

Advanced Materials Modeling

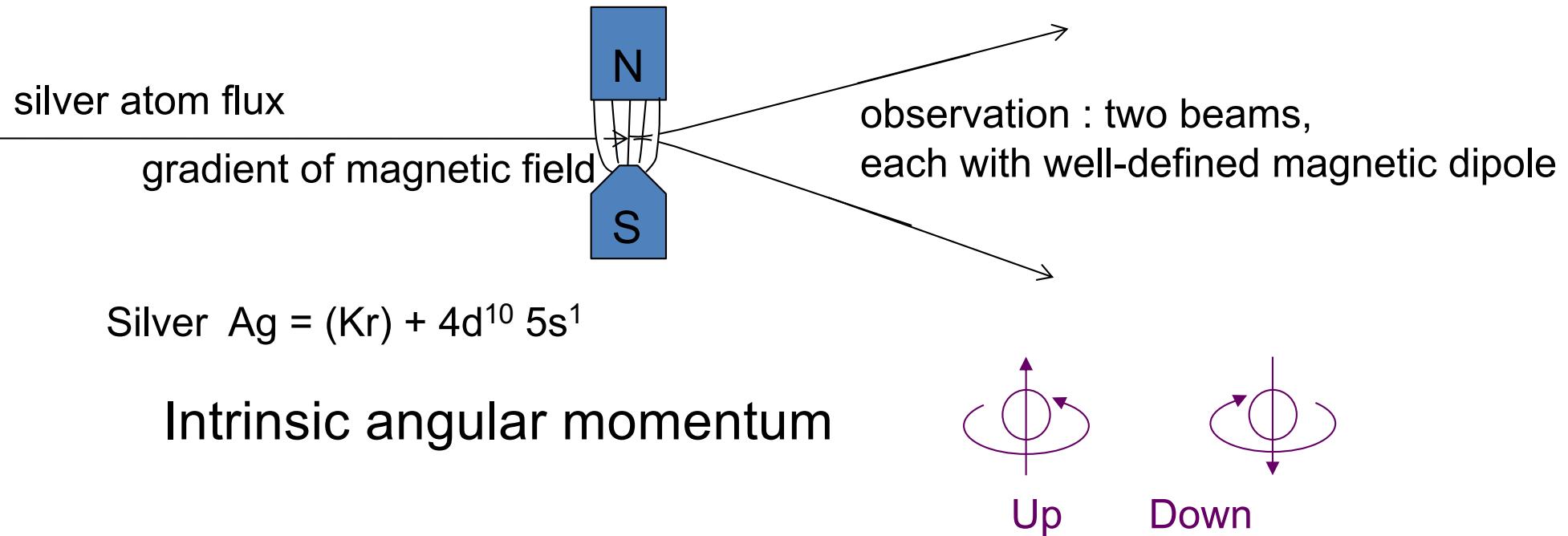
Advanced DFT 3

- + Spin, magnetism, spin-orbit coupling, spin-density functional theory
- + Forces, and brief introduction to density functional perturbation theory
- + Iterative techniques
- + Troubleshooting

Spin, magnetism, spin-orbit coupling, spin-density functional theory

Intrinsic angular momentum

Spin Notion proposed by Uhlenbeck et Goudsmit (1925) to explain Stern and Gerlach experiment



Behaves like an orbital moment, although with only two possible values of the z projection : $m_s \hbar$ $m_s = \pm 1/2$

No intrinsic angular momentum in classical mechanics,
no correspondance principle ...

Reminder : orbital momentum

- Vector property $\vec{\ell} = \vec{r} \times \vec{p} \rightarrow \hat{\vec{\ell}} = -i\hbar(\hat{\vec{r}} \times \hat{\vec{\nabla}})$

$$\hat{\ell}_z = -i\hbar \left(\hat{x} \frac{\partial}{\partial y} - \hat{y} \frac{\partial}{\partial x} \right)$$

Quantum eigenfunctions and eigenvalues

$$\hat{\ell}_z \varphi = \ell_z \varphi \quad -i\hbar \left(x \frac{\partial \varphi}{\partial y} - y \frac{\partial \varphi}{\partial x} \right) = \ell_z \varphi$$

In spherical coordinates

$$\varphi(r, \theta, \Phi) = f(r, \theta) e^{i\ell_z \Phi / \hbar}$$

$$\ell_z = m \cdot \hbar \quad m = 0, \pm 1, \pm 2, \dots \quad \text{Discrete spectrum of } \hat{\ell}_z$$

Representation of spin

Representation based on **vectors** (wavefunction) and **matrices** (operator) : intrinsic angular momentum cannot be represented using spatially-dependent wavefunction

Spinorial wavefunction: 2-component vector, called “spinor”

$$\Psi(\mathbf{r}) = \begin{pmatrix} \psi_{\uparrow}(\mathbf{r}) \\ \psi_{\downarrow}(\mathbf{r}) \end{pmatrix} \text{ or } \psi(\mathbf{r}, \sigma)$$

$\curvearrowleft \quad \sigma \quad \text{can take two “values”: } \uparrow \text{ or } \downarrow$

Other notation : define $|\alpha\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ et $|\beta\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

$$\alpha(\uparrow) = 1 \quad \alpha(\downarrow) = 0 \quad \beta(\uparrow) = 0 \quad \beta(\downarrow) = 1$$

$$\psi(\mathbf{r}, \sigma) = \psi_{\uparrow}(\mathbf{r})\alpha(\sigma) + \psi_{\downarrow}(\mathbf{r})\beta(\sigma)$$

Norm / scalar product

Norm

$$N^2 = \langle \Psi | \Psi \rangle = \int \begin{pmatrix} \psi_{\uparrow}(\mathbf{r}) & \psi_{\downarrow}(\mathbf{r}) \end{pmatrix}^* \begin{pmatrix} \psi_{\uparrow}(\mathbf{r}) \\ \psi_{\downarrow}(\mathbf{r}) \end{pmatrix} d\mathbf{r} = \int |\psi_{\uparrow}(\mathbf{r})|^2 + |\psi_{\downarrow}(\mathbf{r})|^2 d\mathbf{r}$$

Scalar product

$$\begin{aligned} \langle \Phi | \Psi \rangle &= \int \begin{pmatrix} \varphi_{\uparrow}(\mathbf{r}) & \varphi_{\downarrow}(\mathbf{r}) \end{pmatrix}^* \begin{pmatrix} \psi_{\uparrow}(\mathbf{r}) \\ \psi_{\downarrow}(\mathbf{r}) \end{pmatrix} d\mathbf{r} \\ &= \int \varphi_{\uparrow}^*(\mathbf{r}) \psi_{\uparrow}(\mathbf{r}) + \varphi_{\downarrow}^*(\mathbf{r}) \psi_{\downarrow}(\mathbf{r}) d\mathbf{r} \end{aligned}$$

Pure spin-up or -down wavefunctions

Most often : pure spin up wavefunctions, or pure spin down wavefunctions

Explicitly $\Psi(\mathbf{r}) = \begin{pmatrix} \psi_{\uparrow}(\mathbf{r}) \\ 0 \end{pmatrix}$ or $\Psi(\mathbf{r}) = \begin{pmatrix} 0 \\ \psi_{\downarrow}(\mathbf{r}) \end{pmatrix}$

Such pure spin-up or spin-down wavefunctions
are orthogonal to each others.

Correspond to different quantum states.

Simply using spin-up or spin-down wavefunctions
often referred to “scalar wavefunctions”, at variance
with “spinorial wavefunctions” or “spinors”

If not pure-up or -down wavefunctions ?

Computation of electronic density, and also spin-density.

Expectation value of electronic density at a point :

$$n(\mathbf{r}) = |\psi_{\uparrow}(\mathbf{r})|^2 + |\psi_{\downarrow}(\mathbf{r})|^2$$

Expectation value of the spin along z (spin-polarization)

$$s_z(\mathbf{r}) = \left(|\psi_{\uparrow}(\mathbf{r})|^2 - |\psi_{\downarrow}(\mathbf{r})|^2 \right) \frac{\hbar}{2}$$

What about the x - or y - spin-polarization ?
Or expectation along any other direction ?

Pauli matrices : spin along $x, y, z + \dots$

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Warning. Potential notation conflict : σ is for \uparrow or \downarrow , $\sigma_x, \sigma_y, \sigma_z$ are Pauli matrices

Spin-polarization vector (expectation value):

$$s_x(\mathbf{r}) = (\psi_{\downarrow}^*(\mathbf{r})\psi_{\uparrow}(\mathbf{r}) + \psi_{\uparrow}^*(\mathbf{r})\psi_{\downarrow}(\mathbf{r}))\frac{\hbar}{2} = \Psi^+(\mathbf{r})\sigma_x\Psi(\mathbf{r})\frac{\hbar}{2}$$

$$s_y(\mathbf{r}) = (i\psi_{\downarrow}^*(\mathbf{r})\psi_{\uparrow}(\mathbf{r}) - i\psi_{\uparrow}^*(\mathbf{r})\psi_{\downarrow}(\mathbf{r}))\frac{\hbar}{2} = \Psi^+(\mathbf{r})\sigma_y\Psi(\mathbf{r})\frac{\hbar}{2}$$

$$s_z(\mathbf{r}) = \left(|\psi_{\uparrow}(\mathbf{r})|^2 - |\psi_{\downarrow}(\mathbf{r})|^2 \right)\frac{\hbar}{2} = \Psi^+(\mathbf{r})\sigma_z\Psi(\mathbf{r})\frac{\hbar}{2}$$

Spin-density matrix :

$$n_{\alpha\beta}(\mathbf{r}) = \psi_{\alpha}^*(\mathbf{r})\psi_{\beta}(\mathbf{r})$$

$$n(\mathbf{r}) = n_{\uparrow\uparrow}(\mathbf{r}) + n_{\downarrow\downarrow}(\mathbf{r}) \quad s_z(\mathbf{r}) = (n_{\uparrow\uparrow}(\mathbf{r}) - n_{\downarrow\downarrow}(\mathbf{r}))\frac{\hbar}{2}$$

Notation, often :

$$n_{\uparrow}(\mathbf{r}) \triangleq n_{\uparrow\uparrow}(\mathbf{r}) \quad n_{\downarrow}(\mathbf{r}) \triangleq n_{\downarrow\downarrow}(\mathbf{r})$$

Operators acting on spinors

$$\hat{A} = \begin{pmatrix} \hat{A}_{\uparrow\uparrow} & \hat{A}_{\uparrow\downarrow} \\ \hat{A}_{\downarrow\uparrow} & \hat{A}_{\downarrow\downarrow} \end{pmatrix} \quad \text{where } \hat{A}_{\uparrow\uparrow}, \hat{A}_{\uparrow\downarrow}, \hat{A}_{\downarrow\uparrow}, \hat{A}_{\downarrow\downarrow} \text{ are operators acting in usual Hilbert space (=spatial part)}$$

$$\hat{A}|\Psi\rangle = \begin{pmatrix} \hat{A}_{\uparrow\uparrow}|\psi_{\uparrow}\rangle + \hat{A}_{\uparrow\downarrow}|\psi_{\downarrow}\rangle \\ \hat{A}_{\downarrow\uparrow}|\psi_{\uparrow}\rangle + \hat{A}_{\downarrow\downarrow}|\psi_{\downarrow}\rangle \end{pmatrix}$$

Simple example (simple, because spatial part = unit operator) :

$$\hat{S}_z = \begin{pmatrix} \frac{\hbar}{2}\hat{1} & 0 \\ 0 & -\frac{\hbar}{2}\hat{1} \end{pmatrix} = \frac{\hbar}{2}\sigma_z\hat{1} \quad \hat{S}_z|\alpha\rangle = \frac{\hbar}{2}|\alpha\rangle \quad \hat{S}_z|\beta\rangle = -\frac{\hbar}{2}|\beta\rangle$$

Generalization for x and y or other directions using Pauli matrices

Classical treatment of a magnetic field

Now, how does spin relates to magnetic dipole ?

Classical mechanics: for a particle rotating around origin, with mass M and charge q :

$$\text{Angular momentum} \quad \mathbf{L} = M\mathbf{r} \times \mathbf{v}$$

$$\text{Magnetic momentum} \quad \boldsymbol{\mu} = \frac{q}{2}\mathbf{r} \times \mathbf{v}$$

$$\text{yields} \quad \boldsymbol{\mu} = \frac{q}{2M}\mathbf{L}$$

Placed inside a magnetic field, \mathbf{B}
a magnetic dipole acquire an energy $E_{\text{int}} = -\boldsymbol{\mu} \cdot \mathbf{B}$

Quantum treatment of a magnetic field

$$\mu = \frac{q}{2M} \mathbf{L} + \text{Correspondance principle} \Rightarrow \hat{\mu} = \frac{q}{2M} \hat{\mathbf{L}}$$

Suppose a magnetic field along z .

For an electron (charge $q=-e$, mass m_e), observable values for the projection of magnetic moment along z are

$$\mu_z = -\frac{e}{2m_e} L_z = -\frac{e}{2m_e} m\hbar$$

where m = magnetic quantum number (integer)

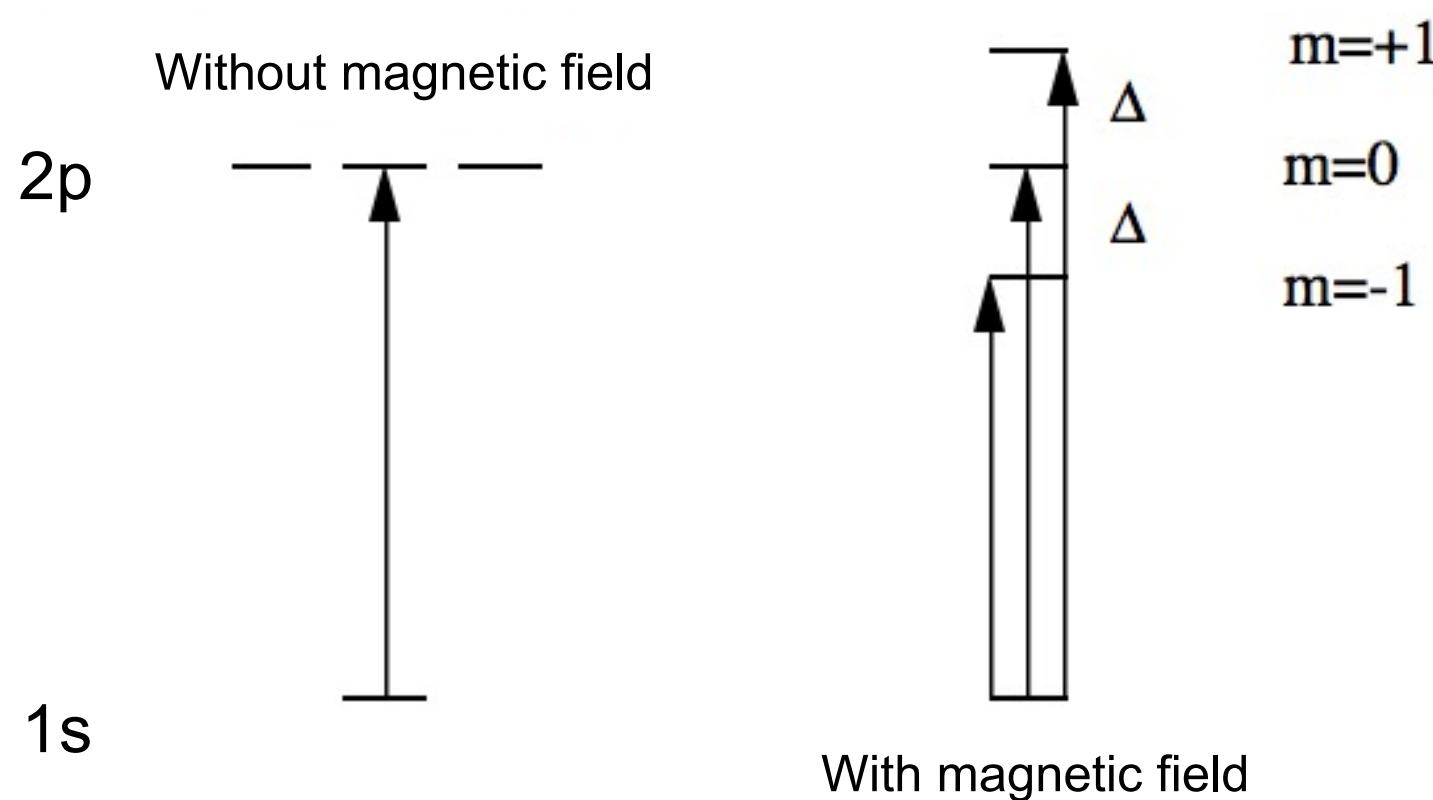
Definition of 'Bohr magneton' $\mu_B = \frac{e\hbar}{2m_e}$

Energy change (first approximation)

$$E_{\text{int}} = -\mu_z \cdot B_z = \frac{e}{2m_e} m\hbar B_z = \frac{e\hbar}{2m_e} m B_z = \mu_B m B_z$$

Zeeman effect : lifting the degeneracy

Atomic levels of same n and ℓ , but different m ,
are not degenerate anymore in a magnetic field



Intrinsic magnetic moment

Classical relation between orbital moment and magnetic moment :

$$\hat{\mu} = \frac{q}{2M} \hat{\mathbf{L}}$$

There exists also a spin magnetic moment, but the proportionality coefficient is not the same :

$$\hat{\mu}_s = g_e \frac{(-e)}{2m_e} \hat{\mathbf{S}}$$

g-factor : $g_e \approx 2$

$$E_{\text{int}} = -\mu_z \cdot B_z = \frac{g_e e}{2m_e} m \hbar B_z = \frac{e \hbar}{2m_e} g_e m B_z = \mu_B g_e m B_z$$

If gradient of B ,
a force appears
=> Stern-Gerlach exp !

Spin-orbit (SOC) coupling (I)

Classically, magnetic dipoles interact with each others.

(magnetic dipole induces a magnetic field, felt by the other dipole)

Magnetic dipole from orbital movement interacts with spin magnetic dipole :

$$E_{SOC} \propto \mathbf{L} \cdot \mathbf{S}$$

Atomic case : degeneracies are lifted.

Example : p levels $m_\ell = -1, 0, +1$ $m_s = -\frac{1}{2}, +\frac{1}{2}$

$$\{|p_{m=-1}, \uparrow\rangle, |p_{m=-1}, \downarrow\rangle, |p_{m=0}, \uparrow\rangle, |p_{m=0}, \downarrow\rangle, |p_{m=+1}, \uparrow\rangle, |p_{m=+1}, \downarrow\rangle\}$$

Due to SOC, 6-fold degeneracy becomes doublet+quadruplet

$$m_j \hbar = m_\ell \hbar + m_s \hbar$$

$$j = \frac{1}{2} \quad \text{with} \quad m_j = -\frac{1}{2}, \frac{1}{2}$$

$$j = \frac{3}{2} \quad \text{with} \quad m_j = -\frac{3}{2}, -\frac{1}{2}, \frac{1}{2}, \frac{3}{2}$$

$p_{1/2} - p_{3/2}$ splitting

Spin-orbit interaction (II)

Theory : originates from relativistic corrections (Schrödinger to Dirac Eq.)

Mass-velocity + Darwin term

= modification of kinetic energy and potential

Spin-orbit interaction

= new type of term $E_{soc} = -\left(\frac{e\hbar}{2m_e c}\right) \sigma \cdot [\nabla V \times \left(\frac{\mathbf{p}}{m_e c}\right)]$ σ = vector of Pauli matrices

Close to a nucleus, $E_{SOC} \propto L.S$ with an energy modif proportional to Z^4

Note : vanishes for s electronic states $\ell = 0$

Often NOT taken into account in first-principles calculations, except for electronic properties. Needs spinorial wavefunctions, increase of CPU time > 4 .

Z=78 ... 83

$6p_{1/2}$ - $6p_{3/2}$ splitting in Bi is about 1.5 eV

es $\ell = 0$					18
13	14	15	16	17	He 4-0026
5 B 10-81	6 C 12-011	7 N 14-0067	8 O 15-9994	9 F 18-9984	10 Ne 20-179
13 Al 26-9815	14 Si 28-0855	15 P 30-9738	16 S 32-06	17 Cl 35-453	18 Ar 39-948
10 Ni 58-69	11 Cu 63-546	12 Zn 65-38	31 Ga 69-72	32 Ge 72-59	33 As 74-9216
28 Pd 106-42	47 Ag 107-868	48 Cd 112-41	49 In 114-82	50 Sn 118-69	51 Sb 121-75
52 Te 126-9045	53 I 131-29	54 Xe 131-29			
78 Pt 195-08	79 Au 196-9665	80 Hg 200-59	81 Tl 204-383	82 Pb 207-2	83 Bi 208-9804
84 Po (209)	85 At (210)	86 Rn (222)			

Incomplete p-shell

Exchange-correlation functionals

In principle, exact functional is still a functional of density only !

$$E_{\text{xc}}^{\text{DFT}}[n] = \int n(\mathbf{r}) \epsilon_{\text{xc}}^{\text{DFT}}(\mathbf{r}; n) d\mathbf{r}$$

However, information on density + spin gives better approximate functionals.
“Spin-density functional theory” (SDFT).

O. Gunnarsson and B.I. Lundqvist, *Phys. Rev. B* **13**, 4274 (1976)

If spin-orbit coupling is ignored, most materials have either no spin-polarization (nowhere in space), or their spin-polarization is aligned with only one direction (chosen as z), everywhere in space.

The latter gives the collinear case :

$$E_{\text{xc}}^{\text{SDFT,coll}}[n_{\uparrow}, n_{\downarrow}] = \int n(\mathbf{r}) \epsilon_{\text{xc}}^{\text{SDFT,coll}}(\mathbf{r}; n_{\uparrow}, n_{\downarrow}) d\mathbf{r}$$

If spin-orbit coupling is taken into account, either there is no spin-polarization (nowhere in space), or spin-polarization is not constrained to be aligned.

The latter gives the non-collinear case :

$$E_{\text{xc}}^{\text{SDFT,non-coll}}[n, \mathbf{s}] = \int n(\mathbf{r}_1) \epsilon_{\text{xc}}^{\text{SDFT,non-coll}}(\mathbf{r}; n, \mathbf{s}) d\mathbf{r}$$

Exchange-correlation potentials

Generalisation of Kohn-Sham equations :

Collinear case : pure spin-up or spin-down wavefunctions $|\psi_i^\alpha\rangle$

$$\left(-\frac{1}{2} \nabla^2 + V_{\text{KS},\alpha}(\mathbf{r}) \right) \psi_{i\alpha}(\mathbf{r}) = \varepsilon_{i\alpha} \psi_{i\alpha}(\mathbf{r}) \quad i \text{ and } \alpha \text{ label the states}$$

$$V_{\text{KS},\alpha}(\mathbf{r}) = V_{\text{ext}}(\mathbf{r}) + \int \frac{n(\mathbf{r}_1)}{|\mathbf{r}_1 - \mathbf{r}|} d\mathbf{r}_1 + \frac{\delta E_{\text{xc}}^{\text{SDFT, coll}}[n_\uparrow, n_\downarrow]}{\delta n_\alpha(\mathbf{r})} \quad \alpha = \uparrow \text{ or } \downarrow$$

Spin-dependent XC potential

Non-collinear case : spinorial wavefunctions $|\Psi_i\rangle$

$$\left(-\frac{1}{2} \nabla^2 + V_{\text{KS}}(\mathbf{r}) + \frac{\hbar}{2} \mathbf{B}_{\text{xc}}(\mathbf{r}) \cdot \boldsymbol{\sigma} \right) \Psi_i(\mathbf{r}) = \varepsilon_i \Psi_i(\mathbf{r}) \quad i \text{ labels the states}$$

$$V_{\text{KS}}(\mathbf{r}) = V_{\text{ext}}(\mathbf{r}) + \int \frac{n(\mathbf{r}_1)}{|\mathbf{r}_1 - \mathbf{r}|} d\mathbf{r}_1 + \frac{\delta E_{\text{xc}}^{\text{SDFT, non-coll}}[n, \mathbf{s}]}{\delta n(\mathbf{r})}$$

$$B_{\text{xc},j}(\mathbf{r}) = \frac{\delta E_{\text{xc}}^{\text{SDFT, non-coll}}[n, \mathbf{s}]}{\delta s_j(\mathbf{r})} \quad j = x, y, z \quad \text{XC magnetic field}$$

Approximate XC functionals

LDA, GGA, hybrids : can easily be generalized.

E.g. Local spin-density approximation for the **collinear case** :

$$E_{\text{xc}}^{\text{LSDA}}[n_{\uparrow}, n_{\downarrow}] = \int n(\mathbf{r}) \epsilon_{\text{xc}}^{\text{LSDA}}(n_{\uparrow}(\mathbf{r}), n_{\downarrow}(\mathbf{r})) d\mathbf{r}$$

$\epsilon_{\text{xc}}^{\text{LSDA}}(n_{\uparrow}(\mathbf{r}), n_{\downarrow}(\mathbf{r}))$ parameterized from homogeneous electron gas

Moreover, exchange energy in **collinear case** has simple general relationship with unpolarized exchange energy :

G.L. Oliver and J.P. Perdew, *Phys. Rev. A* **20**, 397 (1979)

$$E_x^{\text{SDFT,coll}}[n_{\uparrow}, n_{\downarrow}] = \frac{1}{2} (E_x^{\text{DFT}}[2n_{\uparrow}] + E_x^{\text{DFT}}[2n_{\downarrow}])$$

Indeed, comes from separate contribution of spins to exchange energy:

$$E_x = -\frac{1}{2} \sum_{\sigma} \sum_{i,j}^{occ} \int d\mathbf{r} d\mathbf{r}' \frac{\psi_{j\sigma}^*(\mathbf{r}') \psi_{i\sigma}^*(\mathbf{r}') \psi_{j\sigma}(\mathbf{r}) \psi_{i\sigma}(\mathbf{r})}{|\mathbf{r} - \mathbf{r}'|}$$

For **non-collinear case**, neglect directional dependence of local spin-polarization => falls back to collinear case.

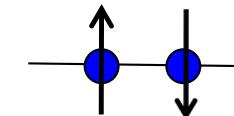
Methodologies : spin+symmetry

In practice ... ?

Materials are treated within one of the following methodologies:

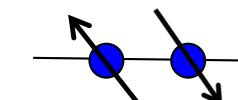
(1) No spin-polarization, scalar wavefunctions

(Many metals/covalent/ionic solids with light nuclei)



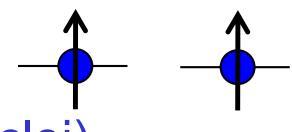
(2) No spin-polarization, spinorial wavefunctions (spin-orbit coupling)

(Many metals/covalent/ionic solids with heavier nuclei ;
the spin-polarization vanishes because of non-broken
time-reversal (TR) symmetry, compensation from TR symmetric states)



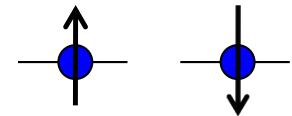
(3) Collinear spin, different spin-up and spin-down wavefunctions

(Ferromagnetic materials with light nuclei, finite systems with net spin,
in particular open shell atoms, and paramagnetic molecules, if light nuclei)



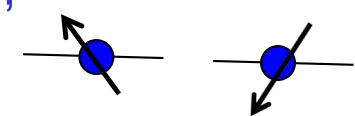
(4) Collinear spin, antiferromagnetic case

(Specific materials, often oxydes – the spin-down density
is space-symmetric to the spin-up density – spin-polarization does not vanish)



(5) Non-collinear spin-polarization

(Heavy nuclei with magnetization, need TR symm. spontaneously broken,
also important for some properties like the magnetocrystalline field,
or spin diffusion)



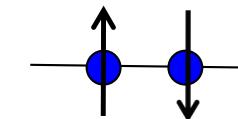
Methodologies : spin+symmetry

In practice ... ?

ABINIT input variables `nspinor`, `nspden`, `nsppol`. Also `spinat(3,natom)`.

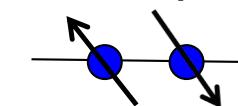
(1) No spin-polarization, scalar wavefunctions

`nspinor=1 nspden=1 nsppol=1`



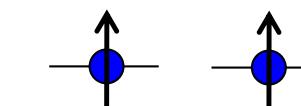
(2) No spin-polarization, spinorial wavefunctions (spin-orbit coupling)

`nspinor=2 nspden=1 nsppol=1`



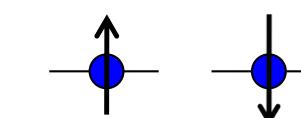
(3) Collinear spin, different spin-up and spin-down wavefunctions

`nspinor=1 nspden=2 nsppol=2`



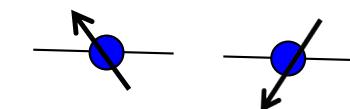
(4) Collinear spin, antiferromagnetic case

`nspinor=1 nspden=2 nsppol=1`

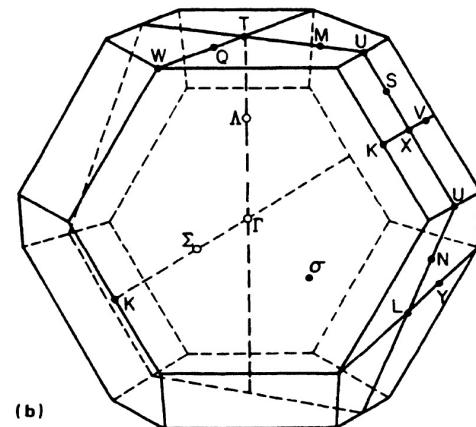
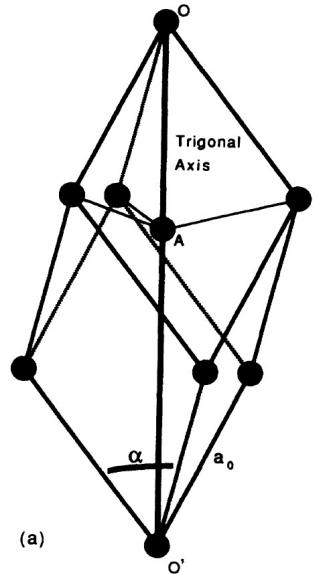
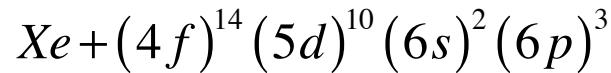


(5) Non-collinear spin-polarization

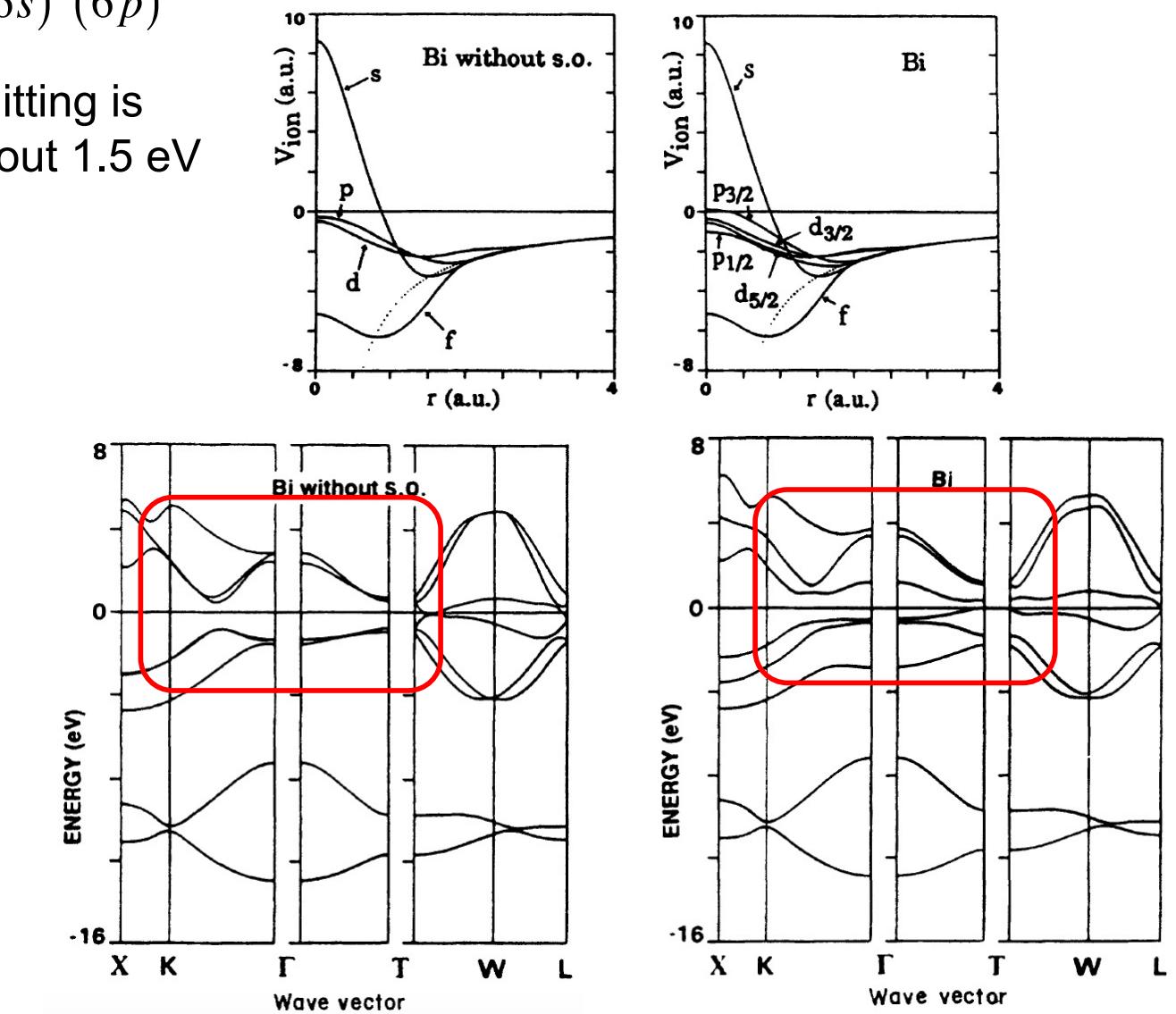
`nspinor=2 nspden=4 nsppol=1`



Example : electronic structure of Bismuth



$6p_{1/2}$ - $6p_{3/2}$ splitting is about 1.5 eV

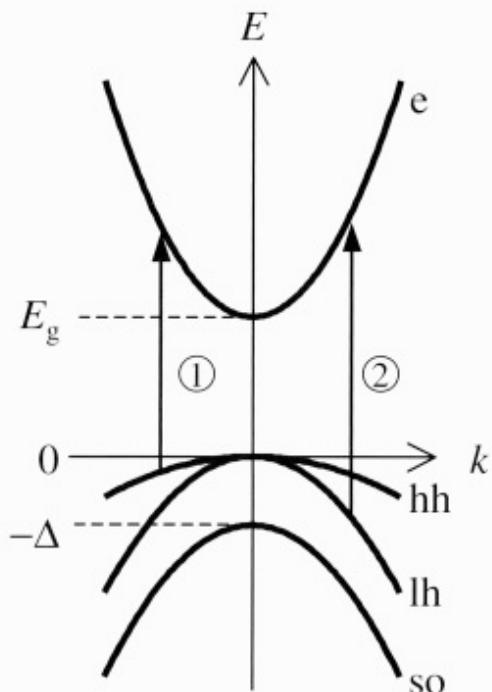


Electronic structure of III-V & II-VI compounds

III-V and II-VI compounds : important for **opto-electronic applications**
(light-emitting diodes, IR and visible photodetectors, solid state lasers)

Crystalline structure : tetrahedrally coordinated
(zinc-blende -AlP, AlAs, AlSb, BN, BP, CdTe, GaAs, HgSe, HgTe, InAs, InP, InSb, ZnS ;
or wurtzite - AlN, GaN, InN, CdS, CdSe, CdTe, ZnO, ZnS)

ZB Band structure in the vicinity of the Γ point (zone center) : Kane model



Kane model

Band structure of several III-V and II-VI compounds near $k=0$.

$E=0$: top of the valence band

$E = E_g$: bottom of the conduction band.

4 bands :

heavy hole (hh) band, light hole (lh) band, spin-orbit hole (so) band, and electron (e) band.

Two optical transitions are indicated.

Transitions can also take place between spin-orbit hole band and conduction band, but are not shown for the sake of clarity.

Δ ranges from <0.1 eV to 0.33 eV for GaAs and 0.8 eV for InSb

Role of XC functional : the case of Iron

Total energy as a function of volume per atom for Iron.

Dashed lines : Spin-LDA

Full lines : GGA

P = paramagnetic (no spin-polarization)

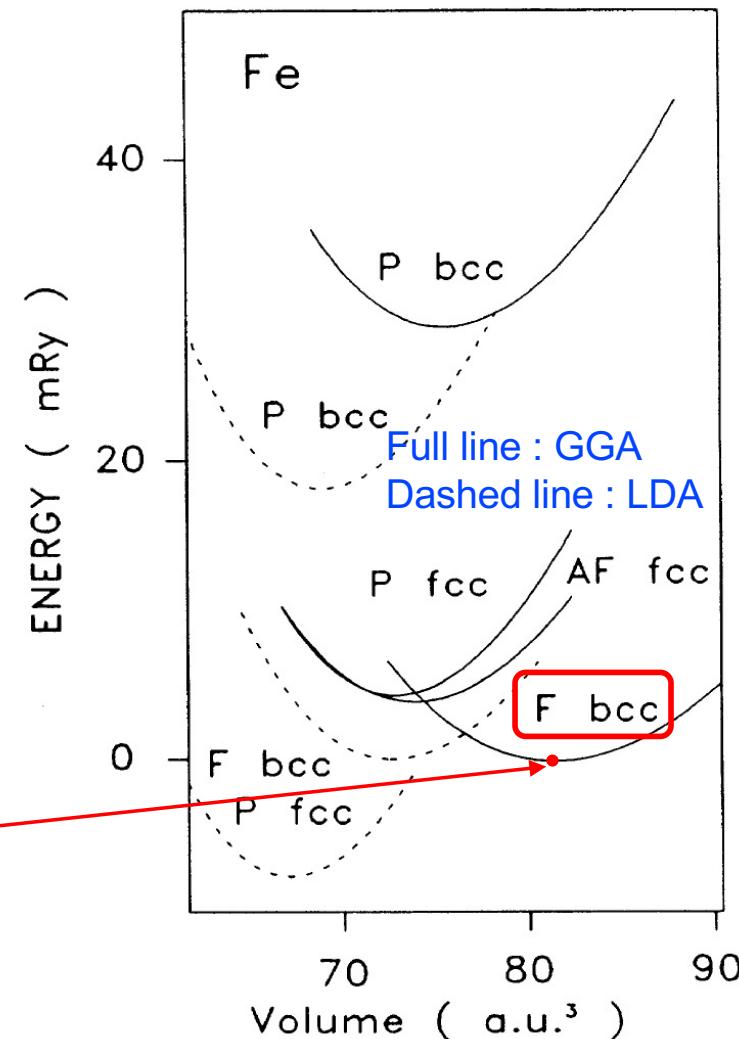
F = ferromagnetic

AF = antiferromagnetic

FCC= face-centered cubic

BCC= body-centered cubic

The GGA gets the **correct phase** :
Experimentally Fe is ferromagnetic BCC.



T.C. Leung, C.T. Chan and B.N. Harmon , *Phys. Rev. B* **44**, 2923 (1991)

Time-reversal operator in the scalar case

Many-body electronic Hamiltonian : invariant upon time-reversal symmetry.

$[\hat{H}_{el}, \hat{T}] = 0$ with \hat{T} time-reversal operator = simply taking the complex conjugate
(Think to time-dependent Schrödinger equation $\hat{H}|\Psi_i(t)\rangle = i\hbar \frac{\partial}{\partial t} |\Psi_i\rangle$)

$$\hat{T}^2 = \hat{1} \quad \hat{T}^{-1} = \hat{T} \quad [\hat{H}_{el}, \hat{T}] = 0 \Rightarrow \hat{T}\hat{H}_{el}\hat{T} = \hat{T}\hat{T}\hat{H}_{el} = \hat{H}_{el}$$

Upon time-reversal: $\hat{\mathbf{r}}$ invariant, $\hat{\mathbf{r}}\hat{T} = \hat{T}\hat{\mathbf{r}}$, while $\hat{\mathbf{p}}$ is reversed $\hat{\mathbf{p}}\hat{T} = -\hat{T}\hat{\mathbf{p}}$

If Kohn-Sham Hamiltonian is invariant upon time-reversal symmetry, and if

$$\hat{H}_{KS}|\Psi_i\rangle = \varepsilon_i|\Psi_i\rangle \quad \text{then} \quad \hat{T}\hat{H}_{KS}\hat{T}|\Psi_i\rangle = \varepsilon_i|\Psi_i\rangle$$
$$\hat{H}_{KS}(\hat{T}|\Psi_i\rangle) = \varepsilon_i(\hat{T}|\Psi_i\rangle)$$

which means $\hat{T}|\Psi_i\rangle$ also eigenfunction, with same eigenenergy than $|\Psi_i\rangle$.

Scalar wavefunctions in periodic boundary conditions and planewave basis

$$\psi_{\mathbf{k}}(\mathbf{r}) = (N\Omega_0)^{-1/2} \sum_{\mathbf{G}} u_{\mathbf{k}}(\mathbf{G}) e^{i(\mathbf{k}+\mathbf{G})\mathbf{r}}$$

Then there exists n' such that

$$\boxed{\begin{aligned} \hat{T}\psi_{n\mathbf{k}}(\mathbf{r}) &= \psi_{n\mathbf{k}}^*(\mathbf{r}) = \psi_{n'-\mathbf{k}}(\mathbf{r}) \\ \varepsilon_{n\mathbf{k}}(\mathbf{r}) &= \varepsilon_{n'-\mathbf{k}}(\mathbf{r}) \end{aligned}}$$

Time-reversal and spin

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

σ_y not invariant upon complex conjugation. Problem with time-reversal symmetry ?

Invariance under complex conjugation AND exchange of spin-up and spin-down.
Time-reversal symmetry exchange spin up and spin down (=angular momentum).

$$\hat{H}_{SOC} = -\left(\frac{e\hbar}{2m_e c}\right) \boldsymbol{\sigma} \cdot \left[\nabla V \times \left(\frac{\hat{\mathbf{p}}}{m_e c} \right) \right]$$

moreover, $\hat{\mathbf{p}}$ changes sign

If Kohn-Sham Hamiltonian including SOC is invariant upon time-reversal symmetry,
and if

$$\hat{H}_{KS} \begin{pmatrix} \psi_{\uparrow n\mathbf{k}}(\mathbf{r}) \\ \psi_{\downarrow n\mathbf{k}}(\mathbf{r}) \end{pmatrix} = \epsilon_{n\mathbf{k}} \begin{pmatrix} \psi_{\uparrow n\mathbf{k}}(\mathbf{r}) \\ \psi_{\downarrow n\mathbf{k}}(\mathbf{r}) \end{pmatrix}$$

Then there exists n' such that $\epsilon_{n\mathbf{k}}(\mathbf{r}) = \epsilon_{n' - \mathbf{k}}(\mathbf{r})$

and
$$\begin{pmatrix} \psi_{\downarrow n\mathbf{k}}^*(\mathbf{r}) \\ -\psi_{\uparrow n\mathbf{k}}^*(\mathbf{r}) \end{pmatrix} = \begin{pmatrix} \psi_{\uparrow n' - \mathbf{k}}(\mathbf{r}) \\ \psi_{\downarrow n' - \mathbf{k}}(\mathbf{r}) \end{pmatrix}$$
 is the corresponding eigenfunction.
= Kramers degeneracy

Time-reversal symmetry forbids spin polarization

Compute expectation value of spin polarization from a pair of degenerate Kramers eigenfunctions :

$$\begin{pmatrix} \psi_{\downarrow n\mathbf{k}}^*(\mathbf{r}) \\ -\psi_{\uparrow n\mathbf{k}}^*(\mathbf{r}) \end{pmatrix} \text{ and } \begin{pmatrix} \psi_{\uparrow n\mathbf{k}}(\mathbf{r}) \\ \psi_{\downarrow n\mathbf{k}}(\mathbf{r}) \end{pmatrix}$$

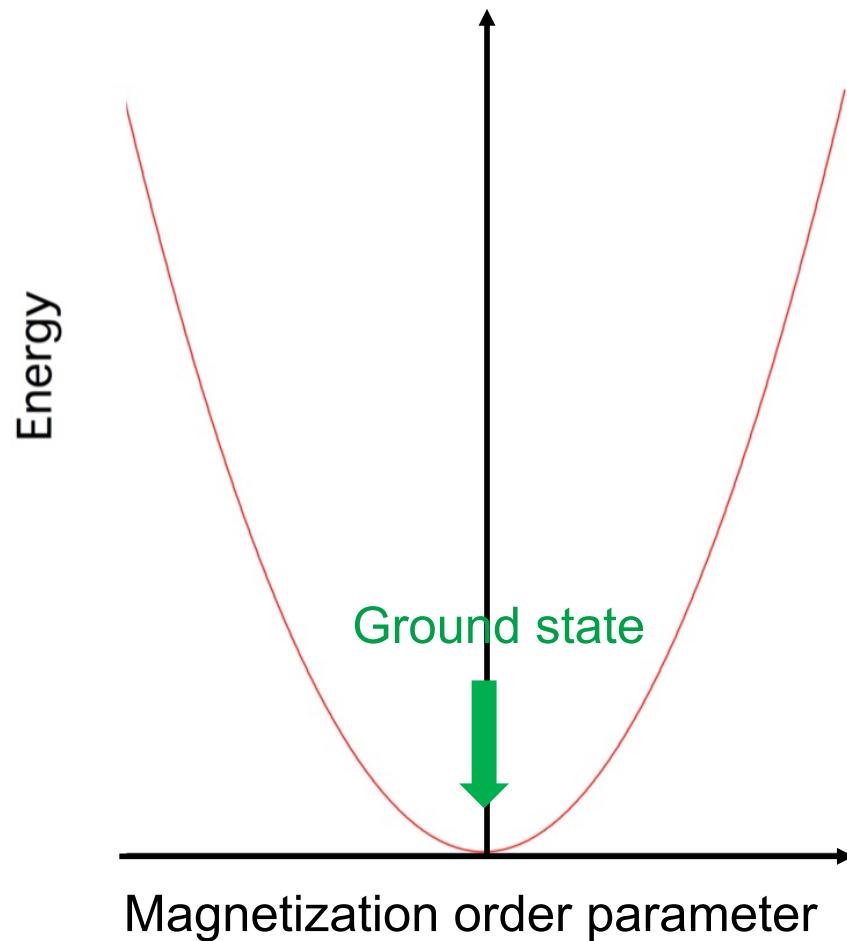
Reminder $s_j(\mathbf{r}) = \Psi^+(\mathbf{r})\sigma_j\Psi(\mathbf{r})\frac{\hbar}{2}$

This gives $s_x(\mathbf{r}) = \left(-\psi_{\uparrow n\mathbf{k}}(\mathbf{r})\psi_{\downarrow n\mathbf{k}}^*(\mathbf{r}) - \psi_{\downarrow n\mathbf{k}}(\mathbf{r})\psi_{\uparrow n\mathbf{k}}^*(\mathbf{r}) \right) \frac{\hbar}{2}$
 $+ \left(\psi_{\uparrow n\mathbf{k}}^*(\mathbf{r})\psi_{\downarrow n\mathbf{k}}(\mathbf{r}) + \psi_{\downarrow n\mathbf{k}}^*(\mathbf{r})\psi_{\uparrow n\mathbf{k}}(\mathbf{r}) \right) \frac{\hbar}{2} = 0$

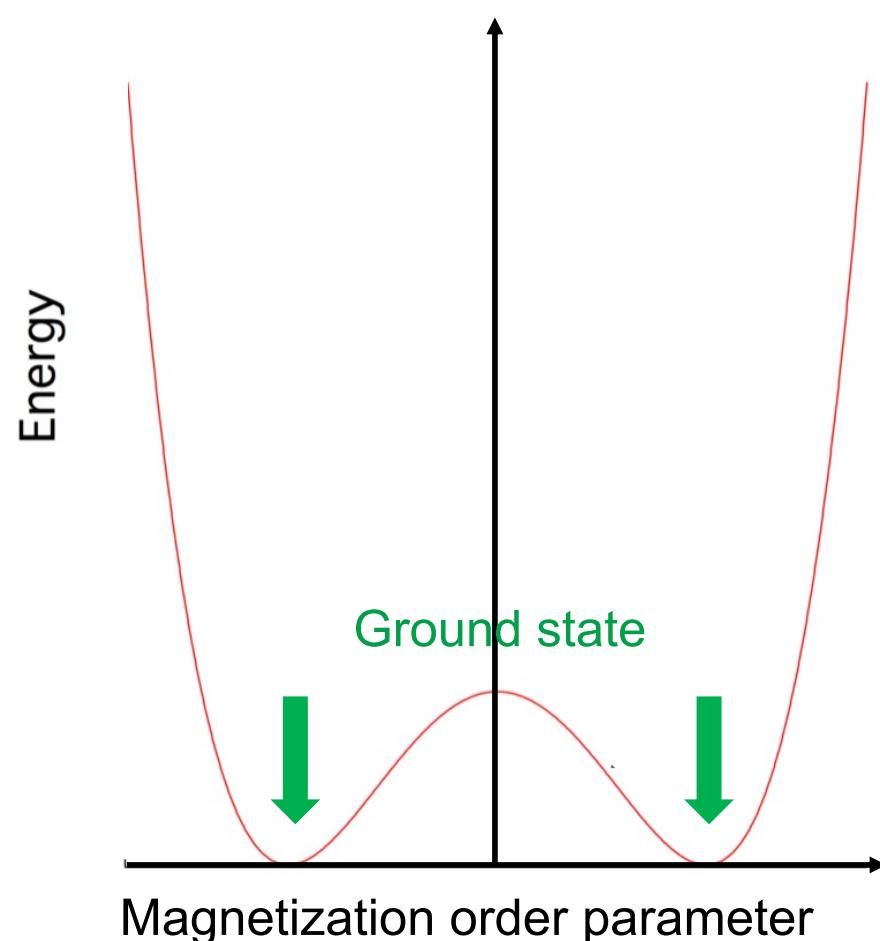
Similarly : $s_y(\mathbf{r}) = 0$ and $s_z(\mathbf{r}) = 0$

So, how can a material be magnetic ?

Spontaneous breaking of time-reversal symmetry



Spin-polarization vanishes,
Time-reversal symmetry not broken
Non-magnetic, possibly spin texture



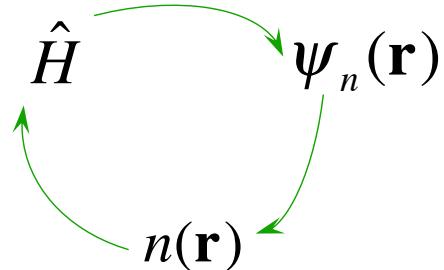
Spin-polarization does not vanish
Time-reversal symmetry broken
Magnetic materials

Computing the forces

Basic equations in DFT

Solve self-consistently the Kohn-Sham equation

$$\left\{ \begin{array}{l} \hat{H} |\Psi_n\rangle = \varepsilon_n |\Psi_n\rangle \\ \hat{H} = \hat{T} + \hat{V} + \hat{V}_{Hxc}[n] \\ n(\mathbf{r}) = \sum_n^{occ} \Psi_n^*(\mathbf{r}) \Psi_n(\mathbf{r}) \end{array} \right.$$



$$\delta_{mn} = \langle \Psi_m | \Psi_n \rangle \text{ for } m, n \in \text{occupied set}$$

or minimize

$$E_{el}\{\Psi\} = \sum_n^{occ} \langle \Psi_n | \hat{T} + \hat{V} | \Psi_n \rangle + E_{Hxc}[n]$$

with

$$\hat{V}(\mathbf{r}) = \sum_{\kappa} -\frac{Z_{\kappa}}{|\mathbf{r} - \mathbf{R}_{\kappa}|}$$

Exchange-correlation functional might be LDA, GGA (e.g. PBE, PBESol), or hybrids, van der waals, etc

Computing the forces (I)

Born - Oppenheimer approx. \Rightarrow find electronic ground state in potential created by nuclei.

A starting configuration of nuclei $\{\mathbf{R}_\kappa\}$ is usually NOT in equilibrium geometry.

$$F_{\kappa,\alpha} = - \frac{\partial E}{\partial R_{\kappa,\alpha}} \Bigg|_{\{\mathbf{R}_\kappa\}} \quad (\text{principle of virtual works})$$

Forces are first derivatives of total energy.
Can be computed by finite differences.

Better approach : compute the response to a perturbation

\Rightarrow What is the energy change ?

$$\{\mathbf{R}_{\kappa,\alpha}\} \rightarrow \{\mathbf{R}_{\kappa,\alpha} + \lambda \delta \mathbf{R}_{\kappa,\alpha}\}$$

Small parameter

Computing the forces (II)

To simplify, let's compute the derivative of an electronic eigenvalue

Perturbation theory : Hellmann - Feynman theorem

$$\hat{H}(\lambda) \quad \varepsilon_n(\lambda) \quad \psi_n(\lambda) \text{ normalized} \quad \hat{H}|\psi_n\rangle = \varepsilon_n|\psi_n\rangle \quad \text{for all } \lambda$$

$$\begin{aligned}\frac{d\varepsilon_n}{d\lambda} &= \frac{d}{d\lambda} \left(\langle \psi_n | \hat{H} | \psi_n \rangle \right) \\ &= \left\langle \frac{d\psi_n}{d\lambda} | \hat{H} | \psi_n \right\rangle + \left\langle \psi_n \left| \frac{d\hat{H}}{d\lambda} \right| \psi_n \right\rangle + \left\langle \psi_n | \hat{H} | \frac{d\psi_n}{d\lambda} \right\rangle \\ &= \left\langle \frac{d\psi_n}{d\lambda} | \varepsilon_n | \psi_n \right\rangle + \left\langle \psi_n \left| \frac{d\hat{H}}{d\lambda} \right| \psi_n \right\rangle + \left\langle \psi_n | \varepsilon_n | \frac{d\psi_n}{d\lambda} \right\rangle \quad = \left\langle \psi_n^{(0)} \left| \frac{d\hat{H}}{d\lambda} \right| \psi_n^{(0)} \right\rangle\end{aligned}$$

Indeed $\left\langle \frac{d\psi_n}{d\lambda} \left| \psi_n \right. \right\rangle + \left\langle \psi_n \left| \frac{d\psi_n}{d\lambda} \right. \right\rangle = \frac{d}{d\lambda} \langle \psi_n | \psi_n \rangle = \frac{d}{d\lambda} (1) = 0$

$\frac{d\psi_n}{d\lambda}$ not needed !

Computing the forces (III)

Application to the derivative of
with respect to an atomic displacement :

$$\hat{H} = \hat{T} + \hat{V}_{ext}\{\hat{R}\} \Rightarrow \frac{\partial \hat{H}}{\partial R_{\kappa,\alpha}} = \frac{\partial \hat{V}_{ext}}{\partial R_{\kappa,\alpha}}$$

$$\frac{\partial \varepsilon_n}{\partial R_{\kappa,\alpha}} = \left\langle \Psi_n \left| \frac{\partial \hat{H}}{\partial R_{\kappa,\alpha}} \right| \Psi_n \right\rangle = \int \frac{\partial \hat{V}_{ext}(\mathbf{r})}{\partial R_{\kappa,\alpha}} n(\mathbf{r}) d\mathbf{r}$$

Computing the forces (IV)

Generalisation to density functional theory

Reminder : $E[\psi] = \sum_n \langle \psi_n | \hat{T} | \psi_n \rangle + \int n(\mathbf{r}) V_{\text{ext}}(\mathbf{r}) d\mathbf{r} + E_{\text{Hxc}}[n]$

If change of atomic positions ...

$$V_{\text{ext}}(\mathbf{r}) = \sum_k -\frac{Z_k}{|\mathbf{r} - \mathbf{R}_k|} \quad (\text{can be generalized to pseudopotential case})$$

$$\frac{\partial V_{\text{ext}}(\mathbf{r})}{\partial R_{k,\alpha}} = +\frac{Z_k}{|\mathbf{r} - \mathbf{R}_k|^2} \cdot \frac{\partial |\mathbf{r} - \mathbf{R}_k|}{\partial R_{k,\alpha}} = -\frac{Z_k}{|\mathbf{r} - \mathbf{R}_k|^3} \cdot (\mathbf{r} - \mathbf{R}_k)_\alpha$$

$$\boxed{\frac{\partial E}{\partial R_{k,\alpha}} = \int n(\mathbf{r}) \frac{\partial V_{\text{ext}}(\mathbf{r})}{\partial R_{k,\alpha}} d\mathbf{r} = - \int \frac{n(\mathbf{r})}{|\mathbf{r} - \mathbf{R}_k|^3} \cdot (\mathbf{r} - \mathbf{R}_k)_\alpha d\mathbf{r}}$$

Forces can be computed directly from the density !

Also, similarly, stresses.

=> Optimization of geometry (lattice parameters and atomic positions)

Density Functional Perturbation Theory

Many physical properties = derivatives of total energy
(or suitable thermodynamic potential) with respect to perturbations.

Consider :

- atomic displacements (phonons)
- dilatation/contraction of primitive cell
- homogeneous external field (electric field, magnetic field ...)

Derivatives of total energy (electronic part + nuclei-nuclei interaction) :

1st order derivatives : forces, stresses, dipole moment ...

2nd order derivatives : dynamical matrix, elastic constants, dielectric susceptibility
atomic polar tensors or Born effective charge tensors
piezoelectricity, internal strains ...

3rd order derivatives : non-linear dielectric susceptibility, Raman susceptibilities
electro-optic effect, phonon - phonon interaction, Grüneisen parameters, ...

Further properties obtained by integration over phononic degrees of freedom :
entropy, thermal expansion, phonon-limited thermal conductivity ...

Iterative algorithms, troubleshooting

Algorithmics : problems to be solved

(1) Kohn - Sham equation

$$\underline{\underline{A}} \underline{x}_i = \lambda_i \underline{x}_i$$

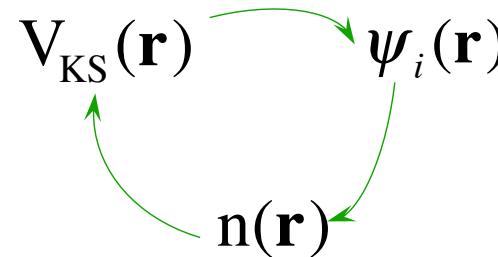
$$\left[-\frac{1}{2} \nabla^2 + V_{KS}(\mathbf{r}) \right] \psi_i(\mathbf{r}) = \epsilon_i \psi_i(\mathbf{r})$$
$$\{\mathbf{G}_j\} \quad \{\mathbf{r}_j\}$$

Size of the system [2 atoms... 600 atoms...] + vacuum ?

Dimension of the vectors \underline{x}_i 300... 100 000... (if planewaves)

of (occupied) eigenvectors 4... 1200...

(2) Self-consistency



(3) Geometry optimization

Find the positions $\{\mathbf{R}_\kappa\}$ of ions such that the forces $\{\mathbf{F}_\kappa\}$ vanish
[= Minimization of energy]

Current practice : iterative approaches

The ‘steepest-descent’ algorithm

Forces are gradients of the energy : moving the atoms along gradients is the steepest descent of the energy surface.

=> Iterative algorithm.

Choose a starting geometry, then a parameter λ , and iterately update the geometry, following the forces :

$$\mathbf{R}_{\kappa,\alpha}^{(n+1)} = \mathbf{R}_{\kappa,\alpha}^{(n)} + \lambda \mathbf{F}_{\kappa,\alpha}^{(n)}$$

Equivalent to the simple mixing algorithm of SCF (see later)

Energy+forces around equilib. geometry

Let us denote the equilibrium geometry as $\underline{\underline{R}}_{\kappa,\alpha}^*$

Analysis of forces close to the equilibrium geometry,
at which forces vanish, thanks to a Taylor expansion :

$$F_{\kappa,\alpha}(R_{\kappa',\alpha'}) = \cancel{F_{\kappa,\alpha}(R_{\kappa',\alpha'}^*)} + \sum_{\kappa',\alpha'} \frac{\partial F_{\kappa,\alpha}}{\partial R_{\kappa',\alpha'}} \Bigg|_{\{R^*\}} (R_{\kappa',\alpha'} - R_{\kappa',\alpha'}^*) + O((R_{\kappa',\alpha'} - R_{\kappa',\alpha'}^*)^2)$$

Moreover, $F_{\kappa,\alpha} = -\frac{\partial E^{BO}}{\partial R_{\kappa,\alpha}}$ $\frac{\partial F_{\kappa',\alpha'}}{\partial R_{\kappa,\alpha}} = -\frac{\partial^2 E^{BO}}{\partial R_{\kappa,\alpha} \partial R_{\kappa',\alpha'}}$

Vector and matrix notation

$$\underline{\underline{R}}_{\kappa,\alpha}^* \rightarrow \underline{\underline{R}}^*$$

$$\underline{\underline{R}}_{\kappa,\alpha} \rightarrow \underline{\underline{R}}$$

$$\underline{\underline{F}}_{\kappa,\alpha} \rightarrow \underline{\underline{F}}$$

$$\frac{\partial^2 E^{BO}}{\partial R_{\kappa,\alpha} \partial R_{\kappa',\alpha'}} \Bigg|_{\{R_{\kappa,\alpha}^*\}} \rightarrow \underline{\underline{H}} \quad (\text{the Hessian})$$

Steepest-descent : analysis (I)

$$\underline{R}_{\kappa,\alpha}^{(n+1)} = \underline{R}_{\kappa,\alpha}^{(n)} + \lambda \underline{F}_{\kappa,\alpha}^{(n)}$$

Analysis of this algorithm, in the linear regime :

$$\underline{F}(\underline{R}) = \underline{F}(\underline{R}^*) - \underline{\underline{H}}(\underline{R} - \underline{R}^*) + O((\underline{R} - \underline{R}^*)^2)$$

$$\underline{R}^{(n+1)} = \underline{R}^{(n)} + \lambda \underline{F}^{(n)} \longrightarrow (\underline{R}^{(n+1)} - \underline{R}^*) = (\underline{R}^{(n)} - \underline{R}^*) - \lambda \underline{\underline{H}}(\underline{R}^{(n)} - \underline{R}^*)$$

$$(\underline{R}^{(n+1)} - \underline{R}^*) = (\underline{\underline{I}} - \lambda \underline{\underline{H}})(\underline{R}^{(n)} - \underline{R}^*)$$

For convergence of the iterative procedure, the "distance" between trial geometry and equilibrium geometry must decrease.

- 1) Can we predict conditions for convergence ?
- 2) Can we make convergence faster ?

Need to understand the action of
the matrix (or operator)

$$\underline{\underline{I}} - \lambda \underline{\underline{H}}$$

Steepest-descent : analysis (II)

What are the eigenvectors and eigenvalues of $\underline{\underline{H}}$?

$\underline{\underline{H}}$ symmetric,
positive definite matrix

$$\left(= \frac{\partial^2 E^{BO}}{\partial R_{\kappa,\alpha} \partial R_{\kappa',\alpha'}} \Bigg|_{\{R_{\kappa,\alpha}^*\}} \right)$$

$\underline{\underline{H}} \underline{f}_i = h_i \underline{f}_i$ where $\{\underline{f}_i\}$ form a complete, orthonormal, basis set

Discrepancy decomposed as $(\underline{R}^{(n)} - \underline{R}^*) = \sum c_i^{(n)} \underline{f}_i$

and $(\underline{R}^{(n+1)} - \underline{R}^*) = (1 - \lambda \underline{\underline{H}}) \sum_i c_i^{(n)} \underline{f}_i = \sum_i c_i^{(n)} (1 - \lambda h_i) \underline{f}_i$

The coefficient of \underline{f}_i is multiplied by $1 - \lambda h_i$

Iteratively : $(\underline{R}^{(n)} - \underline{R}^*) = \sum_i c_i^{(0)} (1 - \lambda h_i)^{(n)} \underline{f}_i$

Steepest-descent : analysis (III)

$$\left(\underline{R}^{(n)} - \underline{R}^* \right) = \sum_i c_i^{(0)} (1 - \lambda h_i)^{(n)} \underline{f}_i$$

The size of the discrepancy decreases if $|1 - \lambda h_i| < 1$

Is it possible to have $|1 - \lambda h_i| < 1$, for all eigenvalues ?

H positive definite \Rightarrow all h_i are positive

Yes ! If λ positive, sufficiently small ...

Steepest-descent : analysis (IV)

$$(\underline{R}^{(n)} - \underline{R}^*) = \sum_i c_i^{(0)} (1 - \lambda h_i)^{(n)} f_i$$

How to determine the optimal value of λ ?

The maximum of all $|1 - \lambda h_i|$ should be as small as possible.

At the optimal value of λ , what will be the convergence rate ?

(= by which factor is reduced the worst component of $(\underline{R}^{(n)} - \underline{R}^*)$?)

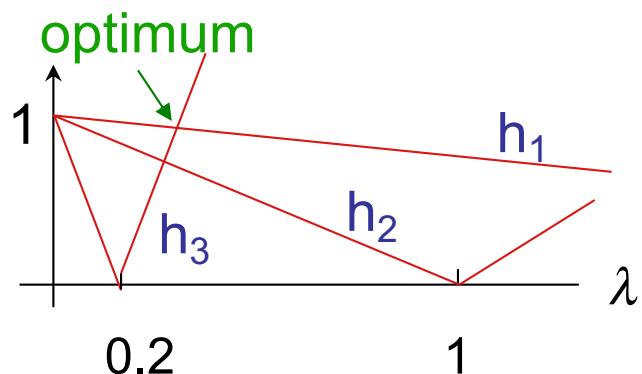
As an exercise : suppose $\begin{cases} h_1 = 0.2 \\ h_2 = 1.0 \\ h_3 = 5.0 \end{cases}$ => what is the best value of λ ?
+ what is the convergence rate ?

Hint : draw the three functions $|1 - \lambda h_i|$ as a function of λ . Then, find the location of λ where the largest of the three curves is the smallest.
Find the coordinates of this point.

Steepest-descent : analysis (V)

Minimise the maximum of $|1 - \lambda h_i|$

$$\begin{array}{ll} h_1 = 0.2 & |1 - \lambda .0.2| \\ h_2 = 1.0 & |1 - \lambda .1| \\ h_3 = 5.0 & |1 - \lambda .5| \end{array}$$



$$\begin{aligned} \text{optimum} &\Rightarrow \lambda = 5 \\ \text{optimum} &\Rightarrow \lambda = 1 \\ \text{optimum} &\Rightarrow \lambda = 0.2 \end{aligned} \quad \left. \right\} ?$$

$$\mu = |1 - \underbrace{\lambda 0.2}_{\text{positive}}| = |1 - \underbrace{\lambda 5}_{\text{negative}}|$$

$$1 - \lambda .0.2 = -(1 - \lambda .5)$$

$$2 - \lambda (0.2+5)=0 \Rightarrow \lambda = 2/5.2$$

$$\mu = 1 - 2 \cdot (0.2 / 5.2)$$

Only $\sim 8\%$ decrease of the error, per iteration ! Hundreds of iterations will be needed to reach a reduction of the error by 1000 or more.

Note : the second eigenvalue does not play any role.

The convergence is limited by the extremal eigenvalues : if the parameter is too large, the smallest eigenvalue will cause divergence, but for that small parameter, the largest eigenvalue lead to slow decrease of the error...

The condition number

In general, $\lambda_{\text{opt}} = 2 / (h_{\min} + h_{\max})$

$$\mu_{\text{opt}} = 2 / [1 + (h_{\max}/h_{\min})] - 1 = [(h_{\max}/h_{\min}) - 1] / [(h_{\max}/h_{\min}) + 1]$$

Perfect if $h_{\max} = h_{\min}$. Bad if $h_{\max} > h_{\min}$.

h_{\max}/h_{\min} called the "condition" number. A problem is "ill-conditioned" if the condition number is large. It does not depend on the intermediate eigenvalues.

Suppose we start from a configuration with forces on the order of 1 Ha/Bohr, and we want to reach the target 1e-4 Ha/Bohr. The mixing parameter is optimal.

How many iterations are needed ?

For a generic decrease factor Δ , with "n" the number of iterations.

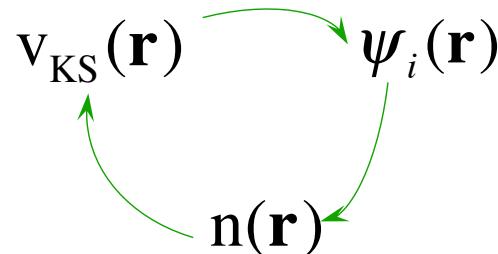
$$\|\underline{F}^{(n)}\| \approx \left(\frac{h_{\max}/h_{\min} - 1}{h_{\max}/h_{\min} + 1} \right)^n \|\underline{F}^{(0)}\|$$

$$\Delta \approx \left(\frac{h_{\max}/h_{\min} - 1}{h_{\max}/h_{\min} + 1} \right)$$

$$n \approx \left[\ln \left(\frac{h_{\max}/h_{\min} + 1}{h_{\max}/h_{\min} - 1} \right) \right]^{-1} \ln \Delta \approx 0.5 \left(h_{\max}/h_{\min} \right) \ln \frac{1}{\Delta}$$

(The latter approximate equality supposes a large condition number)

Analysis of self-consistency



Natural iterative methodology (KS : in => out) :

$$v_{\text{in}}(\mathbf{r}) \rightarrow \psi_i(\mathbf{r}) \rightarrow n(\mathbf{r}) \rightarrow v_{\text{out}}(\mathbf{r})$$

Which quantity plays the role of a force, that should vanish at the solution ?

The difference

$v_{\text{out}}(\mathbf{r}) - v_{\text{in}}(\mathbf{r})$ (generic name : a "residual")

Simple mixing algorithm
(≈ steepest - descent)

$$\underline{v}_{\text{in}}^{(n+1)} = \underline{v}_{\text{in}}^{(n)} + \lambda \left(\underline{v}_{\text{out}}^{(n)} - \underline{v}_{\text{in}}^{(n)} \right)$$

Analysis ... $\underline{v}_{\text{out}} \left[\underline{v}_{\text{in}} \right] = \underline{v}_{\text{out}} \left[\underline{v}^* \right] + \frac{\delta \underline{v}_{\text{out}}}{\delta \underline{v}_{\text{in}}} \left(\underline{v}_{\text{in}} - \underline{v}^* \right)$

$\underline{\underline{H}}$

Like the steepest-descent algorithm, this leads to the requirement to minimize $|1 - \lambda h_i|$ where h_i are eigenvalues of (actually, the dielectric matrix)

$$\frac{\delta \underline{v}_{\text{out}}}{\delta \underline{v}_{\text{in}}}$$

Modify the condition number (II)

$$\underline{R}^{(n+1)} = \underline{R}^{(n)} + \lambda \left(\underline{\underline{H}}^{-1} \right)_{\text{approx}} \underline{F}^{(n)}$$

$$\underline{F}(\underline{R}) = -\underline{\underline{H}} \left(\underline{R} - \underline{R}^* \right) \implies \left(\underline{R}^{(n+1)} - \underline{R}^* \right) = \left(1 - \lambda \left(\underline{\underline{H}}^{-1} \right)_{\text{approx}} \underline{\underline{H}} \right) \left(\underline{R}^{(n)} - \underline{R}^* \right)$$

- Notes :
- 1) If approximate inverse Hessian perfect, optimal geometry is reached in one step, with $\lambda = 1$.
Steepest-descent NOT the best direction.
 - 2) Non-linear effects not taken into account. For geometry optimization, might be quite large. Even with perfect hessian, need 5-6 steps to optimize a water molecule.
 - 3) Approximating inverse hessian by a multiple of the unit matrix is equivalent to changing the λ value.
 - 4) Eigenvalues and eigenvectors of $\left(\underline{\underline{H}}^{-1} \right)_{\text{approx}} \underline{\underline{H}}$ govern the convergence : the condition number can be changed.
 $\left(\underline{\underline{H}}^{-1} \right)_{\text{approx}}$ often called a "pre-conditioner".
 - 5) Generalisation to other optimization problems is trivial.
(The Hessian is referred to as the Jacobian if it is not symmetric.)

Modify the condition number (III)

Approximate Hessian can be generated on a case-by-case basis.

Selfconsistent determination of the Kohn-Sham potential :

Jacobian = dielectric matrix.

Lowest eigenvalue **close to 1**.

Largest eigenvalue :

= **1.5 ... 2.5** for small close-shell molecules, and small unit cell solids

(Simple mixing will sometimes converge with parameter set to 1 !)

= the **macroscopic dielectric constant** (e.g. 12 for silicon),

for larger close-shell molecules and large unit cell insulators,

= **diverge** for large-unit cell metals, or open-shell molecules !

Model dielectric matrices known for rather homogeneous systems.

Knowledge of approx. macroscopic dielectric constant

=> efficient preconditioner

Work in progress for inhomogeneous systems

(e.g. metals/vacuum systems).

Advanced algorithms : using the history

Instead of using only previously computed forces,
take into account past forces for past positions

Large class of methods :

- Broyden (quasi-Newton-type),
- Davidson,
- conjugate gradients,
- Lanczos ...

+ Approximate Hessian can be combined with usage of history

For **geometry optimization** : current practice

Broyden ionmov 22

First without cell relaxation **optcell 0**

then optimizing the cell **optcell 2** (or **1**, or **3...9**)

Iterative algorithms always working ?

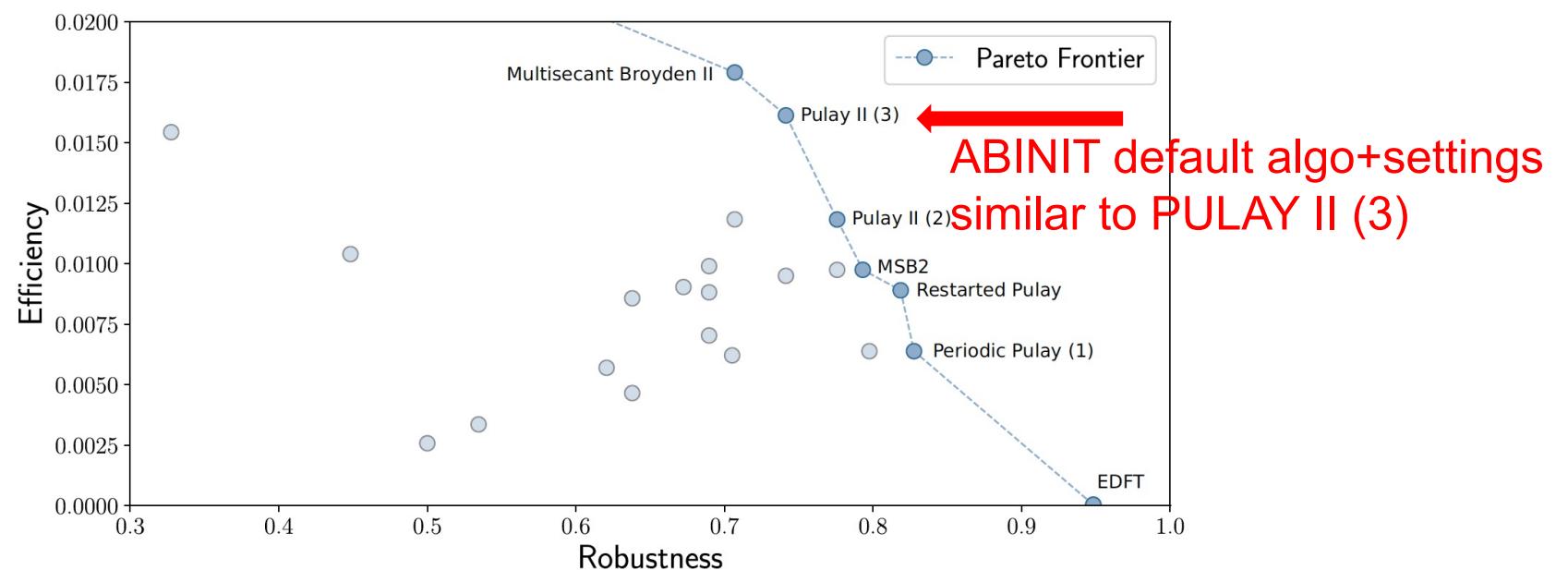
No !

Hypothesis of harmonic energy can be problematic

Many hard cases where iterations go on, without reaching SCF convergence or vanishing forces / stresses.
Trade-off between speed and robustness.

Example : test of 24 iterative algorithms/settings for a set of 56 hard systems

N.D. Woods, M.C. Payne and P.J. Hasnip, *J. Phys.: Condens. Matt.* **31**, 453001 (2019)



Iterative algorithms always working ?

SCF default `iscf` 7 or 17

try changing `iscf`, `npulayit`, `diemac`, `dielng`

is the solution of KS equation accurate ? `nline`, `nnsclo`, `tolrde`

Geometry optimization default `ionmov` 22

try improving underlying SCF accuracy using `tolrff`, `nstep`

try changing `ionmov`

if optimizing primitive lattice vectors, beware about `ecutsm`,

Other troubleshooting ?

Read ABINIT error messages !

Possibly also warnings and comments.

In log file, or in error file, or in the `__ABI_MPIABORTFILE__` file if parallel

Problems with the initial geometry (cell, atom type, nuclei positions)

Many mistakes done by beginners !

Check the units (reminder : ABINIT uses atomic units by default)

Check whether natom is coherent with xred (or xcart)

Check that you have not switched atom types

Relax first the atomic positions **at fixed primitive vectors** before optimizing the cell. Read optcell documentation

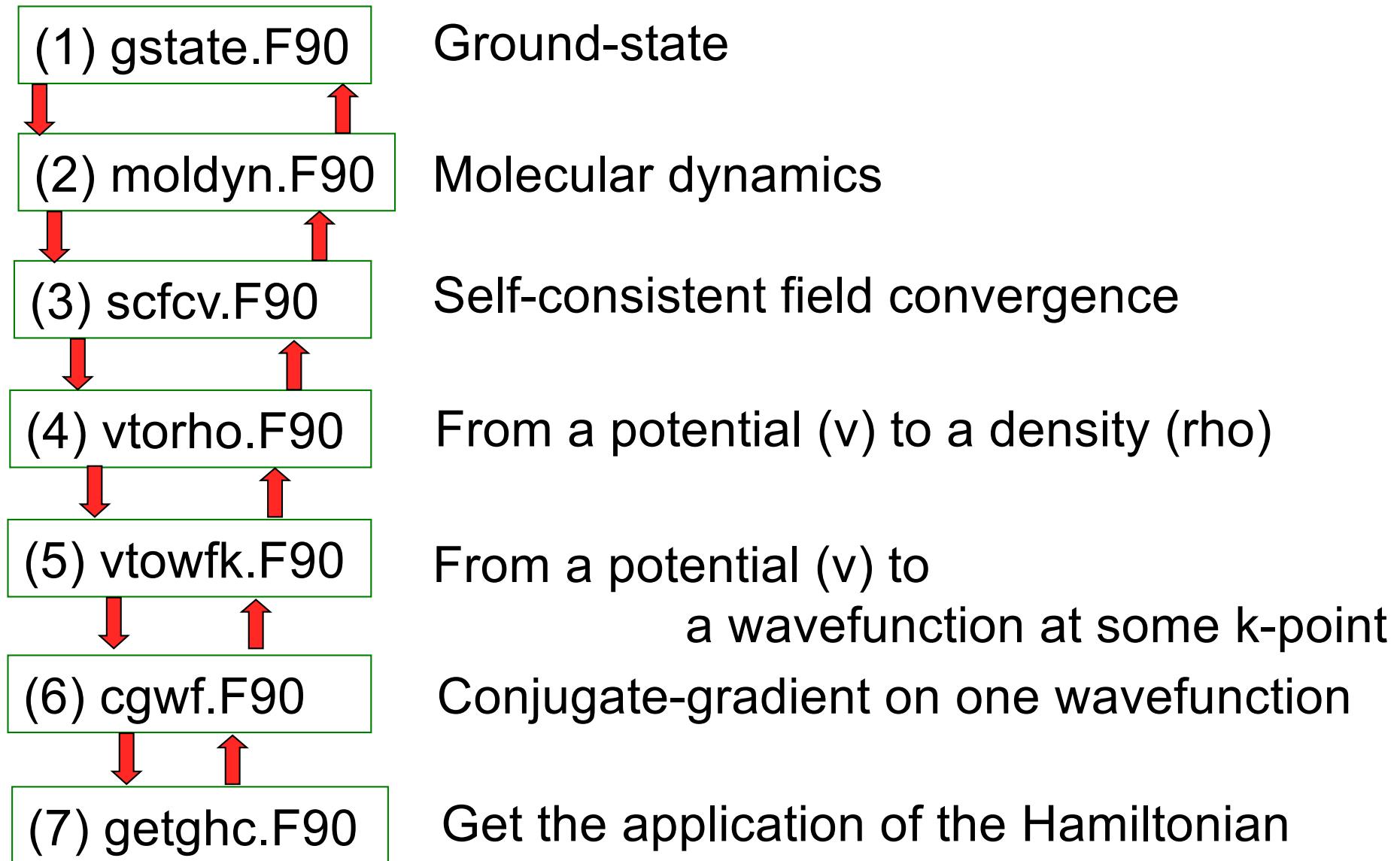
ABINIT can read VASP POSCAR external files containing unit cell parameters and atomic positions. See the input variable structure.

Possibly try to visualize your input geometry using Abipy, VESTA, XCrysDEN, see docs.abinit.org/tutorial/analysis_tools to produce .xsf files for XCrysDEN

More on troubleshooting at

wiki.abinit.org/doku.php?id=howto:troubleshooting

ABINIT : levels in main processing unit



An example run : tbase1_1

- First run in tutorial base1 (<https://docs.abinit.org/tutorial/base1>)
- Very quick (less than 1 sec)
- Look at tbase1_1.abi, O1h.pspgth
- abinit tbase1_1.abi > log
- Look at tbase1_1.abo (such *.abo file will usually be much longer)
- Look at log (such log file is mainly for finding errors or debugging)
- Issue: (-A3 might not be enough for more atoms)

grep -i -B1 -A3 ETOT tbase1_1.abo

grep -i -A3 forces tbase1_1.abo

grep -i -A3 stress tbase1_1.abo

grep -i -A3 energy tbase1_1.abo

grep -i -A3 coordinates tbase1_1.abo

grep -i -A3 density tbase1_1.abo

Additional slides

Exact result for exchange-correlation energy

(without demonstration)

The exchange-correlation energy, functional of the density is the integral over the whole space of the density times the local exchange-correlation energy per particle

$$E_{xc}[n] = \int n(\mathbf{r}_1) \varepsilon_{xc}(\mathbf{r}_1; n) d\mathbf{r}_1$$

while the local exchange-correlation energy per particle is the electrostatic interaction energy of a particle with its DFT exchange-correlation hole.

$$\varepsilon_{xc}(\mathbf{r}_1; n) = \int \frac{1}{2} \frac{\overline{n^{xc}}(\mathbf{r}_2 | \mathbf{r}_1; n)}{|\mathbf{r}_1 - \mathbf{r}_2|} d\mathbf{r}_2$$

$$\text{Sum rule : } \int \overline{n^{xc}}(\mathbf{r}_2 | \mathbf{r}_1; n) d\mathbf{r}_2 = -1$$

Local density approximation (I)

Hypothesis :

- the local XC energy per particle only depend on the local density
- and is equal to the local XC energy per particle of an homogeneous electron gas of same density
(in a neutralizing background - « jellium »)

$$\varepsilon_{xc}^{\text{LDA}}(\mathbf{r}_1; n) = \varepsilon_{xc}^{\text{hom}}(n(\mathbf{r}_1))$$

Gives excellent numerical results ! Why ?

- 1) Sum rule is fulfilled
- 2) Characteristic screening length indeed depend on local density

Local density approximation (II)

Actual function : exchange part + correlation part

$$\varepsilon_x^{\text{hom}}(n) = Cn^{1/3} \quad \text{with} \quad C = -\frac{3}{4\pi}(3\pi^2)^{1/3}$$

for the correlation part, one resorts to accurate numerical simulations beyond DFT (e.g. Quantum Monte Carlo)

Corresponding exchange-correlation potential $V_{\text{xc}}(\mathbf{r}) = \frac{\delta E_{\text{xc}}[n]}{\delta n(\mathbf{r})}$

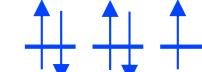
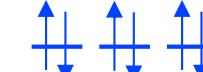
$$V_{\text{xc}}^{\text{approx}}(\mathbf{r}) = \mu_{\text{xc}}(n(\mathbf{r}))$$

$$\mu_{\text{xc}}(n) = \frac{d(n\varepsilon_{\text{xc}}^{\text{approx}}(n))}{dn}$$

$$\mu_x(n) = C \frac{4}{3} n^{1/3} = \frac{4}{3} \varepsilon_x^{\text{hom}}(n)$$

Hund's rule

Example: p levels

						
$ S _{\text{total}}$	1/2	1	3/2	1	1/2	0

Suppose frozen spherically symmetric effective potential created by nuclei and core electrons



Indeed, density identical \Rightarrow Hartree energy identical
However, exchange energy more negative

$$E_H[n] = \frac{1}{2} \int \frac{n(\mathbf{r}_1)n(\mathbf{r}_2)}{|\mathbf{r}_1 - \mathbf{r}_2|} d\mathbf{r}_1 d\mathbf{r}_2$$

$$E_x = -\frac{1}{2} \sum_{\sigma} \sum_{i,j}^{\text{occ}} \int d\mathbf{r} d\mathbf{r}' \frac{\psi_{j\sigma}^*(\mathbf{r}')\psi_{i\sigma}^*(\mathbf{r}')\psi_{j\sigma}(\mathbf{r})\psi_{i\sigma}(\mathbf{r})}{|\mathbf{r} - \mathbf{r}'|}$$



Indeed, density in the same orbital (same spatial location)
increase Hartree energy (also E_x energy less negative !)

ABINIT + python : Abipy, Abitutorials ...

ABINIT organization on GitHub <https://github.com/abinit>

Abipy : python library for launching ABINIT jobs,
and analysing/plotting the results <http://pythonhosted.org/abipy>

=> e.g. connecting ABINIT with tools for high-throughput
calculations developed in the **Materials Project** context
(like Pymatgen, Fireworks).

Abitutorials : tutorial based on Jupyter notebooks ABINIT+python

