

# Wightman function of Klein–Fock–Gordon field

Yi-Fan Wang

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## 1 Covariant differential

$\phi$  a  $\mathbb{C}$ -valued 0-form

$$\mathbb{D}\phi := (\mathfrak{d} - ieA) \wedge \phi, \quad \mathbb{D}\phi^* := (\mathfrak{d} + ieA) \wedge \phi^*, \quad (1.0.1)$$

$A$  the  $\mathfrak{u}(1)$ -valued connection form.

## 2 Covariant codifferential

Define the covariant codifferential of a  $\mathbb{C}$ -valued  $k$ -form  $\zeta$  as follows. Let  $\eta$  be an arbitrary  $\mathbb{C}$ -valued  $(k - 1)$ -form.

$$\int \mathfrak{d}(\eta^* \wedge \star \zeta) \equiv \int \mathfrak{d}\eta^* \wedge \star \zeta - (-)^k \eta^* \wedge \mathfrak{d} \star \zeta =: \int \mathbb{D}\eta^* \wedge \star \zeta - \eta^* \wedge \star \mathbb{D}^\dagger \zeta \quad (2.0.1)$$

$$\begin{aligned} &= \int \mathbb{D}\eta^* \wedge \star \zeta - ieA \wedge \eta^* \wedge \star \zeta - (-)^k \eta^* \wedge \mathfrak{d} \star \zeta \\ &= \int \mathbb{D}\eta^* \wedge \star \zeta + \eta^* \wedge (-)^k ieA \wedge \star \zeta - (-)^k \eta^* \wedge \mathfrak{d} \star \zeta \\ &= \int \mathbb{D}\eta^* \wedge \star \zeta - \eta^* \wedge \star (-)^k \star^{-1} (\mathfrak{d} - ieA) \wedge \star \zeta. \end{aligned} \quad (2.0.2)$$

$$\boxed{\mathbb{D}^\dagger \zeta = (-)^k \star^{-1} (\mathfrak{d} - ieA) \wedge \star \zeta.} \quad (2.0.3)$$

### 3 Maxwell–Klein–Fock–Gordon theory

$$S = \int -\mathbb{D}\phi^* \wedge \star \mathbb{D}\phi - m^2 \phi^* \wedge \star \phi - \frac{1}{2} F \wedge \star F. \quad (3.0.1)$$

$$\delta \mathbb{D}\phi = -ie\delta A\phi + \mathbb{D}\delta\phi. \quad (3.0.2)$$

$$\begin{aligned} \delta(\mathbb{D}\phi^* \wedge \star \mathbb{D}\phi) &= \mathfrak{d}(\delta\phi^* \wedge \star \mathbb{D}\phi + \delta\phi \wedge \star \mathbb{D}\phi^*) \\ &\quad + \delta\phi^* \wedge \star \mathbb{D}^\dagger \mathbb{D}\phi + \delta\phi \wedge \star \mathbb{D}^\dagger \mathbb{D}\phi^* \\ &\quad + \delta A \wedge (ie(\phi^* \star \mathbb{D}\phi - \phi \star \mathbb{D}\phi^*)), \end{aligned} \quad (3.0.3)$$

$$\delta(F \wedge \star F) = 2\mathfrak{d}(\delta A \wedge \star F) - 2\delta A \wedge \mathfrak{d}\star F. \quad (3.0.4)$$

$$\begin{aligned} \delta S &= \int -\mathfrak{d}(\delta\phi^* \wedge \star \mathbb{D}\phi + \delta\phi \wedge \star \mathbb{D}\phi^* + \delta A \wedge \star F) \\ &\quad + \delta\phi^* \wedge \star (\mathbb{D}^\dagger \mathbb{D} - m^2)\phi + \delta\phi \wedge \star (\mathbb{D}^\dagger \mathbb{D} - m^2)\phi^* + \\ &\quad + \delta A \wedge (-\mathfrak{d}\star F + ie(\phi^* \star \mathbb{D}\phi - \phi \star \mathbb{D}\phi^*)). \end{aligned} \quad (3.0.5)$$

#### 3.1 Lorenz gauge

$$\square^2 := (\mathfrak{d} + \mathfrak{d}^\dagger)^2 = \mathfrak{d}\mathfrak{d}^\dagger + \mathfrak{d}^\dagger\mathfrak{d}. \quad (3.1.1)$$

$$\mathfrak{d}\star F = \mathfrak{d} \star \mathfrak{d}A = \star(-)^2 \star^{-1} \mathfrak{d} \star \mathfrak{d}A = \star \mathfrak{d}^\dagger \mathfrak{d}A = \star(\square^2 - \mathfrak{d}\mathfrak{d}^\dagger)A. \quad (3.1.2)$$

One would like to have  $\mathfrak{d}\mathfrak{d}^\dagger A = 0$ , or  $\mathfrak{d}^\dagger A = \text{const.}$  This would be fulfilled if

$$\mathfrak{d}^\dagger A = 0, \quad (3.1.3)$$

which is the Lorenz gauge[1, 2, 4].

### 4 Free Klein–Fock–Gordon equation in $(d+1)$ -dimensions

In the absence of external electromagnetic field and Cartesian coordinates,

$$0 = (\mathbb{D}^\dagger \mathbb{D} - m^2)\phi = (-\partial_t^2 + \partial_x^2 - m^2)\phi. \quad (4.0.1)$$

Linearly independent solutions are

$$\exp(\pm i(-\omega_{\vec{k}} t + \vec{k} \cdot \vec{x})), \quad \omega_{\vec{k}} = \sqrt{\vec{k}^2 + m^2}. \quad (4.0.2)$$

One uses the Noether charge to normalise, which reads

$$(\phi_1, \phi_2) = \int \mathfrak{d}^d x j^0 \quad (4.0.3)$$

gives

$$\begin{aligned}\phi_+(\omega_{\vec{k}}, \vec{k}; t, x) &= \frac{1}{\sqrt{2\omega_{\vec{k}}}} \exp\{+i(-\omega_{\vec{k}}t + \vec{k} \cdot \vec{x})\}, \\ \phi_-(\omega_{\vec{k}}, \vec{k}; t, x) &= \frac{1}{\sqrt{2\omega_{\vec{k}}}} \exp\{-i(-\omega_{\vec{k}}t + \vec{k} \cdot \vec{x})\},\end{aligned}\quad (4.0.4)$$

For  $d \geq 2$ ,

$$\int \frac{\mathbb{d}^d k}{(2\pi)^d} \frac{1}{2\omega_{\vec{k}}} e^{+i(-\omega_{\vec{k}}\Delta t + \vec{k} \cdot \Delta \vec{x})} = \int_0^{+\infty} \frac{\mathbb{d}k}{(2\pi)^d} \frac{k^{d-1} e^{-i\omega_{\vec{k}}\Delta t}}{2\omega_{\vec{k}}} \int e^{ik\Delta x \cos \phi_1} \mathbb{d}\Omega_{d-1}, \quad (4.0.5)$$

where

$$\begin{aligned}& \int \mathbb{d}\Omega_{d-1} \\ &:= \int_0^{2\pi} \mathbb{d}\phi_{d-1} \int_0^\pi \sin^{d-2} \phi_1 \sin^{d-3} \phi_2 \dots \sin^1 \phi_{n-2} \mathbb{d}\phi_1 \mathbb{d}\phi_2 \dots \mathbb{d}\phi_{d-2}.\end{aligned}\quad (4.0.6)$$

Equipped with

$$\int_0^\pi \sin^\nu \phi \mathbb{d}\phi = \pi^{\frac{1}{2}} \frac{\Gamma(\frac{\nu+1}{2})}{\Gamma(\frac{\nu}{2}+1)} \quad \Re \nu > -1, \quad (4.0.7)$$

one has

$$\int \mathbb{d}\Omega_{d-1} = \frac{2\pi^{\frac{d-1}{2}}}{\Gamma(\frac{d-1}{2})} \int_0^\pi \sin^{d-2} \phi_1 \mathbb{d}\phi_1. \quad (4.0.8)$$

The inner integral in eq. (4.0.5) turns to be

$$\begin{aligned}\int e^{ik\Delta x \cos \phi_1} \mathbb{d}\Omega_{d-1} &= \frac{2\pi^{\frac{d-1}{2}}}{\Gamma(\frac{d-1}{2})} \int_0^\pi e^{ik\Delta x \cos \phi_1} \sin^{d-2} \phi_1 \mathbb{d}\phi_1 \\ &= \frac{2\pi^{\frac{d-1}{2}}}{\Gamma(\frac{d-1}{2})} \int_0^\pi \cos(k\Delta x \cos \phi_1) \sin^{d-2} \phi_1 \mathbb{d}\phi_1 \\ &= \frac{2\pi^{\frac{d-1}{2}}}{\Gamma(\frac{d-1}{2})} \frac{\pi^{\frac{1}{2}} \Gamma(\frac{d-1}{2})}{(\frac{k\Delta x}{2})^{\frac{d}{2}-1}} J_{\frac{d}{2}-1}(k\Delta x) \\ &= (2\pi)^{\frac{d}{2}} \frac{J_{\frac{d}{2}-1}(k\Delta x)}{(k\Delta x)^{\frac{d}{2}-1}}\end{aligned}\quad (4.0.9)$$

using [3, eq. (10.9.4)]. Now eq. (4.0.5) turns to be

$$\int_0^{+\infty} \frac{\mathbb{d}k}{(2\pi)^{\frac{d}{2}}} \frac{k^{d-1} e^{-i\omega_{\vec{k}}\Delta t}}{2\omega_{\vec{k}}} \frac{J_{\frac{d}{2}-1}(k\Delta x)}{(k\Delta x)^{\frac{d}{2}-1}}. \quad (4.0.10)$$

For  $\Delta t > \Delta x$ , choose a frame in which  $\Delta x = 0$ . Using [3, eq. (10.2.2)], one has

$$\lim_{x \rightarrow 0} \frac{J_\nu(k\Delta x)}{(k\Delta x)^\nu} = \frac{1}{2^\nu \Gamma(\nu+1)} \quad (4.0.11)$$

## References

## References

- [1] J. van Bladel, “Lorenz or lorentz?”, IEEE Antennas and Propagation Magazine **33**, 69–69 (1991) (cit. on p. 2).
- [2] J. van Bladel, “Lorenz or lorentz? [addendum]”, IEEE Antennas and Propagation Magazine **33**, 56–56 (1991) (cit. on p. 2).
- [3] *NIST DIGITAL LIBRARY OF MATHEMATICAL FUNCTIONS*, <http://dlmf.nist.gov/>, Release 1.0.22 of 2019-03-15, F. W. J. Olver, A. B. Olde Daalhuis, D. W. Lozier, B. I. Schneider, R. F. Boisvert, C. W. Clark, B. R. Miller and B. V. Saunders, eds., <http://dlmf.nist.gov/> (cit. on p. 3).
- [4] L. Lorenz, “XXXVIII. on the identity of the vibrations of light with electrical currents”, The London, Edinburgh, and Dublin Philosophical Magazine and Journal of Science **34**, 287–301 (1867) (cit. on p. 2).