Chapter 5: Simple applications of macroscopic thermodynamics

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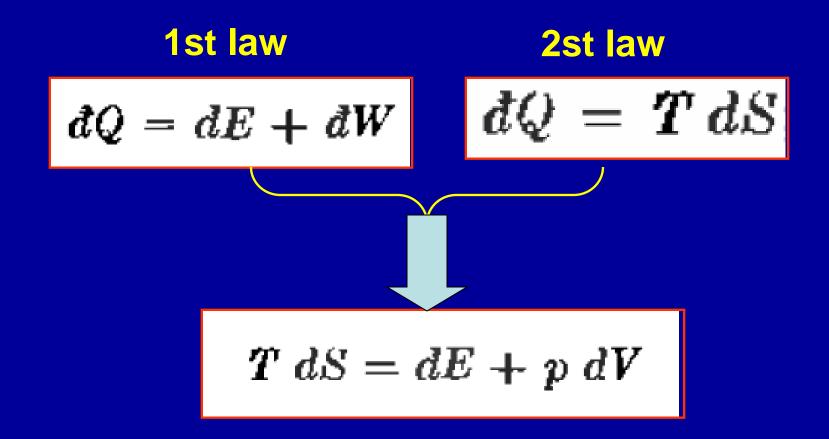
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- Important relationships between macroscopic quantities
- ➤ Based on the fundamental thermodynamic laws (0th, 1st, 2nd, 3rd)
- Quantities such as p, E always are mean values

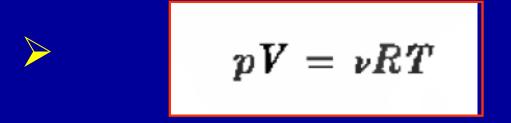
Independent variables:

- Most of systems only have the external para V;
- two independent variables can specify the system, such as E and V; or E and P; or T and V... Other variables can be determined by the two independent variables;

Fundamental equations:



5.1 Equation of state and internal energy



$$\Omega \propto V^{N} \chi(E)$$

$$\beta = \frac{\partial \ln \chi(E)}{\partial E}$$

$$\beta = \beta(E)$$



5.1 Equation of state and internal energy

Generally,

$$E = E(T,V)$$

$$dE = \left(\frac{\partial E}{\partial T}\right)_{V} dT + \left(\frac{\partial E}{\partial V}\right)_{T} dV$$

Fundamentally,

$$T dS = dQ = dE + p dV$$

$$dS = \frac{1}{T} dE + \frac{\nu R}{V} dV$$

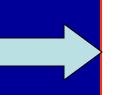
5.1 Equation of state and internal energy

$$dS = \frac{1}{T} \left(\frac{\partial E}{\partial T} \right)_{V} dT + \left[\frac{1}{T} \left(\frac{\partial E}{\partial V} \right)_{T} + \frac{\nu R}{V} \right] dV$$

Because
$$S = S(T, V)$$

$$dS = \left(\frac{\partial S}{\partial T}\right)_{V} dT + \left(\frac{\partial S}{\partial V}\right)_{T} dV$$

5.1 Equation of state and internal energy



$$\begin{pmatrix} \frac{\partial S}{\partial T} \rangle_{V} = \frac{1}{T} \left(\frac{\partial E}{\partial T} \right)_{V} \\
\left(\frac{\partial S}{\partial V} \right)_{T} = \frac{1}{T} \left(\frac{\partial E}{\partial V} \right)_{T} + \frac{\nu R}{V}$$

$$E = E(T)$$



$$\left(\frac{\partial E}{\partial V}\right)_T = 0$$

Because

$$\frac{\partial^2 S}{\partial V \, \partial T} = \frac{\partial^2 S}{\partial T \, \partial V}$$

Thus
$$\left(\frac{1}{T}\left(\frac{\partial^2 E}{\partial V \ \partial T}\right)\right) = \left[-\frac{1}{T^2}\left(\frac{\partial E}{\partial V}\right)_T + \frac{1}{T}\left(\frac{\partial^2 E}{\partial T \ \partial V}\right)\right] + 0$$

5.2 Specific heats

1st law

$$dQ = dE + p dV$$

At constant volume

$$dQ = dE$$

$$c_V \equiv rac{1}{
u} \left(rac{dQ}{dT}
ight)_V = rac{1}{
u} \left(rac{\partial E}{\partial T}
ight)_V$$

$$dE = \left(\frac{\partial E}{\partial T}\right)_V dT$$



 $dE = \nu c_V dT$

5.2 Specific heats 1st law

$$dQ = \nu c \nu dT + p dV$$

For ideal gas at constant pressure

$$p dV = \nu R dT$$

Equation of state

$$dQ = \nu c_V dT + \nu R dT$$

$$c_{p} = \frac{1}{\nu} \left(\frac{dQ}{dT} \right)_{p}$$

$$c_p = c_V + R$$

5.2 Specific heats

$$\gamma \equiv \frac{c_p}{c_V} = 1 + \frac{R}{c_V}$$

Table 5 · 2 · 1	Specific heats of some gases (at 15°C and 1 atm)*				
Gas	Symbol	cv (experimental) (joules mole ⁻¹ deg ⁻¹)	γ (experimental)	(computed by $(5 \cdot 2 \cdot 9)$)	
Hehum Argon Nitrogen Oxygen Carbon dioxide Ethane	He Ar N ₂ O ₂ CO ₂ C ₂ H ₆	12 5 12.5 20.6 21 1 28 2 39.3	1.666 1.666 1.405 1.396 1.302 1.220	1.666 1.666 1.407 1.397 1.298 1.214	

^{*} Experimental values taken from J. R. Partington and W. G. Shilling, "The Specific Heats of Gases," p. 201, Benn, London, 1924

Microscopic calculations of specific heats

For monatomic gas

$$\Omega(E,V) = BV^N E^{8N/2}$$

$$\ln \Omega = \ln B + N \ln V + \frac{3N}{2} \ln E$$

$$eta = \frac{\partial \ln \Omega}{\partial E} = \frac{3N}{2} \frac{1}{E}$$
 $E = \frac{3N}{2\beta} = \frac{3N}{2} kT$

Relation between internal energy and T

Microscopic calculations of specific heats

$$u$$
 is the molar, $E=rac{3}{2}
u(N_ak)T=rac{3}{2}
u RT$

Thus,

$$c_{V} = \frac{1}{\nu} \left(\frac{\partial E}{\partial T} \right)_{V} = \frac{3}{2} R$$

$$c_V = 12.47 \text{ joules deg}^{-1} \text{ mole}^{-1}$$

$$c_p = \frac{3}{2}R + R = \frac{5}{2}R$$

$$\gamma \equiv \frac{c_p}{c_v} = \frac{5}{3} = 1.667$$

Table 5 · 2 · 1	Specific heats of some gases (at 15°C and 1 atm)*				
Gas	Symbol	cv (experimental) (joules mole ⁻¹ deg ⁻¹)	γ (experimental)	(computed by $(5 \cdot 2 \cdot 9)$)	
Hehum	He	12 5	1.666	1.666	
Argon	Ar	12.5	1.666	1.666	
_				1 407	
Nitrogen	N ₂	20.6	1.405		
Oxygen	Oz	21 1	1.396	1.397	
Carbon dioxide	CO_2	2 8 2	1.302	1.298	
Ethane	C_2H_6	39.3	1.220	1.214	

^{*} Experimental values taken from J. R. Partington and W. G. Shilling, "The Specific Heats of Gases," p. 201, Benn, London, 1924

5.3 Adiabatic expansion or compression

Keep a heat contact with a heat reservoir. The gas is allowed to expansion under isothermal. Thus,

$$pV = constant$$

If the gas is thermally insulted,

$$dQ = 0.$$
 $0 = \nu c_V dT + p dV$

Equation of state $p dV + V dp = \nu R dT$

dT is replaced by dV in the former equation

5.3 Adiabatic expansion or compression

$$0 = \frac{c_V}{R} (p \, dV + V \, dp) + p \, dV = \left(\frac{c_V}{R} + 1\right) p \, dV + \frac{c_V}{R} \, V \, dp$$

$$(c_V + R) p \, dV + c_V V \, dp = 0$$

$$\gamma \, \frac{dV}{V} + \frac{dp}{p} = 0$$

$$\gamma = \frac{c_V + R}{c_V} = \frac{c_p}{c_V} \quad function \ of \ T$$

 $\gamma \frac{dV}{V} + \frac{dp}{p} = 0$ γ is T-independent for most gas, and may be a slowly varying

$$\gamma \ln V + \ln p = \text{constant}$$

$$pV^{\gamma} = \text{constant}$$

5.3 Adiabatic expansion or compression

$$pV^{\gamma} = \text{constant}$$

•

$$pV = constant$$

For isothermal process

For adiabatic process

 $\gamma > 1$ p varies more rapidly in adiabatic process

since $p = \nu RT/V$, Can be V, T; or P, T

$$V^{\gamma-1}T = \text{constant}$$

$$dQ = dE + p dV \quad dE = \nu c_V dT$$

$$dE = \nu c_V dT$$

$$T dS = \nu c_V(T) dT + \frac{\nu RT}{V} dV$$

$$dS = \nu c_V(T) \frac{dT}{T} + \nu R \frac{dV}{V}$$

From (T_0, V_0) to (T, V), through quasi-static process

$$S(T,V;\nu) - \nu \varepsilon_0 = \nu \int_{T_0}^T \frac{c_V(T')}{T'} \frac{dT'}{T'} + \nu R \int_{V_0(\nu/\nu_0)}^V \frac{dV'}{V'}$$

5.4 entropy

The last integration is immediate:

$$\int_{V_0(\nu/\nu_0)}^{V} \frac{dV'}{V''} = [\ln V']_{V_0(\nu/\nu_0)}^{V} = \ln V - \ln \left(V_0 \frac{\nu}{\nu_0}\right) = \ln \frac{V}{\nu} - \ln \frac{V_0}{\nu_0}$$

Hence $(5 \cdot 4 \cdot 2)$ becomes

$$S(T,V;\nu) = \nu \left[\int_{T_0}^T \frac{c_V(T')}{T'} dT' + R \ln \frac{V}{\nu} - R \ln \frac{V_0}{\nu_0} + s_0 \right]$$
 (5-4.3)

or $S(T,V;\nu) = \nu \left[\int \frac{c_V(T')}{T'} dT' + R \ln V - R \ln \nu + \text{constant} \right]$

if
$$c_V$$
 is constant.
$$\int \frac{c_V}{T'} dT' = c_V \ln T$$

5.5 Derivation of general relations

Start from
$$dQ = T dS = dE + p dV$$

* Independent variables: S and V

$$dE = T dS - p dV$$

$$E = E(S, V)$$

$$dE = \left(\frac{\partial E}{\partial S}\right)_{V} dS + \left(\frac{\partial E}{\partial V}\right)_{S} dV$$

Combining,

$$\begin{pmatrix} \frac{\partial E}{\partial S} \end{pmatrix}_{V} = T$$

$$\begin{pmatrix} \frac{\partial E}{\partial V} \end{pmatrix}_{S} = -p$$

5.5 Derivation of general relations

Start from

$$dQ = T dS = dE + p dV$$

* Independent variables: S and V

$$\frac{\partial^2 E}{\partial V \partial S} = \frac{\partial^2 E}{\partial S \partial V}$$

$$\left(\frac{\partial}{\partial V}\right)_{\mathcal{S}} \left(\frac{\partial E}{\partial S}\right)_{\mathcal{V}} = \left(\frac{\partial}{\partial S}\right)_{\mathcal{V}} \left(\frac{\partial E}{\partial V}\right)_{\mathcal{S}}$$

$$\left(\frac{\partial E}{\partial S}\right)_{V} = T$$

$$\left(\frac{\partial E}{\partial V}\right)_{S} = -p$$

$$\left(\frac{\partial T}{\partial V}\right)_{\mathcal{S}} = -\left(\frac{\partial p}{\partial S}\right)_{V}$$

5.5 Derivation of general relations

Start from
$$dQ = T dS = dE + p dV$$

* Independent variables: S and p

$$p \ dV = d(pV) - V \ dp$$

$$dE = T dS - p dV = T dS - d(pV) + V dp$$

$$d(E + pV) = T dS + V dp$$

Define
$$H \equiv E + pV$$

5.5 Derivation of general relations

* Independent variables: S and p

$$dH = T dS + V dp$$

$$dH = H(S,p)$$

$$dH = \left(\frac{\partial H}{\partial S}\right)_{p} dS + \left(\frac{\partial H}{\partial p}\right)_{B} dp$$

$$\begin{pmatrix} \frac{\partial H}{\partial \bar{S}} \end{pmatrix}_{p} = T$$

$$\begin{pmatrix} \frac{\partial H}{\partial p} \end{pmatrix}_{S} = V$$

$$\frac{\partial^2 H}{\partial p \, \partial S} = \frac{\partial^2 H}{\partial S \, \partial p}$$

$$\left(\frac{\partial T}{\partial p}\right)_{S} = \left(\frac{\partial V}{\partial S}\right)_{p}$$

5.5 Derivation of general relations

* Independent variables: T and V

Define
$$F \equiv E - TS$$

$$\frac{dE = T dS - p dV = d(TS) - S dT - p dV}{dF = -S dT - p dV}$$

$$F = F(T, V)$$

$$dF = \left(\frac{\partial F}{\partial T}\right)_{V} dT + \left(\frac{\partial F}{\partial V}\right)_{T} dV$$

$$\begin{pmatrix} \frac{\partial F}{\partial T} \rangle_{V} = -S \\
\left(\frac{\partial F}{\partial V} \right)_{T} = -p$$

Properties of ideal gas 5.5 Derivation of general relations

* Independent variables: T and V

$$\frac{\partial^2 F}{\partial V \partial T} = \frac{\partial^2 F}{\partial T \partial V}$$

$$\left(\frac{\partial F}{\partial T}\right)_{V} = -S$$

$$\left(\frac{\partial F}{\partial V}\right)_{T} = -p$$

$$\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial p}{\partial T}\right)_V$$

5.5 Derivation of general relations

* Independent variables: T and p

Define

$$G \equiv E - TS + pV$$

$$dE = T dS - p dV = d(TS) - S dT - d(pV) + V dp$$

$$dG = -S dT + V dp$$

$$G = G(T,p)$$

$$dG = \left(\frac{\partial G}{\partial T}\right)_{p} dT + \left(\frac{\partial G}{\partial p}\right)_{T} dp$$

$$\left(\frac{\partial G}{\partial T}\right)_{p} = -S$$

Properties of ideal gas 5.5 Derivation of general relations

* Independent variables: T and p

$$\begin{pmatrix} \frac{\partial G}{\partial T} \end{pmatrix}_{p} = -S \\
\left(\frac{\partial G}{\partial p} \right)_{T} = V$$

$$\frac{\partial^2 G}{\partial p \ \partial T} = \frac{\partial^2 G}{\partial T \ \partial p}$$

$$-\left(\frac{\partial S}{\partial p}\right)_{T} = \left(\frac{\partial V}{\partial T}\right)_{p}$$

5.6 Summary of Maxwell relations and thermodynamic functions

Maxwell relations

$$dE = T dS - p dY$$

$$\begin{pmatrix} \frac{\partial T}{\partial V} \end{pmatrix}_{S} = -\left(\frac{\partial p}{\partial S}\right)_{V} \\
\begin{pmatrix} \frac{\partial T}{\partial p} \end{pmatrix}_{S} = \left(\frac{\partial V}{\partial S}\right)_{p} \\
\begin{pmatrix} \frac{\partial S}{\partial V} \end{pmatrix}_{T} = \left(\frac{\partial p}{\partial T}\right)_{V} \\
\begin{pmatrix} \frac{\partial S}{\partial p} \end{pmatrix}_{T} = -\left(\frac{\partial V}{\partial T}\right)_{p}
\end{pmatrix}$$

T, p, V, S are not completely in-dependent

???

One can use Ω to give a complete macroscopic description on system

5.6 Summary of Maxwell relations and thermodynamic functions

Maxwell relations

??? One can use Ω to give a complete macroscopic description on system

$$\beta = \frac{\partial \ln \Omega}{\partial E} \begin{cases} \frac{\partial S}{\partial E} \\ \frac{\partial S}{\partial E} \\ \frac{\partial S}{\partial E} \\ \frac{\partial E}{\partial E} \\ \frac{\partial E}{\partial$$

Number of states

generalized force

Properties of ideal gas 5.7 thermodynamic functions

$$E$$
 $H \equiv E + pV$
 $F \equiv E - TS$
 $G \equiv E - TS + pV$
 $E = E(S,V)$
 $H = H(S,p)$
 $F = F(T,V)$
 $G = G(T,p)$

$$dE = T dS - p dV$$

$$dH = T dS + V dp$$

$$dF = -S dT - p dV$$

$$dG = -S dT + V dp$$

5.8 Specific heats

at constant volume

$$C_{V} = \left(\frac{dQ}{dT}\right)_{V} = T\left(\frac{\partial S}{\partial T}\right)_{V}$$

at constant pressure

$$C_{p} = \left(\frac{dQ}{dT}\right)_{p} = T\left(\frac{\partial S}{\partial T}\right)_{p}$$



$$dQ = T dS = T \left[\left(\frac{\partial S}{\partial T} \right)_{p} dT + \left(\frac{\partial S}{\partial p} \right)_{T} dp \right]$$

$$dQ = T dS = C_p dT + T \left(\frac{\partial S}{\partial p}\right)_T dp$$

5.8 Specific heats

dp

$$dQ = T dS = C_{p} dT + T \left(\frac{\partial S}{\partial p} \right)_{T} \left[\left(\frac{\partial p}{\partial T} \right)_{V} dT + \left(\frac{\partial p}{\partial V} \right)_{T} dV \right]$$

when V is constant

$$dQ = C_V dT$$

$$C_{V} = T \left(\frac{\partial S}{\partial T} \right)_{V} = C_{P} + T \left(\frac{\partial S}{\partial P} \right)_{T} \left(\frac{\partial P}{\partial T} \right)_{V}$$

$$\left(\frac{\partial S}{\partial p}\right)_{T} = -\left(\frac{\partial V}{\partial T}\right)_{p}$$

The quantities in the right of equation is easily to be measured

5.8 Specific heats

$$lpha \equiv rac{1}{V} igg(rac{\partial V}{\partial T} igg)_{p}$$

 $\alpha \equiv \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_{T}$ volume coefficient of expansion

$$\left(\frac{\partial S}{\partial p}\right)_{T} = -V_{\alpha}$$

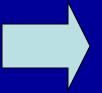
$$dV = \left(\frac{\partial V}{\partial T}\right)_{p} dT + \left(\frac{\partial V}{\partial p}\right)_{T} dp$$

$$dV = \left(\frac{\partial V}{\partial T}\right)_p dT + \left(\frac{\partial V}{\partial p}\right)_T dp \left(\frac{\partial P}{\partial T}\right)_V = -\frac{\left(\frac{\partial V}{\partial T}\right)_p}{\left(\frac{\partial V}{\partial p}\right)_T}$$

5.8 Specific heats
$$C_V = T \left(\frac{\partial S}{\partial T} \right)_V = C_P + T \left(\frac{\partial S}{\partial P} \right)_T \left(\frac{\partial P}{\partial T} \right)_V$$

$$\kappa \equiv -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_T$$

isothermal compressibility



$$\left(\frac{\partial p}{\partial T}\right)_{V} = \frac{\alpha}{\kappa}$$

$$C_V = C_P + T(-V\alpha) \left(\frac{\alpha}{\kappa}\right)$$

$$C_P - C_V = VT \frac{\alpha^2}{\kappa}$$

Class-work

P 192 5.1

Homework

P 192 5.2, 5.3, 5.5, 5.6, 5.7