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TALENT School on Nuclear Quantum Monte Carlo Methods North Carolina State University (NCSU), July 11-29 2016. Let's consider a many-body general Hamiltonian describing a system of ${\it N}$ particles.

$$H = -\frac{\hbar^2}{2m} \sum_{i}^{N} \nabla_i^2 + \sum_{i < j} v(r_{ij})$$

and we want to calculate the energy:

$$E_0 \leq E = \frac{\langle \Psi_T | H | \Psi_T \rangle}{\langle \Psi_T | \Psi_T \rangle} = \frac{\int dr_1 \dots dr_N \, \Psi_T^*(r_1 \dots r_N) H \Psi_T(r_1 \dots r_N)}{\int dr_1 \dots dr_N \, \Psi_T^*(r_1 \dots r_N) \Psi_T(r_1 \dots r_N)}$$

where $\Psi_T(r_1 \dots r_N)$ is a variational wave function that depends on the coordinates (positions) of the particles, with variational parameters $\{\alpha\}$.

For a given form of $\Psi_T(r_1...r_N)$ we want to find the best value(s) of α giving the lower energy, i.e. we need to study $E[\alpha]$. How do we solve the above integral?

Some definition:

- $r_i = (x_i, y_i, z_i, ...)$ has coordinate(s) of particle i
- when used, $R = r_1...r_N$ (and $dR = dr_1...dr_N$)
- $\Psi(R) = \langle R | \Psi \rangle$



Remember:

$$\int F(x)dx = \int \frac{F(x)}{P(x)}P(x)dx = \int f(x)P(x)dx$$

So we can rewrite the integral as:

$$E = \frac{\int dr_1 \dots dr_N \, \Psi_T^*(r_1 \dots r_N) H \Psi_T(r_1 \dots r_N)}{\int dr_1 \dots dr_N \, \Psi_T^*(r_1 \dots r_N) \Psi_T(r_1 \dots r_N)}$$
$$= \int dR \, P(R) \frac{\Psi_T^*(R) H \Psi_T(R)}{\Psi_T^*(R) \Psi_T(R)}$$

where

$$P(R) = \frac{\Psi_T^*(r_1 \dots r_N) \Psi_T(r_1 \dots r_N)}{\int dr_1 \dots dr_N \, \Psi_T^*(r_1 \dots r_N) \Psi_T(r_1 \dots r_N)} = \frac{|\Psi_T(R)|^2}{\int dR \, |\Psi_T(R)|^2}$$

Easy case: 1-dimensional Harmonic oscillator ($\hbar = m = 1$).

$$H = -\frac{1}{2}\frac{d^2}{dx^2} + \frac{1}{2}\Omega^2 x^2$$

Let's start from the following variational wave function:

$$\Psi_T(x) = e^{-\alpha x^2/2}$$

($\Omega = \alpha$ gives the exact solution).

We will call walker one configuration that will be used to generate new configurations according to the Metropolis/Markov chain algorithm.

In the 1D harmonic oscillator case, walker i is simply

$$W_i = \{x_i\}$$

but in general (3D, multiparticles, spins, etc.) each walker will contain many more information, i.e. $\{\vec{r}_1, \dots, \vec{r}_N, \dots\}$.



Following Metropolis, the VMC algorithm can be summarized in the following steps:

- Generate a set of N walkers randomly distributed
- ② Loop over the N walkers, and for each i-th walker:
- **3** Calculate $P(x_i) = |\langle \Psi_T(x) | x_i \rangle|^2$
- Make a "trial" step: $x_{\text{test}} = x + (\xi 0.5)\Delta$
- Calculate $P(x_{\text{test}}) = |\langle \Psi_T(x) | x_{\text{test}} \rangle|^2$
- **©** Calculate the ratio $P(x_{\text{test}})/P(x_i)$ and do Metropolis
- **1** If the move is accepted, set $x_i = x_{\text{test}}$.
- Calculate averages
- Iterate with 2) until the equilibration is reached, then reset estimators and iterate until the error is small enough

Note: in VMC, N can be even 1. But using more configurations will be very convenient when parallel machines are available.



The calculation of observables is carried out over a total of M walkers (uncorrelated points in the Markov chain) as

$$\langle O \rangle = \frac{\sum_{i}^{M} \omega_{i} \frac{\langle \Psi_{T}(x) | W_{i} \rangle \langle W_{i} | \hat{O} | \Psi_{T}(x) \rangle}{\langle \Psi_{T}(x) | W_{i} \rangle \langle W_{i} | \Psi_{T}(x) \rangle}}{\sum_{i}^{M} \omega_{i}}$$

where ω_i is the weight of the configuration (usually 1).

For example (remember that $\Psi_T = \exp(-\alpha x^2/2)$

$$\begin{split} &\frac{\langle \Psi_{\mathcal{T}}(x)|W_{i}\rangle\langle W_{i}|\hat{H}|\Psi_{\mathcal{T}}(x)\rangle}{\langle \Psi_{\mathcal{T}}(x)|W_{i}\rangle\langle W_{i}|\Psi_{\mathcal{T}}(x)\rangle} = \\ &= -\frac{1}{2}\frac{\langle \Psi_{\mathcal{T}}(x)|x_{i}\rangle\langle x_{i}|\frac{d^{2}}{dx^{2}}|\Psi_{\mathcal{T}}(x)\rangle}{\langle \Psi_{\mathcal{T}}(x)|x_{i}\rangle\langle x_{i}|\Psi_{\mathcal{T}}(x)\rangle} + \frac{1}{2}\Omega^{2}\frac{\langle \Psi_{\mathcal{T}}(x)|x_{i}\rangle\langle x_{i}|x^{2}|\Psi_{\mathcal{T}}(x)\rangle}{\langle \Psi_{\mathcal{T}}(x)|x_{i}\rangle\langle x_{i}|\Psi_{\mathcal{T}}(x)\rangle} \\ &= -\frac{1}{2}(\alpha^{2}x_{i}^{2} - \alpha) + \frac{1}{2}\Omega^{2}x_{i}^{2} \end{split}$$

Note: if Ψ_T is the exact solution, then $\frac{\Psi_T^*H\Psi_T}{\Psi_T^*\Psi_T}=E=\frac{\Omega}{2} \quad \forall x!$

VMC: the first piece of code

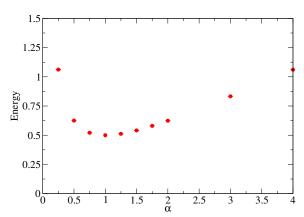
```
do j=1,nstep
do n=1,nwalk
  psiold=exp(-0.5*alpha*xold(n)**2)
  call ran(csi.irn)
  xtest=xold(n)+sigma*(csi-0.5)
  psinew=exp(-0.5*alpha*xtest**2)
  prob=(psinew/psiold)**2
  call ran(csi,irn)
  if (csi.gt.prob) xtest=xold(n)
  eloc=-0.5*((alpha*xtest)**2-alpha)+0.5*omega**2*xtest**2
  xold(n)=xt.est
 enddo
enddo
. . .
```

- Loop over the N walkers, and for each i-th walker: do j=1,nstep do n=1,nwalk
- ② Calculate $P(x_i) = |\langle \Psi_T(x) | x_i \rangle|^2$ psiold=exp(-0.5*alpha*xold(n)**2)
- Make a "trial" step: $x_{\text{test}} = x + (\xi 0.5)\Delta$ call ran(csi,irn) xtest=xold(n)+sigma*(csi-0.5)
- Calculate $P(x_{\text{test}}) = |\langle \Psi_T(x) | x_{\text{test}} \rangle|^2$ psinew=exp(-0.5*alpha*xtest**2)

- Calculate the ratio P(x_{test})/P(x_i) and do Metropolis prob=(psinew/psiold)**2 call ran(csi,irn)
- ② If the move is accepted, set x_i = x_{test}, otherwise reject the move
 if (csi.gt.prob) xtest=xold(n)
 xold(n)=xtest
- Ocalculate averages
 eloc=-0.5*((alpha*xtest)**2-alpha)+0.5*omega**2*xtest**2
- Iterate with 2) until the equilibration is reached, then reset estimators and iterate until the error is small enough

An example: 1D harmonic oscillator, VMC calculations.

Energy as a function of the variational parameter α :



VMC is a tool to calculate the energy of a many-body system for a given wave function as a function of its variational parameters.

This requires to solve multi-dimensional integrals, that are solved using Monte Carlo integration.

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... for now :-)

Backup: Harmonic Oscillator

$$\begin{split} \Psi_T(x) &= e^{-\alpha x^2/2} \\ |\Psi_T|^2 &= \int dx \, e^{-\alpha x^2} = \sqrt{\frac{\pi}{\alpha}} \\ \langle H \rangle &= \int dx \, e^{-\alpha x^2/2} \left(-\frac{1}{2} \frac{d^2}{dx^2} + \frac{1}{2} \Omega^2 x^2 \right) e^{-\alpha x^2/2} \\ &= \int dx \, e^{-\alpha x^2/2} \left[-\frac{1}{2} (\alpha^2 x^2 - \alpha) + \frac{1}{2} \Omega^2 x^2 \right] e^{-\alpha x^2/2} \\ &= \left[\frac{1}{4\alpha} (\Omega^2 - \alpha^2) + \frac{\alpha}{2} \right] \sqrt{\frac{\pi}{\alpha}} \\ \frac{\langle \Psi_T | H | \Psi_T \rangle}{\langle \Psi_T | \Psi_T \rangle} &= \frac{1}{4\alpha} (\Omega^2 - \alpha^2) + \frac{\alpha}{2} \end{split}$$