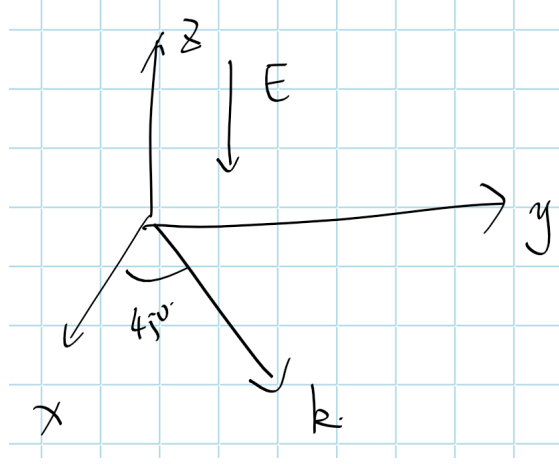


横向应用



电场施加方向与光束传播方向垂直，如上图所示

$$k_x = \sin \theta \quad k_y = \cos \theta \quad (11)$$

以 $\bar{4}3m$ 晶类为例，有介电张量和二阶极化率张量为

$$\begin{pmatrix} \varepsilon_{xx} & 0 & 0 \\ 0 & \varepsilon_{xx} & 0 \\ 0 & 0 & \varepsilon_{xx} \end{pmatrix} \quad \begin{bmatrix} 0 & 0 & 0 & xyz & xyz & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & xyz & xyz & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & xyz & xyz \end{bmatrix} \quad (12)$$

根据电光效应下相对介电常数张量变换， $(\varepsilon_{\mu\alpha})_{eff} = \left[\varepsilon_{r\mu\alpha} + 2\vec{\chi}_{\mu\alpha\beta}^{(2)}(\omega, 0) \cdot E_{0\beta} \right]$,

有施加电场后，有相对介电常数变化，

$$\begin{aligned} (\varepsilon_{\mu\alpha})_{eff} &= \left[\varepsilon_{r\mu\alpha} + 2\vec{\chi}_{\mu\alpha\beta}^{(2)}(\omega, 0) \cdot E_{0\beta} \right] \\ &= \begin{pmatrix} \varepsilon_{xx} + 2\chi_{xxz}^{(2)}E_{0z} & 0 + 2\chi_{xyz}^{(2)}E_{0z} & 0 + 2\chi_{xzz}^{(2)}E_{0z} \\ 0 + 2\chi_{yxz}^{(2)}E_{0z} & \varepsilon_{xx} + 2\chi_{yyz}^{(2)}E_{0z} & 0 + 2\chi_{yzz}^{(2)}E_{0z} \\ 0 + 2\chi_{zxz}^{(2)}E_{0z} & 0 + 2\chi_{zyz}^{(2)}E_{0z} & \varepsilon_{xx} + 2\chi_{zzz}^{(2)}E_{0z} \end{pmatrix} \\ &= \begin{pmatrix} \varepsilon_{xx} & 2\chi_{xyz}^{(2)}E_{0z} & 0 \\ 2\chi_{yxz}^{(2)}E_{0z} & \varepsilon_{xx} & 0 \\ 0 & 0 & \varepsilon_{xx} \end{pmatrix} \end{aligned} \quad (13)$$

代入晶体光学基本方程 $D = \frac{n^2}{\mu_0 c^2} [\mathbf{E} - \mathbf{k}(\mathbf{k} \cdot \mathbf{E})] = \varepsilon_{eff} \cdot \mathbf{E}$ ，有

$$\frac{n^2}{\mu_0 c^2} \begin{pmatrix} E_x(\omega) \\ E_y(\omega) \\ E_z(\omega) \end{pmatrix} - \frac{n^2}{\mu_0 c^2} \hat{k} \cdot (\hat{k} \cdot \vec{E}) = \varepsilon_0 \begin{pmatrix} \varepsilon_{xx} & 2\chi_{xxy}^{(2)}E_{0z} & 0 \\ 2\chi_{zyx}^{(2)}E_{0z} & \varepsilon_{xx} & 0 \\ 0 & 0 & \varepsilon_{xx} \end{pmatrix} \begin{pmatrix} E_x(\omega) \\ E_y(\omega) \\ E_z(\omega) \end{pmatrix} \quad (14)$$

$$\begin{pmatrix} \varepsilon_{xx} & 2\chi_{xyz}^{(2)}E_{0z} & 0 \\ 2\chi_{xyx}^{(2)}E_{0z} & \varepsilon_{xx} & 0 \\ 0 & 0 & \varepsilon_{xx} \end{pmatrix} \begin{pmatrix} E_x(\omega) \\ E_y(\omega) \\ E_z(\omega) \end{pmatrix} + n^2 \cdot \begin{pmatrix} k_x^2 & k_x k_y & 0 \\ k_y k_x & k_y^2 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix} = \begin{pmatrix} n^2 & 0 & 0 \\ 0 & n^2 & 0 \\ 0 & 0 & n^2 \end{pmatrix} \begin{pmatrix} E_x(\omega) \\ E_y(\omega) \\ E_z(\omega) \end{pmatrix} \quad (15)$$

$$\begin{pmatrix} \varepsilon_{xx} + k_x^2 n^2 & 2\chi_{xyz}^{(2)}E_{0z} + k_x k_y n^2 & 0 \\ 2\chi_{xyx}^{(2)}E_{0z} + k_y k_x n^2 & \varepsilon_{xx} + k_y^2 n^2 & 0 \\ 0 & 0 & \varepsilon_{xx} \end{pmatrix} \begin{pmatrix} E_x(\omega) \\ E_y(\omega) \\ E_z(\omega) \end{pmatrix} = \begin{pmatrix} n^2 & 0 & 0 \\ 0 & n^2 & 0 \\ 0 & 0 & n^2 \end{pmatrix} \begin{pmatrix} E_x(\omega) \\ E_y(\omega) \\ E_z(\omega) \end{pmatrix} \quad (16)$$

求久期方程，有

$$\begin{vmatrix} \varepsilon_{xx} + k_x^2 n^2 - n^2 & 2\chi_{xyz}^{(2)}E_{0z} + k_x k_y n^2 & 0 \\ 2\chi_{xyx}^{(2)}E_{0z} + k_y k_x n^2 & \varepsilon_{xx} + k_y^2 n^2 - n^2 & 0 \\ 0 & 0 & \varepsilon_{xx} - n^2 \end{vmatrix} = 0 \quad (17)$$

可以看出有一个解为

$$n_1^2 = \varepsilon_{xx} \quad (18)$$

代入 $k_x = \sin \theta$ ， $k_y = \cos \theta$ ，取久期方程前两阶有

$$\begin{aligned} (\varepsilon_{xx} + n^2 \sin^2 \theta - n^2) (\varepsilon_{xx} + n^2 \cos^2 \theta - n^2) &= \left(2\chi_{xyz}^{(2)}E_{0z} + n^2 \sin \theta \cos \theta \right)^2 \\ n_2^2 &= \frac{\sec^2(\theta) \sqrt{(\varepsilon (\text{Cod}(\theta)^2 + \sin^2(\theta) - 2) - 2a\chi \sin(2\theta))^2 - 2 \cos^2(\theta) (\varepsilon^2 - 4a^2 \chi^2) (-2 \text{Cod}(\theta) - 2 \cos^2(\theta))}}{2 (\text{Cod}(\theta)^2 - \cos^2(\theta))} \end{aligned}$$

代回基本方程，有本征方向

