2.2

用电像法讨论在内半径为a的中空接地导电球内的点电荷a的问题,试求:

- (a) 球内的势
- (b) 感生面电荷密度
- (c) q所受作用力的大小与方向
- (d) 如果球体保持在固定电位V,则解是否有变化?如果球体内外表面带有总电荷Q,解是否有变换?

(a)

记球心为坐标原点,

球内情形相当于球外情形调换像电荷和源电荷,即

$$r' = \frac{a^2}{\tilde{r}'}, \quad q = -\frac{a}{\tilde{r}'}\tilde{q}$$

$$\downarrow \qquad \qquad \downarrow$$

$$\tilde{r}' = \frac{a^2}{r'}, \quad \tilde{q} = -\frac{a}{r'}q$$

$$(1)$$

有

$$R = |\vec{r} - \vec{r}'| = \sqrt{r^2 + r'^2 - 2rr'\cos\alpha}$$

$$\tilde{R} = |\vec{r} - \tilde{r}'|$$

$$= \sqrt{r^2 + \tilde{r}'^2 - 2r\tilde{r}'\cos\alpha}$$

$$= \frac{1}{r'}\sqrt{r^2r'^2 + R_0^4 - 2R_0^2rr'\cos\alpha}$$
(2)

所以

$$\varphi(\vec{r}) = \frac{q}{4\pi\varepsilon_0} \left(\frac{1}{R} - \frac{R_0}{r'} \frac{1}{\tilde{R}^2} \right)
= \frac{q}{4\pi\varepsilon_0} \left(\frac{1}{\sqrt{r^2 + r' - 2rr'\cos\alpha}} - \frac{R_0}{\sqrt{r^2r^2 + R_0^4 - 2R_0^2rr'\cos\alpha}} \right)$$
(3)

(b)

$$\sigma = -\varepsilon_0 \frac{\partial \varphi}{\partial r} \bigg|_{r=a} = -\frac{q}{4\pi a^2} \left(\frac{a}{r'}\right) \frac{1 - \frac{a^2}{r'^2}}{\left(1 + \frac{a^2}{r'^2} - 2\frac{a}{r'}\cos\alpha\right)^{3/2}} \tag{4}$$

(c)

$$|\mathbf{F}| = \frac{1}{4\pi\epsilon_0} \frac{q^2}{a^2} \left(\frac{a}{r'}\right)^3 \left(1 - \frac{a^2}{r'^2}\right)^{-2} \tag{5}$$

方向指向像电荷

(d)

保持固定电位,则解加一常数

表面带有电荷,则在球心加一源电荷

2.7

- (a) 写出适当的格林函数 $G(\vec{x},\vec{x'})$
- (b) 如果在边界上z=0处,以原点为圆心、半径为a的圆内电势 $\Phi=V$,圆外电势为0。求以柱坐标 $(
 ho,\phi,z)$ 指定的P点处的电势的积分表达式
- (c) 证明,沿圆轴 $(\rho=0)$ 的电势为

$$\Phi = V \left(1 - \frac{z}{\sqrt{a^2 + z^2}} \right) \tag{6}$$

(d) 证明,在远距离处 $(
ho^2+z^2\gg a^2)$,电势可以对 $(
ho^2+z^2)^{-1}$ 幂级数展开,首项为

$$\Phi = \frac{Va^2}{2} \frac{z}{(\rho^2 + z^2)^{3/2}} \left[1 - \frac{3a^2}{4(\rho^2 + z^2)} + \frac{5(3\rho^2 a^2 + a^4)}{8(\rho^2 + z^2)^2} + \cdots \right]$$
(7)

验证(c)、(d)的结果在共同的定义域内是一致的

(a)

$$G\left(\vec{x}, \vec{x}'\right) = \left(\frac{1}{\sqrt{\left(x - x'\right)^2 + \left(y - y'\right)^2 + \left(z - z'\right)^2}} - \frac{1}{\sqrt{\left(x - x'\right)^2 + \left(y - y'\right)^2 + \left(z + z'\right)^2}}\right) \tag{8}$$

(b)

边界条件即为

$$\varphi|_{\rho \le a, z=0} = V \quad \varphi|_{\rho > a, z=0} = 0 \tag{9}$$

有Dirichlet形式解

$$\varphi(\vec{r}) = \frac{1}{4\pi\varepsilon_0} \int_{V'} \rho\left(\vec{r'}\right) G dV' - \frac{1}{4\pi} \oint_{S'} \left[\varphi \frac{\partial G}{\partial n'} \right] dS'
= -\frac{1}{4\pi} \oint_{S'} \left[\varphi \frac{\partial G}{\partial n'} \right] dS'$$
(10)

其中,

$$\frac{\partial G}{\partial n'} = -\frac{\partial G}{\partial z'}\Big|_{z'=0}
= -\frac{2z}{\left((x-x')^2 + (y-y')^2 + z^2\right)^{3/2}}$$
(11)

有

$$\varphi(\vec{r}) = \frac{1}{4\pi} \oint_{S'} \left[\frac{2\varphi z}{\left((x - x')^2 + (y - y')^2 + z^2 \right)^{3/2}} \right] dS'
= \frac{1}{4\pi} \oint_{S'} \left[\frac{2\varphi z}{\left(-2\rho\rho'\cos(\phi - \phi') + \rho^2 + \rho'^2 + z^2 \right)^{3/2}} \right] dS'
= \frac{V}{2\pi} \oint_{S'} \left[\frac{z}{\left(-2\rho\rho'\cos(\phi - \phi') + \rho^2 + \rho'^2 + z^2 \right)^{3/2}} \right] dS'$$
(12)

下边界上有面元?

$$dS' = \rho' d\phi' d\rho' \tag{13}$$

$$\varphi(\vec{r}) = \frac{V}{2\pi} \int_0^a \int_0^{2\pi} \left[\frac{z}{(-2\rho\rho'\cos(\phi - \phi') + \rho^2 + \rho'^2 + z^2)^{3/2}} \right] \rho' d\phi' d\rho'$$
(14)

(c)

$$\varphi|_{\rho=0} = \frac{V}{2\pi} \int_0^a \int_0^{2\pi} \frac{z}{(\rho'^2 + z^2)^{3/2}} \rho' d\phi' d\rho'$$

$$= \frac{V}{2\pi} \cdot 2\pi \cdot \left(1 - \frac{z}{\sqrt{a^2 + z^2}}\right)$$
(15)

$$= V \cdot \left(1 - \frac{\tilde{}}{\sqrt{a^2 + z^2}}\right)$$

(d)

当 $ho^2+z^2\gg a^2$ 时,利用, $(1+x)^{-3/2}=O\left(x^5
ight)+rac{315x^4}{128}-rac{35x^3}{16}+rac{15x^2}{8}-rac{3x}{2}+1$

$$\varphi(\vec{r}) = \frac{V}{2\pi} \int_{0}^{a} \int_{0}^{2\pi} \left[\frac{z}{(-2\rho\rho'\cos(\phi - \phi') + \rho^{2} + \rho'^{2} + z^{2})^{3/2}} \right] \rho' d\phi' d\rho'
= \frac{V}{2\pi} \cdot \frac{z}{(\rho^{2} + z^{2})^{3/2}} \cdot \iint \left(1 - \frac{\rho'^{2} - 2\rho\rho'\cos(\phi - \phi')}{\rho^{2} + z^{2}} \right)^{-3/2} \rho' d\phi' d\rho'
= \frac{V}{2\pi} \cdot \frac{z}{(\rho^{2} + z^{2})^{3/2}} \cdot \iint \left[1 - \frac{3}{2} \cdot \frac{\rho'^{2} - 2\rho\rho'\cos(\phi - \phi')}{\rho^{2} + z^{2}} + \cdots \right] \rho' d\phi' d\rho'
= \frac{V}{2\pi} \cdot \frac{z}{(\rho^{2} + z^{2})^{3/2}} \cdot \left(\frac{a^{2}}{2} \cdot 2\pi - \frac{3}{2} \cdot \frac{1}{\rho^{2} + z^{2}} \cdot \frac{\pi a^{4}}{2} + \cdots \right)
= \frac{Va^{2}}{2} \cdot \frac{z}{(\rho^{2} + z^{2})^{3/2}} \cdot \left(1 - \frac{3a^{2}}{4(\rho^{2} + z^{2})} + \cdots \right)$$
(16)

QED

考虑z轴上 $(\rho = 0)$, 上式化为

$$\varphi(\vec{r})|_{z\to\infty} = \frac{Va^2}{2} \cdot \frac{z}{(\rho^2 + z^2)^{3/2}} \cdot \left(1 - \frac{3a^2}{4(\rho^2 + z^2)} + \cdots\right)$$
(17)

对于 (c) , 有

$$\varphi = V \left(1 - \frac{z}{\sqrt{a^2 + z^2}} \right)
= V \left(1 - \left(1 + \frac{a^2}{z^2} \right)^{-1/2} \right)
\approx \frac{Va^2}{2z^2} \left(1 - \frac{3a^2}{4z^2} + \frac{5a^4}{8z^4} \right)$$
(18)

可见(c)、(d)在共同定义域内结果一致

2.9

半径为a的绝缘、球形、导体壳处于均匀电场 E_0 中。如果球体被垂直于电场的平面切割成两个半球,求阻止两球分离的力。

- (a) 如果壳层不带电
- (b) 如果壳层总电荷量为Q

(a)

考虑电场沿z负方向,球体被xOy平面切开

根据课本(2.15),有感应电荷密度

$$\sigma = -\epsilon_0 \frac{\partial \Phi}{\partial r} \bigg|_{r=a} = 3\epsilon_0 E_0 \cos \theta \tag{19}$$

因此球面上面元 $dec{S}$ 受力,利用 $dec{F}=ec{E}dq=rac{\sigma^2dec{S}}{2\epsilon_0}$

$$\mathbf{E}_{ind} = rac{\sigma}{2\epsilon_0}\hat{\mathbf{r}}$$
 $\mathbf{F} = rac{1}{2\epsilon_0}\int \sigma^2\hat{\mathbf{r}}d\sigma$

W.J.Duffin.Electricity and Magnetism.McGraw-Hill Book Company,fourth edition,1990 P51

$$d\vec{F}_z = d\vec{F}\cos\theta = \frac{9}{2}\epsilon_0 E_0^2 \cos^3\theta d\vec{S}$$
 (20)

左右两半球积分有

$$\vec{F}_{z} = 2 \int_{0}^{2\pi} \int_{0}^{\pi/2} \frac{9}{2} \epsilon_{0} E_{0}^{2} \cos^{3}\theta a^{2} \sin\theta d\theta d\phi \hat{e_{z}}$$

$$= 9\pi a^{2} \epsilon_{0} E_{0}^{2} \int_{0}^{\pi/2} \cos^{3}\theta \sin\theta d\theta \hat{e_{z}}$$

$$= 9\pi a^{2} \epsilon_{0} E_{0}^{2} \int_{0}^{\pi/2} \cos^{3}\theta \sin\theta d\theta \hat{e_{z}}$$

$$= 9\pi a^{2} \epsilon_{0} E_{0}^{2} \int_{0}^{\pi/2} \cos^{4}\theta e^{2} \hat{e_{z}}$$
(21)

$$= \sin a \; \epsilon_0 \mathcal{L}_0 \left[-1/4 \cos \; \sigma \right]_0 \; \; \epsilon_z$$

$$= \frac{9}{4} \pi a^2 \epsilon_0 E_0^2 \hat{e_z}$$

(b)

考虑壳层带电Q,均匀分布于球面,有面密度

$$\sigma_a = rac{Q}{4\pi a^2}$$
 (22)

再考虑感应电荷, 有总电荷密度

$$\sigma = 3\epsilon_0 E_0 \cos \theta + \frac{Q}{4\pi a^2} \tag{23}$$

面源 $d\vec{S}$ 受力

$$d\vec{F}_z = d\vec{F}\cos\theta = \left(\frac{Q^2}{32\pi^2 a^4 \epsilon} + \frac{3eQ\cos(\theta)}{4\pi a^2} + \frac{9}{2}e^2\epsilon\cos^2(\theta)\right)d\vec{S}$$
 (24)

同(a)积分有

$$\vec{F}_{\vec{z}} = \left(\frac{9}{4}\pi a^2 \epsilon_0 E_0^2 + \frac{Q^2}{32\pi\epsilon_0 a^2} + \frac{E_0 Q}{2}\right) \hat{e_z}$$
 (25)

2.12

从二维电势问题的的级数解((2.71))出发,在半径为(0.71)的圆柱体表面指定电势,正式评估系数,将其代入级数,然后求和,以泊松积分的形式得到圆柱体内的电势。

$$\Phi(\rho,\phi) = \frac{1}{2\pi} \int_0^{2\pi} \Phi(b,\phi') \frac{b^2 - \rho^2}{b^2 + \rho^2 - 2b\rho\cos(\phi' - \phi)} d\phi'$$
 (26)

如果想得到在圆柱体和无穷大的空间所围区域内的电势,需要做什么修改?

(2.71)

$$\Phi(\rho, \phi) = a_0 + b_0 \ln \rho + \sum_{n=1}^{\infty} a_n \rho^n \sin(n\phi + \alpha_n)$$

$$+ \sum_{n=1}^{\infty} b_n \rho^{-n} \sin(n\phi + \beta_n)$$
(27)

柯西积分公式

设 Ω 是复平面 $\mathbb C$ 的一个单连通的开子集。 $f:\Omega\to\mathbb C$ 是一个 Ω 上的全纯函数。设 γ 是 Ω 内的一个简单闭合的可求长曲线(即连续而不自交并且能定义长度的闭合曲线),那么函数 f在 γ 内部的点a上的值是:

$$f(a) = \frac{1}{2\pi i} \oint_{\gamma} \frac{f(z)}{z - a} dz \tag{28}$$

有边界条件

①:
$$\Phi(\rho=0)=\infty$$

②:
$$\Phi(\rho = b) = V$$

利用边界条件①有

$$b_0 = b_n = 0 \tag{29}$$

有解

$$\Phi(\rho,\phi) = a_0 + \sum_{n=1}^{\infty} a_n \rho^n \sin(n\phi - \alpha_n)$$

$$= a_0 + \sum_{n=1}^{\infty} a_n \left(\frac{\rho}{b}\right)^n \sin(n\phi - \alpha_n)$$
(30)

利用, $\sin heta = ig(e^{i heta} - e^{-i heta}ig)/(2i)$, 有

$$\Phi(
ho,\phi) = a_0 + \sum_{n=1}^{\infty} rac{a_n}{2i} \Big(rac{
ho}{b}\Big)^n \left[e^{-ilpha_n}e^{in\phi} - e^{ilpha_n}e^{-in\phi}
ight]$$

$$= a_0 + \sum_{n=1}^{\infty} \left(\frac{\rho}{b}\right)^n \left[c_n e^{in\phi} + d_n e^{-in\phi}\right]$$

$$= \sum_{n=-\infty}^{\infty} c_n \left(\frac{\rho}{b}\right)^{|n|} e^{in\phi}$$
(31)

利用边界条件②有

$$V = \sum_{n=-\infty}^{\infty} c_n e^{in\phi}$$

$$\int_0^{2\pi} V(\phi) e^{-in'\phi} d\phi = \sum_{n=-\infty}^{\infty} c_n \int_0^{2\pi} e^{i(n-n')\phi} d\phi$$

$$\int_0^{2\pi} V(\phi) e^{-in'\phi} d\phi = \sum_{n=-\infty}^{\infty} c_n 2\pi \delta_{nn'}$$
(32)

有

$$c_n = \frac{1}{2\pi} \int_0^{2\pi} V(\phi) e^{-in\phi} d\phi \tag{33}$$

有圆柱内区域电势

$$\Phi(\rho,\phi) = \frac{1}{2\pi} \int_{0}^{2\pi} d\phi' V(\phi') \sum_{n=-\infty}^{\infty} \left(\frac{\rho}{b}\right)^{|n|} e^{in(\phi-\phi')}
= \frac{1}{2\pi} \int_{0}^{2\pi} d\phi' V(\phi') \left[-1 + \sum_{n=0}^{\infty} \left(\frac{\rho}{b}\right)^{n} e^{in(\phi-\phi')} + \sum_{n=0}^{\infty} \left(\frac{\rho}{b}\right)^{n} e^{-in(\phi-\phi')} \right]
= \frac{1}{2\pi} \int_{0}^{2\pi} d\phi' V(\phi') \left[-1 + \sum_{n=0}^{\infty} \left[\left(\frac{\rho}{b}\right) e^{i(\phi-\phi')} \right]^{n} + \sum_{n=0}^{\infty} \left[\left(\frac{\rho}{b}\right) e^{-i(\phi-\phi')} \right]^{n} \right]$$
(34)

利用, $\sum_0^\infty r^n = rac{1}{1-r}$, 有

$$\Phi(\rho,\phi) = \frac{1}{2\pi} \int_{0}^{2\pi} d\phi' V\left(\phi'\right) \left[-1 + \frac{1}{1 - \left(\frac{\rho}{b}\right) e^{i(\phi - \phi')}} + \frac{1}{1 - \left(\frac{\rho}{b}\right) e^{-i(\phi - \phi')}} \right] \\
= \frac{1}{2\pi} \int_{0}^{2\pi} d\phi' V\left(\phi'\right) \left[\frac{\left(\left(1 - \left(\frac{\rho}{b}\right) e^{-i(\phi - \Phi')}\right) + \left(1 - \left(\frac{\rho}{b}\right) e^{i(\phi - \phi')}\right) - \left(1 - \left(\frac{\rho}{b}\right) e^{i(\phi - \Phi')}\right) \left(1 - \left(\frac{\rho}{b}\right) e^{-i(\phi'\Phi')}\right)}{\left(1 - \left(\frac{\rho}{b}\right) e^{i(\phi - \phi')}\right) \left(1 - \left(\frac{\rho}{b}\right) e^{-i(\phi - \phi')}\right)} \right] (35) \\
= \frac{1}{2\pi} \int_{0}^{2\pi} \Phi\left(b, \phi'\right) \frac{b^{2} - \rho^{2}}{b^{2} + \rho^{2} - 2b\rho\cos\left(\phi' - \phi\right)} d\phi'$$

对于圆柱外区域,交换 b, ρ 即可,有

$$\Phi(\rho,\phi) = \frac{1}{2\pi} \int_0^{2\pi} \Phi(b,\phi') \frac{\rho^2 - b^2}{b^2 + \rho^2 - 2b\rho\cos(\phi' - \phi)} d\phi'$$
 (36)