

Graduate Cosmology Spring 2025

Homework 2 solutions

How to self-grade: full points for correct answer with correct reasoning; half-points for correct reasoning but incorrect answer due to algebra error; zero point for incorrect reasoning (even if final answer is correct out of luck).

Problem 1

1. We rewrite $dt/da = 1/\dot{a} = \frac{1}{aH}$. Using Friedmann's equation, and the boundary condition $t = 0$ at $a = 0$, we get

$$t(a) = H_0^{-1} \int_0^a \frac{db}{b(\Omega_r b^{-4} + \Omega_m b^{-3} + \Omega_k b^{-2} + \Omega_\Lambda)^{1/2}}.$$

2. (a) For $a \ll a_{m\Lambda}$, since $b < a$ in the integral, we may neglect Ω_Λ (as well as Ω_k) in the integrand, so we have

$$t \approx H_0^{-1} \int_0^a \frac{db}{b\sqrt{\Omega_r b^{-4} + \Omega_m b^{-3}}} = H_0^{-1} \int_0^a \frac{bdb}{\sqrt{\Omega_r + \Omega_m b}} = \frac{a_{\text{eq}}^{3/2}}{H_0 \sqrt{\Omega_m}} \int_0^{a/a_{\text{eq}}} \frac{x dx}{\sqrt{1+x}}.$$

Using the hint, we get

$$t \approx \frac{a_{\text{eq}}^{3/2}}{H_0 \sqrt{\Omega_m}} \frac{2}{3} \left(2 + (a/a_{\text{eq}} - 2) \sqrt{1 + a/a_{\text{eq}}} \right),$$

where we set the constant adequately so that $t(0) = 0$.

(b) Let us now compute the result in the limit $a \ll a_{\text{eq}}$. We first Taylor-expand $\sqrt{1 + a/a_{\text{eq}}} \approx 1 + \frac{1}{2}(a/a_{\text{eq}}) - \frac{1}{8}(a/a_{\text{eq}})^2$. We see that the constant term and the term linear in a/a_{eq} cancel out, so the leading order is quadratic in a/a_{eq} :

$$t(a \ll a_{\text{eq}}) \approx \frac{a_{\text{eq}}^{3/2}}{2H_0 \sqrt{\Omega_m}} (a/a_{\text{eq}})^2.$$

The $a \gg a_{\text{eq}}$ expansion is easier:

$$t(a_{\text{eq}} \ll a \ll a_{m\Lambda}) \approx \frac{2}{3} \frac{a^{3/2}}{H_0 \sqrt{\Omega_m}}.$$

Both of these solutions match the $a(t)$ results we found in class, using $\Omega_r = a_{\text{eq}}\Omega_m$.

- (c) We plug in $a = a_{\text{eq}}$ and find

$$t_{\text{eq}} = \frac{2}{3} \frac{a_{\text{eq}}^{3/2}}{H_0 \sqrt{\Omega_m}} (2 - \sqrt{2})$$

With $H_0^{-1} \approx 14$ Gyr, we obtain $t_{\text{eq}} \approx 50$ kyr.

3. (a) For $a \gg a_{\text{eq}}$, the integral is dominated by $b \approx a$, i.e. we may neglect matter and radiation:

$$\begin{aligned} t(a \gg a_{\text{eq}}) &\approx H_0^{-1} \int_0^a \frac{db}{b\sqrt{\Omega_m b^{-3} + \Omega_\Lambda}} = \frac{1}{H_0 \sqrt{\Omega_\Lambda}} \int_0^{a/a_{m\Lambda}} \frac{dx}{x\sqrt{1+1/x^3}} = \frac{1}{H_0 \sqrt{\Omega_\Lambda}} \int_0^{a/a_{m\Lambda}} \frac{x^{1/2} dx}{\sqrt{1+x^3}} \\ &= \frac{2}{3} \frac{1}{H_0 \sqrt{\Omega_\Lambda}} \text{Arcsinh}((a/a_{m\Lambda})^{3/2}), \end{aligned}$$

where we recall that $a_{m\Lambda} = (\Omega_m/\Omega_\Lambda)^{1/3}$. It is straightforward to see that this is precisely the inverse of (2.177) in Baumann.

(b) We now plug in $a = 1$ in this result to find the age of the Universe:

$$t_0 = \frac{2}{3} \frac{1}{H_0 \sqrt{\Omega_\Lambda}} \text{Arcsinh}(\sqrt{\Omega_\Lambda/\Omega_m}) \approx 0.95 H_0^{-1} \approx 13.7 \text{ Gyr}$$

where we used $H_0 \approx 68 \text{ km/s/Mpc}$.

4) The conformal time η satisfies the ODE $d\eta = dt/a$. So it can be expressed as an integral in a similar to $t(a)$, but with an additional factor of $1/a$. In the matter- Λ era, this is

$$\eta(a) \approx \frac{1}{H_0} \int_0^a \frac{db}{b^2 \sqrt{\Omega_m b^{-3} + \Omega_\Lambda}}.$$

This integral is not analytic, so we just do it numerically, for $a = 1$, and find $\eta_0 \approx 3.24 H_0^{-1} \approx 47 \text{ Gyr} \approx 14 \text{ Gpc}$.

Problem 2:

1) By definition $d\eta = dt/a(t) = dt/da \, da/a = da/a\dot{a} = da/(a^2 H(a))$. Since we are asked to express things in terms of redshift $z = 1/a - 1$, we get $dz = -da/a^2$, hence $d\eta = -dz/H(z)$. So, putting it together, we get

$$\begin{aligned} \chi(z) &= \int_{\eta(z)}^{\eta_0} d\eta = \int_0^z \frac{dz'}{H(z')} = \frac{1}{H_0} \int_0^z \frac{dz'}{\sqrt{\Omega_r(1+z')^4 + \Omega_m(1+z')^3 + \Omega_k(1+z')^2 + \Omega_\Lambda}} \\ &= \frac{1}{100 \text{ km/s/Mpc}} \int_0^z \frac{dz'}{\sqrt{\omega_r(1+z')^4 + \omega_m(1+z')^3 + \omega_k(1+z')^2 + \omega_\Lambda}}. \end{aligned}$$

2. and 3.: see attached Python notebook.

4. For a spatially flat Universe, $S_k(\chi) = \chi$, and the luminosity distance is therefore just $d_L(z) = (1+z)\chi(z)$. Neglecting radiation, and since the Universe is spatially flat, $\Omega_m = 1 - \Omega_\Lambda$, so that

$$d_L(z) = H_0^{-1}(1+z) \int_0^z \frac{dz'}{\sqrt{\Omega_\Lambda + (1-\Omega_\Lambda)(1+z')^3}}.$$

This can be expressed in terms of hypergeometric functions but that is not particularly useful, so we compute it numerically.

5. (a) We first find the parameter K that minimizes the χ^2 by taking the derivative of χ^2 with respect to K and setting it to zero. This gives

$$K_{\min}(\Omega_\Lambda) = \left(\sum_i \frac{\mu_i - 5 \log_{10}(F(z_i, \Omega_\Lambda))}{\sigma_i^2} \right) / \left(\sum_i 1/\sigma_i^2 \right).$$

We then plug this back into the original χ^2 to find the desired expression.

(b) See attached notebook. The best-fit is $\Omega_\Lambda = 0.715$.