribution:

$$P(\mathcal{R}) = Z^{-1} e^{-U_{eff}(\mathcal{R})}$$

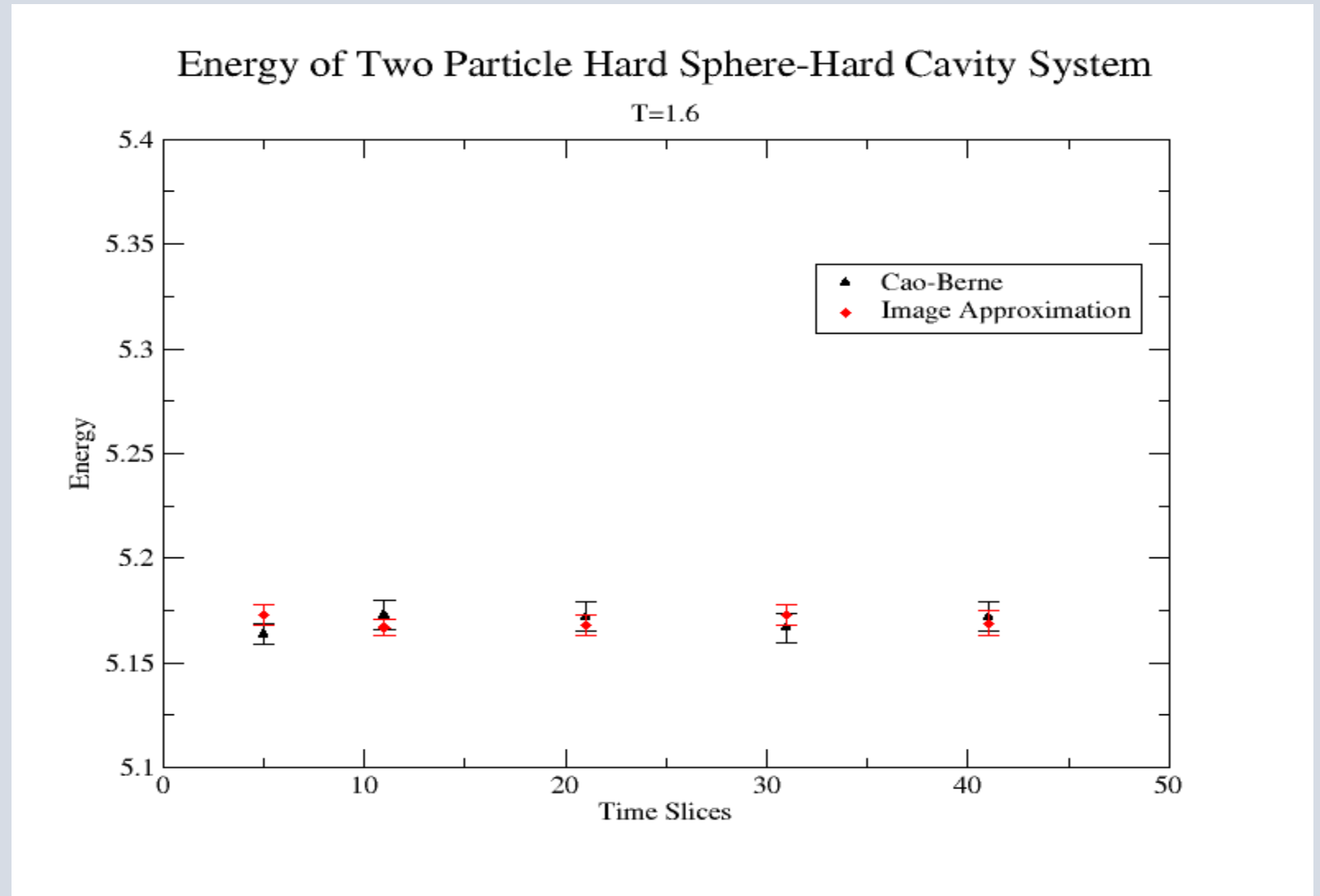
# Calculating Observables

- PIMC simulation of two particles at a given Monte Carlo Step
- Particles are represented as ring polymers



# Convergence in Energy of Two Hard Sphere Particles Inside Hard Cavity

- $V$  is an infinite spherical well in separation distance
- Calculation performed with both propagators
- Convergence seems to be similar between propagators

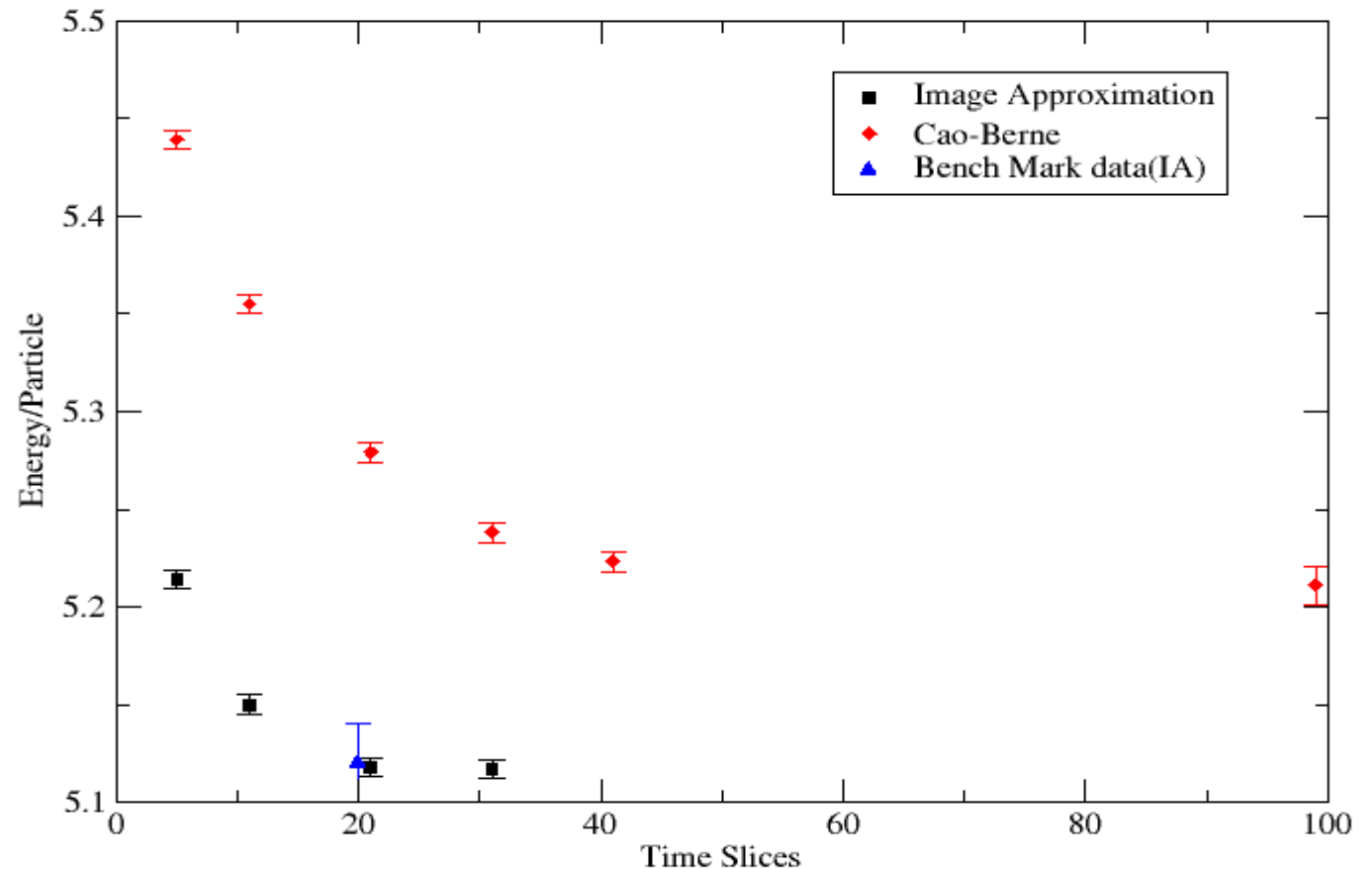


# Convergence in Energy of 'N' Hard Sphere System

- 108 Hard sphere particles
- Convergence in  $E/N$  occurs faster with image approximation propagator

PIMC Simulation of 108 Hard Sphere Particles, Energy/ $N$  Vs. Time Slices

$T=1.6, \text{density}=.15$

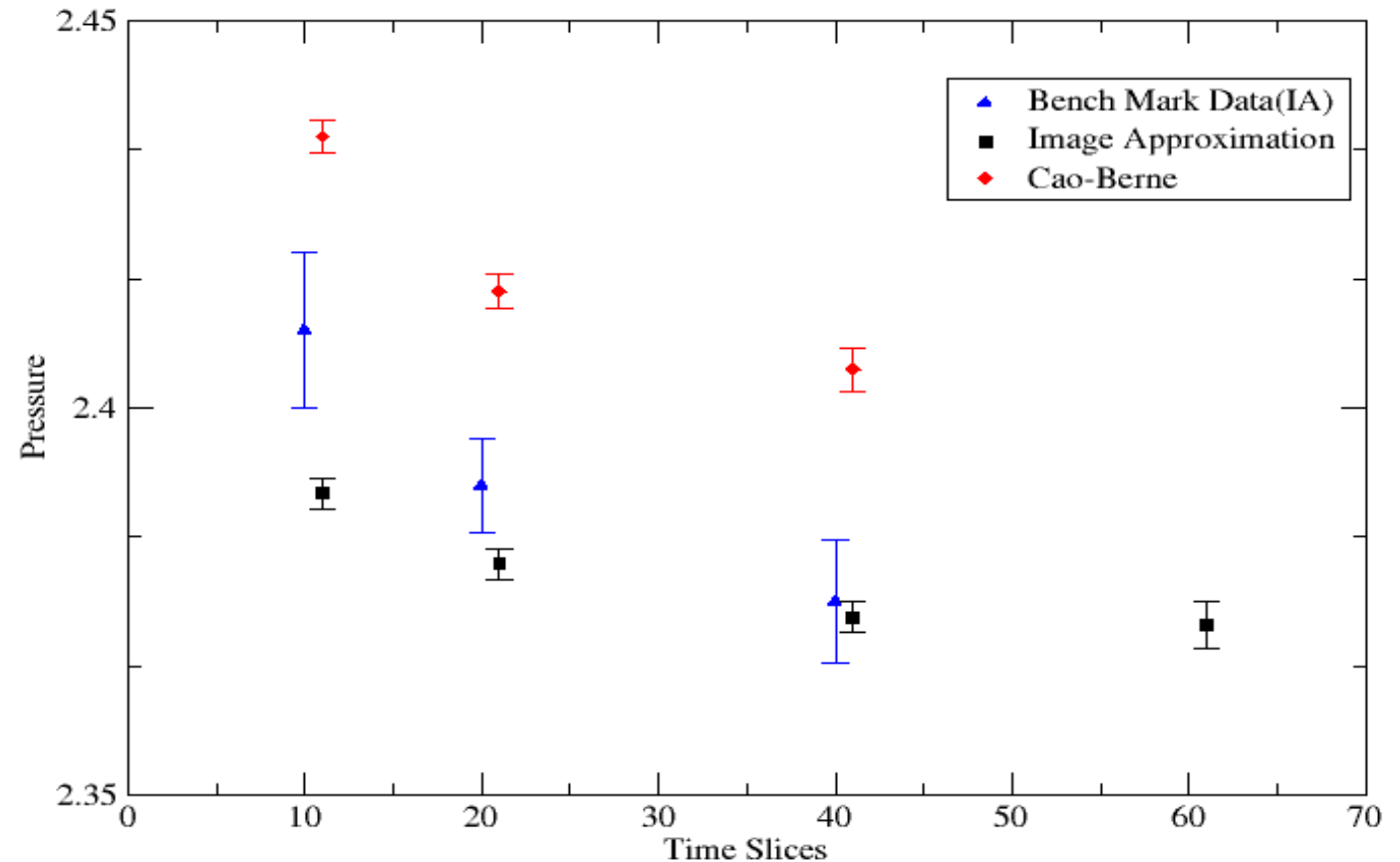


# Convergence in Pressure of 'N' Hard Sphere System

- 108 Hard sphere particles
- Convergence in pressure appears to be equal between propagators

PIMC Simulation of 108 Hard Sphere Particles, Pressure Vs. Time Slices

$T=1.6, \text{Density}=.2063$



# Conclusions

- Path integral Monte Carlo is a technique used to calculate thermodynamic averages on finite temperature systems
- Hard Sphere systems are a popular choice due to being a realistic model that is computationally straight forward
- Varying propagators yield varying convergence/accuracy