

# §5. 球外(内)的 Dirichlet 问题

1. 球外(无源区) Green 函数、势的表达式

$$G(\vec{r}, \vec{r}') = \frac{1}{|\vec{r} - \vec{r}'|} - \frac{R_0}{r' |\vec{r} - \vec{r}'|}$$

$$\text{其中 } \vec{r}' \cdot \vec{r} = R_0^2$$

$$= \frac{1}{\sqrt{r^2 + r'^2 - 2rr' \cos \gamma}} - \frac{R_0}{r' \sqrt{r^2 + \tilde{r}'^2 - 2r\tilde{r}' \cos \gamma}}$$

$$= \frac{1}{\sqrt{r^2 + r'^2 - 2rr' \cos \gamma}} - \frac{R_0}{\sqrt{r^2 r'^2 + R_0^4 - 2R_0^2 rr' \cos \gamma}}$$

$$\frac{\partial G}{\partial n'} = - \frac{\partial}{\partial r'} \left( \frac{1}{\sqrt{\dots}} - \frac{R_0}{\sqrt{\dots}} \right) \Big|_{r'=R_0}$$

$$= - \frac{\partial}{\partial r'} \left( \frac{1}{\sqrt{\dots}} - \frac{R_0}{\sqrt{\dots}} \right) \Big|_{r'=R_0}$$

由 Gauss 公式 (Green 定理) 中,  $V$  与  $S$  的  
正方向约定: 研究对象  $V$  在  $S$  内,  $\vec{n}'$   
由内指向外。若  $\dots \dots S$  外,  $\vec{n}'$   
由外向内。

$$= - \left[ \left( -\frac{1}{2} \right) \frac{r' - r \cos \gamma}{(r^2 + r'^2 - 2rr' \cos \gamma)^{\frac{3}{2}}} - \left( -\frac{1}{2} \right) \frac{R_0 (r^2 r' - R_0^2 r \cos \gamma)}{(r^2 r'^2 + R_0^4 - 2R_0^2 rr' \cos \gamma)^{\frac{3}{2}}} \right] \Big|_{r'=R_0}$$

$$= \left[ \frac{r' - r \cos \gamma}{(r^2 + r'^2 - 2rr' \cos \gamma)^{\frac{3}{2}}} - \frac{R_0 (r^2 r' - R_0^2 r \cos \gamma)}{(r^2 r'^2 + R_0^4 - 2R_0^2 rr' \cos \gamma)^{\frac{3}{2}}} \right] \Big|_{r'=R_0}$$

$$= \frac{R_0 - r \cos \gamma}{(r^2 + R_0^2 - 2R_0 r \cos \gamma)^{\frac{3}{2}}} - \frac{R_0^2 (r^2 - R_0 r \cos \gamma)}{(r^2 R_0^2 + R_0^4 - 2R_0^3 r \cos \gamma)^{\frac{3}{2}}}$$

$$= \dots - \frac{r^2 - R_0 r \cos \gamma}{R_0 (r^2 + R_0^2 - 2R_0 r \cos \gamma)^{\frac{3}{2}}}$$

$$= \frac{1}{R_0 (r^2 + R_0^2 - 2R_0 r \cos \gamma)^{\frac{3}{2}}} (R_0^2 - R_0 r \cos \gamma - r^2 + R_0 r \cos \gamma)$$

$$= \frac{R_0^2 - r^2}{R_0 (R_0^2 + r^2 - 2R_0 r \cos \gamma)^{\frac{3}{2}}}$$

故 Dirichlet 问题 (特别对于无源区) 球外

$$\varphi(\vec{r}) = - \frac{1}{4\pi} \oint \varphi(\vec{r}') \frac{\partial G(\vec{r}, \vec{r}')}{\partial n'} dS'$$

$$= \frac{R_0^2 - r^2}{4\pi R_0} \oint \varphi(\vec{r}') \frac{dS'}{(R_0^2 + r^2 - 2R_0 r \cos \gamma)^{\frac{3}{2}}}$$

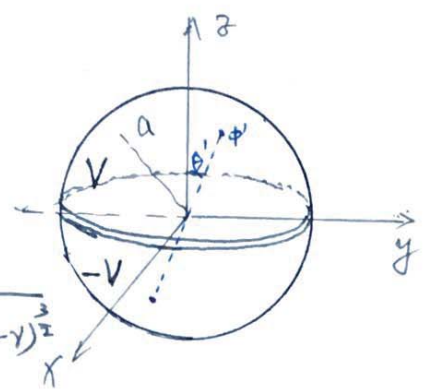
$$\left. \begin{aligned} \text{其中 } dS' &= R_0^2 d\Omega' \\ &= R_0^2 \sin \theta' d\theta' d\phi' \end{aligned} \right\}$$

$$= \frac{R_0 (R_0^2 - r^2)}{4\pi} \oint \frac{\varphi(\vec{r}') d\Omega'}{(R_0^2 + r^2 - 2R_0 r \cos \gamma)^{\frac{3}{2}}}$$

## 2. 双球导体外的电势

$$\varphi(\vec{r}) = \frac{a(a^2-r^2)V}{4\pi} \times$$

$$\int_{\Omega^+} \frac{d\Omega'}{(a^2+r^2-2ar\cos\gamma)^{\frac{3}{2}}} - \int_{\Omega^-} \frac{d\Omega'}{(a^2+r^2-2ar\cos\gamma)^{\frac{3}{2}}}$$



$$= \frac{Va(a^2-r^2)}{4\pi} \int_0^{2\pi} d\phi' \left[ \int_0^{\frac{\pi}{2}} \frac{\sin\theta' d\theta'}{(a^2+r^2-2ar\cos\gamma)^{\frac{3}{2}}} - \int_{\frac{\pi}{2}}^{\pi} \frac{\sin\theta' d\theta'}{(a^2+r^2-2ar\cos\gamma)^{\frac{3}{2}}} \right]$$

其中  $\cos\gamma = \cos\theta\cos\theta' + \sin\theta\sin\theta'\cos(\phi-\phi')$

$$= \frac{Va(a^2-r^2)}{4\pi} \int_0^{2\pi} d\phi' \left[ \int_0^{\frac{\pi}{2}} \frac{d(\cos\theta')}{(a^2+r^2-2ar\cos\gamma)^{\frac{3}{2}}} - \int_{-1}^0 \frac{d(\cos\theta')}{(a^2+r^2-2ar\cos\gamma)^{\frac{3}{2}}} \right]$$

第二项作代换 (反演)  $\theta' = \pi - \theta''$ ,  $\phi' = \phi'' + \pi$

$$\cos\gamma = -\cos\theta\cos\theta'' - \sin\theta\sin\theta''\cos(\phi-\phi''-\pi)$$

$$= -\cos\theta\cos\theta'' - \sin\theta\sin\theta''\cos(\phi-\phi'') = -\cos\gamma'$$

且  $d(\cos\theta') = -d(\cos\theta'')$

$$\int_{-1}^0 \dots = \int_1^0 \dots \quad \int_0^{\frac{\pi}{2}} \dots = \int_{\frac{\pi}{2}}^{\pi} \dots$$

$$\therefore \varphi(\vec{r}) = \frac{Va(a^2-r^2)}{4\pi} \int_0^{2\pi} d\phi' \int_0^1 d(\cos\theta') \left[ \frac{1}{(a^2+r^2-2ar\cos\gamma)^{\frac{3}{2}}} - \frac{1}{(a^2+r^2+2ar\cos\gamma)^{\frac{3}{2}}} \right]$$

$$= \frac{V}{4\pi} \frac{a(a^2-r^2)}{(a^2+r^2)^{\frac{3}{2}}} \int_0^{2\pi} d\phi' \int_0^1 d(\cos\theta') \left[ \frac{1}{(1-2\alpha\cos\gamma)^{\frac{3}{2}}} - \frac{1}{(1+2\alpha\cos\gamma)^{\frac{3}{2}}} \right]$$

其中  $\alpha = \frac{ar}{a^2+r^2}$

展开  $(1 \pm x)^{-\frac{3}{2}} = 1 \mp \frac{3}{2}x + \frac{1}{2!} \cdot \frac{15}{4}x^2 \mp \frac{1}{3!} \cdot \frac{105}{8}x^3 + \dots$

$$(1-2\alpha\cos\gamma)^{-\frac{3}{2}} = (1+2\alpha\cos\gamma)^{-\frac{3}{2}}$$

$$= 6\alpha\cos\gamma + 35\alpha^3\cos^3\gamma$$

$$\therefore \varphi(\vec{r}) = \frac{V}{4\pi} \frac{a(a^2-r^2)}{(a^2+r^2)^{\frac{3}{2}}} \int_0^{2\pi} d\phi' \int_0^1 d(\cos\theta') (6\alpha\cos\gamma + 35\alpha^3\cos^3\gamma)$$

其中  $\int_0^{2\pi} d\phi' \int_0^1 d(\cos\theta') \cos\gamma = \int_0^{2\pi} d\phi' \int_0^1 d(\cos\theta') \cos\theta' \cdot \cos\theta$

$$+ \int_0^{2\pi} d\phi' \cos(\phi-\phi') \int_0^1 d(\cos\theta') \sin\theta' \sin\theta$$

$$= 2\pi \cdot \cos\theta \cdot \frac{1}{2} \cos^2\theta' \Big|_0^1 = \pi\cos\theta$$

§6. 分离变量法之极坐标

1. 极坐标下 Laplace 方程通解

$$\nabla^2\varphi = 0 \rightarrow \text{极}$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial \varphi}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \varphi}{\partial \phi^2} = 0$$

取  $\varphi = R(\rho) \Phi(\phi)$

$$\frac{1}{R} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial R}{\partial \rho} \right) = - \frac{1}{\Phi} \frac{\partial^2 \Phi}{\partial \phi^2} \equiv \nu^2$$

作代: 2.2, 2.7, 2.9, 2.12