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Cosmology Homework 5

$$a = \frac{1}{1+z} \quad da = -(1+z)^{-2} dz$$

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code on GitHub at github.com/cmhainje/cosmo-hw

Problem 1: 5/6

part 1. see notebook 2/2

part 2. $\Theta(\eta_0, \vec{x}=0, \hat{p}) = \int_0^{\eta_0} d\eta \, g(\eta) \, S(\eta, \vec{x}=(\eta-\eta_0)\hat{p}, \hat{p})$

1/1 where $g(\eta) \equiv a n_e \sigma_T \exp\left[-\int_{\eta}^{\eta_0} d\eta' a' n_e' \sigma_T\right] \quad S \equiv \Theta_0 + \vec{V}_e \cdot \hat{p} + \frac{3}{20} Q_{ij} \hat{p}_i \hat{p}_j$

for $z \lesssim z_{\text{reio}}$, we assume S to be negligible

$$\Rightarrow \Theta(\eta_0, \vec{x}=0, \hat{p}) \simeq \int_0^{\eta_{\text{reio}}} d\eta \, g(\eta) \, S(\eta, \vec{x}=(\eta-\eta_0)\hat{p}, \hat{p})$$

define $\tau_{\text{reio}} = \int_{\eta_{\text{reio}}}^{\eta_0} dt \, n_e \sigma_T = \int_{\eta_{\text{reio}}}^{\eta_0} d\eta \, a n_e \sigma_T$

$$g(\eta_{\text{reio}}) = a n_e \sigma_T e^{-\tau_{\text{reio}}}$$

$$g(\eta < \eta_{\text{reio}}) = a n_e \sigma_T e^{-\tau_{\text{reio}}} \exp\left[-\int_{\eta}^{\eta_{\text{reio}}} d\eta' a' n_e' \sigma_T\right]$$

$$\equiv e^{-\tau_{\text{reio}}} \tilde{g}(\eta), \text{ defining } \tilde{g} \text{ to be the value of } g \text{ ignoring reio}$$

$$\Rightarrow \Theta(\eta_0, \vec{x}=0, \hat{p}) = e^{-\tau_{\text{reio}}} \int_0^{\eta_{\text{reio}}} d\eta \, \tilde{g}(\eta) \, S(\eta, \vec{x}=(\eta-\eta_0)\hat{p}, \hat{p}) \quad \checkmark \quad \square$$

part 3. given Y_{He} and assuming H, He are fully ionized, $n_p = \frac{1-Y_{\text{He}}}{m_p} \rho_b$ and $n_{\text{He}^{++}} = \frac{Y_{\text{He}}}{m_{\text{He}^{++}}} \rho_b$.

2/3 assume charge neutrality $\Rightarrow n_e = n_H + 2n_{\text{He}} = \left(\frac{1-Y_{\text{He}}}{m_p} + \frac{2Y_{\text{He}}}{m_{\text{He}^{++}}}\right) \rho_b$

$$\rho_b = \rho_{b,0} a^{-3}, \quad \rho_{b,0} = \rho_{\text{crit},0} \Omega_{b,0} = \rho_{\text{crit},0} \omega_b h^2 = \frac{3(100 \text{ km/s/Mpc})^2}{8\pi G} \omega_b$$

$$\begin{aligned} \text{then } \tau_{\text{reio}} &= \int_{\eta_{\text{reio}}}^{\eta_0} dt \, n_e \sigma_T \\ &= \left(\frac{1-Y_{\text{He}}}{m_p} + \frac{2Y_{\text{He}}}{m_{\text{He}^{++}}}\right) \rho_{b,0} \sigma_T \int_{\eta_{\text{reio}}}^{\eta_0} dt \, a^{-3} \\ &= \left(\frac{1-Y_{\text{He}}}{m_p} + \frac{2Y_{\text{He}}}{m_{\text{He}^{++}}}\right) \rho_{b,0} \sigma_T \int_{a_{\text{reio}}}^1 \frac{da}{\dot{a} a^3} \end{aligned}$$

assuming $\omega_k = 0$ and $\omega_r \ll \omega_m, \omega_\Lambda$ then $\dot{a} a^3 = (100 \frac{\text{km/s}}{\text{Mpc}}) \sqrt{\omega_\Lambda} a_{\text{m}\Lambda}^4 x^2 \sqrt{x+x^4} \quad (x \equiv a_{\text{m}\Lambda})$

$$\begin{aligned} &\approx \left(\frac{1-Y_{\text{He}}}{m_p} + \frac{2Y_{\text{He}}}{m_{\text{He}^{++}}}\right) \frac{\rho_{b,0} \sigma_T}{(100 \text{ km/s/Mpc}) \sqrt{\omega_\Lambda} a_{\text{m}\Lambda}^3} \int_{a_{\text{reio}}/a_{\text{m}\Lambda}}^{1/a_{\text{m}\Lambda}} \frac{dx}{x^2 \sqrt{x+x^4}} \quad \checkmark \quad a_{\text{m}\Lambda} = (\omega_m/\omega_\Lambda)^{1/3} \\ &= \left(\frac{1-Y_{\text{He}}}{m_p} + \frac{2Y_{\text{He}}}{m_{\text{He}^{++}}}\right) \frac{\rho_{b,0} \sigma_T}{(100 \text{ km/s/Mpc}) \sqrt{\omega_\Lambda} a_{\text{m}\Lambda}^3} \frac{2}{3} \left[\sqrt{1 + \left(\frac{a_{\text{m}\Lambda}}{a_{\text{reio}}}\right)^3} - \sqrt{1 + \left(\frac{a_{\text{m}\Lambda}}{1}\right)^3} \right] \quad \checkmark \\ &= \left(\frac{1-Y_{\text{He}}}{m_p} + \frac{2Y_{\text{He}}}{m_{\text{He}^{++}}}\right) \frac{\rho_{b,0} \sigma_T}{(100 \text{ km/s/Mpc}) \sqrt{\omega_m}} \frac{2}{3} \left[\sqrt{a_{\text{m}\Lambda}^{-3} + a_{\text{reio}}^{-3}} - \sqrt{a_{\text{m}\Lambda}^{-3} + 1} \right] \quad \checkmark \\ &= \alpha \left[\sqrt{a_{\text{m}\Lambda}^{-3} + (1+z_{\text{reio}})^3} - \sqrt{a_{\text{m}\Lambda}^{-3} + 1} \right] \quad \text{where } \alpha = 2.415 \cdot 10^{-3} \text{ for Planck 18 parameters} \end{aligned}$$

inverting: $z_{\text{reio}} = \left[\left(\frac{\tau_{\text{reio}}}{\alpha}\right)^2 + 2 \frac{\tau_{\text{reio}}}{\alpha} \sqrt{a_{\text{m}\Lambda}^{-3} + 1} + 1 \right]^{1/3} - 1 = 11.88 \quad (\text{using } \tau_{\text{reio}} = 0.054 \pm 0.007)$

error prop: $\sigma_z = \frac{\partial z}{\partial \tau} \sigma_\tau = \frac{1}{3} \left[\left(\frac{\tau_{\text{reio}}}{\alpha}\right)^2 + 2 \frac{\tau_{\text{reio}}}{\alpha} \sqrt{a_{\text{m}\Lambda}^{-3} + 1} + 1 \right]^{-2/3} \left[\frac{2\tau_{\text{reio}}}{\alpha^2} + \frac{2}{\alpha} \sqrt{a_{\text{m}\Lambda}^{-3} + 1} \right] \sigma_\tau = 0.43$

I had a small bug in my code to compute z_{reio} : accidentally wrote $a_{\text{m}\Lambda} = (\Omega_m/\Omega_\Lambda)^{1/3}$. -1

after fix, I got $z_{\text{reio}} = 9.8 \pm 0.29$. I am dubious about the assumption of matter domination, though.

Problem 2: 4/8

part 1. $\Theta(\eta, \vec{x}=0, \hat{p}) = \Theta_0(\eta_*, -(\eta-\eta_*)\hat{p}) + \hat{p} \cdot \vec{V}_e(\eta_*, -(\eta-\eta_*)\hat{p})$
 0/3 Fourier: $\Theta_0(\eta_*, \vec{x}) = \Theta_* e^{i\vec{k} \cdot \vec{x}}$, scalar V_e : $\vec{V}_e = i\hat{k} V_* e^{i\vec{k} \cdot \vec{x}}$
 then $\Theta(\eta, \vec{0}, \hat{p}) = (\Theta_* + i\hat{k} \cdot \hat{p} V_*) \exp[-i k(\eta-\eta_*)(\hat{k} \cdot \hat{p})]$

I was confused and didn't see that I needed to integrate over \hat{p} .

so $\Theta_0 = \Theta_* \exp[-i k(\eta-\eta_*)(\hat{k} \cdot \hat{p})]$ and $V_g = V_* \exp[-i k(\eta-\eta_*)(\hat{k} \cdot \hat{p})]$?
 can be written $\Theta_0 = \Theta_* (e^{-i\hat{k} \cdot \hat{p}})^{k(\eta-\eta_*)}$, $V_g = V_* (e^{-i\hat{k} \cdot \hat{p}})^{k(\eta-\eta_*)}$
 $e^{-i\hat{k} \cdot \hat{p}} < 1$ except when $\hat{k} \perp \hat{p}$, so $k(\eta-\eta_*) \gg 1$ suppresses

part 2. $\frac{1}{H(\eta_*)}$ is a physical length, $\frac{1}{\partial \ell(\eta_*)}$ is the corresponding comoving length
 1/2 taking $z_* = 1100$, comoving wavenumber $k_* = \frac{1}{c \text{ length}} = a H(\eta_*) = 0.0048 \text{ Mpc}^{-1} = (208 \text{ Mpc})^{-1}$

didn't calculate suppression

part 3. $\Theta = \sum_{\ell} (2\ell+1) \Theta_{\ell}(\eta) P_{\ell}(\hat{p} \cdot \hat{k})$
 but also $\Theta = (\Theta_* + i\hat{k} \cdot \hat{p} V_*) \exp[i k \chi_* (\hat{k} \cdot \hat{p})]$ defining $\chi_* = -(\eta-\eta_*)$
 3/3 $= (\Theta_* + V_* \frac{d}{d(k\chi_*)}) \sum_{\ell} (2\ell+1) i^{\ell} j_{\ell}(k\chi_*) P_{\ell}(\hat{k} \cdot \hat{p})$
 $= \sum_{\ell} (2\ell+1) i^{\ell} [\Theta_* j_{\ell}(k\chi_*) + V_* j'_{\ell}(k\chi_*)] P_{\ell}(\hat{k} \cdot \hat{p})$

then $\Theta_{\ell} = i^{\ell} [\Theta_* j_{\ell}(k\chi_*) + V_* j'_{\ell}(k\chi_*)]$ ✓

↑ sign wrong due to $(-i)^{\ell}$, I'm counting it

Problem 3: 5/5

part 1. $\tilde{T}^\mu{}_\nu = \frac{\partial \tilde{x}^\mu}{\partial x^\alpha} \frac{\partial x^\beta}{\partial \tilde{x}^\nu} T^\alpha{}_\beta(x^\lambda)$

$$\begin{aligned} &= (\delta^\mu{}_\alpha + \partial_\alpha \xi^\mu) (\delta^\beta{}_\nu - \partial_\nu \xi^\beta) T^\alpha{}_\beta(\tilde{x}^\lambda - \xi^\lambda) \\ &= (\delta^\mu{}_\alpha \delta^\beta{}_\nu - \delta^\mu{}_\alpha \partial_\nu \xi^\beta + \partial_\alpha \xi^\mu \delta^\beta{}_\nu) [T^\alpha{}_\beta(\tilde{x}^\lambda) - \xi^\lambda \partial_\lambda T^\alpha{}_\beta] \\ &= T^\mu{}_\nu(\tilde{x}^\lambda) - (\partial_\nu \xi^\beta T^\mu{}_\beta - \partial_\alpha \xi^\mu T^\alpha{}_\nu + \xi^\lambda \partial_\lambda T^\mu{}_\nu) \end{aligned}$$

$$\tilde{T}^\mu{}_\nu = \bar{T}^\mu{}_\nu(\tilde{x}^\lambda) + \delta \tilde{T}^\mu{}_\nu(\tilde{x}^\lambda) \quad \text{and} \quad T^\mu{}_\nu(\tilde{x}^\lambda) = \bar{T}^\mu{}_\nu(\tilde{x}^\lambda) + \delta T^\mu{}_\nu(\tilde{x}^\lambda)$$

so $\delta \tilde{T}^\mu{}_\nu - \delta T^\mu{}_\nu|_{\tilde{x}^\lambda} = -(\partial_\nu \xi^\beta \bar{T}^\mu{}_\beta - \partial_\alpha \xi^\mu \bar{T}^\alpha{}_\nu + \xi^\lambda \partial_\lambda \bar{T}^\mu{}_\nu)$ ✓

part 2. $T^0{}_0 = -(\bar{\rho} + \delta\rho) \quad T^i{}_0 = -(\bar{\rho} + \bar{P}) V^i \quad T^i{}_j = (\bar{P} + \delta P) \delta^i{}_j + \Pi^i{}_j$

3/3 $\Rightarrow \delta T^0{}_0 = -\delta\rho \quad \delta T^i{}_0 = -(\bar{\rho} + \bar{P}) V^i \quad \delta T^i{}_j = \delta P \delta^i{}_j + \Pi^i{}_j$

thus

$${}^0{}_0: (-\tilde{\delta\rho}) - (-\delta\rho) = -(\partial_0 \xi^\beta \bar{T}^0{}_\beta - \partial_\alpha \xi^0 \bar{T}^\alpha{}_0 + \xi^\lambda \partial_\lambda \bar{T}^0{}_0)$$

$$\delta\rho - \tilde{\delta\rho} = -(-\partial_0 \xi^0 \bar{\rho} + \partial_0 \xi^0 \bar{\rho} - \xi^\lambda \partial_\lambda \bar{\rho})$$

$$\tilde{\delta\rho} = \delta\rho - \xi^0 \partial_0 \bar{\rho} \quad \checkmark$$

$${}^i{}_0: \delta \tilde{T}^i{}_0 - \delta T^i{}_0 = -(\partial_0 \xi^\beta \bar{T}^i{}_\beta - \partial_\alpha \xi^i \bar{T}^\alpha{}_0 + \xi^\lambda \partial_\lambda \bar{T}^i{}_0)$$

$$-(\bar{\rho} + \bar{P})(\tilde{V}^i - V^i) = -\partial_0 \xi^i (\bar{\rho} + \bar{P})$$

$$\tilde{V}^i = V^i + \partial_0 \xi^i \quad \checkmark$$

$${}^i{}_j: \delta \tilde{T}^i{}_j - \delta T^i{}_j = -(\partial_j \xi^\beta \bar{T}^i{}_\beta - \partial_\alpha \xi^i \bar{T}^\alpha{}_j + \xi^\lambda \partial_\lambda \bar{T}^i{}_j)$$

$$\tilde{\delta P} \delta^i{}_j + \tilde{\Pi}^i{}_j - \delta P \delta^i{}_j - \Pi^i{}_j = -\xi^\lambda \partial_\lambda \bar{P} \delta^i{}_j$$

grouping terms proportional to $\delta^i{}_j$ and \perp to it gives:

$$\tilde{\delta P} = \delta P - \xi^0 \partial_0 \bar{P} \quad \checkmark \quad \text{and} \quad \tilde{\Pi}^i{}_j = \Pi^i{}_j \quad \checkmark$$

part 3. $\mathcal{R} = -C - \frac{1}{3} k^2 E + \mathcal{H} \frac{\delta\rho}{\bar{\rho}}$

1/1 $\tilde{\mathcal{R}} = -\tilde{C} - \frac{1}{3} k^2 \tilde{E} + \mathcal{H} \frac{\delta\tilde{\rho}}{\bar{\rho}}$

$$= -[C - \mathcal{H} \xi^0 + \frac{1}{3} k^2 L] - \frac{1}{3} k^2 [E - L] + \mathcal{H} \frac{1}{\bar{\rho}} [\delta\rho - \xi^0 \partial_0 \bar{\rho}]$$

$$= -C + \cancel{\mathcal{H} \xi^0} - \cancel{\frac{1}{3} k^2 L} - \frac{1}{3} k^2 E + \cancel{\frac{1}{3} k^2 L} + \mathcal{H} \frac{\delta\rho}{\bar{\rho}} - \cancel{\mathcal{H} \xi^0}$$

$$= \mathcal{R}$$

□

Problem 4: 4/9

part 1. $g_{\mu\nu} = a^2 [\eta_{\mu\nu} + h_{\mu\nu}]$

1/1

consider $A^{\mu\nu} = a^{-2} [\eta^{\mu\nu} - h^{\mu\nu}] = a^{-2} [\eta^{\mu\alpha} - \eta^{\mu\alpha} \eta^{\nu\beta} h_{\alpha\beta}]$

if $A^{\mu\nu} g_{\nu\rho} = \delta^\mu_\rho$, then $A^{\mu\nu}$ is the inverse metric i.e. $g^{\mu\nu} = A^{\mu\nu}$.

let's evaluate it:

$$\begin{aligned} A^{\mu\nu} g_{\nu\rho} &= a^{-2} [\eta^{\mu\nu} - \eta^{\mu\alpha} \eta^{\nu\beta} h_{\alpha\beta}] a^2 [\eta_{\nu\rho} + h_{\nu\rho}] \\ &= [\eta^{\mu\nu} - \eta^{\mu\alpha} \eta^{\nu\beta} h_{\alpha\beta}] [\eta_{\nu\rho} + h_{\nu\rho}] \\ &= \delta^\mu_\rho + \eta^{\mu\nu} h_{\nu\rho} - \eta^{\mu\alpha} \delta^\beta_\rho h_{\alpha\beta} - \eta^{\mu\alpha} \eta^{\nu\beta} h_{\alpha\beta} h_{\nu\rho} \\ &= \delta^\mu_\rho + \cancel{\eta^{\mu\nu} h_{\nu\rho}} - \cancel{\eta^{\mu\alpha} h_{\alpha\rho}} + O(h^2) \\ &= \delta^\mu_\rho + O(h^2) \end{aligned}$$

so indeed $g^{\mu\nu} = a^{-2} [\eta^{\mu\nu} - \eta^{\mu\alpha} \eta^{\nu\beta} h_{\alpha\beta}]$ to linear order in perturbation h . □

part 2. note that $h^{\mu\nu} = -h_{\mu\nu}$ for our metric

2/2 $G_{00} = 3 \left(\frac{a'}{a} \right)^2 - 2 \frac{a'}{a} \partial_i h_{0i} = 3 \left(\frac{a'}{a} \right)^2$
 $= 3 \mathcal{H}^2$ ✓

$$\begin{aligned} G_{0i} &= \left(\frac{a'}{a} \right)^2 h_{0i} - 2 \frac{a''}{a} h_{0i} + \partial_j \partial_{(i} h_{j)0} = \left(\frac{a'}{a} \right)^2 h_{0i} - 2 \frac{a''}{a} h_{0i} - \frac{1}{2} \partial_j \partial_j h_{0i} \\ &= - \left(2\mathcal{H}' + 2\mathcal{H}^2 + \frac{1}{2} \vec{\nabla}^2 \right) h_{0i} \quad \checkmark \end{aligned}$$

$$\begin{aligned} G_{ij} &= \delta_{ij} \left[\left(\frac{a'}{a} \right)^2 - 2 \frac{a''}{a} + 2 \frac{a'}{a} \partial_k h_{0k} + \partial_0 \partial_k h_{0k} \right] \\ &\quad - \frac{a'}{a} \partial_j h_{0i} - \frac{a'}{a} \partial_i h_{0j} - \frac{1}{2} \partial_0 \partial_j h_{0i} - \frac{1}{2} \partial_0 \partial_i h_{0j} \\ &= \delta_{ij} \left[-2\mathcal{H}' - \mathcal{H}^2 \right] - \mathcal{H} \partial_j h_{0i} - \mathcal{H} \partial_i h_{0j} - \frac{1}{2} \partial_0 \partial_j h_{0i} - \frac{1}{2} \partial_0 \partial_i h_{0j} \quad \checkmark \end{aligned}$$

part 3. EFEs: $G_{\mu\nu} = 8\pi G T_{\mu\nu} = 8\pi G T_{\nu\mu} = 8\pi G g_{\nu\alpha} T^\alpha_\mu$

2/3 EFE_{0i}: $G_{0i} = 8\pi G g_{i\alpha} T^\alpha_0 = 8\pi G [g_{i0} T^0_0 + g_{ij} T^j_0]$
 $= -8\pi G [h_{0i} (\bar{\rho} + \delta\rho) + (\bar{\rho} + \bar{P}) V_i] a^2 \quad (-1)$
 $- \left(2\mathcal{H}' + 2\mathcal{H}^2 + \frac{1}{2} \vec{\nabla}^2 \right) h_{0i} = -8\pi G (\bar{\rho} h_{0i} + (\bar{\rho} + \bar{P}) V_i)$

re-arranging: $(\mathcal{H}^2 + 2\mathcal{H}' + \frac{1}{2} \vec{\nabla}^2 - 8\pi G \bar{\rho}) h_{0i} = 8\pi G (\bar{\rho} + \bar{P}) V_i$

this \bar{P} term is not in the solution: why?

isn't $T_{i0} = g_{i\mu} T^\mu_0$ not there

$$= g_{i0} T^0_0 + g_{ij} T^j_0$$

$$= (a^2 h_{0i}) (-(\bar{\rho} + \delta\rho)) + (a^2 \delta_{ij}) (-(\bar{\rho} + \bar{P}) V^j)$$

$$= -a^2 (h_{0i} \bar{\rho} + (\bar{\rho} + \bar{P}) V^i) ?$$

part 4.

0/3

$$\begin{aligned} 0 &= \nabla_{\mu} T^{\mu}_{i} = \nabla_0 T^0_{i} + \nabla_j T^j_{i} \\ &= -\nabla_0 (\bar{\rho} + \bar{p}) v_i + \nabla_j (\bar{p} + \delta p) \delta^j_i \\ &= -\nabla_0 (\bar{\rho} + \bar{p}) v_i + \nabla_i \delta p \end{aligned}$$

$$\nabla_0 \left[\left(\mathcal{H}^2 + 2\mathcal{H}' + \frac{1}{2} \vec{\nabla}^2 - 8\pi G \bar{\rho} \right) h_{0i} \right] = 8\pi G \nabla_i \delta p \quad ?$$

1) couldn't figure this out