

# Notes on Lagrangian Singular Dynamics

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## 1 Lagrangian formalism

Lagrangian with velocity

$$L^\vee := L|_{\dot{q}=v} \quad (1.1)$$

Equations of motion

$$\sum_j M_{ij} \dot{v}_j = K_i^\vee, \quad \dot{q}_i = v_i. \quad (1.2)$$

where

$$M_{ij}(q, v) := \frac{\partial^2 L}{\partial v_i \partial v_j}. \quad (1.3)$$

## 2 Examples

### 2.1 Toy examples

Example 0

Gitman and Tyutin 1990, sec. 1.2

$$L = \frac{1}{2}(\dot{x} - y)^2 \quad (2.1)$$

Example 1

$$L = \frac{1}{2}\dot{x}^2 + \dot{x}y - \frac{1}{2}(x - y)^2. \quad (2.2)$$

Example 2

$$L = \frac{1}{2}\dot{x}^2 + \dot{x}y + \frac{1}{2}(x - y)^2 \quad (2.3)$$

Example 3

$$L = \frac{1}{2}(\dot{q}_2 - e^{q_1})^2 + \frac{1}{2}(\dot{q}_3 - q_2)^2. \quad (2.4)$$

## 2.2 Parametrised systems

### Non-relativistic point particle

Kiefer 2012, sec. 3.1.1

$$S[q(t)] := \int_{t_1}^{t_2} \mathbb{d}t L\left(q, \frac{\mathbb{d}q}{\mathbb{d}t}\right) \quad (2.5)$$

### Relativistic charged point particle

Landau and Lifshitz 1975, sec. 15, Kiefer 2012, sec. 3.1.2

$$S := \int -m \mathbb{d}s + e A_\mu(x) \mathbb{d}x^\mu =: \int \mathbb{d}\tau L, \quad (2.6)$$

where the Lagrangian reads

$$L = -m \sqrt{-\eta_{\mu\nu} \dot{x}^\mu \dot{x}^\nu} + q \dot{x}^\mu A_\mu(x). \quad (2.7)$$

$$M_{\mu\nu} := \frac{\partial^2 L}{\partial \dot{x}^\mu \partial \dot{x}^\nu} = m \frac{-\eta_{\mu\nu} \eta_{\alpha\beta} + \eta_{\mu\alpha} \eta_{\nu\beta} \dot{x}^\alpha \dot{x}^\beta}{(-\eta_{\rho\sigma} \dot{x}^\rho \dot{x}^\sigma)^{3/2}}, \quad (2.8)$$

which has one and only one eigenvector with null eigenvalue

$$\dot{x}^\mu M_{\mu\nu} = 0. \quad (2.9)$$

#### 2.2.1 Neutral scalar field

ibid., sec. 3.3

## 2.3 Maxwell–Proca theory

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} m^2 A_\mu A^\mu + \frac{1}{2} F_{\mu\nu} S^{\mu\nu} + A_\mu J_f^\mu, \quad (2.10)$$

where  $m > 0$  corresponds to the Proca theory Gitman and Tyutin 1990, sec. 2.3, and  $m = 0$  the Maxwell theory H. J. Rothe and K. D. Rothe 2010, sec. 3.3.3 Gitman and Tyutin 1990, sec. 2.4.

## 2.4 String theories

### Relativistic point particle with auxiliary coordinate

Blumenhagen, Lüst, and Theisen 2013, sec. 2.1

$$L := \frac{1}{2} (e^{-1} \eta_{\mu\nu} \dot{x}^\mu \dot{x}^\nu - m^2 e) \quad (2.11)$$

### Nambu–Gotō action

Generalising the kinetic part of (2.6), one has

$$S_{\text{NG}} := -T \int_{\Sigma} \mathrm{d}A =: -T \int_{\Sigma} \mathrm{d}^2\sigma \mathcal{L}, \quad (2.12)$$

where the Lagrangian density

$$\mathcal{L} = \sqrt{-\Gamma}, \quad \Gamma := \det \Gamma_{\alpha\beta}, \quad \Gamma_{\alpha\beta} := \frac{\partial X^\nu}{\partial \sigma^\alpha} \frac{\partial X_\nu}{\partial \sigma^\beta}. \quad (2.13)$$

Historically Gotō 1971; Nambu 1970; Reference e.g. Blumenhagen, Lüst, and Theisen 2013 Kiefer 2012, sec. 3.2

### Polyakov action

Generalising (2.11)

$$S_{\text{P}}[X^\mu, h_{\alpha\beta}] = -\frac{T}{2} \int_{\Sigma} \mathcal{L}, \quad (2.14)$$

where

$$\mathcal{L} := \sqrt{-h} h^{\alpha\beta} \Gamma_{\alpha\beta}. \quad (2.15)$$

Historically Brink, Di Vecchia, and Howe 1976; Deser and Zumino 1976; Polyakov 1981; Reference Kiefer 2012, sec. 3.2

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