Black Hole Thermodynamics

And Hawking Radiation

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Advanced Seminar on Relativity and Cosmology



Outline

1. Introduction

- Tim
- 2. Laws of black hole mechanics and thermodynamics
 - 2.0 The second law Tim
 - 2.1 The zeroth law Sebastiàn
 - 2.2 The first law Sebastiàn
 - 2.3 The third law Sebastiàn
- 3. Hawking radiation Yi-Fan

Outline

- 1. Motivations
- 2. Particle creation in curved space-time
- 3. Gravitational collapse
- 4. Eternal black hole
- 5. Information loss problem



Motivations for the Radiation

Developments in classical physics

- (Speculative) entropy and temperature of black holes [Bek73; Bek80]
- Penrose process and super-radiance



Motivations for the Radiation

Developments in quantum physics

- Arguments from fundamental principles
 - Uncertainty principle (black holes "might split up" [Sus08])
 - Vacuum polarisation (e.g. Schwinger Effect, where electron-positron pairs are pulled out from vacuum in strong external electric field), see e.g. [FN98]
- Arguments from field theory in curved space-time: definition of particle ambiguous [Ful73], leading to particle creations
 - in dynamical universes, see e.g. [PT09]
 - for accelerated observers [Ful73; Dav75; Unr76]



Scalar Field in Special Relativity

Warm-up

Diagonalising Hamiltonian: Fourier kernel & 'momentum' p

$$\tilde{\phi}(p) \propto \int \mathrm{d} x \, \mathrm{e}^{-\mathrm{i} p x} \phi(x), \quad \tilde{\pi}(p) \propto \int \mathrm{d} x \, \mathrm{e}^{+\mathrm{i} p x} \pi(x).$$

- Trsfing field between inertial obs. A and B: Poincaré $\mathcal P$

$$\begin{split} x \to \mathcal{P} x &= \Lambda x + x_0, \quad p \to \mathcal{P} p = p + p_0; \\ \tilde{\phi}_A(p) \to \tilde{\phi}_A\big(\mathcal{P}^{-1} p\big) \propto \tilde{\phi}_B\big(p'\big), \quad \text{etc.} \end{split}$$

Canonical quantisation: ladder operators

$$\hat{a}_\phi(p) \propto E_p^{+\frac{1}{2}} \hat{\hat{\phi}}(p) + \mathbb{i} E_p^{-\frac{1}{2}} \hat{\tilde{\pi}}^\dagger(p), \quad \text{etc.}$$

• Trsfing ladder optrs btw. inertial obs.: particle remains p.

$$\hat{b}(p') = \widehat{U}_{\mathcal{P}} \hat{a}(p) \widehat{U}_{\mathcal{P}}^{\dagger} \propto \hat{a}(\mathcal{P}p), \quad \text{etc.}$$



Scalar Field in General Relativity

The technical advance we just mentioned

• Diagonalising Hamiltonian: general kernel & parameter k

$$\tilde{\phi}(k) \propto \int \mathrm{d}x \, K(k;x) \phi(x), \quad \tilde{\pi}(k) \propto \int \mathrm{d}x \, K^*(k;x) \pi(x)$$

- Trsfing field between (time-like) obs.: general diffeom. $\mathcal D$

$$\begin{split} x \to \chi = f(x), \quad k \to \kappa = g(k); \\ \tilde{\phi}_A(k) \to \tilde{\phi}_A\big(g^{-1}(\kappa)\big) = \int \mathrm{d}\kappa \, T(k;\kappa) \tilde{\phi}_B(\kappa), \quad \text{etc.} \end{split}$$

Canonical quantisation: ladder operators

$$\hat{a}_{\phi}(k) \propto E_k^{+\frac{1}{2}} \hat{\tilde{\phi}}(k) \pm \mathbb{i} E_k^{-\frac{1}{2}} \hat{\tilde{\pi}}^{\dagger}(k), \quad \text{etc.}$$

• Transforming ladder optrs: Bogolyubov, mixing p. & anti-p.

$$\widehat{b}(\kappa) = \widehat{U}_{\mathcal{D}}\widehat{a}(k)\widehat{U}_{\mathcal{D}}^{\dagger} \propto \int \mathrm{d}k\, \alpha(\kappa;k)\widehat{a}(k) + \beta(\kappa;k)\widehat{a}^{\dagger}(k), \quad \text{etc.}$$

Scalar Field in Special and General Relativity

Particle 'creation'

 \bullet Defining the $\hat{a}\text{-vacuum}$ as being annihilated by all $\hat{a}\text{'s}$

$$\hat{a}(p \text{ or } k) \left| \Omega \right\rangle \equiv 0$$

Number operators transform as

$$\begin{split} \hat{a}^{\dagger}(p)\hat{a}(p) &\to \hat{b}^{\dagger}(p')\hat{b}(p') \\ &\propto \hat{a}^{\dagger}(\mathcal{P}p)\hat{a}(\mathcal{P}p); &= \dots \hat{a}^{\dagger}\hat{a} + \dots \hat{a}\hat{a} + \dots \hat{a}^{\dagger}\hat{a}^{\dagger} \\ &+ \int \mathrm{d}k \left|\beta(\kappa;k)\right|^2 \hat{a}(k)\hat{a}^{\dagger}(k). \end{split}$$

• $\langle \hat{n}_a \rangle_{\Omega}$'s transform as

$$\langle \hat{a}^{\dagger}(p)\hat{a}(p)\rangle \to \langle \hat{b}^{\dagger}(p')\hat{b}(p')\rangle \quad \langle \hat{a}^{\dagger}(k)\hat{a}(k)\rangle \to \langle \hat{b}^{\dagger}(\kappa)\hat{b}(\kappa)\rangle$$

$$= 0; \qquad \qquad = \int dk \left|\beta(\kappa;k)\right|^{2}.$$

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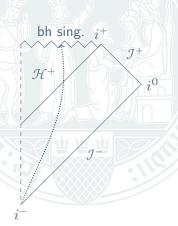
Spherically gravitational collapse

The model and its quantisation

 Massless neutral scalar field on the given background space-time w/ inand out-going eigen-modes

$$\begin{split} f_{\omega'lm} &\propto r^{-1} F_{\omega'} \mathrm{e}^{\mathrm{i}\omega'v} Y_{lm}(\theta,\phi), \\ p_{\omega lm} &\propto r^{-1} P_{\omega} \mathrm{e}^{\mathrm{i}\omega u} Y_{lm}(\theta,\phi) \end{split}$$

- Quantising at Cauchy surfaces
 - 1. Early time: \hat{a} 's defined on \mathcal{I}^-
 - 2. Late time: \hat{b} 's on \mathcal{I}^+ . \hat{c} 's on \mathcal{H}^+
- Physical vacuum: $\hat{a} | H \rangle \equiv 0$
- Concerned with $\langle H \mid \hat{n}_b \mid H \rangle$



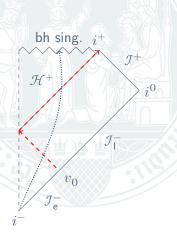
Spherically gravitational collapse

Hawking's construction [Haw74; Haw75]

- $\hbox{ Consider p_ω on \mathcal{I}^+ propagating backwardly to \mathcal{I}^-, ending up with }$
 - The same freq. $p_{\omega}^{(1)}$: only gives a $\delta(\omega \omega')$ term in $\alpha(\omega; \omega')$
 - The rest $p_{\omega}^{(2)}$: contributing to β as well; of interest
- \bullet On \mathcal{I}^- , approximately $p_\omega^{(2)} \propto$

$$\begin{cases} 0, & v > v_0; \\ r^{-1} P_{\omega} \left(\frac{v_0 - v}{CD} \right)^{-\frac{\delta}{\kappa} \frac{\omega}{\kappa_{\mathsf{S}}}}, & v \lesssim v_0 \end{cases}$$

where v_0 is the latest time a null geodesic could leave $\mathcal{I}^-,\ C$ and D const.





Results and Interpretation

Hawking temperature

Expectation value of particle number

$$\left\langle \hat{n}_b(\omega) \right\rangle_H = \int d\omega' \left| \beta(\omega; \omega') \right|^2 \approx \Gamma_\omega \left(e^{\frac{2\pi\omega}{\kappa_S}} - 1 \right)^{-1} \tag{1}$$

• Comparing eq. (1) with Bose–Einstein dist. (black body) $\left\langle \hat{n}(\omega) \right\rangle_{\mathsf{BE}} = \left(\mathrm{e}^{\frac{\omega}{T}} - 1 \right)^{-1} \text{, one may conclude that the physical system concerned is a } \textit{grey} \text{ body, with a temperature of}$

$$T_{\rm H} = rac{\kappa_{
m S}}{2\pi} pprox \left(rac{M_{
m O}}{M}
ight) \cdot 6.169 imes 10^{-8} \, {
m K}.$$



Results and Interpretation

Hawking temperature

Expectation value of particle number

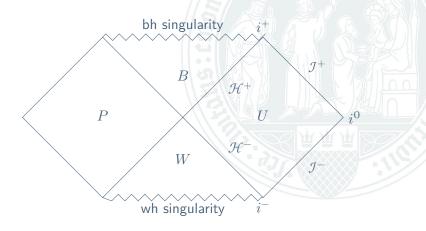
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• $\hat{
ho}_H=|H
angle\,\langle H|$ pure, while the thermal $\hat{
ho}_{
m BE}=rac{1}{Z}{
m e}^{-\widehat{H}/T}$ mixed?

The conformal diagram

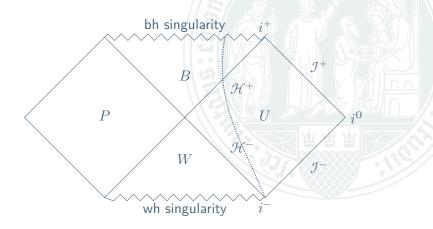


[Can80; FN98]

- Clean calculation; $\left\langle \hat{\phi}^2 \right\rangle_{\rm ren}$, $\left\langle \widehat{T}_{\mu
 u} \right\rangle_{\rm ren}$ etc. obtainable
- Different 'vacua' can be defined
 - ${\color{red}\bullet}$ Boulware: No flux, $\left<\widehat{\phi}^2\right>_{\rm ren}$ blows up at $\mathcal{H}^-\cup\mathcal{H}^+$
 - Miracle near the horizons
 - Israel–Hartle–Hawking: I/O flux, $\left\langle \widehat{\phi}^2 \right\rangle_{\rm ren}$ finite at $\mathcal{H}^- \cup \mathcal{H}^+$
 - Black hole in a heat bath
 - Mathematically more desirable
 - Unruh: O flux, $\left\langle \widehat{\phi}^2 \right\rangle_{\mathrm{ren}}$ blows up at \mathcal{H}



The conformal diagram, with cut



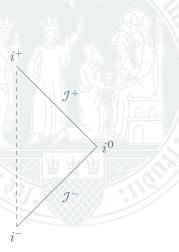
[Can80; FN98]

- Comparing the space-time with that of the collapsing body:
 cut along the dots and stick the right to an interior solution
- Clean calculation; $\left\langle \hat{\phi}^2 \right\rangle_{\rm ren}$, $\left\langle \widehat{T}_{\mu\nu} \right\rangle_{\rm ren}$ etc. obtainable
- Different 'vacua' can be defined
 - Boulware: No flux, $\left\langle \widehat{\phi}^2 \right\rangle_{\mathrm{ren}}$ blows up at $\mathcal{H}^- \cup \mathcal{H}^+$
 - Israel–Hartle–Hawking: I/O flux, $\left\langle \widehat{\phi}^2 \right\rangle_{\mathrm{ren}}$ finite at $\mathcal{H}^- \cup \mathcal{H}^+$
 - \bullet Unruh: O flux, $\left\langle \widehat{\phi}^{2}\right\rangle _{\mathrm{ren}}$ blows up at \mathcal{H}
 - ullet \mathcal{H}^- and the divergence therein can be hidden
 - Similar to the (physical) collapsing case



See e.g. [Mat09; Man15]

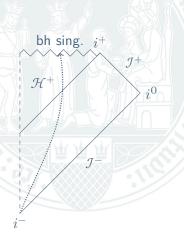
• Minkowski: particles from \mathcal{I}^- to \mathcal{I}^+





See e.g. [Mat09; Man15]

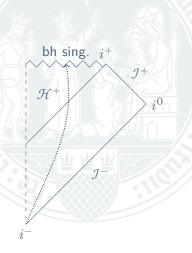
- Minkowski: particles from \mathcal{I}^- to \mathcal{I}^+
- Collapsing body: can also go to \mathcal{H}^+
- Obs. able to set up input at J⁻ and collect output at J⁺, but not at H⁺





See e.g. [Mat09; Man15]

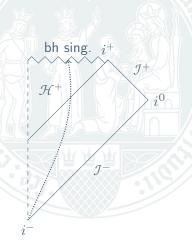
- Minkowski: particles from \mathcal{I}^- to \mathcal{I}^+
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- Obs. able to set up input at \mathcal{I}^- and collect output at \mathcal{I}^+ , but not at \mathcal{H}^+
- $|\alpha\rangle$ on \mathcal{I}^- evolves to $|\beta\rangle$ on $\mathcal{H}^+\cup\mathcal{I}^+$
- Pure $\hat{\rho}_{l} = |\alpha\rangle\langle\alpha|$ turns mixed $\hat{\rho}_{O} = \operatorname{tr}_{\mathcal{H}^{+}} |\beta\rangle\langle\beta|!$ Something is lost





See e.g. [Mat09; Man15]

- Minkowski: particles from \mathcal{I}^- to \mathcal{I}^+
- Collapsing body: can also go to \mathcal{H}^+
- Obs. able to set up input at \$\mathcal{I}^-\$ and collect output at \$\mathcal{I}^+\$, but not at \$\mathcal{H}^+\$
- $|\alpha\rangle$ on \mathcal{I}^- evolves to $|\beta\rangle$ on $\mathcal{H}^+\cup\mathcal{I}^+$
- Pure $\hat{\rho}_{\rm I} = |\alpha\rangle \langle \alpha|$ turns mixed $\hat{\rho}_{\rm O} = {\rm tr}_{\mathcal{H}^+} |\beta\rangle \langle \beta|!$ Something is lost
- Non-conservation arguments [Haw76]
- Conservation arguments [Pag93]



Summary

- Robust calculation of outgoing particle flux for the space-time of collapsing body
- Controversial interpretations and extrapolations
- Back-reaction to the metric: no time
- Seeks combination with
 - Quantum gravitation
 - Quantum information

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