

热力学与统计物理-第八次作业

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Problem 7.1

Answer:

(a)

We label the positions and momenta such that r_{ij} and p_{ij} refer to the j th molecule of type i . There are N_i molecules of species i . Then the classical partition function for the mixture of ideal gas is:

$$z' = \int \exp \left[-\frac{\beta}{2m_1} (p_{11}^2 + \dots + p_{1N_1}^2) \dots - \frac{\beta}{2m_k} (p_{k1}^2 + \dots + p_{kN_k}^2) \right] \frac{d^3r_{11} \dots d^3r_{kN_k} d^3p_{11} \dots d^3p_{kN_k}}{h_0^{3N_1} \dots h_0^{3N_k}} \quad (1.1)$$

This integration over r yield the volume, V , while the p integrals are identical. Since there are $N_1 + N_2 + \dots + N_k$ integration, we have:

$$z' = V^{(N_1 + \dots + N_k)} \left[\int e^{-\frac{\beta p^2}{2m}} \frac{d^3p}{h_0^3} \right]^{(N_1 + \dots + N_k)} \quad (1.2)$$

The term in brackets is independent of volume, consequently:

$$\ln z' = (N_1 + \dots + N_k) \ln V + \ln(\text{constant}) \quad (1.3)$$

And:

$$\bar{p} = \frac{1}{\beta} \frac{\partial}{\partial V} \ln z' = (N_1 + \dots + N_k) \frac{1}{\beta V} \quad (1.4)$$

$$\bar{p}V = (N_1 + \dots + N_k) kT = (\nu_1 + \dots + \nu_k) RT \quad (1.5)$$

(b)

For the i^{th} gas, $p_i V = \nu_i RT$, from (1.5):

$$\bar{p} = \sum_i \bar{p}_i \quad (1.6)$$

Problem 7.3

Answer:

Before the partition is removed, we have on the left, $pV = \nu RT$. After removal the pressure is:

$$p_f = \frac{2\nu}{(1+b)V} = \frac{2p}{1+b} \quad (2.1)$$

(b)

The initial and final entropies of the system for different gases are:

$$S_i = \nu R \left[\ln \frac{V}{N_A \nu} + \frac{3}{2} \ln T + \sigma_1 \right] + \nu R \left[\ln \frac{bV}{N_A \nu} + \frac{3}{2} \ln T + \sigma_2 \right] \quad (2.2)$$

$$S_f = \nu R \left[\ln \frac{(1+b)V}{N_A \nu} + \frac{3}{2} \ln T + \sigma_1 \right] + \nu R \left[\ln \frac{(1+b)V}{N_A \nu} + \frac{3}{2} \ln T + \sigma_2 \right] \quad (2.3)$$

Here one adds the entropies of the gases in the left and right compartments for S_i while S_f is the entropy of two different gases in volume $(1+b)V$. Then:

$$\Delta S = S_f - S_i = \nu R \left[2 \ln \frac{V(1+b)}{N_A \nu} - \ln \frac{V}{N_A \nu} - \ln \frac{bV}{N_A \nu} \right] = \nu R \ln \frac{(1+b)^2}{b} \quad (2.4)$$

(c)

In the case of identical gases, S_i is again the sum of the entropies of the left and right compartments. S_f is the entropy of 2ν moles in a volume $(1+b)V$

$$S_i = \nu R \left[\ln \frac{V}{N_A \nu} + \frac{3}{2} \ln T + \sigma_0 \right] + \nu R \left[\ln \frac{bV}{N_A \nu} + \frac{3}{2} \ln T + \sigma_0 \right] \quad (2.5)$$

$$S_f = 2\nu R \left[\ln \frac{(1+b)V}{2N_A \nu} + \frac{3}{2} \ln T + \sigma_0 \right] \quad (2.6)$$

Thus:

$$\Delta S = \nu R \left[\ln \frac{(1+b)V}{2N_A \nu} - \ln \frac{V}{N_A \nu} - \ln \frac{bV}{N_A \nu} \right] = \nu R \ln \frac{(1+b)^2}{4b} \quad (2.7)$$

Problem 7.4

Answer:

(a)

The system is isolated so its total energy is constant, and since the energy of an ideal gas depends only on temperature, we have

$$\Delta E_1 + \Delta E_2 = C_V (T_f - T_1) + C_V (T_f - T_2) = 0 \quad (3.1)$$

Or

$$T_f = \frac{T_1 + T_2}{2} \quad (3.2)$$

The total volume is found from the equation of state

$$V = \frac{\nu_1 R T_1}{p_1} + \frac{\nu_2 R T_2}{p_2} \quad (3.3)$$

Thus the final pressure is

$$p_f = \frac{(\nu_1 + \nu_2) R T_f}{V} = \frac{(\nu_1 + \nu_2)}{2} \left(\frac{T_1 + T_2}{(\nu_1 T_1/p_1) + (\nu_2 T_2/p_2)} \right) \quad (3.4)$$

(b)

Using $\frac{V}{N} = \frac{kT}{p}$, we have for the initial and final entropies of different gases

$$S_i = \nu_1 R \left[\ln \frac{kT_1}{p_1} + \frac{3}{2} \ln T_1 + \sigma_1 \right] + \nu_2 R \left[\ln \frac{kT_2}{p_2} + \ln T_2 + \sigma_2 \right] \quad (3.5)$$

$$\begin{aligned} S_f &= \nu_1 R \left[\ln \frac{k}{\nu_1} \left(\frac{\nu_1 T_1}{p_1} + \frac{\nu_2 T_2}{p_2} \right) + \frac{3}{2} \ln \frac{T_1 + T_2}{2} + \sigma_1 \right] \\ &+ \nu_2 R \left[\ln \frac{k}{\nu_2} \left(\frac{\nu_1 T_1}{p_1} + \frac{\nu_2 T_2}{p_2} \right) + \frac{3}{2} \ln \frac{T_1 + T_2}{2} + \sigma_2 \right] \end{aligned} \quad (3.6)$$

So:

$$\begin{aligned} \Delta S &= S_f - S_i = \nu_1 R \left[\ln \left(1 + \frac{\nu_2 T_2 p_1}{\nu_1 T_1 p_2} \right) + \frac{3}{2} \ln \frac{T_1 + T_2}{2T_1} \right] \\ &+ \nu_2 R \left[\ln \left(1 + \frac{\nu_1 T_1 p_2}{\nu_2 T_2 p_1} \right) + \frac{3}{2} \ln \frac{T_1 + T_2}{2T_2} \right] \end{aligned} \quad (3.7)$$

(c)

For identical gases:

$$S_i = \nu_1 R \left[\ln \frac{kT_1}{p_1} + \frac{3}{2} \ln T_1 + \sigma_0 \right] + \nu_2 R \left[\ln \frac{kT_2}{p_2} + \frac{3}{2} \ln T_2 + \sigma_0 \right] \quad (3.8)$$

$$S_f = (\nu_1 + \nu_2) R \left[\ln \frac{k}{(\nu_1 + \nu_2)} \left(\frac{\nu_1 T_1}{p_1} + \frac{\nu_2 T_2}{p_2} \right) + \frac{3}{2} \ln \frac{T_1 + T_2}{2} + \sigma_0 \right] \quad (3.9)$$

So:

$$\begin{aligned} \Delta S &= \nu_1 R \left[\ln \frac{1}{(\nu_1 + \nu_2)} \left(\nu_1 + \frac{\nu_2 T_2 p_1}{T_1 p_2} \right) + \frac{3}{2} \ln \frac{T_1 + T_2}{2T_1} \right] \\ &+ \nu_2 R \left[\ln \frac{1}{(\nu_1 + \nu_2)} \left(\nu_2 + \frac{\nu_1 T_1 p_2}{T_2 p_1} \right) + \frac{3}{2} \ln \frac{T_1 + T_2}{2T_2} \right] \end{aligned} \quad (3.9)$$

Problem 7.5

Answer:

We take the zero of potential energy so that if a segment is oriented parallel to the vertical it contributes energy Wa and if antiparallel it contributes $-Wa$ to the total energy of the rubber band (Thus if the rubber band were fully extended, the total energy would be $-NWa$). Since the segments are non-interacting,

$$\bar{l} = N \frac{ae^{-Wa\beta} - ae^{Wa\beta}}{e^{-Wa\beta} + e^{Wa\beta}} = Na \tanh \frac{Wa}{kT} \quad (4.1)$$

Problem 7.6

Answer:

Since the total energy is additive:

$$E_i = \epsilon_i(p_i) + U(q_1 \cdots q_n) = \frac{p^2}{2m} + U(q_1 \cdots q_n) \quad (5.1)$$

where U is the energy of interaction, the equipartition theorem still applies and

$$\bar{\epsilon} = \frac{3}{2}kT \quad (5.2)$$

Problem 7.7

Answer:

If the gas is ideal, its mean energy per particle is

$$\bar{\epsilon} = \frac{\overline{p_x^2}}{2m} + \frac{\overline{p_y^2}}{2m} = kT \quad (6.1)$$

and the mean energy per mole becomes $\bar{E} = N_A kT$:

$$C = \frac{\partial \bar{E}}{\partial T} = N_A k = R \quad (6.2)$$

Problem 7.10

Answer:

(a)

Let the restoring force be $-\alpha x$. Then the mean energy of N particles is

$$\bar{E} = N \left(\frac{1}{2} m \overline{x^2} + \frac{1}{2} \alpha \overline{x^2} \right) \quad (7.1)$$

Then:

$$\bar{E} = N \left(\frac{1}{2}kT + \frac{1}{2}kT \right) = NkT \quad (7.2)$$

So:

$$c = \frac{\partial \bar{E}}{\partial T} = Nk \quad (7.3)$$

(b)

If the restoring force is $-\alpha x^3$, the mean energy per particle is

$$\bar{\epsilon} = \frac{1}{2}m\bar{x}^2 + \frac{1}{4}\alpha\bar{x} \quad (7.4)$$

So, similarly with (a), we can get:

$$C = \frac{\partial \bar{E}}{\partial T} = \frac{3}{4}Nk \quad (7.5)$$

Problem 7.12

Answer:

(a)

Consider a cube of side a . The force necessary to decrease the length of a side by Δa is $\kappa_0 A \Delta a$ and therefore the pressure is $\Delta p = \kappa_0 \Delta a / a^2$. The change in volume is $\Delta V = -a^2 \Delta a$

$$\kappa = -\frac{1}{V} \left(\frac{\Delta V}{\Delta p} \right) = -\frac{1}{a^3} \left(-\frac{a^2 \Delta a}{\kappa_0 \Delta a / a^2} \right) = \frac{a}{\kappa_0} \quad (8.1)$$

(b)

The Einstein temperature is $\theta_E = \hbar\omega/k$. Since $\omega = \sqrt{\kappa_0/m}$ where m is mass, we have from (8.1)

$$\theta_E = \frac{\hbar}{k} \sqrt{\frac{\kappa}{m}} = \frac{\hbar}{k} \sqrt{\frac{a}{m\kappa}} \quad (8.2)$$

So:

$$a = \left(\frac{\mu}{\rho N_A} \right)^{1/3} = \left(\frac{63 \cdot 5}{(8.9)(6 \times 10^{23})} \right)^{1/3} = 2.3 \times 10^{-8} \text{cm} \quad (8.3)$$

So:

$$\theta_E \approx 150^\circ \text{K} \quad (8.4)$$