

热力学与统计物理-第十次作业

吴远清-2018300001031

2020 年 5 月 21 日

Problem 7.14

Answer:

The energy of the magnetic moment in the field H is:

$$E = -\mu \cdot H = -\mu H \cos \theta \quad (1.1)$$

where θ is the angle between the magnetic moment and the direction of the field or z-axis. Then the probability that the magnetic moment lies in the range θ to $\theta + d\theta$ is proportional to the Boltzmann factor and the solid angle $2\pi \sin \theta d\theta$

Thus:

$$P(\theta)d\theta \propto e^{\beta\mu H \cos \theta} \sin \theta d\theta \quad (1.2)$$

And:

$$\bar{M}_Z = N_0 \bar{\mu}_Z = \frac{N_0 \int_0^\pi e^{\beta\mu H \cos \theta} \cos \theta \sin \theta d\theta (\mu \cos \theta)}{\int_0^\pi e^{\beta\mu H \cos \theta} \sin \theta d\theta} \quad (1.3)$$

$$\begin{aligned} \bar{M}_Z &= \frac{N_0}{H} \frac{\partial}{\partial \beta} \ln \int_0^\pi e^{\beta\mu H \cos \theta} \sin \theta d\theta = \frac{N_0}{H} \frac{\partial}{\partial \beta} \ln \frac{e^{\beta\mu H} - e^{-\beta\mu H}}{\beta\mu H} \\ &= \frac{N_0}{H} \frac{\partial}{\partial \beta} \ln \frac{2 \sinh \beta\mu H}{\beta\mu H} = \frac{N_0}{H\beta \sinh \beta\mu H} [\mu H \beta \cosh \beta\mu H - \sinh \beta\mu H] \end{aligned} \quad (1.4)$$

So:

$$\bar{M}_Z = N_0 \mu \left[\coth \beta\mu H - \frac{1}{\beta\mu H} \right] \quad (1.5)$$

Problem 7.15

Answer:

We have:

$$\bar{M}_Z = N_0 g \mu_0 J B_J(\eta) \quad (2.1)$$

Then:

$$B_J(\eta) = \frac{1}{J} \left[\left(J + \frac{1}{2} \right) \coth \left(J + \frac{1}{2} \right) \eta - \frac{1}{2} \coth \frac{1}{2} \eta \right] \quad (2.2)$$

If $\eta \ll 1$ and $J \gg 1$ in such a way that $J\eta \gg 1$, $B_J(\eta)$ becomes:

$$B_J(\eta) = \frac{1}{J} \left[J \coth J\eta - \frac{1}{2} \left(\frac{2}{\eta} \right) \right] = \coth J\eta - \frac{1}{J\eta} \quad (2.3)$$

Let $\mu H \beta = J\eta = Jg\mu_0 H \beta$ where $\mu = g\mu_0 J$ is by (7.8.2) the classical magnetic moment. , then (2.1) becomes:

$$\overline{M}_Z = N_0 \mu \left[\coth \beta \mu H - \frac{1}{\beta \mu H} \right] \quad (2.4)$$

Problem 7.17

Answer:

To find the fraction, ξ , of molecules with x component of velocity between $-\tilde{v}$ and \tilde{v} , we must integrate the distribution between these limits, i.e.:

$$\xi = \frac{1}{n} \int_{-\tilde{v}}^{\tilde{v}} g(v_x) dv_x = \int_{-\sqrt{\frac{2kT}{m}}}^{\sqrt{\frac{2kT}{m}}} \left(\frac{m}{2\pi kT} \right)^{\frac{1}{2}} e^{-(mv_x^2/2kT)} dv_x \quad (3.1)$$

Making the change of variable, $y = \sqrt{\frac{m}{kT}} v_x$, we have:

$$\xi = \frac{1}{\sqrt{2\pi}} \int_{-\sqrt{2}}^{\sqrt{2}} e^{-(y^2/2)} dy = \frac{2}{\sqrt{2\pi}} \int_0^2 e^{-(y^2/2)} dy = 2 \operatorname{erf} \sqrt{2} \quad (3.2)$$

Problem 7.18

Answer:

In problem 5.9 we found that the velocity of sound is:

$$u = \left(\frac{\gamma RT}{\mu} \right)^{\frac{1}{2}} \quad (4.1)$$

where $\gamma = C_p/C_v$ and μ is the atomic weight. Since $\mu = N_A m$, we have:

$$u = \left(\frac{\gamma RT}{N_A m} \right)^{\frac{1}{2}} = \left(\frac{\gamma kT}{m} \right)^{\frac{1}{2}} \quad (4.2)$$

The most probable speed is $\tilde{v} = (2kT/m)^{\frac{1}{2}}$. Thus:

$$u = \left(\frac{\gamma}{2} \right)^{\frac{1}{2}} \tilde{v} \quad (4.3)$$

For helium, $\gamma = 1.66$ so that $u = 0.91\tilde{v}$, and the fraction of molecules with speeds less than u is:

$$\xi = \frac{1}{n} \int_0^{0.91\tilde{v}} F(v) dv = 4\pi \int_0^{0.91(2kT/m)^{\frac{1}{2}}} \left(\frac{m}{2\pi kT} \right)^{\frac{3}{2}} v^2 e^{-\frac{mv^2}{2kT}} dv \quad (4.4)$$

Making the change of variable $y = (m/kT)^{\frac{1}{2}}v$, we have:

$$4\xi = \frac{4\pi}{(2\pi)^{\frac{3}{2}}} \int_0^{0.91\sqrt{2}} y^2 e^{-y^2/2} dy \quad (4.5)$$

This integral may be evaluated by noticing that integration by parts of $\int_0^a e^{-y^2/2} dy$ yields:

$$\int_0^a e^{-y^2/2} dy = ye^{-y^2/2} \Big|_0^a + \int_0^a y^2 e^{-y^2/2} dy \quad (5.6)$$

Then we find:

$$\xi = 4\pi^{-\frac{1}{2}}(2)^{-\frac{3}{2}} \int_0^{0.91\sqrt{2}} y^2 e^{-y^2/2} dy = 4\pi^{-\frac{1}{2}}(2)^{-\frac{3}{2}} \int_0^{0.91\sqrt{2}} e^{-y^2/2} dy - 0.91\sqrt{2} \exp[-(0.912)^2/2] \approx 0.37 \quad (5.7)$$

Problem 7.19

Answer:

$$(a) \bar{v}_x = 0$$

$$(b) \overline{v_x^2} = \frac{kT}{m}$$

$$(c) \overline{(v^2 v_x)} = (\overline{v_x^3} + \overline{v_y^2 v_x} + \overline{v_z^2 v_x}) = 0$$

$$(d) \overline{(v_x^3 v_y)} = \overline{v_x^3} \bar{v}_y = 0$$

$$(e) \overline{(v_x + bv_y)^2} = \overline{v_x^2} + 2b\bar{v}_x \bar{v}_y + b^2 \overline{v_y^2} = (1+b)^2 \frac{kT}{m}$$

$$(f) \overline{v_x^2 v_y^2} = \left(\frac{kT}{m} \right)^2$$

Problem 7.21

Answer:

The most probable energy is given by the condition $\frac{dF(\epsilon)}{d\epsilon} = 0$:

$$\frac{1}{2} \epsilon^{-\frac{1}{2}} e^{-\frac{\epsilon}{kT}} - \frac{\epsilon^{\frac{1}{2}}}{kT} e^{-\frac{\epsilon}{kT}} = 0 \quad (7.1)$$

Then:

$$\tilde{\epsilon} = \frac{1}{2}kT \quad (7.2)$$

The most probable speed is:

$$\tilde{v} = (2kT/m)^{\frac{1}{2}} \quad (7.3)$$

So:

$$\frac{1}{2}mv^2 = kT \quad (7.4)$$

Problem 7.23

Answer:

(a)

The number of molecules which leave the source slit per second is:

$$\Phi_0 A = \frac{1}{4}n\bar{v}A = \frac{\bar{p}_s A}{\sqrt{2\pi mkT}} = 1.1 \times 10^{18} \text{ molecules/sec} \quad (8.1)$$

where A is the area of slit.

(b)

Approximating the slit as a point source, we have by (7.11.7), the number of molecules with speed in the range between v and $v + dv$ which emerge into solid angle $d\Omega$ is:

$$A\Phi(v)d^3v = A[f(v)v \cos \theta] [v^2 dv d\Omega] \quad (8.2)$$

$\cos \theta \approx 1$ for molecules arriving at the detector slit; hence the total number which reach the detector is

$$A \int_{\Omega_d} d\Omega \int_0^\infty n \left(\frac{m}{2\pi kT} \right)^{\frac{3}{2}} v^3 e^{-\frac{\beta mv^2}{2}} dv = \frac{A\Omega_d \bar{p}_s}{\pi \sqrt{2\pi mkT}} \quad (8.2)$$

From (8.2), we have:

$$\frac{A^2 \bar{p}_s}{\pi L^2 \sqrt{2\pi mkT}} = \phi_d A = \frac{A \bar{p}_d}{\sqrt{2\pi mkT}} \quad (8.3)$$

Or:

$$\bar{p}_d = \bar{p}_s \frac{A}{\pi L^2} = 2.4 \times 10^{-8} \text{ mm of Hg.} \quad (8.4)$$

Problem 7.27

Answer:

The rate of change of the number of particles inside the container is

$$\frac{dN}{dt} = -\frac{1}{4}n\bar{v}A = -\frac{NA}{4V}\sqrt{\frac{8}{T}}\frac{kT}{m} = -\frac{\lambda N}{\sqrt{m}} \quad (9.1)$$

thus defining λ . since pressure is proportional to the number of particles, we find after integrating:

$$p/p_0 = \exp\left[-\frac{\lambda t}{\sqrt{m}}\right] \quad (9.2)$$

For Helium gas $p/p_0 = 1/2$ at $t = 1$ hour. Substituting we have:

$$\lambda = \sqrt{\frac{m}{\text{He}}} \ln 2 \quad (9.3)$$

So:

$$n_{\text{Ne}}/n_{\text{He}} = 2^{(1-\sqrt{m_{\text{He}}/m_{\text{Ne}}})} = 2^{(1-\sqrt{\mu_{\text{He}}/\mu_{\text{Ne}}})} \quad (9.4)$$

Problem 7.29

Answer:

(a)

Inside the container $\bar{v}_z = 0$ by symmetry.

(b)

The velocity distribution, $\phi(\bar{v})$, of the molecules which have effused into the vacuum is:

$$\Phi(\bar{v})d^3\bar{v} = f(\tilde{v})v_z d^3\bar{v} \quad (10.1)$$

where $f(\tilde{v})$ is Maxwell distribution.

So:

$$\bar{v}_z = \frac{\frac{\sqrt{\pi}}{4} \left(\frac{2k\pi}{m}\right)^{\frac{3}{2}}}{\frac{1}{2} \left(\frac{2k\pi}{m}\right)} = \sqrt{\frac{\pi}{2}} \frac{kT}{m} \quad (10.2)$$