## On Algorithmic Universality in F-theory Compactifications

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We study universality of geometric gauge sectors in the string landscape in the context of F-theory compactifications. A finite time construction algorithm for  $\frac{4}{3} \times 1.3998 \times 10^{755}$  F-theory geometries is presented. High probability geometric assumptions uncover universal structures in the ensemble without explicitly constructing it. For example, non-Higgsable clusters of seven-branes with intricate gauge sectors occur with probability above  $1 - 1.07 \times 10^{-755}$ , and the geometric gauge group rank is above 160 with probability .999995. In the latter case there are at least 10  $E_8$  factors, the structure of which fixes the gauge groups on certain nearby seven-branes. Visible sectors may arise from  $E_6$  or SU(3) seven-branes, which occur in certain random samples with probability  $\simeq 1/200$ .

I. Introduction. String theory is a consistent theory of quantum gravity that naturally gives rise to interesting gauge and cosmological sectors. As such, it is a promising candidate for a unified theory. However, there is a vast landscape of four-dimensional metastable vacua that may realize different physics, making predictions difficult.

A possible way forward, as in many areas of physics, is to demonstrate universality in large ensembles. For string vacua, such as the oft-quoted  $O(10^{500})$  type IIb flux vacua [1], studying universality via explicit construction is complex and impractical [2]. However, it may be possible to derive universality from a precise construction algorithm, rather than from the constructed ensemble. We refer to this as algorithmic universality, and find it a promising way forward in the string landscape.

We present such an algorithm in the context of 4d F-theory [3] compactifications. The ensemble is a collection of  $4/3 \times 1.3998 \times 10^{755}$  six-manifolds, perhaps the largest set of string geometries to date, that serve as the extra spatial dimensions. Their topological structure determines the 4d gauge group that arises geometrically from configurations of seven-branes that form a network of so-called non-Higgsable clusters (NHC) [4]. We establish that non-Higgsable clusters arise with probability above  $1-1.07\times 10^{-755}$  in this ensemble, and demonstrate that a rich minimal gauge structure arises with high probability. We also present results from random sampling that are potentially relevant for visible sectors.

A number of recent results suggest that NHC are very important in the 4d F-theory landscape. These configurations exist for generic vacuum expectation values of scalar fields (complex structure moduli), and therefore gauge symmetry does not require stabilization on subloci in moduli space [5], which can have high codimension [6]. Standard model structures may arise naturally [5], strong coupling is generic [7], and 4d NHC may exhibit features [8] (such as loops and branches) not present in 6d. NHC arise in the F-theory geometry with the largest number of flux vacua [9], and they arise universally in known ensembles [10] that are closely related to ours. NHC in 6d have been studied extensively [4, 11].

In Sec. II we review non-Higgsable clusters. In Sec. III we present our ensemble. In Sec. IV we exhibit universality. In Sec. V we discuss our results.

II. Seven-Branes and Non-Higgsable Clusters. A 4d F-theory geometry is a Calabi-Yau elliptic fibration X over six extra spatial dimensions described by a complex threefold base space B defined by the equation

$$y^2 = x^3 + fx + g \tag{1}$$

where f and g are polynomials in the coordinates of B; technically,  $f \in \Gamma(\mathcal{O}(-4K_B))$ ,  $g \in \Gamma(\mathcal{O}(-6K_B))$ , with  $K_B$  the canonical class. Seven-branes are localized on the discriminant locus  $\Delta = 4f^3 + 27g^2 = 0 \subset B$ .

Upon compactification the gauge group structure of seven-branes gives rise to four-dimensional gauge sectors. It is controlled by f and g, and for a typical B the most general f,g take the form  $f=\tilde{f}\prod_i x_i^{t_i}, g=\tilde{g}\prod_i x_i^{m_i}$ , so

$$\Delta = \tilde{\Delta} \prod_{i} x_i^{\min(3l_i, 2m_i)} =: \tilde{\Delta} \prod_{i} x_i^{n_i}, \qquad (2)$$

and therefore f, g and  $\Delta$  vanish along  $x_i = 0$  to  $mult_{x_i=0}(f,g,\Delta) = (l_i,m_i,n_i)$ . This seven-brane carries a gauge group  $G_i$  given in Table I according to the Kodaira classification. In some cases further geometric data is necessary to uniquely specify  $G_i$  (see e.g. [10] for conditions) but this data always exists for fixed B. For generic f and g, a seven-brane on  $x_i = 0$  requires  $(l_i, m_i) \geq (1, 1)$ .

Such a seven-brane is called a geometrically non-Higgsable seven-brane (NH7) because it carries a gauge group that cannot be removed by deforming f or g. A NH7 may have geometric gauge group

$$G \in \{E_8, E_7, E_6, F_4, SO(8), SO(7), G_2, SU(3), SU(2)\},\$$

which could be broken by turning on particular fluxes. A typical base B, as we will show in the strongest generality to date, has many non-Higgsable seven-branes that often intersect in pairs, giving rise to jointly charged matter. Such a cluster of seven-branes is a geometrically non-Higgsable cluster (NHC). For brevity, we henceforth drop the adjectives geometric and geometrically.

$F_i$	$l_i$	$m_i$	$n_i$	Sing.	$G_i$
$I_0$	$\geq 0$	$\geq 0$	0	none	none
$I_n$	0	0	$n \ge 2$	$A_{n-1}$	$SU(n)$ or $Sp(\lfloor n/2 \rfloor)$
II	$\geq 1$	1	2	none	none
III	1	$\geq 2$	3	$A_1$	SU(2)
IV	$\geq 2$	2	4	$A_2$	SU(3) or $SU(2)$
$I_0^*$	$\geq 2$	$\geq 3$	6	$D_4$	$SO(8)$ or $SO(7)$ or $G_2$
$I_n^*$	2	3	$n \geq 7$	$D_{n-2}$	SO(2n-4) or $SO(2n-5)$
$IV^*$	$\geq 3$	4	8	$E_6$	$E_6 \text{ or } F_4$
$III^*$	3	$\geq 5$	9	$E_7$	$E_7$
$II^*$	$\geq 4$	5	10	$E_8$	$E_8$

TABLE I. Kodaira fiber  $F_i$ , singularity, and gauge group  $G_i$  on the seven-brane at  $x_i = 0$  for given  $l_i$ ,  $m_i$ , and  $n_i$ .

## III. Large Landscapes of Geometries from Trees.

We now introduce our construction, which utilizes building blocks in toric varieties that we call trees to systematically build up F-theory geometries. After describing the geometric setup and defining terms that simplify the discussion, we will present a criterion, classify all trees satisfying it, and build the F-theory geometries.

Our construction begins with a smooth weak-Fano toric threefold  $B_i$ , and then builds structure on top. These geometries  $B_i$  are determined by fine regular star triangulations (FRST) of one of the 4319 3d reflexive polytopes [12]; there are an estimated  $O(10^{15})$  such geometries [6]. The 2d faces of the 3d polytope are known as facets, and a triangulated polytope will have triangulated facets. Such  $B_i$  do not support NHC.

Consider such a  $B_i$  determined by an FRST of a 3d reflexive polytope  $\Delta^{\circ}$ , a triangulated facet F in  $\Delta^{\circ}$ , and an edge between two points  $v_1$  and  $v_2$  in F with associated homogeneous coordinates  $x_1$  and  $x_2$ . Since  $v_{1,2}$  are connected by an edge,  $x_1 = x_2 = 0$  defines a Riemann surface (algebraic curve) in  $B_i$ , which can be "blown up" using a new ray  $v_e = v_1 + v_2$  and subdividing cones using standard toric techniques. This is a topological transition that introduces a new ("exceptional") divisor e = 0 in B, where e is the coordinate associated to  $v_e$ . This process can be iterated, for example blowing up along  $e = x_1 = 0$ , which would add a new ray  $v_e + v_1 = 2v_1 + v_2$ .

After a number of iterations the associated toric variety will have a collection of exceptional divisors with associated rays  $v_{e_i} = a_i v_1 + b_i v_2$ , which will appear to have formed a tree above the ground that connects  $v_1$  and  $v_2$  in F. Each  $v_{e_i}$  is a leaf with height  $h_{e_i} = a_i + b_i$ , and we will refer to trees built on edges within F as edge trees. The height of a tree is the height of its highest leaf. As an example,  $\{v_1 + v_2, 2v_1 + v_2, v_1 + 2v_2\}$  appears as



where the  $v_1$  to  $v_2$  line is the edge (ground) in F, dashed

green lines are above the ground, 0 is the origin of  $\Delta^{\circ}$ , and the new rays are labelled by their heights.

Similarly, one can also build face trees by beginning with a face on F, with vertices  $v_1, v_2, v_3$  associated to  $x_1, x_2, x_3$ . Adding  $v_e = v_1 + v_2 + v_3$  and subdividing appropriately blows up the point  $x_1 = x_2 = x_3 = 0$  and produces a new toric variety. Again such blowups can be iterated. This process builds a collection of leaves  $v_{e_i} = a_i v_1 + b_i v_2 + c_i v_3$  with  $a_i, b_i, c_i > 0$  of height  $h_{e_i} = a_i + b_i + c_i$  that comprise a face tree. Face trees are built above the interior of the face due to the strict inequality in the definition. Note if one leaf coefficient was zero the associated leaf would be above an edge of the face, not above the face interior.

Geometries can be systematically constructed by adding a face tree to each face in each triangulated facet of  $\Delta^{\circ}$ , and then an edge tree to each edge. The associated smooth toric threefold B has a collection of rays v, each of which can be written  $v = av_1 + bv_2 + cv_3$  with  $v_i$  3d cone vertices in  $B_i$ . If (a,b,c) = (1,0,0) or some permutation thereof,  $v \in \Delta^{\circ}$  and this height  $h_v = 1$  "leaf" is more appropriately a root, since it is on the ground.

A natural question in systematically building up geometries is whether there is a maximal tree height. For a toric variety B to be an allowed F-theory base it must not have any so-called (4,6) divisors (see Appendix), which we ensure by a simple height criterion proven in Prop. 1:

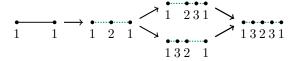
If 
$$h_v \leq 6$$
 for all leaves  $v \in B$ , then there are no  $(4,6)$  divisors.

This condition is simple and sufficient, but not necessary for the absence of (4,6) divisors. Nevertheless, it will allow us to build a large class of geometries.

The task is now clear: we must systematically build all topologically distinct edge trees and face trees of height  $\leq$  6. Since the combinatorics are daunting, let us exemplify the problem for  $h \leq 3$  trees. Viewing the facet head on, an edge in F appears as

$$v_1$$
  $v_2$ 

with the vertices and their heights labelled. Adding  $v_1 + v_2$  subdivides the edge, and further subdivision gives

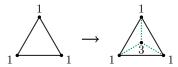


where we have dropped the vertex labels and kept the heights. The trees emerge out of the page, but visualization is made easier by projecting on to the edge; the right-most tree is the one previously presented vertically. There are five edge trees with height < 3. Similarly,

N	# Edge Trees	# Face Trees
3	5	2
4	10	17
5	50	4231
6	82	41,439,964

$h_v$	Probability
3	.99999998
4	.999995
5	.999997
6	.9999898

TABLE II. Left: The number of edge trees and face trees with height  $h \leq N$ . Right: The probability that a face tree with  $h \leq 6$  has a leaf v with a given height  $h_v$ .



shows that there are 2 face trees of height  $\leq 3$ . Here we have denoted the new edges by green lines since they do not sit in the facet. With our definitions, edge trees are built above an edge in the facet, whereas higher leaves in face trees may be built on new edges that do not sit in the facet. For example, a height 4 leaf could be added on any of the green lines above. A (tedious) straightforward calculation shows that the number of edge or face trees with  $h \leq N$  grows rapidly, as in Table II.

Having classified the number of  $h \leq 6$  face trees and edge trees, we now give a lower bound for the number of F-theory geometries that arise from building trees on an FRST of  $\Delta^{\circ}$ . We construct an ensemble  $S_{\Delta^{\circ}}$  of geometries by systematically putting  $h \leq 6$  face trees on all faces  $\tilde{F}$  of  $\Delta^{\circ}$  and then putting  $h \leq 6$  edge trees on all edges  $\tilde{E}$  of  $\Delta^{\circ}$ . Using Table II, the size of  $S_{\Delta^{\circ}}$  is

$$|S_{\Delta^{\circ}}| = \sum_{F} 82^{\#\tilde{E} \text{ on } F} (4.2 \times 10^{6})^{\#\tilde{F} \text{ on } F} - \sum_{F} 82^{\#\tilde{E} \text{ on } E},$$

which may be simplified [13] since

$$\#\tilde{E} = 2n_B + 3n_I - 3 \qquad \#\tilde{F} = n_B + 2n_I - 2 \qquad (3)$$

on a triangulated facet F, where  $n_I$  and  $n_B$  are the number of facet interior and boundary points, respectively.

Two 3d reflexive polytopes give a far larger number  $|S_{\Delta^{\circ}}|$  than the others. They are the convex hulls  $\Delta_i^{\circ} := \operatorname{Conv}(S_i), i = 1, 2$  of the vertex sets

$$S_1 = \{(-1, -1, -1), (-1, -1, 5), (-1, 5, -1), (1, -1, -1)\},\$$
  

$$S_2 = \{(-1, -1, -1), (-1, -1, 11), (-1, 2, -1), (1, -1, -1)\}.$$

After triangulation  $\Delta_1^{\circ}$  and  $\Delta_2^{\circ}$  have the same number of edges and faces. Their largest facets are displayed in Fig. 1 and have  $\#\tilde{E}=63$  and  $\#\tilde{F}=36$ . We compute

$$|S_{\Delta_1^{\circ}}| = \frac{1.3998}{3} \times 10^{755} \qquad |S_{\Delta_2^{\circ}}| = 1.3998 \times 10^{755}, (4)$$

where the factor of 1/3 is due to the symmetries discussed in the Appendix. All other polytopes  $\Delta^{\circ}$  contribute negligibly:  $|S_{\Delta^{\circ}}| \leq 1.65 \times 10^{692}$  configurations. This gives

# 4d F-theory Geometries 
$$\geq \frac{4}{3} \times 1.3998 \times 10^{755}$$
, (5)

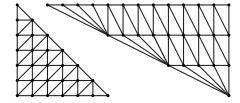


FIG. 1. The largest facets in the two 3d reflexive polytopes  $\Delta_1^{\circ}$  and  $\Delta_2^{\circ}$  with the most number of interior points. Presented is one triangulation of each, from which we see  $\#\tilde{E}=63$  edges and  $\#\tilde{F}=36$  faces in both facets.

which undercounts due to the facts that we choose to do face blowups followed by edge blowups to simplify the subdivision combinatorics, and that we have not taken into account the  $O(10^{15})$  FRS of  $\Delta_2^{\circ}$  and  $\Delta_1^{\circ}$ .

## IV. Universality and Non-Higgsable Clusters.

We now study universality in the dominant sets of F-theory geometries  $S_{\Delta_1^{\circ}}$  and  $S_{\Delta_2^{\circ}}$ . We prove non-Higgsable cluster universality, minimal gauge group universality, and discuss results from random sampling.

Algorithmic Universality and Gauge Groups. We wish to establish the likelihood that an F-theory base in  $S_{\Delta_1^\circ}$  or  $S_{\Delta_2^\circ}$  give rise to non-Higgsable seven-branes. The result arises from Prop. 2: if there is a tree anywhere on F, even a single leaf, there is a non-Higgsable seven-brane on all divisors associated to interior points of F. For any  $S_{\Delta^\circ}$  only one configuration has no trees, and therefore

$$P(\text{NHC in } S_{\Delta^{\circ}}) \ge 1 - \frac{1}{|S_{\Delta^{\circ}}|}.$$
 (6)

This is always very close to one, and in particular

$$P(\text{NHC in } S_{\Delta_1^{\circ}}) \ge 1 - 2.14 \times 10^{-755}$$
  
 $P(\text{NHC in } S_{\Delta_2^{\circ}}) \ge 1 - .71 \times 10^{-755}$ . (7)

We see that NHC are universal in these ensembles.

We now wish to study physics in our ensemble. Consider a geometric assumption  $A_i$  and a physical property  $P_i$  such that  $A_i \Longrightarrow P_i$ . Our goal is to determine high probability assumptions that lead to interesting physical properties, computing  $P(A_i)$  since  $A_i \Longrightarrow P_i$  ensures  $P(P_i) \ge P(A_i)$ . We will focus on  $S_{\Delta_1^\circ}$  and  $S_{\Delta_2^\circ}$  since these dominate the ensemble.

Consider first  $S_{\Delta_1^\circ}$  and let  $A_1$  be the assumption that any simplex in an FRST of  $\Delta_1^\circ$  containing a vertex of  $\Delta_1^\circ$  has an  $h \geq 3$  face tree on it. For the 3 symmetric facets of  $\Delta_1^\circ$  there are 17 ways to choose simplices containing the vertices, and 1796 ways for its largest facet. The maximum number of simplices containing vertices is 24. Using  $P(h \geq 3 \text{ tree on simplex})$  from Table II,

$$P(A_1 \text{ in } S_{\Delta_1^\circ}) \ge .9999998^{24} = .999995.$$
 (8)

There are  $17^3 \times 1796$  ways to choose simplices that contain the vertices, all of which yield  $G \geq F_4^{18} \times E_6^{10} \times U^9$ 

**A.** Appendix: Technical Subtleties. We now address technical subtleties that are important for establishing, but not understanding, results in the main text.

Polytope Symmetries. In equation (4) we have included a factor of 1/3 relative to the count one would obtain directly from the algorithm. This takes into account an overcounting of geometries due to toric equivalences, which arise when there is a  $GL(3,\mathbb{Z})$  transformation on the toric rays that preserves the cone structure of the fan. In general, there may be many such equivalences between elements of two ensembles  $S_{\Delta_i^{\circ}}$  and  $S_{\Delta_i^{\circ}}$ , where  $\Delta_{i,j}^{\circ}$  are any two 3d reflexive polytopes. However, to ensure that the count (4) is accurate, we only need to consider whether there are equivalences between two elements in  $S_{\Delta_1^{\circ}}$ , two in  $S_{\Delta_2^{\circ}}$ , or one in  $S_{\Delta_1^{\circ}}$  to  $S_{\Delta_2^{\circ}}$ . It is sufficient to consider  $GL(3,\mathbb{Z})$  actions on the ground, i.e. on the facets. This follows from the fact that rays of different height cannot be exchanged under automorphisms of the fan. First note that points in a hyperplane remain in a hyperplane after a  $GL(3,\mathbb{Z})$  transformation, and points in a line remain in a line by linearity. Facets must therefore map to facets. The big facets in  $\Delta_1^{\circ}$  and  $\Delta_2^{\circ}$  cannot map to other facets by point counting, and therefore they must map to themselves. There is no nontrivial map taking the big facet in  $\Delta_2^{\circ}$  to itself, but there is a  $\mathbb{Z}_3$  rotation taking the big facet in  $\Delta_1^{\circ}$  to itself, giving a factor of 1/3 in  $S_{\Delta_1^{\circ}}$ . There is no non-trivial map between the big facets in  $\Delta_1^{\circ}$  and  $\Delta_2^{\circ}$ , and therefore  $S_{\Delta_2^{\circ}} \cap S_{\Delta_1^{\circ}} = \emptyset$ . Together, these establish (4).

Multiplicites of Vanishing and Resolutions. In discussing what constitutes an allowed 4d F-theory geometry  $X \to B$ , we mentioned certain criteria on multiplicities of vanishing that we now elaborate on. In [14, 15] it was shown that if a Calabi-Yau variety has at worst canonical singularities, then it is at finite distance from the bulk of the moduli space in the Weil-Petersson metric. This criterion is general and therefor applies to elliptic fibrations such as X. The reason that it is physically relevant is that if X has worse singularities than a nearby Calabi-Yau X' that is known to represent a physical configuration, and X is at finite distance in the moduli space from X', we should expect that X is also a physical configuration. This criterion, which we refer to as the Hayakawa-Wang criterion, gives a related criterion by studying elliptic fibrations [16– 18]: if  $mult_D(f,g) < (4,6)$ ,  $mult_C(f,g) < (8,12)$ , and  $ord_p(f,g) < (12,18)$  for all divisors  $D \subset B$ , curves  $C \subset B$ , and points  $p \subset B$ , respectively, then X has at worst canonical singularities and is at finite distance in the moduli space due to the Hayakawa-Wang criterion<sup>1</sup>. Here "less than" means that at least one of the multiplicities or orders is strictly less than the given multiplicity

or order. We now translate this multiplicity of vanishing (MOV) condition to constraints on the height of the tree.

**Proposition 1.** Suppose each leaf  $v \in B$  has height  $h_v \leq 6$ . Then B has no (4,6) divisors.

Proof. Consider a facet F, which has a unique associated point  $m_F$  satisfying  $(m_F, \tilde{v}) = -1 \ \forall \tilde{v} \in F$ ; furthermore since  $m_F \in \Delta$ ,  $(m_F, v) \geq -1 \ \forall v \in \Delta^{\circ}$ . Now suppose  $h_v \leq 6/n$ ,  $n \in \mathbb{N}$  for all rays  $v = av_1 + bv_2 + cv_3$  in B, with  $v_i$  3d cone vertices in  $B_i$ . Then  $(nm_F, v) \geq -n(a+b+c) = -nh_v \geq -6$  for all rays v and therefore  $nm_F \in \Delta_g$ . Here we denote  $\Delta_f, \Delta_g$  as the polytopes corresponding to  $\Gamma(\mathcal{O}(-4K_B)), \Gamma(\mathcal{O}(-6K_B))$ , respectively. If  $h_v \leq 6 \ \forall v$ , then  $m_F \in \Delta_g$ . This monomial has multiplicity of vanishing  $(v, m_F) + 6 = 5$  for any v in or above F, which protects v from being a (4,6) divisor. If  $h_v \leq 6 \ \forall v$  then  $m_F \in \Delta_g \ \forall F$  and there is a monomial that prevents each divisor from being (4,6).

It is also simple to see that in our ensemble, f and g can only vanish to multiplicities less than (8,12) along curves and orders less than (12,18) at points, respectively. Consider any toric curve  $C = D_s \cdot D_t \subset B$ . Take  $v_s = \sum_i a_{i,s} v_i$  and  $v_t = \sum_i a_{i,t} v_i$  and define  $a := \sum_i a_{i,s}$  and  $b := \sum_i a_{i,t}$ . Let F be a facet on which or above which  $v_s$  and  $v_t$  sit, with  $m_F$  the dual facet. As an element of  $\Delta_g$  the associated monomial may be written  $s^{\langle m,v_s\rangle+6}t^{\langle m,v_t\rangle+6}\times\ldots$ , and the monomial vanishes to multiplicity  $\langle m,v_s\rangle+\langle m,v_t\rangle+12=-a-b+12$  along C. For g to vanish to multiplicity 12 along a curve, this requires a+b<0, which cannot happen. A similar argument shows that g cannot vanish to order 18 or higher at points.

On the other hand, our ensemble is generated by a series of repeated blowups along curves and points, and a crepant resolution only exists if the MOV is  $\geq (4,6)$  for a curve, and  $\geq (8,12)$  for a point. One can achieve the required MOV by tuning in complex structure moduli space, but one has to ensure that no infinite distance singularities (as in the above) are introduced in the process. However, it is simple to see that the desired MOV can be achieved, without introducing any disallowed singularities, by simple tuning without turning off the monomial corresponding to  $m_F$ , for all F.

**7-Branes and Gauge Enhancement.** We now prove some useful results that allow us to determine a universal gauge sector in our ensemble, as well as show that NH 7-branes are ubiquitous.

**Proposition 2.** Suppose that there exists v in or above a facet F, i.e.  $v = av_1 + bv_2 + cv_3$  with  $v_i$  simplex vertices in F, such that  $h_v \geq 2$ . Then there is a non-Higgsable seven-brane on the divisor associated to each interior point of F.

Proof. Then  $(6m_F, v) = -6h_v \le -12$  implies  $6m_F \notin \Delta_g$ . Similarly,  $4m_F \notin \Delta_f$ . Since any point p interior to F has

<sup>&</sup>lt;sup>1</sup> We thank D. Morrison for discussions on this and related points.

 $(m,p)=-1 \iff m=m_F$  and reflexive polytopes of dimension three are normal, i.e. any  $m_f \in \Delta_f$   $(m_f \in \Delta_g)$  has  $m_f = \sum_i m_i, m_i \in \Delta$   $(m_g = \sum_i m_i, m_i \in \Delta)$ , it follows that  $(m_f,p)=-4 \iff m_f = 4m_F$  and  $(m_g,p)=-6 \iff m_g = 6m_F$ . Therefore, if there is any tree on F then  $4m_F \notin \Delta_f$  and  $6m_F \notin \Delta_g$ . By normality, for any p interior to F this implies  $\nexists m_f \in \Delta_f | (m_f,p) = -4$  and  $\nexists m_g \in \Delta_g | (m_g,p) = -6$ , and therefore  $ord_p(f,g) > (0,0)$ , which implies there is a non-Higgsable seven-brane on the divisor associated to p.

**Proposition 3.** Let v be a leaf  $v = av_1 + bv_2 + cv_3$  with  $v_i$  simplex vertices in F. If the associated divisors  $D_{1,2,3}$  carry a non-Higgsable  $E_8$  seven-brane, and if v has height  $h_v = 1, 2, 3, 4, 5, 6$  it also has Kodaira fiber  $F_v = II^*, IV_{ns}^*, I_{0,ns}^*, IV_{ns}, II, -$  and gauge group  $G_v = E_8, F_4, G_2, SU(2), -, -$ , respectively.

Proof. Recall that the height criterion gives  $mult_v(g) \leq 6-h_v$ . If  $v = av_1 + bv_2 + cv_3$  with  $v_i$  each carrying  $E_8$ , then  $(m_f, v_i) \geq 0$ ,  $(m_g, v_i) \geq -1$ ,  $\forall m_f \in \Delta_f$  and  $\forall m_g \in \Delta_g$ . This gives  $(m_f, v) \geq 0$ ,  $(m_g, v) \geq -(a + b + c) = -h_v$ . Together, we see  $mult_v(f) \geq 4$ ,  $mult_v(g) = 6 - h_v$ . For  $h_v = 1, 5, 6$  this fixes  $G_v$ , but to determine  $G_v$  for  $h_v = 2, 3, 4$  we must study the split condition. A necessary condition is that there is one monomial  $m_g \in \Delta_g$  such that  $(m_g, v) + 6 = 6 - h_v$ , and since  $m_F \in \Delta_g$  always, where F is the facet in which  $v_i$  lie, then  $m_g = m_F$ . Morever, the monomial m in g associated to  $m_F$  must be a perfect square; since  $(m_F, v_i) + 6 = 5$ ,  $m \sim x_i^5$  and m is not a perfect square. Therefore the fibers are all non-split. This establishes the result.

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