Notes on Canonical Singular Dynamics

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1 Classical formalism

Lagrangian with velocity

$$L^{\mathrm{v}} := L\big|_{\dot{q}=v} \tag{1.1}$$

Equations of motion

$$\sum_j M_{ij}\dot{v}_j = K_i^{\mathrm{v}}, \quad \dot{q}_i = v_i. \tag{1.2}$$

where

$$M_{ij}(q,v) := \frac{\partial^2 L^{\mathrm{v}}}{\partial v_i\,\partial v_j}. \tag{1.3}$$

Adding

$$p_i := \frac{\partial L^{\mathrm{v}}}{\partial v_i}. \tag{1.4}$$

Variation of

$$S[q,p;v] := \int \mathrm{d}t \left[L^{\mathrm{v}} + \sum_i p_i(\dot{q}_i - v_i) \right]. \tag{1.5}$$

gives the *extended Euler–Lagrange equations*

$$\dot{q}_i = v_i, \quad \dot{p}_i = \frac{\partial L^{\mathrm{v}}}{\partial q_i}, \quad p_i = \frac{\partial L^{\mathrm{v}}}{\partial v_i}. \tag{1.6}$$

Hamiltonian with velocity

$$H^{\mathrm{v}}(q,p;v) := \sum_i p_i v_i - L^{\mathrm{v}}. \tag{1.7}$$

Identities

$$\frac{\partial H^{\mathrm{v}}}{\partial q_i} \equiv -\frac{\partial L^{\mathrm{v}}}{\partial q_i}, \quad \frac{\partial H^{\mathrm{v}}}{\partial p_i} \equiv v_i, \quad \frac{\partial H^{\mathrm{v}}}{\partial v_i} \equiv p_i - \frac{\partial L^{\mathrm{v}}}{\partial v_i}. \tag{1.8}$$

Variation of

$$S[q,p;v] := \int \mathrm{d}t \left[\sum_i p_i \dot{q}_i - H^{\mathrm{v}} \right] \tag{1.9}$$

gives the *extended canonical equations*

$$\dot{q}_i = [q_i, H^{\mathrm{v}}]_{\mathrm{P}}, \quad \dot{p}_i = [p_i, H^{\mathrm{v}}]_{\mathrm{P}}, \quad \frac{\partial H^{\mathrm{v}}}{\partial v_i} = 0, \tag{1.10}$$

where the *Poisson bracket* is defined as

$$[f^{\mathrm{v}}, g^{\mathrm{v}}]_{\mathrm{P}} := \sum_i \left(\frac{\partial f^{\mathrm{v}}}{\partial q_i} \frac{\partial g^{\mathrm{v}}}{\partial p_i} - \frac{\partial f^{\mathrm{v}}}{\partial p_i} \frac{\partial g^{\mathrm{v}}}{\partial q_i} \right). \tag{1.11}$$

$v_a = \overline{v}_a(q, p; \{v_\alpha\})$ can be solved, $a = 1, 2, \dots, r_M$; v_α cannot be solved, $\alpha = r_M + 1, \dots, n$, where $r_M = \operatorname{rank} M$.

(need to show $v_a = \overline{v}_a(q, p_a)$)

Primary constraints in the standard form

$$\Phi_\alpha(q, p) := \left. \frac{\partial H^{\mathrm{v}}}{\partial v_\alpha} \right|_{\{v_\alpha = \overline{v}_\alpha\}} \equiv p_\alpha - \overline{p}_\alpha(q, \{p_a\}), \tag{1.12}$$

where

$$\overline{p}_\alpha(q, \{p_a\}) := \left. \frac{\partial L^{\mathrm{v}}}{\partial v_\alpha} \right|_{\{v_a = \overline{v}_a\}}. \tag{1.13}$$

Total Hamiltonian

$$H^{\mathrm{t}} := \left. H^{\mathrm{v}} \right|_{\{v_a = \overline{v}_a\}} \equiv H^{\mathrm{v}}(q, p; \{\overline{v}^a(q, p_a; \{v_\alpha\}), v_\alpha\}). \tag{1.14}$$

Subspace of primary constraints

$$\Gamma_{\mathrm{P}} = \{(q, p) \mid \Phi_\alpha(q, p) = 0, \forall \alpha\} \tag{1.15}$$

Since

$$\frac{\partial H^{\mathrm{t}}}{\partial v_\alpha} = \left. \frac{\partial H^{\mathrm{v}}}{\partial v_\alpha} \right|_{\{v_a = \overline{v}_a\}} = \Phi_\alpha \equiv p_\alpha - \overline{p}_\alpha(q, \{p_a\}), \tag{1.16}$$

H^{t} is linear in v_α. One writes

$$H^{\mathrm{t}}(q, \{p_a\}; \{p_\alpha\}, \{v_\alpha\}) = H(q, \{p_a\}) + \sum_\alpha v_\alpha \Phi_\alpha, \tag{1.17}$$

where H^{c} is the *canonical Hamiltonian* or simply *Hamiltonian.*

Proposition H^{c} is independent of $\{p_\alpha\}$.

Proposition Canonical equations with primary constraints

$$\dot{q}_i = [q_i, H]_{\mathrm{P}} + \sum_\beta v_\beta [q_i, \phi_\beta]_{\mathrm{P}}, \tag{1.18}$$

$$\dot{p}_i = [p_i, H]_{\mathrm{P}} + \sum_\beta v_\beta [p_i, \phi_\beta]_{\mathrm{P}}, \tag{1.19}$$

$$\Phi_\alpha(q, p) = 0, \tag{1.20}$$

where v_β's are undetermined. Note that eq. (1.18) for $i = \alpha$ holds identically: $\dot{q}_\alpha = \dot{q}_\alpha$.

Weak equality: $f_1 \approx f_2$ iff $\left. f_1 \right|_{\Gamma_{\mathrm{P}}} = \left. f_2 \right|_{\Gamma_{\mathrm{P}}}$.

Proposition if f and g are two functions over the phase space Γ, and $f \approx h$, then

$$\frac{\partial}{\partial q_i}\left(f - \sum_\beta \phi_\beta \frac{\partial f}{\partial p_\beta}\right) \approx \frac{\partial}{\partial q_i}\left(h - \sum_\beta \phi_\beta \frac{\partial h}{\partial p_\beta}\right), \tag{1.21}$$

$$\frac{\partial}{\partial p_i}\left(f - \sum_\beta \phi_\beta \frac{\partial f}{\partial p_\beta}\right) \approx \frac{\partial}{\partial p_i}\left(h - \sum_\beta \phi_\beta \frac{\partial h}{\partial p_\beta}\right). \tag{1.22}$$

Corollary $\forall H_1 \approx H$,

$$\dot{q}_i \approx [q_i, H]_{\mathrm{P}}, \qquad \dot{p}_i \approx [p_i, H]_{\mathrm{P}}. \tag{1.23}$$

Primary and second constraints $\phi_\mu^{(1,)}$, $\phi_\omega^{(2,)}$; first and second class constraints $\phi_u^{(,1)}$, $\phi_w^{(,2)}$.

2 Examples

2.1 Toy examples

Example 0

Gitman and Tyutin 1990, sec. 1.2

$$L = \frac{1}{2}(\dot{x} - y)^2 \tag{2.1}$$

Example 1

$$L = \frac{1}{2}\dot{x}^2 + \dot{x}y - \frac{1}{2}(x - y)^2. \tag{2.2}$$

One has

$$L^{\mathrm{v}} = \frac{1}{2}v_x^2 + v_x y - \frac{1}{2}(x - y)^2, \tag{2.3}$$

so that

$$p_x = \frac{\partial L^{\mathrm{v}}}{\partial v_x} = v_x + y, \qquad p_y = 0, \tag{2.4}$$

thus

$$\overline{v}_x = p_x - y. \tag{2.5}$$

So that v_y is the primary inexpressible velocity.

The Hamiltonian with velocity reads

$$H^{\mathrm{v}}(q, p; v) = v_x p_x + v_y p_y - \frac{1}{2}v_x^2 - v_x y + \frac{1}{2}(x - y)^2, \tag{2.6}$$

whilst the total and canonical Hamiltonians are

$$H^{\mathrm{t}}\big(q, p; \overline{v}_x, v_y\big) = H^{\mathrm{c}} + v_y \Phi_1, \tag{2.7}$$

$$H^{\mathrm{c}}(q,p) = \frac{1}{2}(p_x - y)^2 + \frac{1}{2}(x - y)^2, \tag{2.8}$$

where the only primary constraint $\Phi_1 = p_y$.

Persistence condition of Φ_1 leads to

$$0 \approx [\Phi_1, H^{\mathrm{t}}]_{\mathrm{P}} = p_x + x - 2y =: \Phi_2. \tag{2.9}$$

Note that $[\Phi_1, \Phi_2] = p_x - x$ does not vanish on Γ, thus $\Phi_{1,2}$ are second-class constraints, and no more constraint can be generated. To solve for v_y one evaluates

$$0 =: [\Phi_1, H^{\mathrm{t}}]_{\mathrm{P}} = p_x - x - 2v_y, \tag{2.10}$$

so that $v_y := (p_x - x)/2$ solves the constraint.

One also has

$$Q_{\alpha\beta} = \big[\Phi_\alpha, \Phi_\beta\big]_{\mathrm{P}} = \begin{pmatrix} 0 & +2 \\ -2 & 0 \end{pmatrix}, \qquad \big(Q^{-1}\big)_{\alpha\beta} = \begin{pmatrix} 0 & -1/2 \\ +1/2 & 0 \end{pmatrix}, \tag{2.11}$$

so that the Dirac brackets are defined as

$$[f,g]_{\mathrm{D}} := [f,g]_{\mathrm{P}} + \frac{1}{2}\Big([f, p_y]_{\mathrm{P}} [p_x + x - 2y, g]_{\mathrm{P}} - [f, p_x + x - 2y]_{\mathrm{P}} \big[p_y, g\big]_{\mathrm{P}} \Big). \tag{2.12}$$

The fundamental ones different from Possion brackets are

$$[x,y]_{\mathrm{D}} = [y, p_x]_{\mathrm{D}} = \frac{1}{2}, \qquad \big[y, p_y\big]_{\mathrm{D}} = 0. \tag{2.13}$$

Last one different from book?

Example 2

$$L = \frac{1}{2}\dot{x}^2 + \dot{x}y + \frac{1}{2}(x-y)^2 \tag{2.14}$$

Primary constraint

$$p_y = 0; \tag{2.15}$$

total Hamiltonian

$$H^{\mathrm{t}} = \frac{1}{2}p_x^2 - p_x y - \frac{1}{2}x^2 + xy + v_y p_y. \tag{2.16}$$

Example 3

$$L = \frac{1}{2}\big(\dot{q}_2 - \mathrm{e}^{q_1}\big)^2 + \frac{1}{2}\big(\dot{q}_3 - q_2\big)^2. \tag{2.17}$$

2.2 Parametrised systems

Non-relativistic point particle

Kiefer 2012, sec. 3.1.1

$$S[q(t)] := \int_{t_1}^{t_2} \mathrm{d}t\, L\Big(q, \frac{\mathrm{d}q}{\mathrm{d}t}\Big) \tag{2.18}$$

Relativistic charged point particle

Landau and Lifshitz 1975, sec. 16, Kiefer 2012, sec. 3.1.2

$$S := \int -m\, \mathrm{d}s + eA_\mu(x)\, \mathrm{d}x^\mu =: \int \mathrm{d}\tau\, L, \tag{2.19}$$

where the Lagrangian reads

$$L = -m\sqrt{-\eta_{\mu\nu}\dot{x}^\mu \dot{x}^\nu} + q\dot{x}^\mu A_\mu(x). \tag{2.20}$$

$$M_{\mu\nu} := \frac{\partial^2 L^{\mathrm{v}}}{\partial v^\mu\, \partial v^\nu} = m\frac{-\eta_{\mu\nu}\eta_{\alpha\beta} + \eta_{\mu\alpha}\eta_{\nu\beta}}{\big(-\eta_{\rho\sigma}v^\rho v^\sigma\big)^{3/2}} v^\alpha v^\beta, \tag{2.21}$$

which has one and only one eigenvector with null eigenvalue

$$v^\mu M_{\mu\nu} = 0. \tag{2.22}$$

Momenta

$$p_\mu = \frac{\partial L^{\mathrm{v}}}{\partial v^\mu} = \frac{m\eta_{\mu\nu}v^\nu}{\sqrt{-\eta_{\rho\sigma}v^\rho v^\sigma}} + qA_\mu. \tag{2.23}$$

If one chooses v^0 to be the primary inexpressible velocity, then eliminating p_0 in eq. (2.23) yields

$$v^i = \frac{\xi\eta^{ij}(p_j - qA_j)v^0}{\sqrt{m^2 + \eta^{kl}(p_k - qA_k)(p_l - qA_l)}}, \tag{2.24}$$

where $\xi = \operatorname{sgn} v^0$. In the following $\xi = +1$ will be chosen.

Inserting eq. (2.24) into the Hamiltonian with velocity

$$H^{\mathrm{v}} = v^\mu p_\mu - L^{\mathrm{v}} = m\sqrt{-\eta_{\mu\nu}v^\mu v^\nu} + v^\mu\big(p_\mu - qA_\mu(x)\big), \tag{2.25}$$

one obtains the total Hamiltonian

$$H^{\mathrm{t}} = v^0\Big(p_0 - qA_0 + \sqrt{m^2 + \eta^{kl}(p_k - qA_k)(p_l - qA_l)}\Big), \tag{2.26}$$

where only a primary constraint survives, which is obviously a first-class constraint

$$\phi^{(1,1)} = p_0 - qA_0 + \sqrt{m^2 + \eta^{kl}(p_k - qA_k)(p_l - qA_l)}, \tag{2.27}$$

and the canonical Hamiltonian vanishes

$$H^{\mathrm{c}} = 0. \tag{2.28}$$

To compare, note in the non-covariant formalism (Landau and Lifshitz 1975, sec. 8)

$$S = \int \mathbb{d}t\, L, \qquad L = -m\sqrt{1 - \dot{\vec{x}}^2} - q\phi + q\dot{\vec{x}} \cdot \vec{A}, \tag{2.29}$$

the system is regular, and the canonical Hamiltonian reads

$$H^{\mathrm{c}} = \sqrt{m^2 + \left(\vec{p} - q\vec{A}\right)^2} + q\phi, \tag{2.30}$$

which corresponds to setting $\phi^{(1,1)} = 0$, $p_0 \to -H^{\mathrm{c}}$ ($p_\mu = (-E, \vec{p})$), and noting $A_\mu = \left(-\phi, \vec{A}\right)$.

Relativistic point particle with einbein

Blumenhagen, Lüst, and Theisen 2013, sec. 2.1

$$L \coloneqq \frac{1}{2}\left(e^{-1}\eta_{\mu\nu}\dot{x}^\mu\dot{x}^\nu - m^2 e\right) \tag{2.31}$$

$$p_\mu = \frac{\partial L^{\mathrm{v}}}{\partial v^\mu} = e^{-1}\eta_{\mu\nu}v^\nu, \qquad p_e = 0. \tag{2.32}$$

Choosing v^e to be the primary inexpressible velocity, one has

$$v^\mu = e\eta^{\mu\nu}p_\nu. \tag{2.33}$$

Hamiltonian with velocity

$$H^{\mathrm{v}} = v^\mu p_\mu + v^e p_e + \frac{1}{2}\left(-e^{-1}\eta_{\mu\nu}v^\mu v^\nu + m^2 e\right); \tag{2.34}$$

total Hamiltonian

$$H^{\mathrm{t}} = \frac{e}{2}\left(\eta^{\mu\nu}p_\mu p_\nu + m^2\right) + v^e p_e; \tag{2.35}$$

canonical Hamiltonian

$$H^{\mathrm{c}} = \frac{e}{2}\left(\eta^{\mu\nu}p_\mu p_\nu + m^2\right). \tag{2.36}$$

The only primary constraint

$$\Phi^{(1,)} = p_e; \tag{2.37}$$

its time evolution

$$\begin{aligned}\left[\Phi^{(1,)}, H^{\mathrm{t}}\right]_{\mathrm{P}} &= [p_e, e]_{\mathrm{P}}\frac{1}{2}\left(\eta^{\mu\nu}p_\mu p_\nu + m^2\right) \\ &= -\frac{1}{2}\left(\eta^{\mu\nu}p_\mu p_\nu + m^2\right).\end{aligned} \tag{2.38}$$

Choose

$$\Phi^{(2,)} = \eta^{\mu\nu} p_\mu p_\nu + m^2, \tag{2.39}$$

whose Possion bracket with H^{t} vanishes; furthermore,

$$\left[\Phi^{(1,)}, \Phi^{(2,)}\right]_{\mathrm{P}} \equiv 0. \tag{2.40}$$

Thus one ends up with two first-class constraints.

2.2.1 Neutral scalar field

Kiefer 2012, sec. 3.3

2.3 Maxwell–Proca theory

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} m^2 A_\mu A^\mu + A_\mu J^\mu, \tag{2.41}$$

where $m > 0$ corresponds to the Proca theory Gitman and Tyutin 1990, sec. 2.3, and $m = 0$ the Maxwell theory H. J. Rothe and K. D. Rothe 2010, sec. 3.3.3, Gitman and Tyutin 1990, sec. 2.4.

Lagrangian density with velocity

$$\mathcal{L}^{\mathrm{v}} = \frac{1}{2}(V_i - \partial_i A_0)^2 - \frac{1}{4} F_{ij}^2 + \frac{m^2}{2}(A_0^2 - A_i^2) + A_0 J^0 + A_i J^i; \tag{2.42}$$

momenta density

$$B^0 := \frac{\partial \mathcal{L}^{\mathrm{v}}}{\partial V_0} = 0, \qquad B^i := \frac{\partial \mathcal{L}^{\mathrm{v}}}{\partial V_i} = V^i + \partial_i A^0; \tag{2.43}$$

total and canonical Hamiltonians

$$\mathcal{H}^{\mathrm{t}} = \mathcal{H}^{\mathrm{c}} + V_0 \Phi_1, \tag{2.44}$$

$$\mathcal{H}^{\mathrm{c}} = \frac{1}{2}(B^i)^2 + B^i \partial_i A_0 + \frac{1}{4} F_{ij}^2 + \frac{m^2}{2}(-A_0^2 + A_i^2) - A_0 J^0 - A_i J^i, \tag{2.45}$$

$$\tag{2.46}$$

where

$$\Phi_1 = B^0 \tag{2.47}$$

is the only primary constraint.

$$\begin{aligned} \left[\Phi_{1\,1}, \mathcal{H}^{\mathrm{t}}{}_2\right]_{\mathrm{P}} &= B^i{}_1 \partial_{i_2} \left[B^0{}_2, A_{0_2}\right]_{\mathrm{P}} + \frac{m^2}{2} \left[B^0{}_1, A_{0_2}^2\right]_{\mathrm{P}} - \left[B^0{}_1, A_{0_2}\right]_{\mathrm{P}} J^0{}_2 \\ &= \left(-B^i{}_2 \partial_{i_2} - m^2 A_{0_2} + J^0{}_2\right) \delta(\vec{x}_1 - \vec{x}_2). \end{aligned} \tag{2.48}$$

Integration with $\mathbb{d}^3 x_2$ yields the secondary constraint

$$[\Phi_1, H^{\mathrm{t}}]_{\mathrm{P}} = \partial_i B^i - m^2 A_0 + J^0 =: \Phi_2, \tag{2.49}$$

so that

$$[\Phi_1, \Phi_2]_{\mathrm{P}} = m^2 \delta(\vec{x}_1 - \vec{x}_2). \tag{2.50}$$

One may further compute

$$\left[{\Phi_2}_1, {\mathcal{H}^{\mathrm{c}}}_2\right]_{\mathrm{P}} = \left[\partial_{i_1} {B^i}_1, \frac{1}{4} F^2_{jk_2} + \frac{m^2}{2} A^2_{j_2} - A_{j_2} {J^j}_2\right]_{\mathrm{P}}, \tag{2.51}$$

in which

$$\begin{aligned} \left[\partial_{i_1} {B^i}_1, \frac{1}{4} F^2_{jk_2}\right]_{\mathrm{P}} &= \left(\partial_{j_2} A_{k_2} - \partial_{k_2} A_{j_2}\right) \partial_{i_1} \left[{B^i}_1, \left(\partial_{j_2} A_{k_2}\right)\right]_{\mathrm{P}} \\ &= -F_{ij_2} \partial_{i_1} \partial_{j_2} \delta(\vec{x}_1 - \vec{x}_2). \end{aligned} \tag{2.52}$$

The Poisson bracket can be evaluated to be

$$\left[{\Phi_2}_1, {\mathcal{H}^{\mathrm{c}}}_2\right]_{\mathrm{P}} = -\left(F_{ij_2} \partial_{j_2} + m^2 {A^i}_2 + {J^i}_2\right) \partial_{i_1} \delta(\vec{x}_1 - \vec{x}_2). \tag{2.53}$$

Integration with $\mathbb{d}^3 x_2$ yields

$$[\Phi_2, H^{\mathrm{c}}]_{\mathrm{P}} = \partial_i (m^2 A^i + J^i). \tag{2.54}$$

For Proca theory $m > 0$, then the algorithm terminates, and one obtains a second-class system.

For Maxwell theory $m = 0$.

2.4 String theories

Nambu–Gotō action

Generalising the kinetic part of (2.19), one has

$$S_{\mathrm{NG}} := -T \int_\Sigma \mathbb{d}A =: -T \int_\Sigma \mathbb{d}^2 \sigma \mathcal{L}, \tag{2.55}$$

where the Lagrangian density

$$\mathcal{L} = \sqrt{-\Gamma}, \quad \Gamma := \det \Gamma_{\alpha\beta}, \quad \Gamma_{\alpha\beta} := \frac{\partial X^\nu}{\partial \sigma^\alpha} \frac{\partial X_\nu}{\partial \sigma^\alpha}. \tag{2.56}$$

Historically Gotō 1971; Nambu 1970; Reference e.g. Blumenhagen, Lüst, and Theisen 2013 Kiefer 2012, sec. 3.2

Polyakov action

Generalising (2.31)

$$S_{\mathrm{P}}[X^\mu, h_{\alpha\beta}] = -\frac{T}{2}\int_\Sigma \mathcal{L}, \tag{2.57}$$

where

$$\mathcal{L} := \sqrt{-h}h^{\alpha\beta}\Gamma_{\alpha\beta}. \tag{2.58}$$

Historically Brink, Di Vecchia, and Howe 1976; Deser and Zumino 1976; Polyakov 1981; Reference Kiefer 2012, sec. 3.2

2.5 Gravitation theories

Closed Friedmann universe

This part adapts ibid., sec. 8.1.2.

The total action reads

$$S := S_{\mathrm{EG}} + S_\phi, \tag{2.59}$$

where S_{EG} follows (2.77), and

$$S_\phi := \int_{\mathcal{M}} \mathbb{d}^4x\,\sqrt{-g}\left(-\frac{1}{2}g^{\mu\nu}(\nabla_\mu\phi)(\nabla_\nu\phi) - m^2\phi^2\right). \tag{2.60}$$

Adapting

$$\mathbb{d}s^2 = -N^2(t)\,\mathbb{d}t^2 + a^2(t)\,\mathbb{d}\Omega_3^2, \tag{2.61}$$

where

$$\Omega_3^2 = \mathbb{d}\chi^2 + \sin^2\chi\,(\mathbb{d}\theta^2 + \sin^2\theta\,\mathbb{d}\phi^2). \tag{2.62}$$

One has

$$\sqrt{-g} = Na^3\sin^2\chi\,\sin\theta, \qquad \sqrt{h} = a^3\sin^2\chi\,\sin\theta; \tag{2.63}$$

whereas

$$R = \frac{6}{N^2}\left(-\frac{\dot{N}\dot{a}}{Na} + \frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^2\right) + \frac{6}{a^2}, \qquad K = \frac{3\dot{a}}{Na}. \tag{2.64}$$

$$S_{\mathrm{EG}} = \frac{A_3}{16\pi G}\left(\int_{t_1}^{t_2}\mathbb{d}t Na^3(R-2\Lambda) - \left[\frac{6\dot{a}a^2}{N}\right]_{t_1}^{t_2}\right), \tag{2.65}$$

where

$$A_3 = \int \sin^2\chi\,\sin\theta\,\mathbb{d}\chi\,\mathbb{d}\theta\,\mathbb{d}\phi = 2\pi^2. \tag{2.66}$$

The term proportional to $\ddot{a}/a$ in the integrand can be integrated by parts

$$\int_{t_1}^{t_2}\mathbb{d}t\,Na^3\frac{6}{N^2}\frac{\ddot{a}}{a} = 6\left(\left[\frac{\dot{a}a^2}{N}\right]_{t_1}^{t_2} - \int_{t_1}^{t_2}\mathbb{d}t\,\dot{a}\frac{\mathbb{d}}{\mathbb{d}t}\frac{a^2}{N^2}\right), \tag{2.67}$$

in which the first term cancels the Gibbons–Hawking–York term. One has

$$S_{\mathrm{EG}} = \frac{3\pi}{4G}\int_{t_1}^{t_2} \mathrm{d}t \left(-\frac{a}{N}\dot{a}^2 + Na - \frac{\Lambda}{3}Na^3\right). \tag{2.68}$$

The matter part of the action reads

$$S_\phi = \pi^2 \int_{t_1}^{t_2} \mathrm{d}t\, a^3 \left(\frac{1}{N}\dot{\phi}^2 - m^2 N\phi^2\right). \tag{2.69}$$

Lagrangian with velocity

$$L^{\mathrm{v}} = \frac{3\pi}{4G}\left(-\frac{a}{N}{v^a}^2 + Na - \frac{\Lambda}{3}Na^3\right) + \pi^2 a^3\left(\frac{1}{N}{v^\phi}^2 - m^2 N\phi^2\right). \tag{2.70}$$

Canonical momenta

$$p_N := \frac{\partial L^{\mathrm{v}}}{\partial v^N} = 0, \quad p_a := \frac{\partial L^{\mathrm{v}}}{\partial v^a} = -\frac{3\pi}{2G}\frac{a}{N}v^a, \quad p_\phi := \frac{\partial L^{\mathrm{v}}}{\partial v^\phi} = 2\pi^2\frac{a^3}{N}v^\phi. \tag{2.71}$$

Choosing v^N to be the primary inexpressible velocity, one obtains the total and canonical Hamiltonians

$$H^{\mathrm{t}} = H^{\mathrm{c}} + v^N\Phi, \tag{2.72}$$

$$H^{\mathrm{c}} = -NH_\perp, \tag{2.73}$$

where

$$H_\perp := \frac{G}{3\pi}\frac{p_a^2}{a} - \frac{1}{4\pi^2}\frac{p_\phi^2}{a^3} - \frac{3\pi}{4G}\left(\frac{\Lambda}{3}a^2 - 1\right)a - \pi^2 m^2 a^3\phi^2, \tag{2.74}$$

$$\Phi = p_N \tag{2.75}$$

are the *Hamiltonian constraint* and the primary constraint, respectively.

Evaluating the time evolution of Φ yields

$$[\Phi, H^t]_{\mathrm{P}} = H_\perp, \tag{2.76}$$

so that the Hamiltonian constraint is indeed a constraint. There is no further constraint, and $[\Phi, H_\perp]_{\mathrm{P}}$ vanishes identically. Therefore there are two and only two first-class constraints.

2.5.1 Einstein–Hilbert action

$$S_{\mathrm{EG}} = S_{\mathrm{EH}} + S_{\mathrm{GHY}}, \tag{2.77}$$

$$S_{\mathrm{EH}} = \frac{1}{16\pi G}\int_{\mathcal{M}} \mathrm{d}^4x\,\sqrt{-g}(R - 2\Lambda), \tag{2.78}$$

and

$$S_{\mathrm{GHY}} = -\frac{1}{8\pi G}\int_{\partial\mathcal{M}} \mathrm{d}^3x\,\sqrt{h}K, \tag{2.79}$$

which is named after Gibbons and Hawking 1977; York 1972 but actually already mentioned in Einstein 1916. See Dyer and Hinterbichler 2009 for a brief review.

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