## 电动力学-第四次作业

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## 2020年3月26日

## Problem3.7

Answer:

For this question, we want to solve the electric field above the plane(z>0), and the Dirichlet Boundary condition is:

$$\begin{cases} V=0, where z=0\\ V=0, where \sqrt{x^2+y^2+z^2} \to \infty \end{cases} \tag{1.1}$$

If we remove the grounded conductor, and place a charge -2q at (0,0,-d), and a charge +q at (0,0,-3d), it's easy to prove that it has the same Dirichlet boundary condition with the origin question. So, the force exert on +q is:

$$F = \frac{1}{4\pi\epsilon_0} \left( -\frac{2q^2}{4d^2} - \frac{2q^2}{16d^2} + \frac{q^2}{36d^2} \right) = -\frac{1}{4 pi\epsilon_0} \frac{43q^2}{72d^2}$$
 (1.2)

Problem3.15

Answer:

(a).

According to the symmetry of the pipe, it's easy to know that the potential inside the pipe is only the function of x and y.

And we can determined the Dirichlet boundary condition:

$$\begin{cases} V(x = 0, y) = 0 \\ V(x = b, y) = V_0(y) \\ V(x, y = 0) = 0 \\ V(x, y = b) = 0 \end{cases}$$
(2.1)

Assume that the V(x,y) could be writen in following form:

$$V(x,y) = \mathcal{X}(x)\mathcal{Y}(y) \tag{2.2}$$

So, the laplace's equation turn to:

$$\mathcal{Y}\frac{\partial^2 \mathcal{X}}{\partial x^2} + \mathcal{X}\frac{\partial^2 \mathcal{Y}}{\partial y^2} = 0 \tag{2.3}$$

Which equal to:

$$\frac{1}{\mathcal{X}}\frac{\partial^2 \mathcal{X}}{\partial x^2} + \frac{1}{\mathcal{Y}}\frac{\partial^2 \mathcal{Y}}{\partial y^2} == 0$$
 (2.4)

The two terms in the left side of the equation is only the function of x and of y. If we want (2.4) always correct, the following equation must be correct.

$$\begin{cases} \frac{\partial^2 \mathcal{X}}{\partial x^2} = k^2 \mathcal{X} \\ \frac{\partial^2 \mathcal{Y}}{\partial y^2} = -k^2 \mathcal{Y} \end{cases}$$
 (2.5)

Solve (2.5) we get:

$$V(x,y) = (C_1 e^{kx} + C_2 e^{-kx})(C_3 \sin(ky) + C_4 \cos(ky))$$
 (2.6)

By consider the boundary condition, we can determined the coefficient:

$$V(x,y) = C_1 \left(e^{\frac{n\pi x}{a}} - e^{-\frac{n\pi x}{a}}\right) \left(C_4 \sin(\frac{n\pi y}{a})\right)$$
$$= 2C_1 C_4 \sinh(\frac{n\pi x}{a}) \sin(\frac{n\pi y}{a}) \tag{2.7}$$

(2.7) could be writen in a sum form:

$$V(x,y) = \sum_{n=1}^{\infty} C_n \sinh(\frac{n\pi x}{a}) \sin(\frac{n\pi y}{a})$$
 (2.8)

And solve  $C_n$  with  $V_0(y)$ :

$$C_n = \frac{2}{a \sinh(n\pi b/a)} \int_0^a V_0(y) \sin(\frac{n\pi y}{a}) dy$$
 (2.9)

(b)

$$C_n = \frac{2V_0}{a \sinh(n\pi b/a)} \int_0^a \sin(\frac{n\pi y}{a}) dy$$
 (2.10)

If n is even:

$$C_n = 0 (2.11)$$

If n in odd:

$$C_n = \frac{4V_0}{n\pi \sinh(n\pi b/a)} \tag{2.12}$$

The potential is:

$$V(x,y) = \frac{4V_0}{\pi} \sum_{n=1,3,5} \frac{\sinh(n\pi x/a)\sin(n\pi y/a)}{n \sinh(n\pi b/a)}$$
(2.13)

Problem3.19

Answer:

$$V_0(\theta) = k\cos(3\theta) \tag{3.1}$$

Express (3.1) in the combination of Legendre polynomials:

$$V_0(\theta) = k[\alpha P_3(\cos\theta) + \beta P_1(\cos\theta)] \tag{3.2}$$

Determined the coefficient  $\alpha$  and  $\beta$ :

$$\cos(3\theta) = 4\cos^3\theta - 3\cos\theta = \alpha \left[\frac{1}{2}(5\cos^3\theta - 3\cos\theta)\right] + \beta\cos\theta \tag{3.3}$$

So:

$$\begin{cases} \alpha = \frac{8}{5} \\ \beta = -\frac{3}{5} \end{cases}$$
 (3.4)

Therefore:

$$V_0(\theta) = \frac{k}{5} [8P_3(\cos\theta) - 3P_1(\cos\theta)] \tag{3.5}$$

For the inside space:

$$V(r,\theta) = \sum_{l=0}^{\infty} A_l \tau^l P_l(\cos\theta) \qquad r \le R$$
 (3.6)

 $A_l$  is determined by:

$$A_{l} = \frac{(2l+1)}{2R^{l}} \int_{0}^{\pi} V_{0}(\theta) P_{l}(\cos\theta) \sin\theta d\theta$$

$$= \frac{(2l+1)}{2R^{l}} \frac{k}{5} \left\{ 8 \int_{0}^{\pi} P_{3}(\cos\theta) P_{l}(\cos\theta) \sin\theta d\theta - 3 \int_{0}^{\pi} P_{1}(\cos\theta) P_{l}(\cos\theta) \sin\theta d\theta \right\}$$

$$= \frac{(2l+1)}{2R^{l}} \frac{k}{5} \left\{ 8 \frac{2}{(2l+1)} \delta_{l3} - 3 \frac{2}{(2l+1)\delta_{l1}} \right\}$$

$$= \frac{k}{5R^{l}} \left[ 8\delta_{l3} - 3\delta_{l1} \right]$$
(3.7)

So:

$$V(r,\theta) = \frac{k}{5} \left[ 8\left(\frac{r}{R}\right)^3 P_3(\cos\theta) * 3\left(\frac{r}{R}\right) P_1(\cos\theta) \right]$$
(3.8)

Then, look outside:

$$V(r,\theta) = \sum_{l=0}^{\infty} \frac{B_l}{r^{l+1}} P_l(\cos\theta) \qquad r \ge R$$
 (3.9)

And we can determined  $B_l$ :

$$B_l = \frac{k}{5} \frac{1}{R^{l+1}} [8\delta_{l3} - 3\delta_{l1}] \tag{3.10}$$

So:

$$V(r,\theta) = \frac{k}{5} \left[ 8(\frac{R}{r})^4 P_3(\cos\theta) - 3(\frac{R}{r})^2 P_1(\cos\theta) \right]$$
 (3.11)

For the charge density:

$$\sigma(\theta) = \epsilon_0 \sum_{l=0}^{\infty} (2l+1)A_l R^l (l-1)P_l(\cos\theta) = \frac{\epsilon_0 k}{5R} [-9P_1(\cos\theta) + 56P_3(\cos\theta)]$$
(3.12)

Problem3.45

Answer:

(a)

$$\frac{1}{2} \sum_{i,j=1}^{3} \hat{r}_i \hat{r}_j Q_{ij} = \frac{1}{2} \int 3 \sum_{i=1}^{3} \hat{r}_i r'_j \sum_{j=1}^{3} \hat{r}_j r'_i - (r')^2 \sum_{i,j} \hat{r}_i \hat{r}_j \delta_{ij} \rho dr'$$
(4.1)

So, the potential of the quadrupole is:

$$V_{quad} = \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} \int \frac{1}{2} (r'^2 \cos\theta' - r') \rho d\tau' = \frac{1}{4\pi\epsilon_0} \int r'^2 P_2(\cos\theta') \rho d\tau' \quad (4.2)$$

(b)

Since  $x^2 = y^2 = (a/2)^2$  for all four charges:

$$Q_{xx} = Q_{yy} = 0 (4.3)$$

And z=0:

$$Q_{zz} = Q_{(xz)} = Q_{(yz)} = Q_{(zx)} = Q_{(zy)} = 0 (4.4)$$

For  $Q_{xy}$  and  $Q_{yx}$ :

$$Q_{xy} = Q_{yx} = 3a^2q (4.5)$$