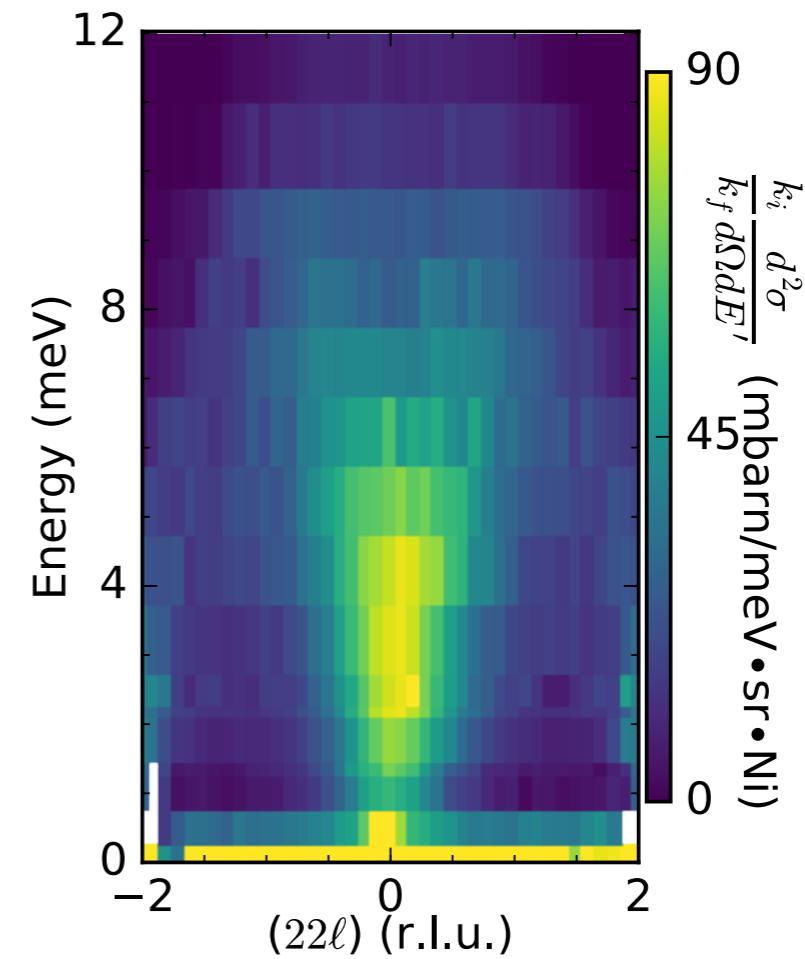
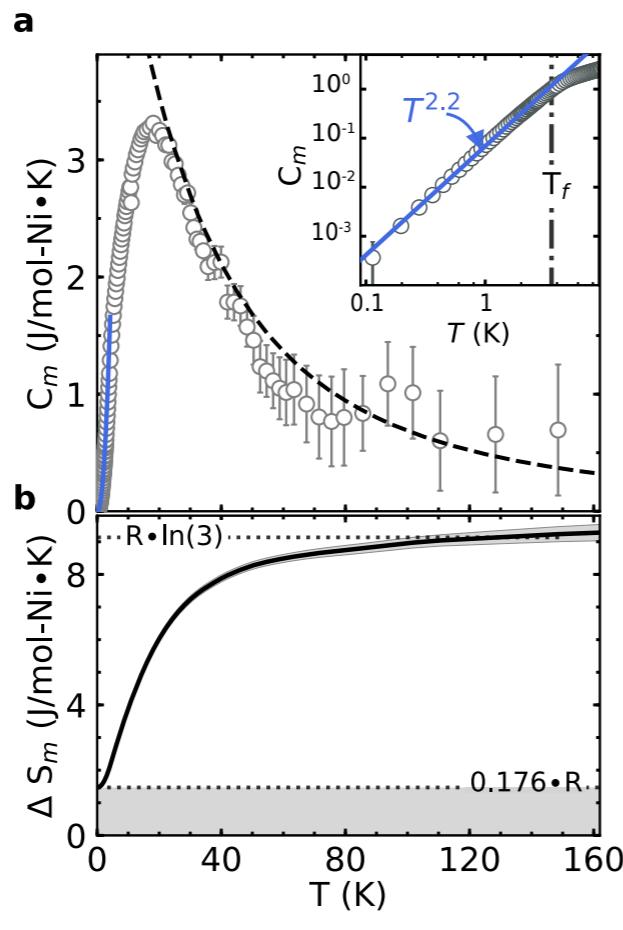


CONTINUUM OF MAGNETIC EXCITATIONS IN $\text{NaCaNi}_2\text{F}_7$

Kemp Plumb
HFM, July 2018



COLLABORATORS



Collin Broholm

Allen Scheie

Hitesh Changlani

Shu Zhang



Bob Cava

Jason Krizan



Jose Rodriguez-Rivera

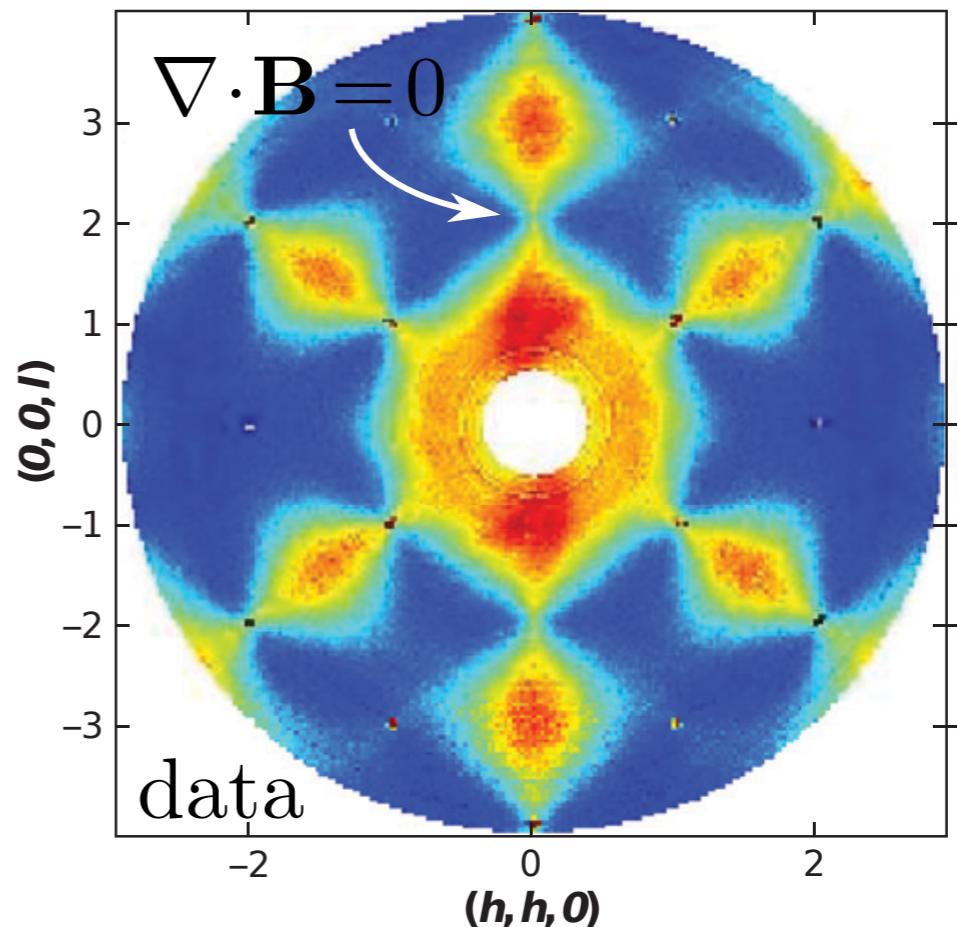
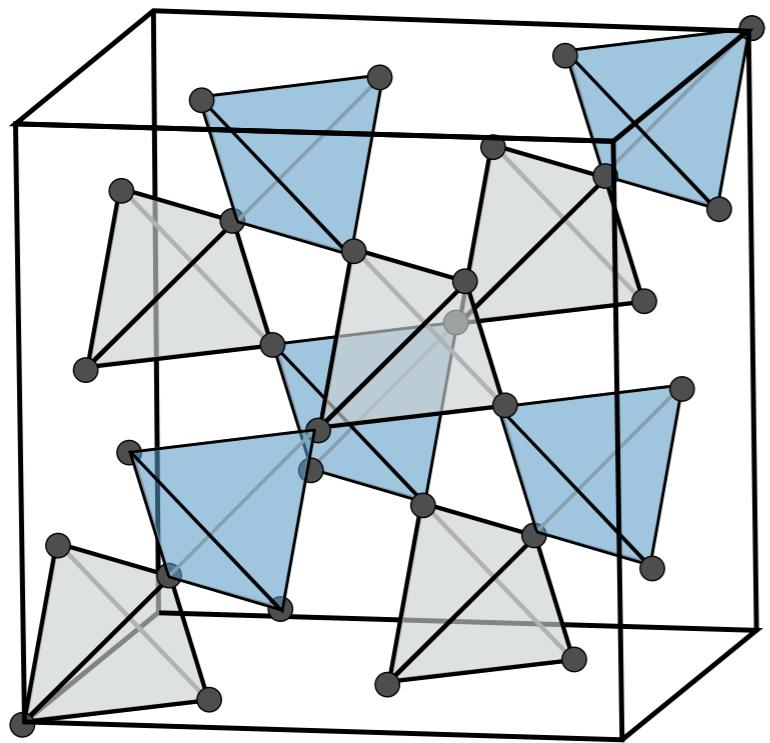
Yiming Qiu



Barry Winn

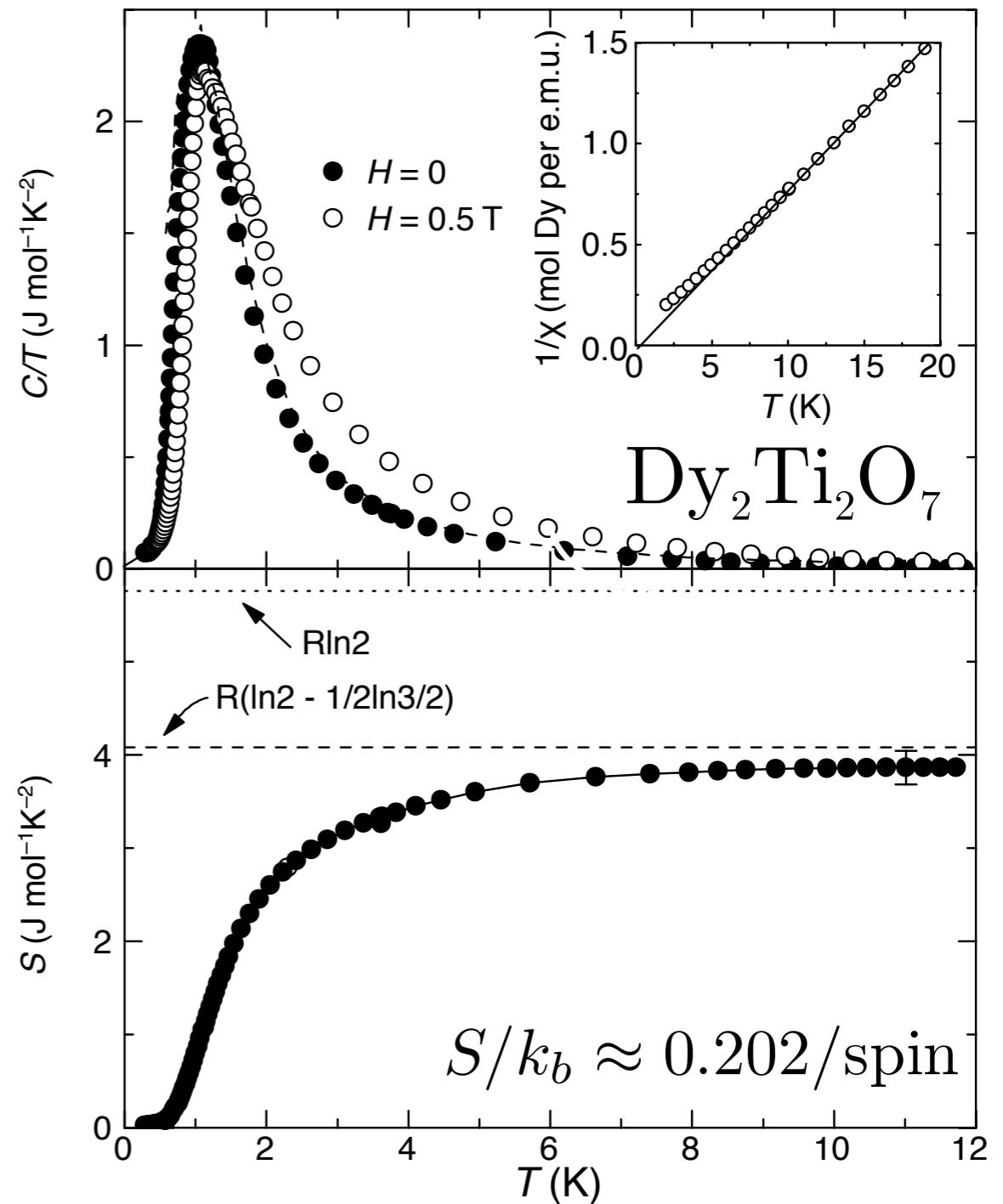
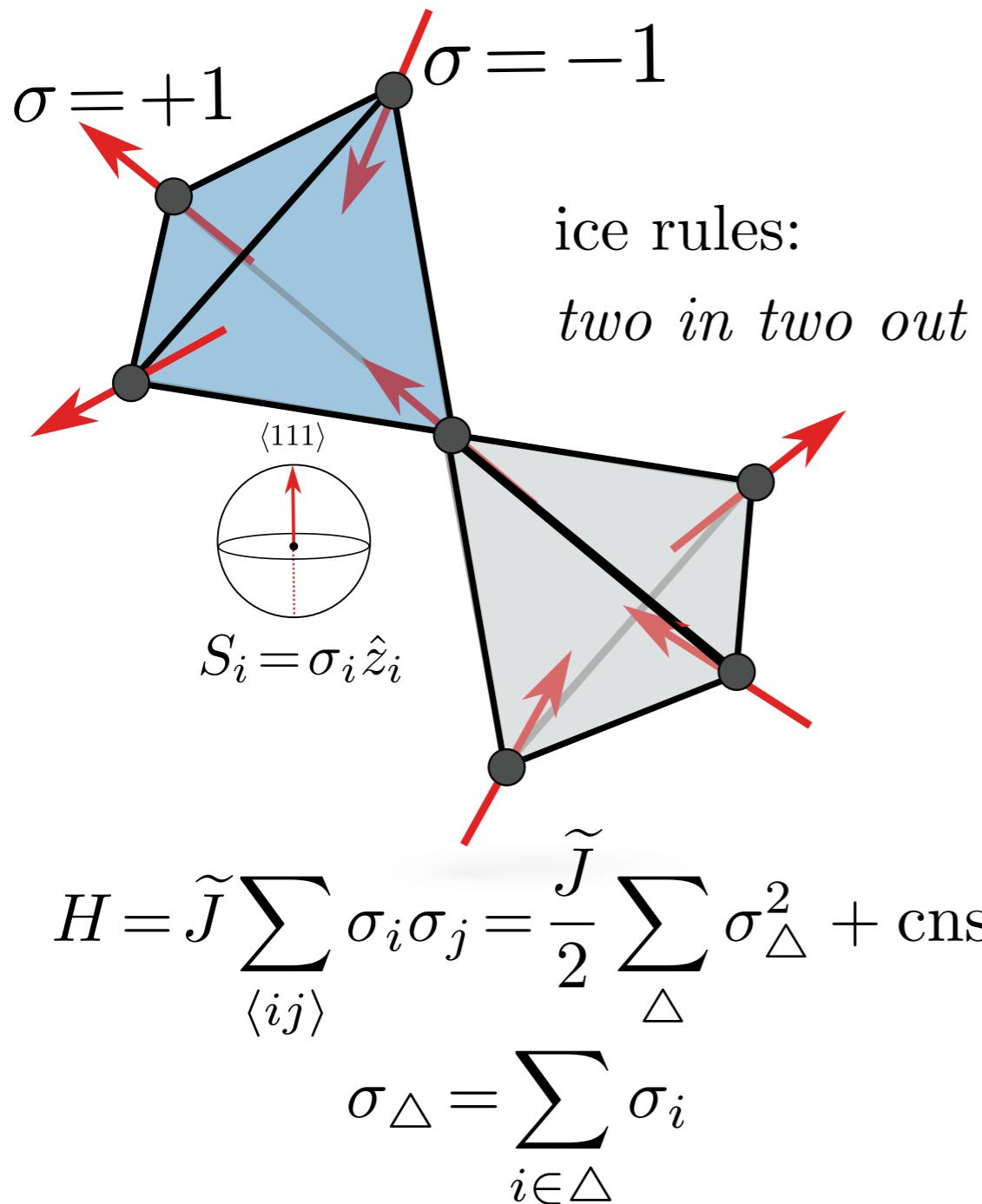
Feng Ye

THE PYROCHLORE LATTICE



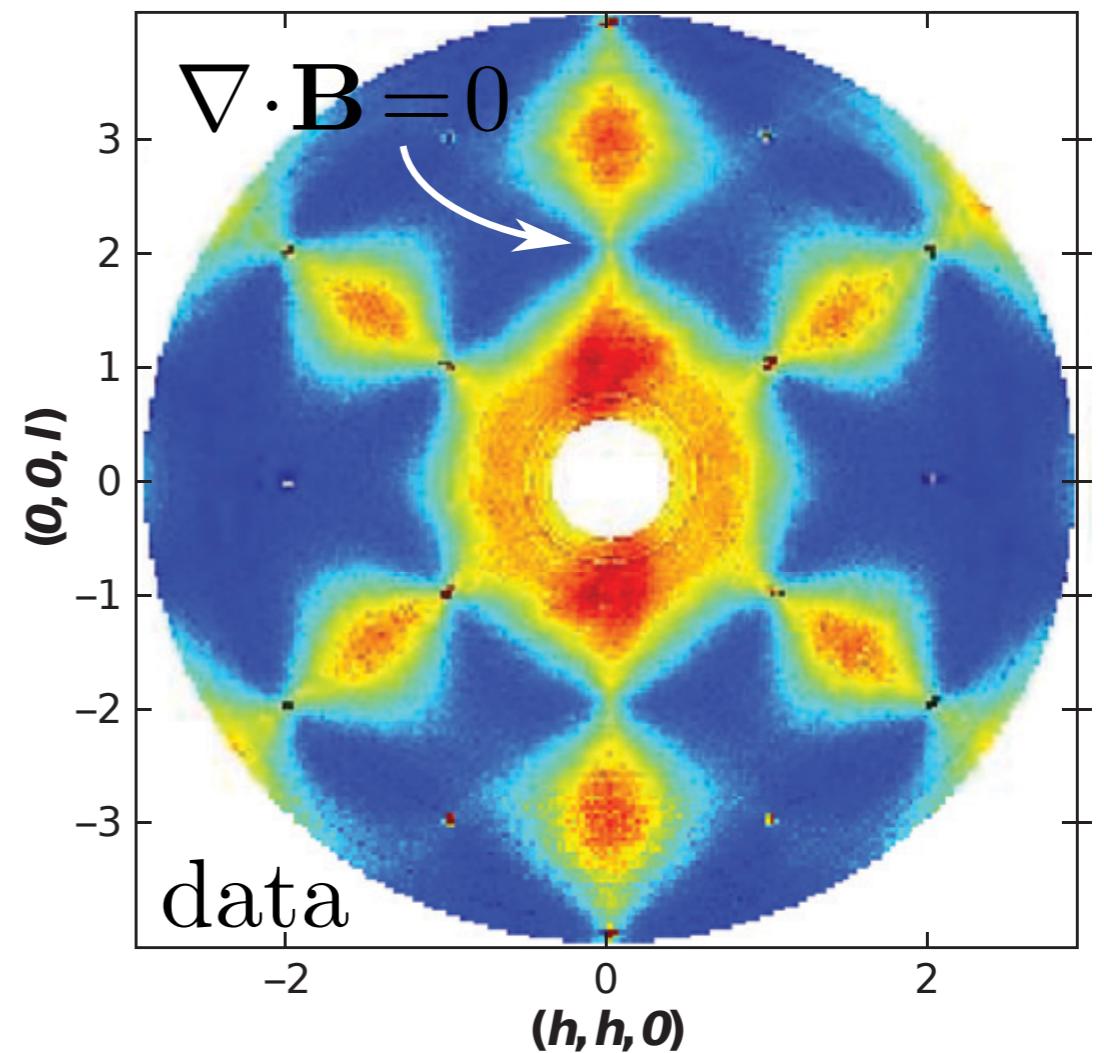
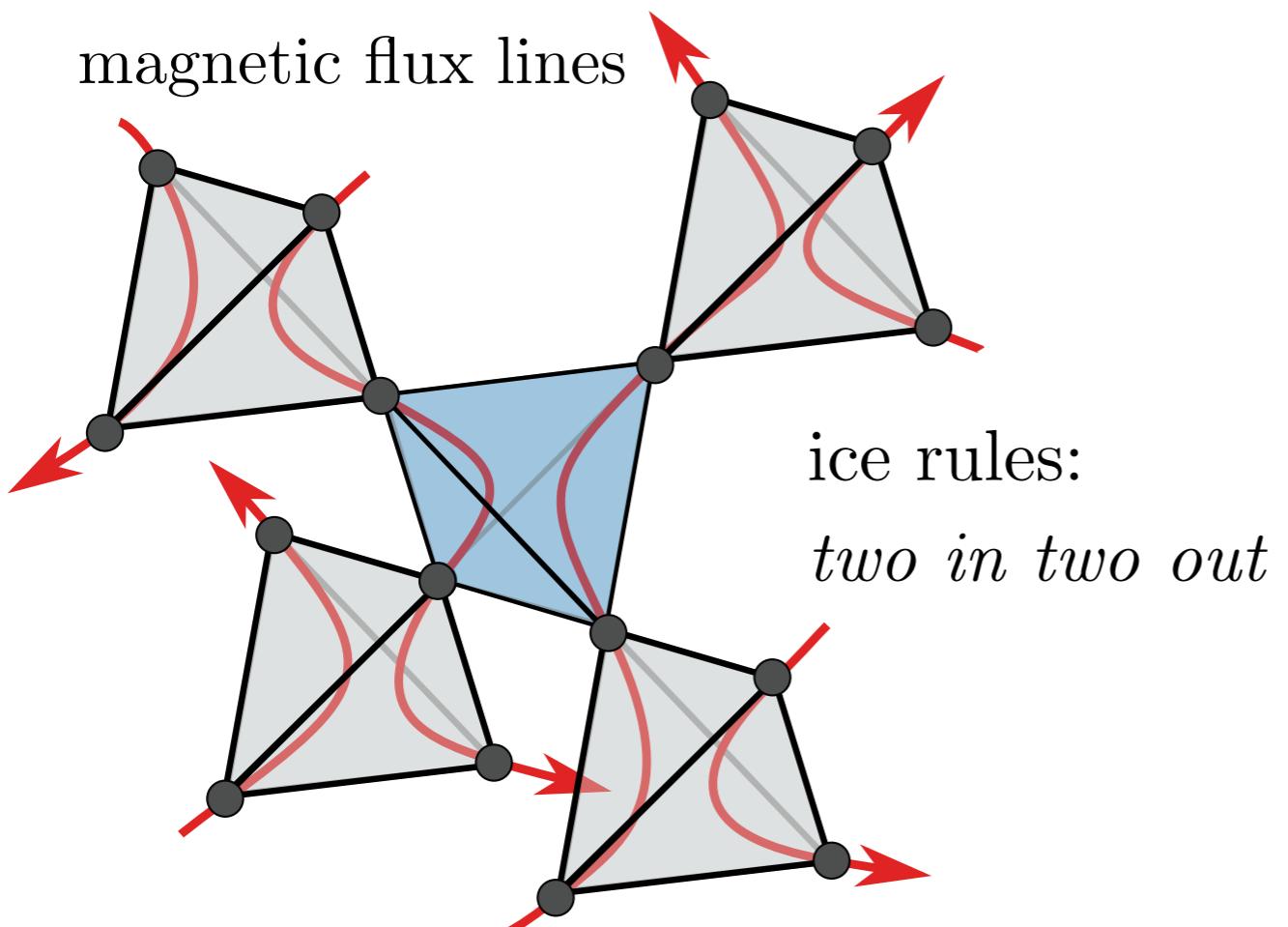
- A paradigm for magnetic frustration: degenerate low energy states subject to local constraints, leads to
 - Spin ice: Ising ferromagnet.
 - Order-by-disorder: xy models.
 - Spin liquid: Heisenberg antiferromagnet. Quantum?

COULOMB PHASE: SPIN-ICE



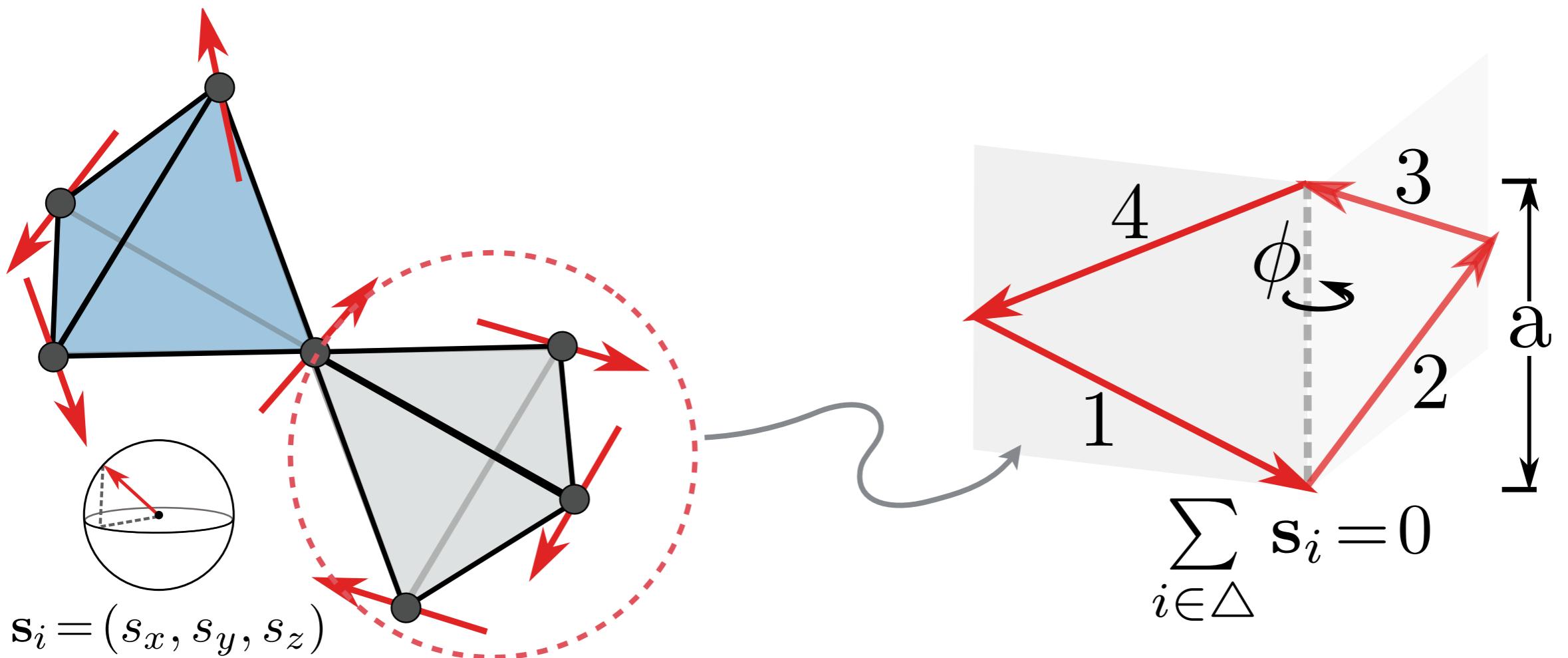
- Lowest energy states have: $\sigma_{\Delta} = 0 \rightarrow$ ice rules.
- Macroscopic degeneracy, $W = (3/2)^{(N/2)}$ states with $\sigma_{\Delta} = 0$.

COULOMB PHASE: SPIN-ICE



- Highly degenerate, like a liquid, but correlated.
- Imagine spins defining flux lines, then $\nabla \cdot \mathbf{B} = 0$.
- Consequence is power law correlations: $\langle \mathbf{B}(0)\mathbf{B}(\mathbf{r}) \rangle \propto \frac{1}{r^3}$.
- Spins are frozen in time.

COULOMB PHASE: CONTINUOUS SPINS



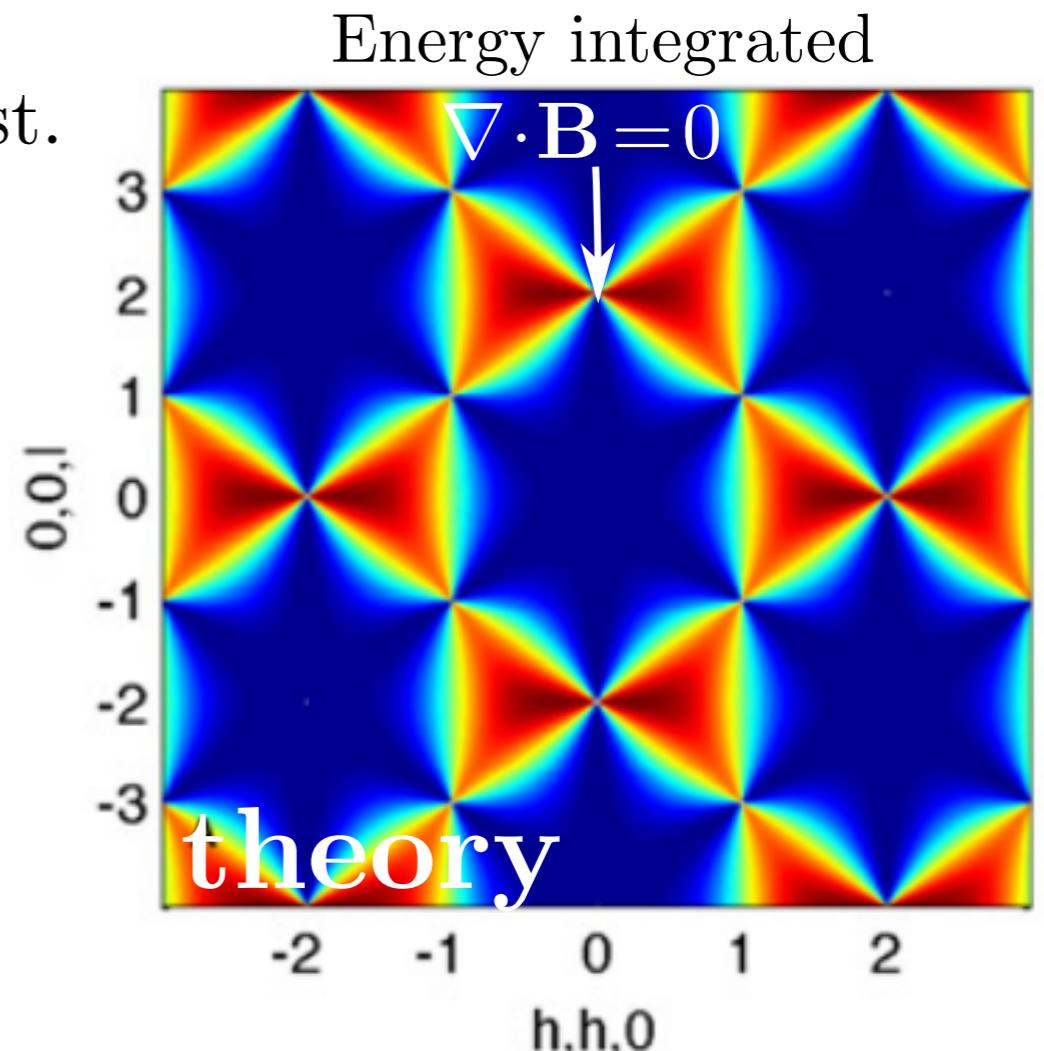
- In lattice: Total degrees of freedom - constraints $\propto N$.
- No internal energy barriers, spins are dynamic.

COULOMB PHASE: CONTINUOUS SPINS

$$H = J \sum_{\langle ij \rangle} \mathbf{s}_i \cdot \mathbf{s}_j = \frac{J}{2} \sum_{\Delta} |\mathbf{L}_{\Delta}|^2 + \text{cnst.}$$

$$L_{\Delta} = \sum_{i \in \Delta} \mathbf{s}_i$$

$$\mathbf{s}_i = (s_x, s_y, s_z)$$

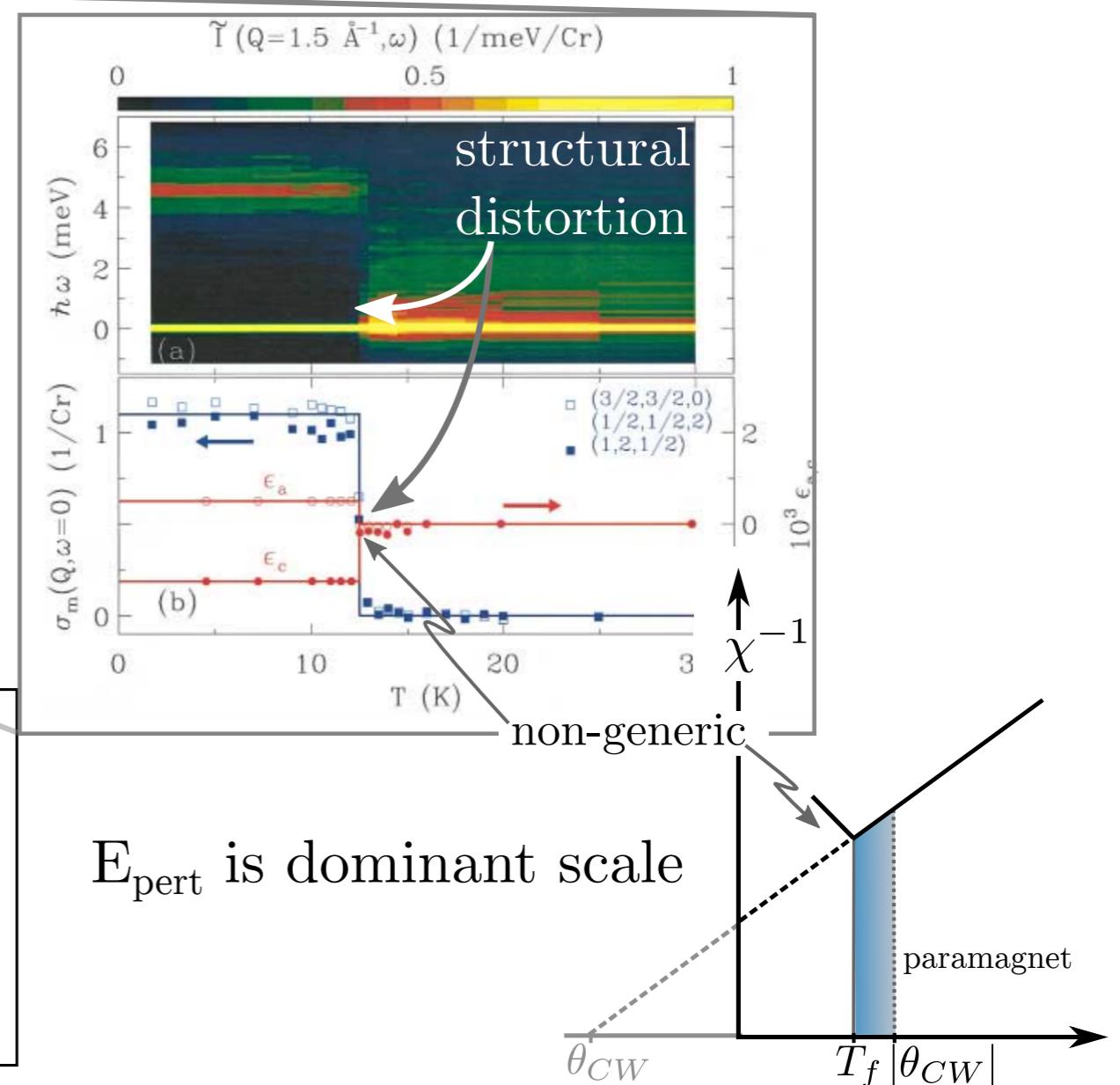
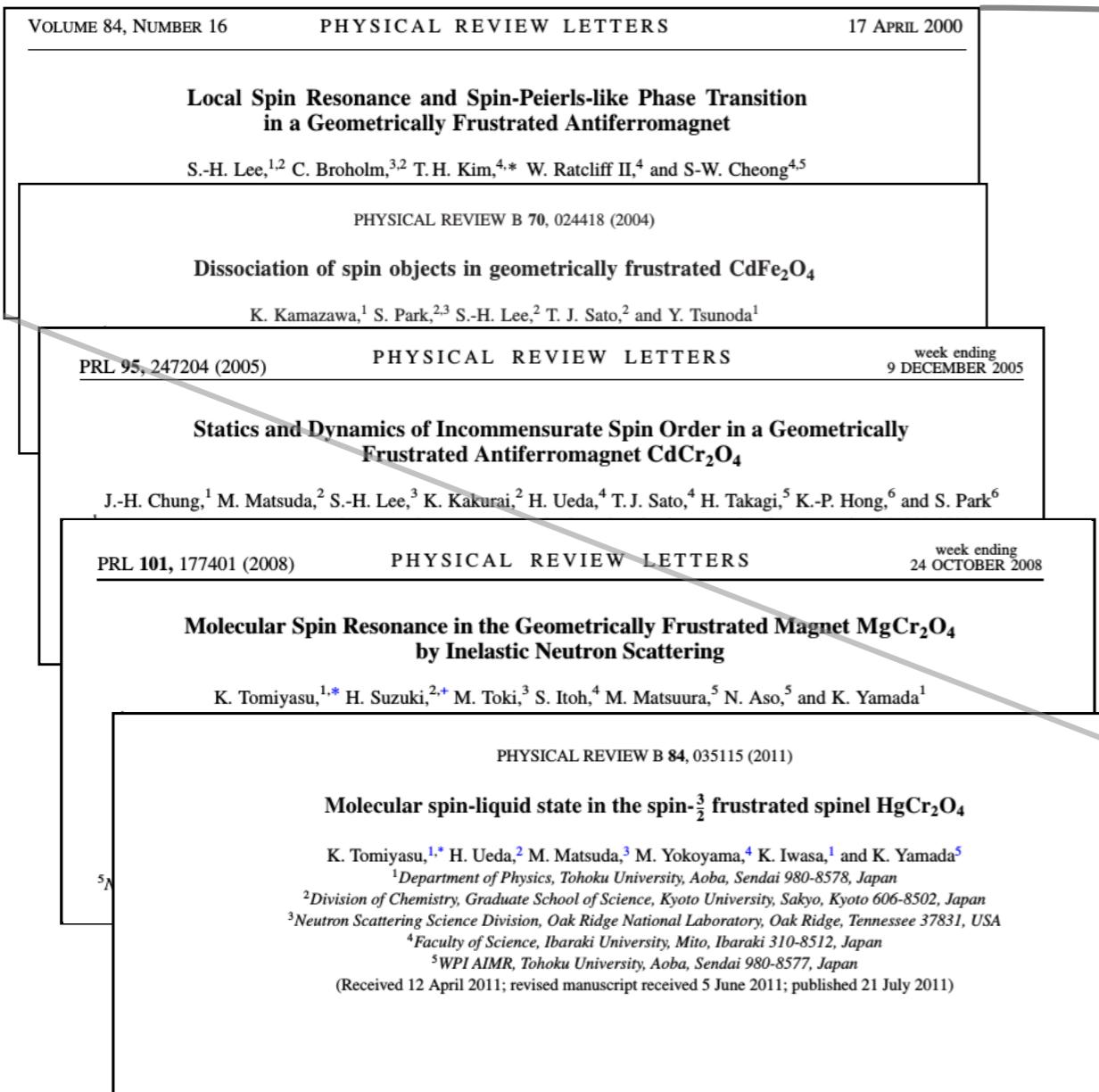


- Can map spin configurations to solenoidal field.
- $\nabla \cdot \mathbf{B} = 0$ implies dipolar correlations $\propto 1/r^3$.
- A (classical) U(1) spin-liquid in three dimensions.
- Unsolved problem in quantum limit.

[†]Henley, PRB (2005)

Villain, Z. Phys. B. (1979)

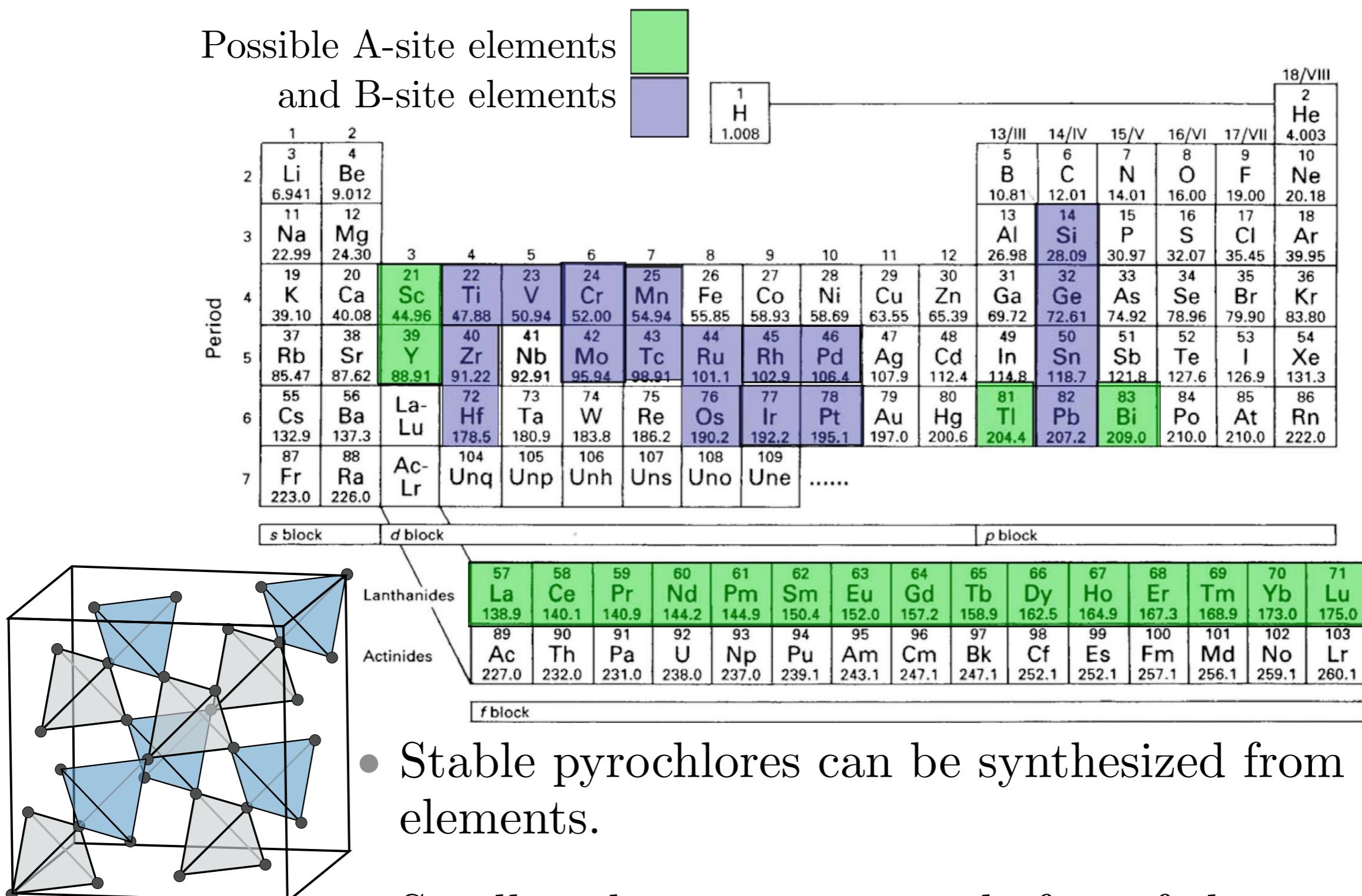
EXPERIMENTS ON HEISENBERG PYROCHLORES



- B-sites of spinel lattice AB₂O₄ have pyrochlore structure.
- In all of these compounds, lattice deformations impact magnetism.
- Perturbations are always there can we access regime where E_p < k_bT < E_J.

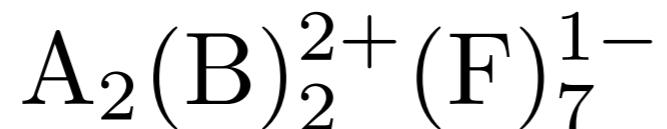
WHY DO WE NEED MORE PYROCHLORES?

Possible A-site elements and B-site elements

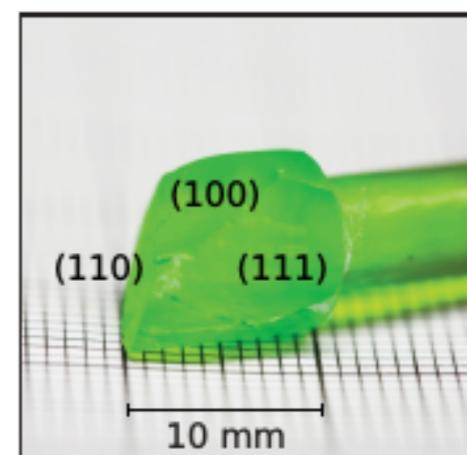


- $\text{A}_2\text{B}_2\overset{\text{O}_7^{2-}}{\text{O}_7}$ • Small exchange energy scale from f-electrons.

WHY NaCaNi₂F₇?

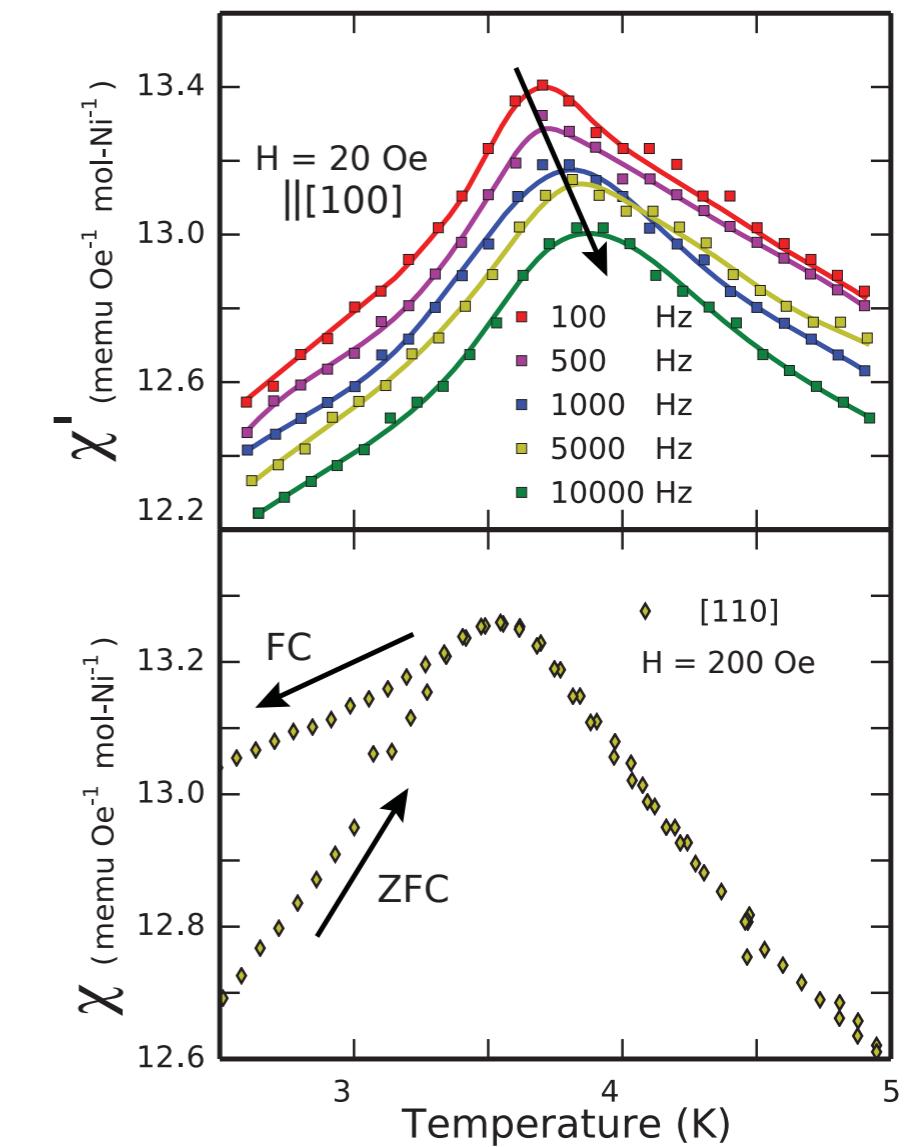
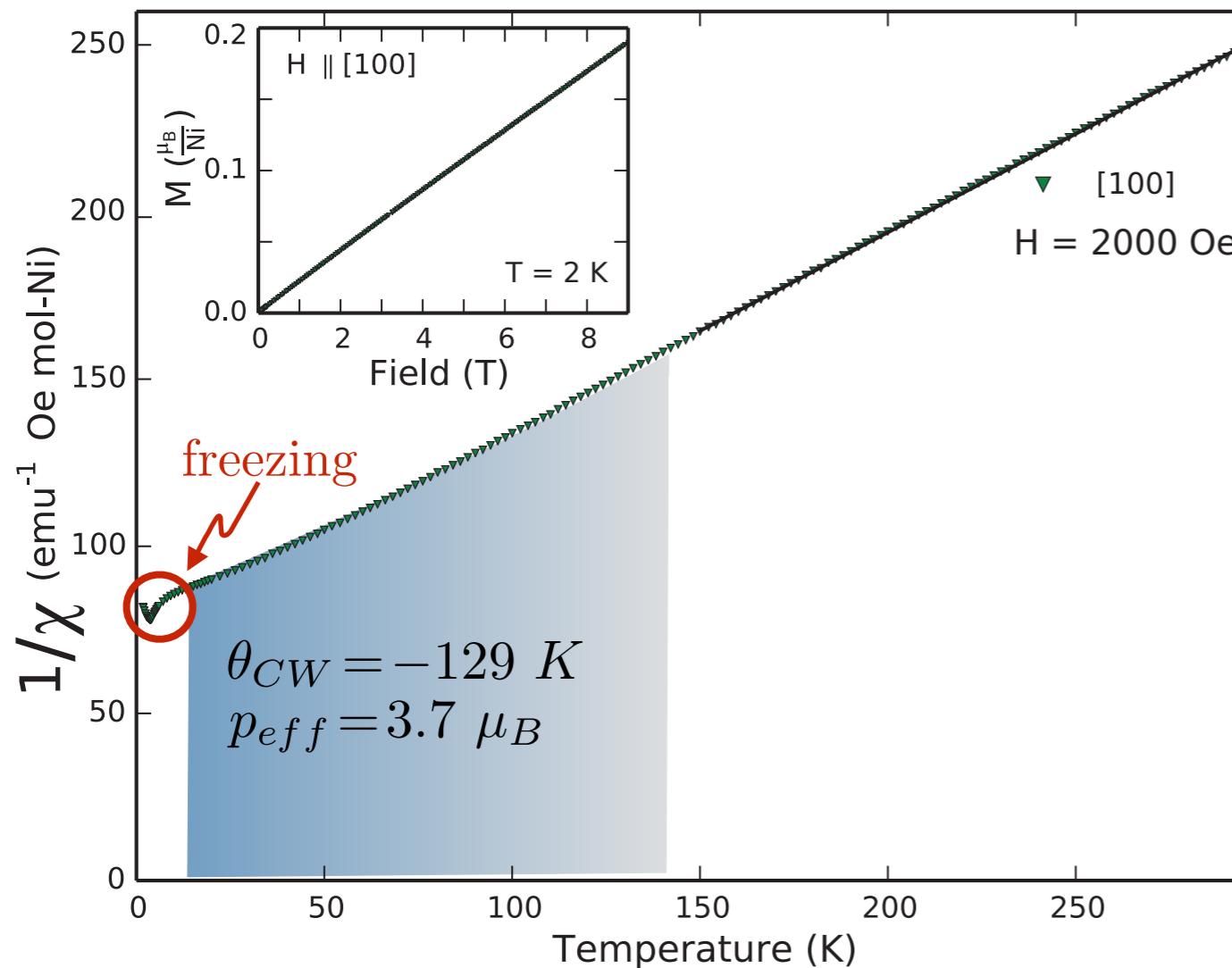


Possible A-site elements and B-site elements



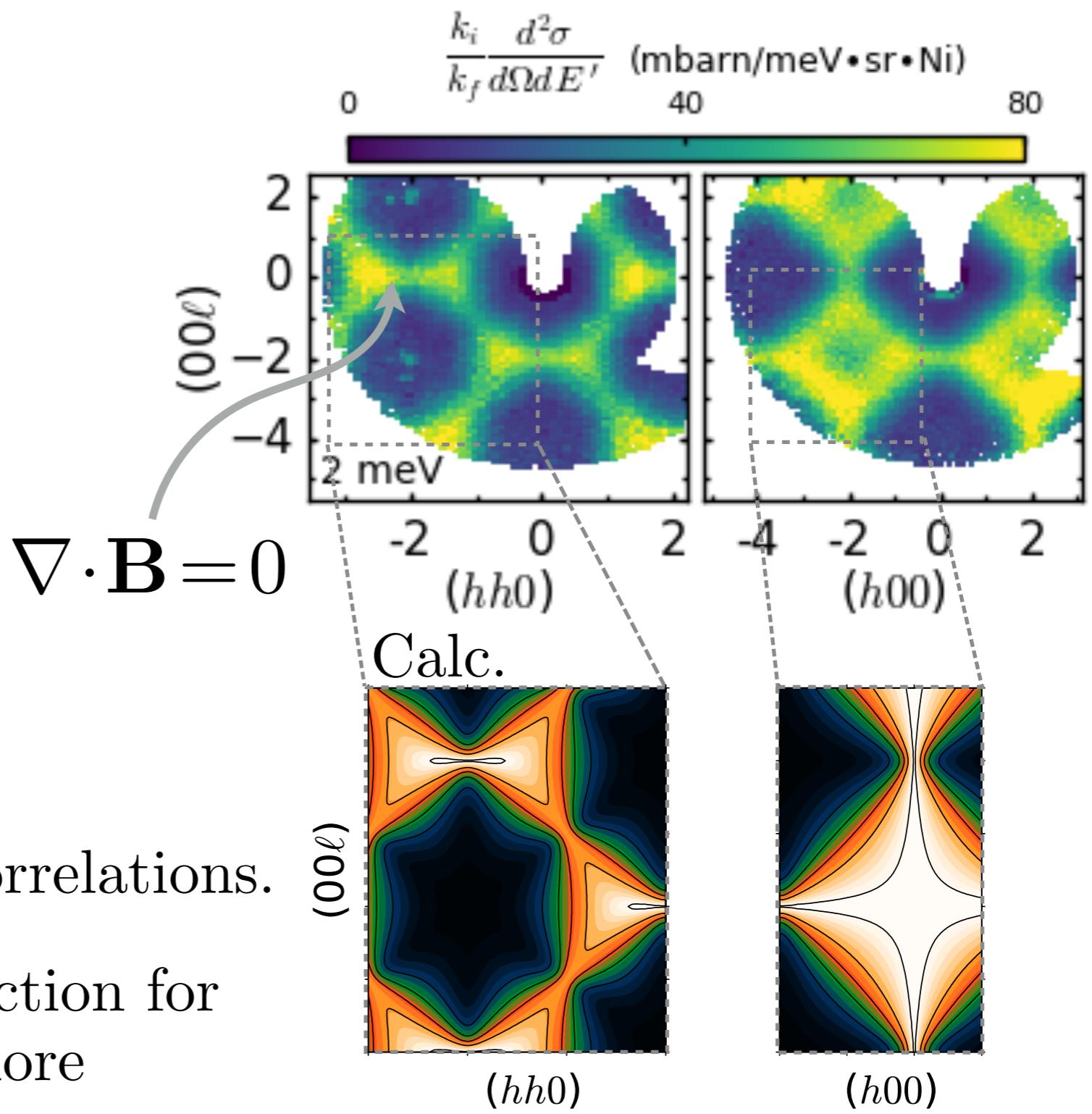
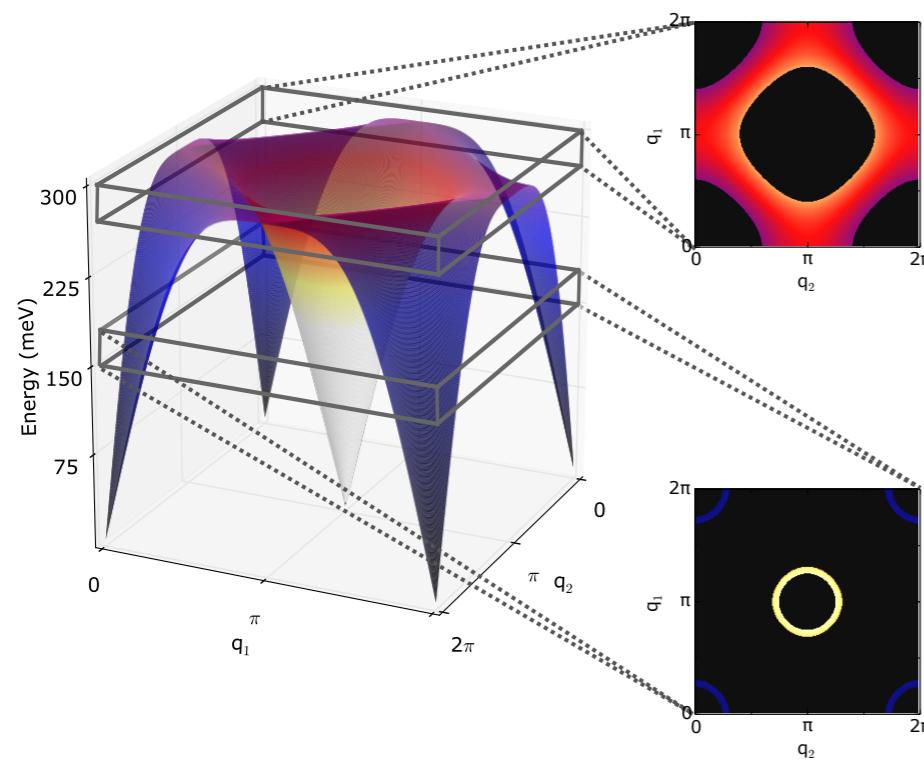
- Stronger magnetic exchange for 3d transition metals.
 - Low oxidation states of 3d TM's rules out oxides.
 - Charge balance requires +1.5 on A-site → mix Na^{1+} and Ca^{2+} .

FLOURIDE PYROCHLORES: $\text{NaCaNi}_2\text{F}_7$



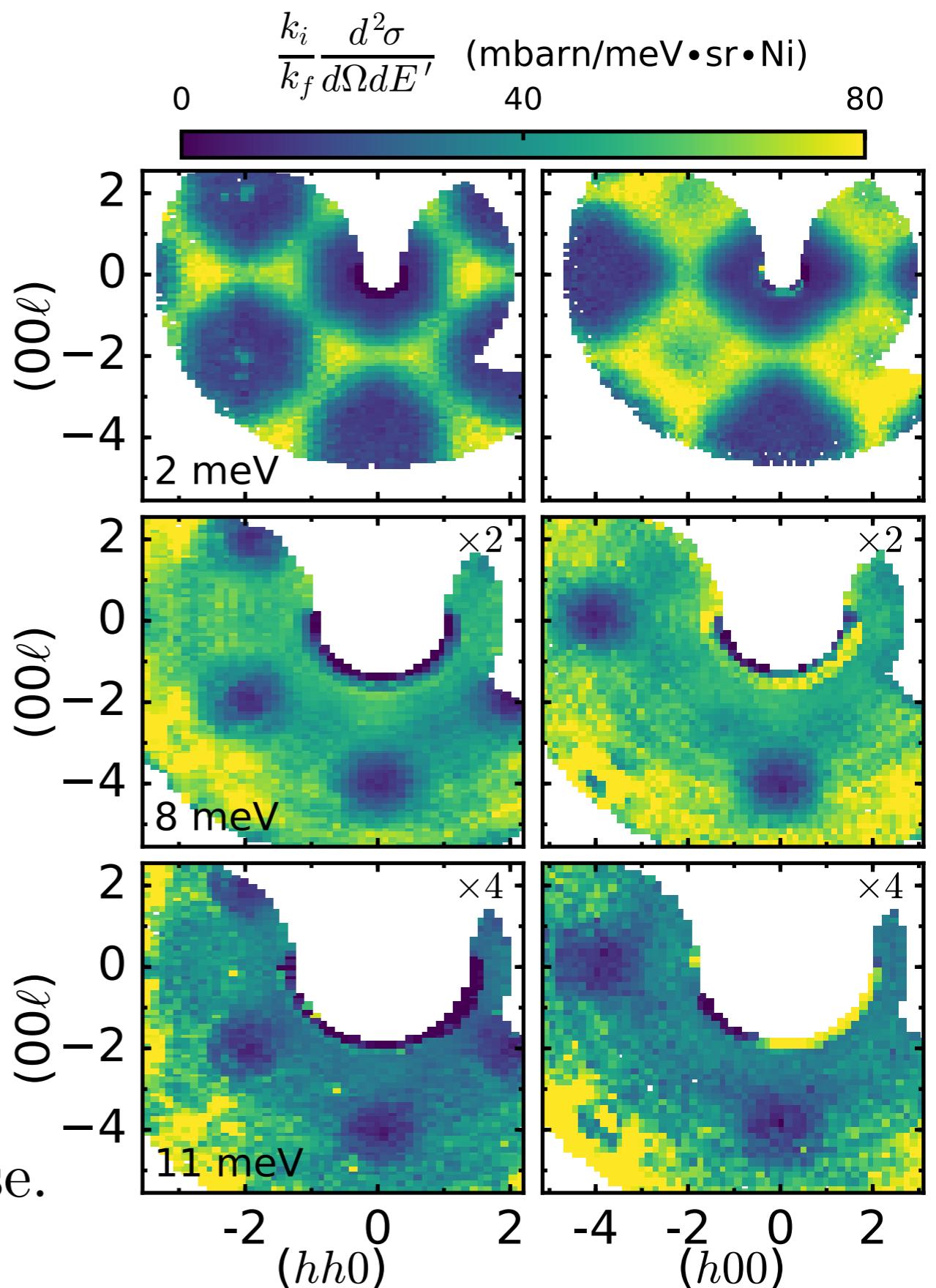
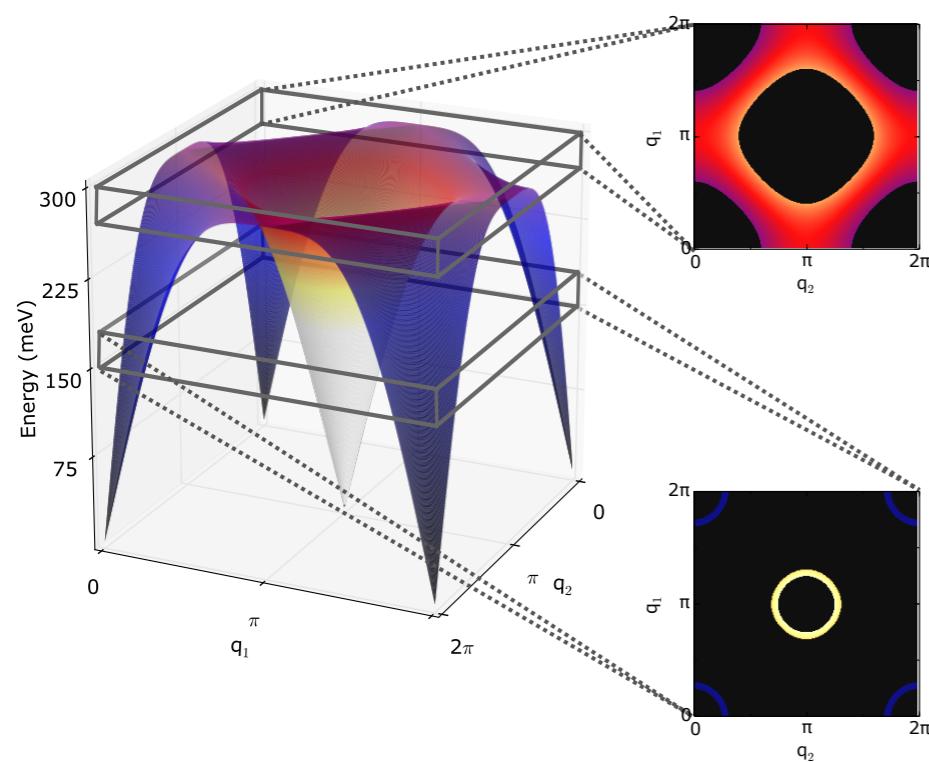
- Magnetic interactions are strong: $\Theta_{CW} = -129 \text{ K}$.
- $\text{Ni}^{2+} \rightarrow S=1$, in the quantum limit.
- Spin freezing at $T_f = 3.5 \text{ K}$. Signatures of a spin glass.
- Estimate bond disorder from T_f : $\delta J \approx 0.2 \text{ meV}$.

MAGNETIC EXCITATIONS IN NaCaNi₂F₇



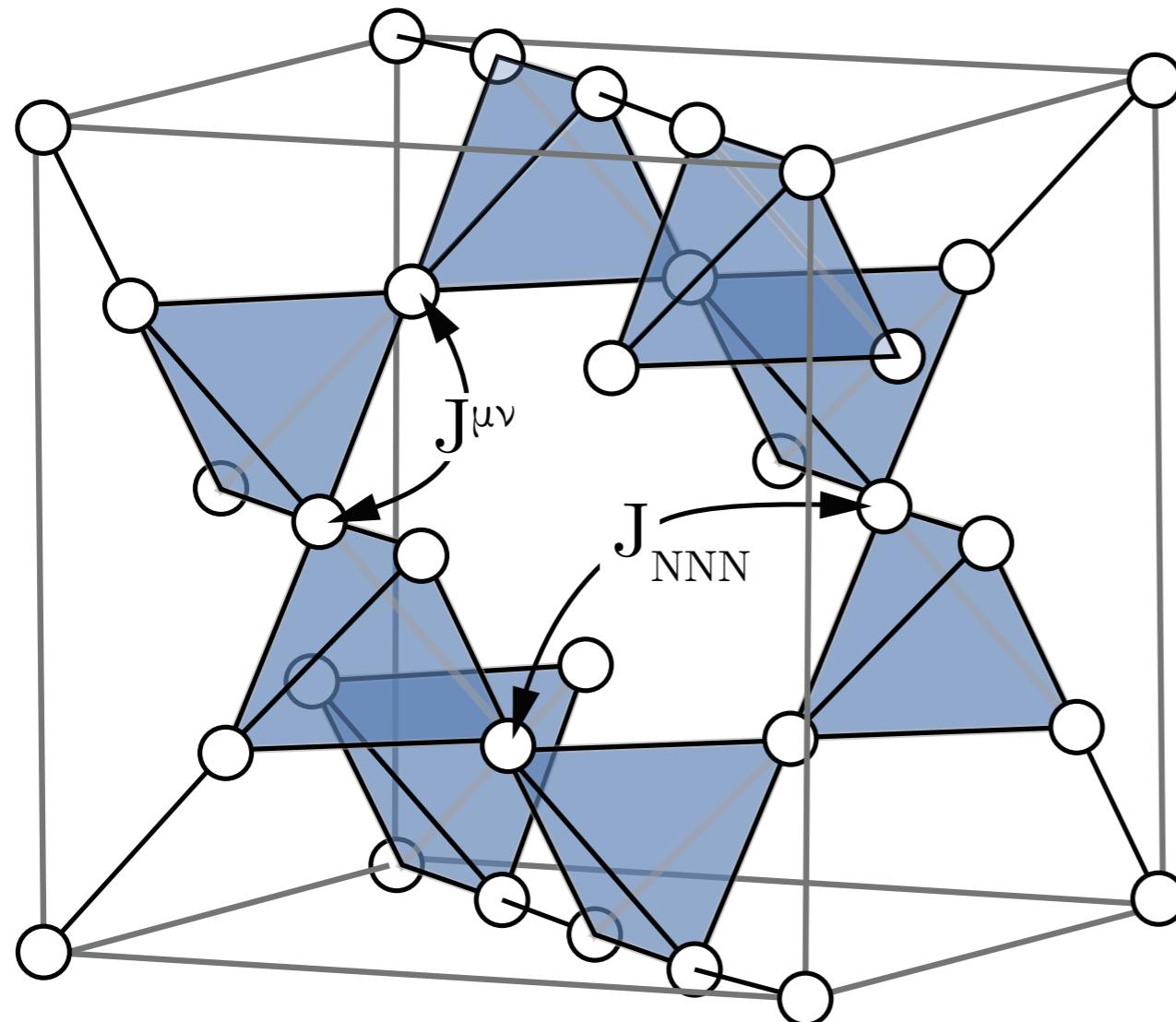
- Excitations show dipolar correlations.
- Looks like theoretical prediction for the spin-liquid on a pyrochlore lattice.

MAGNETIC EXCITATIONS IN $\text{NaCaNi}_2\text{F}_7$



- No LRO and no spin waves.
- Dispersion perpendicular to bow-ties.
- Excitations of the Coulomb phase.

DETERMINING THE MAGNETIC INTERACTIONS



$$H = \sum \mathbf{S}_i \cdot \mathbf{J} \cdot \mathbf{S}_j$$

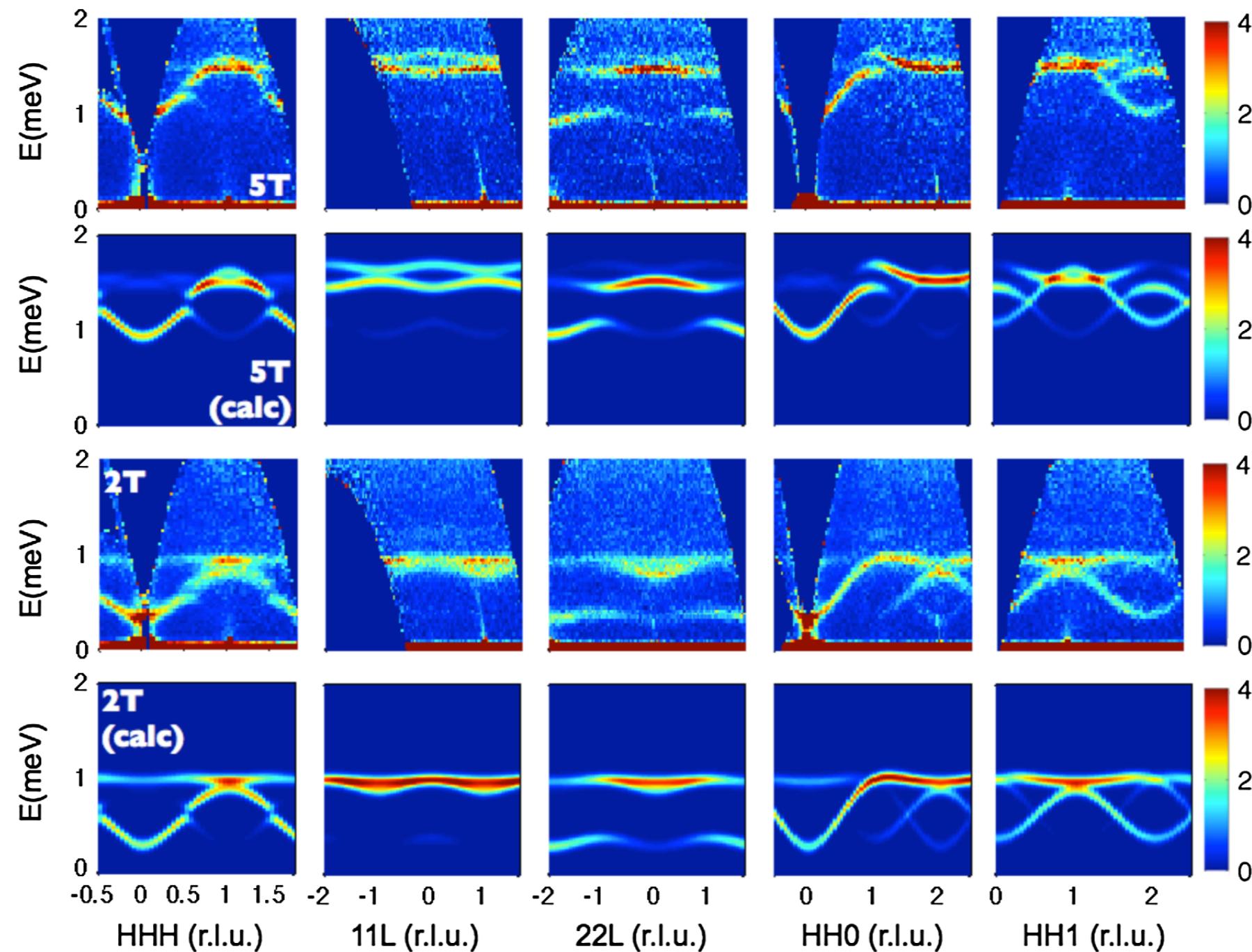
$$\mathbf{J} = \begin{pmatrix} J_2 & J_4 & J_4 \\ -J_4 & J_1 & J_3 \\ -J_4 & J_3 & J_1 \end{pmatrix}$$

- Must consider symmetry allowed exchange matrix.
- Further neighbor exchange also possible.

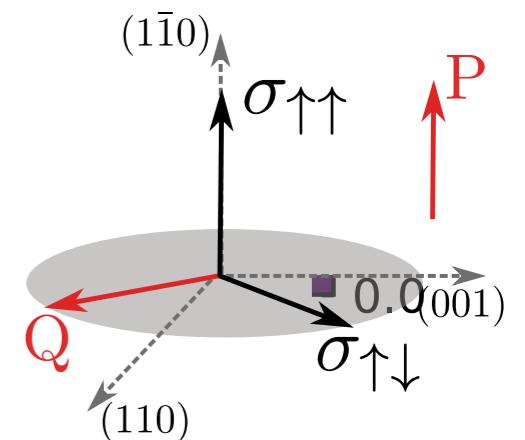
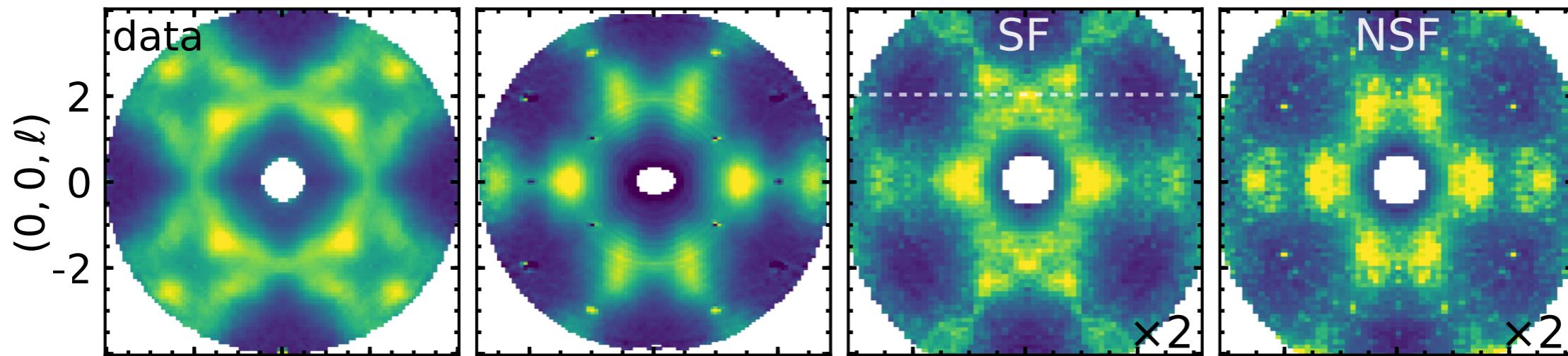
DETERMINING THE MAGNETIC INTERACTIONS

$$H = \sum \mathbf{S}_i \cdot \mathbf{J} \cdot \mathbf{S}_j$$

$$\mathbf{J} = \begin{pmatrix} J_2 & J_4 & J_4 \\ -J_4 & J_1 & J_3 \\ -J_4 & J_3 & J_1 \end{pmatrix}$$

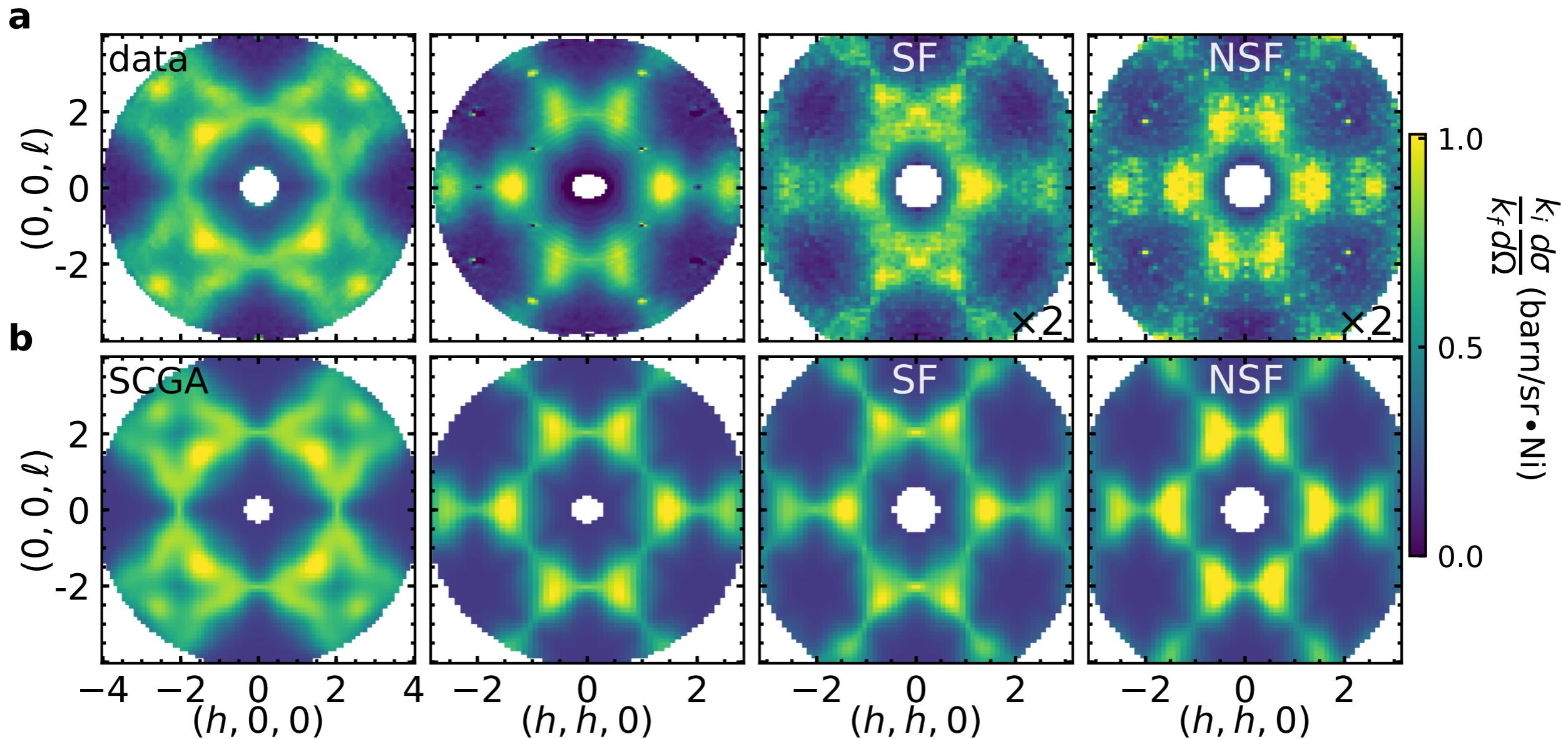


DETERMINING THE MAGNETIC INTERACTIONS



- Measure equal time correlation function $S(Q)$.
- Resolve spin components with polarized neutrons, large anisotropies are ruled out.
- Refine interactions using Gaussian approximation to calculate $S(Q)$.

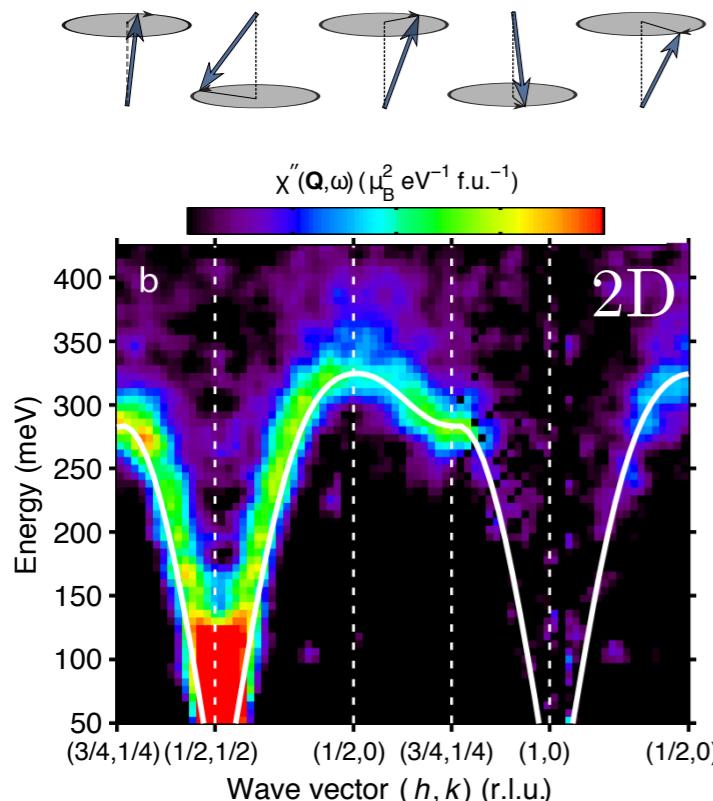
DETERMINING THE MAGNETIC INTERACTIONS



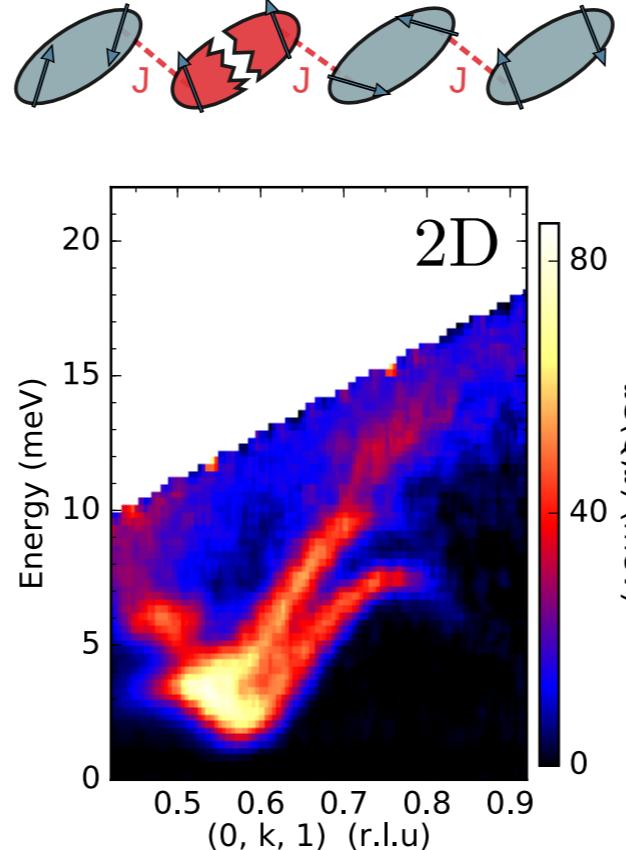
- Global refinement of Gaussian approximation to $S(Q)$.
- Nearest neighbor: $J_1 = J_2 = 3.2(1)$ meV, $J_3 = 0.019(3)$ meV, $J_4 = -0.070(4)$ meV.
- Next nearest neighbor: $J_{NNN} = -0.025(5)$ meV.

COLLECTIVE EXCITATIONS IN MAGNETS

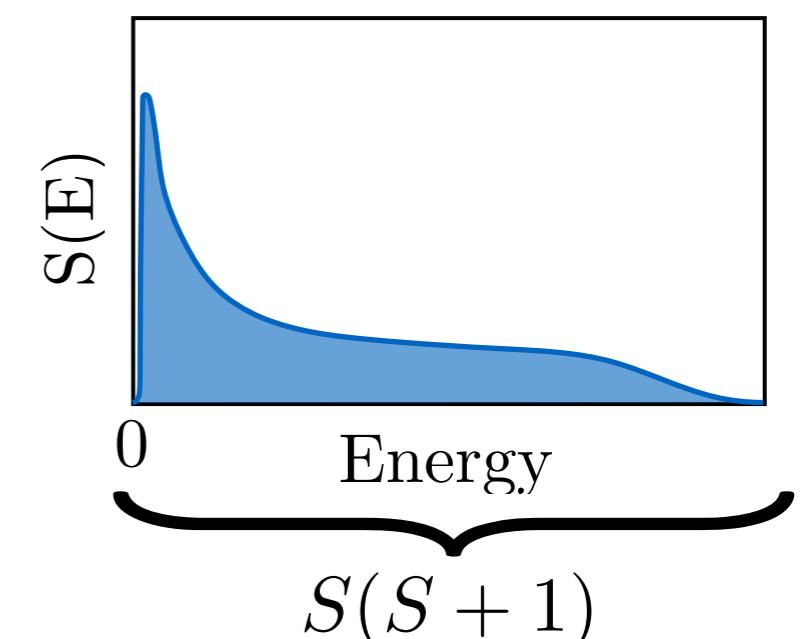
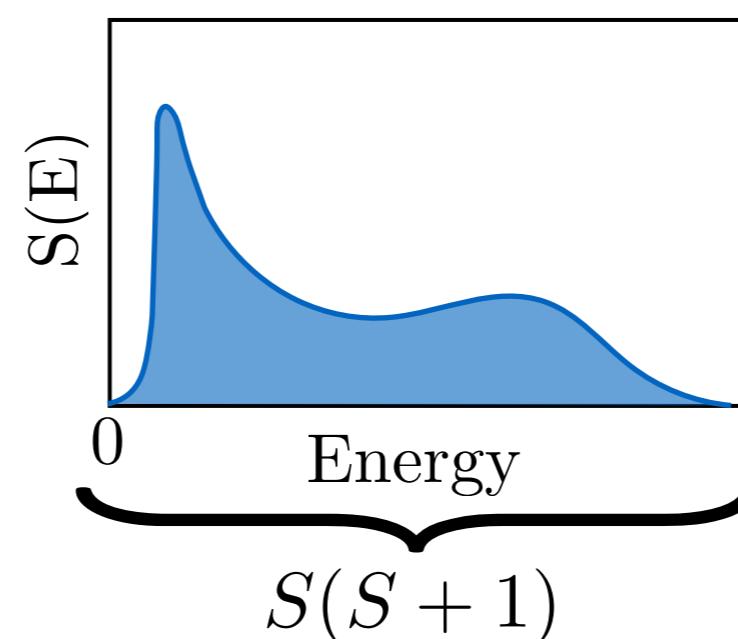
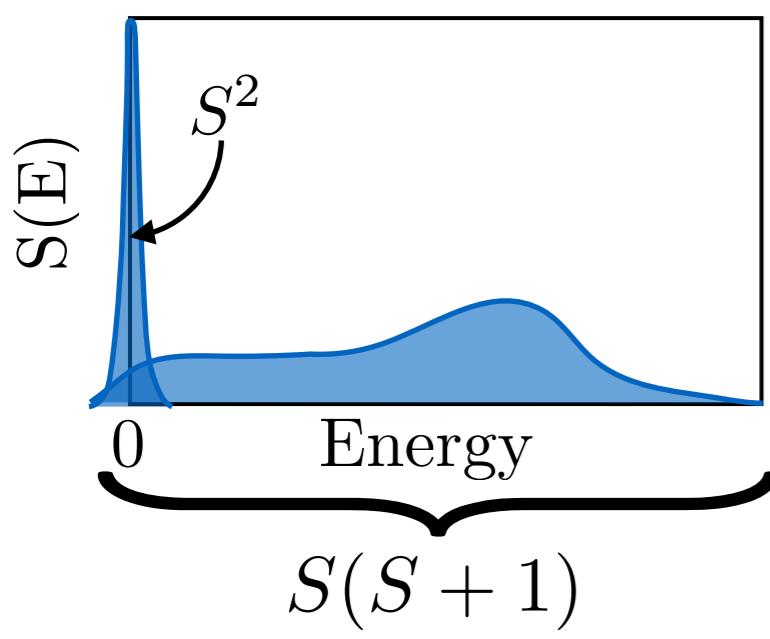
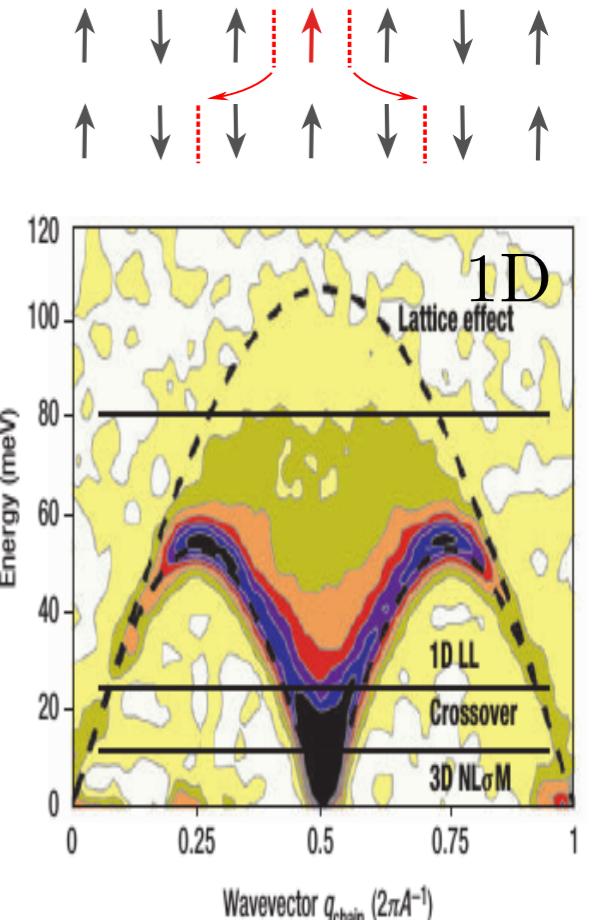
magnetic order



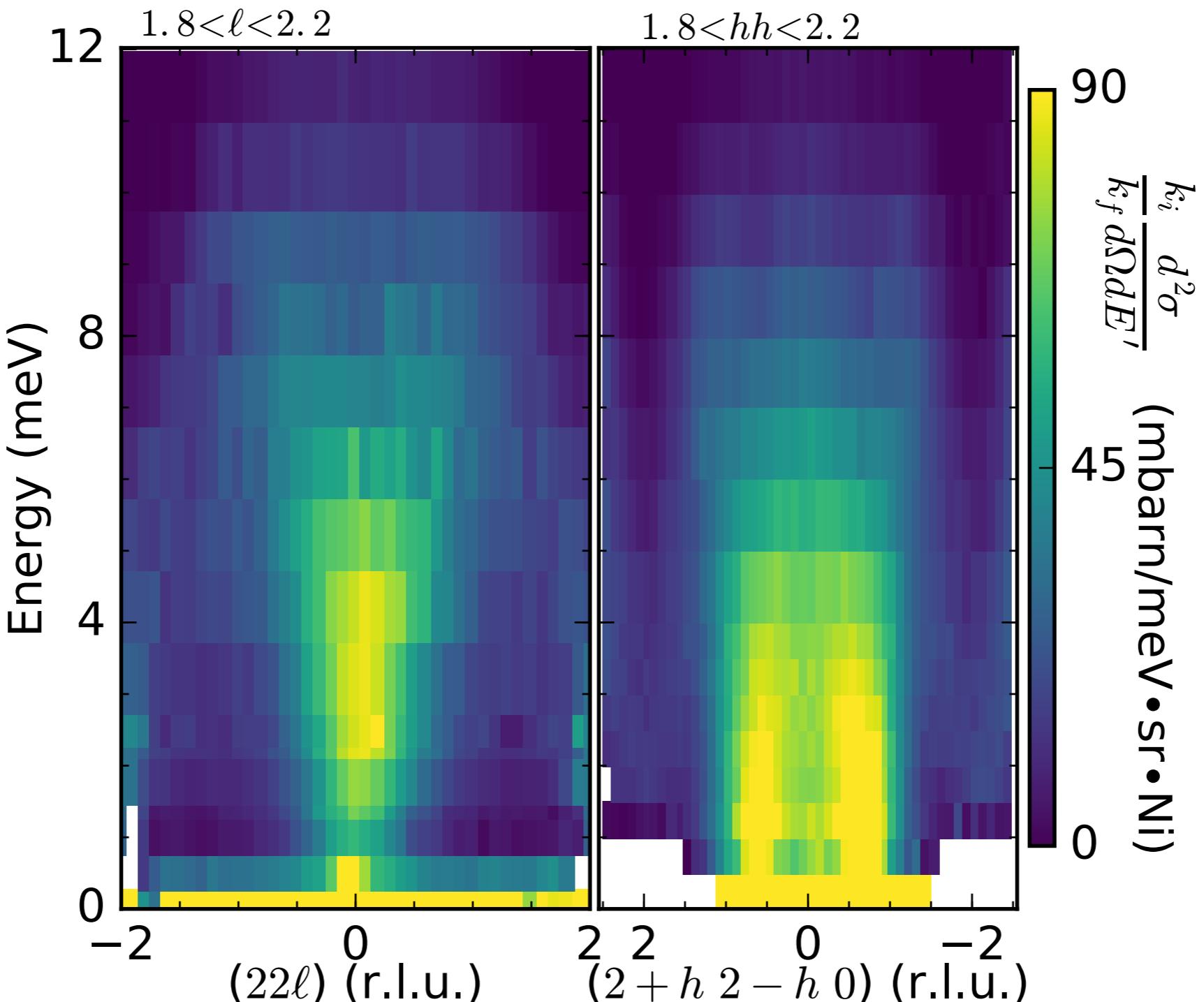
valence bond order



fractionalization

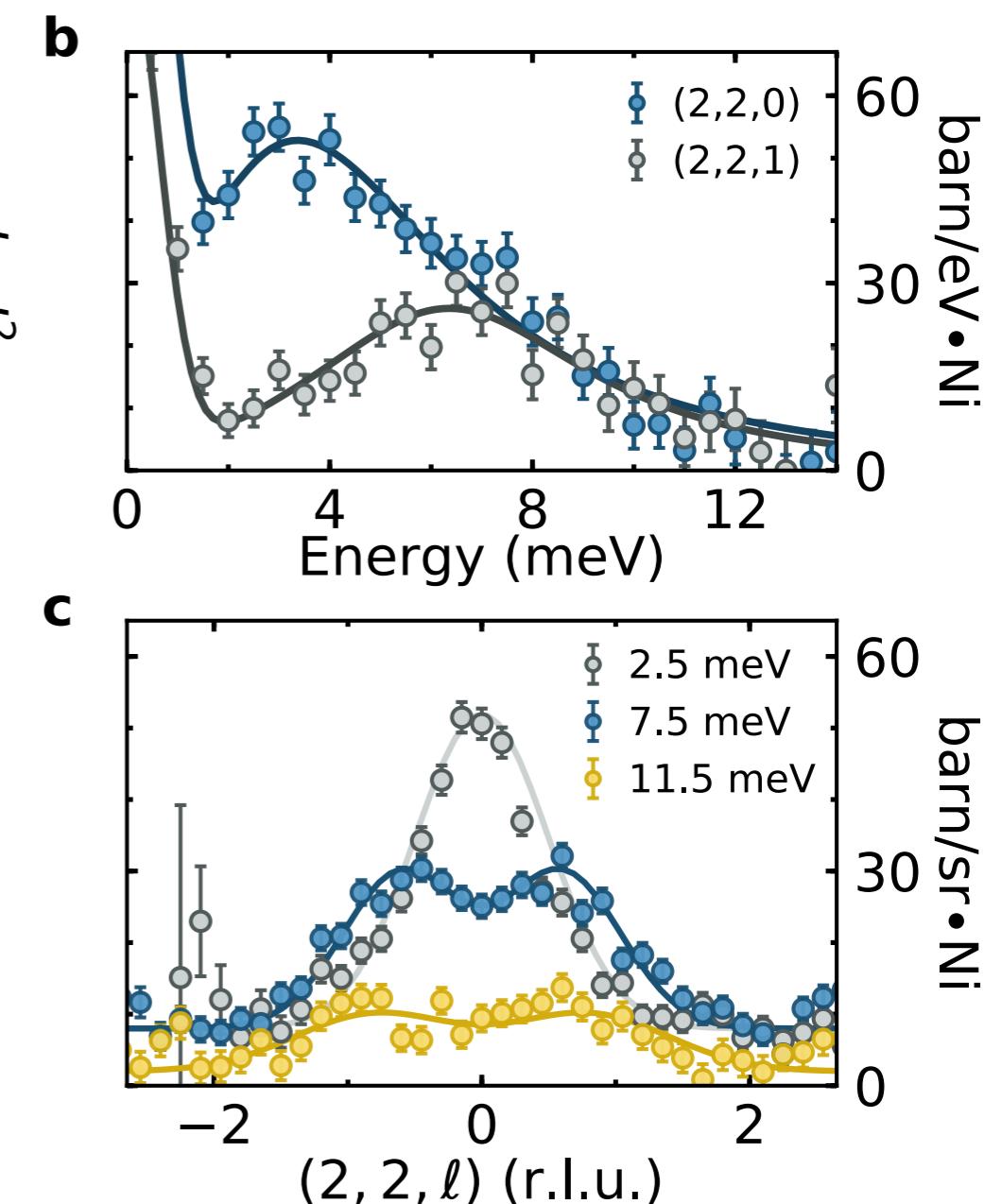
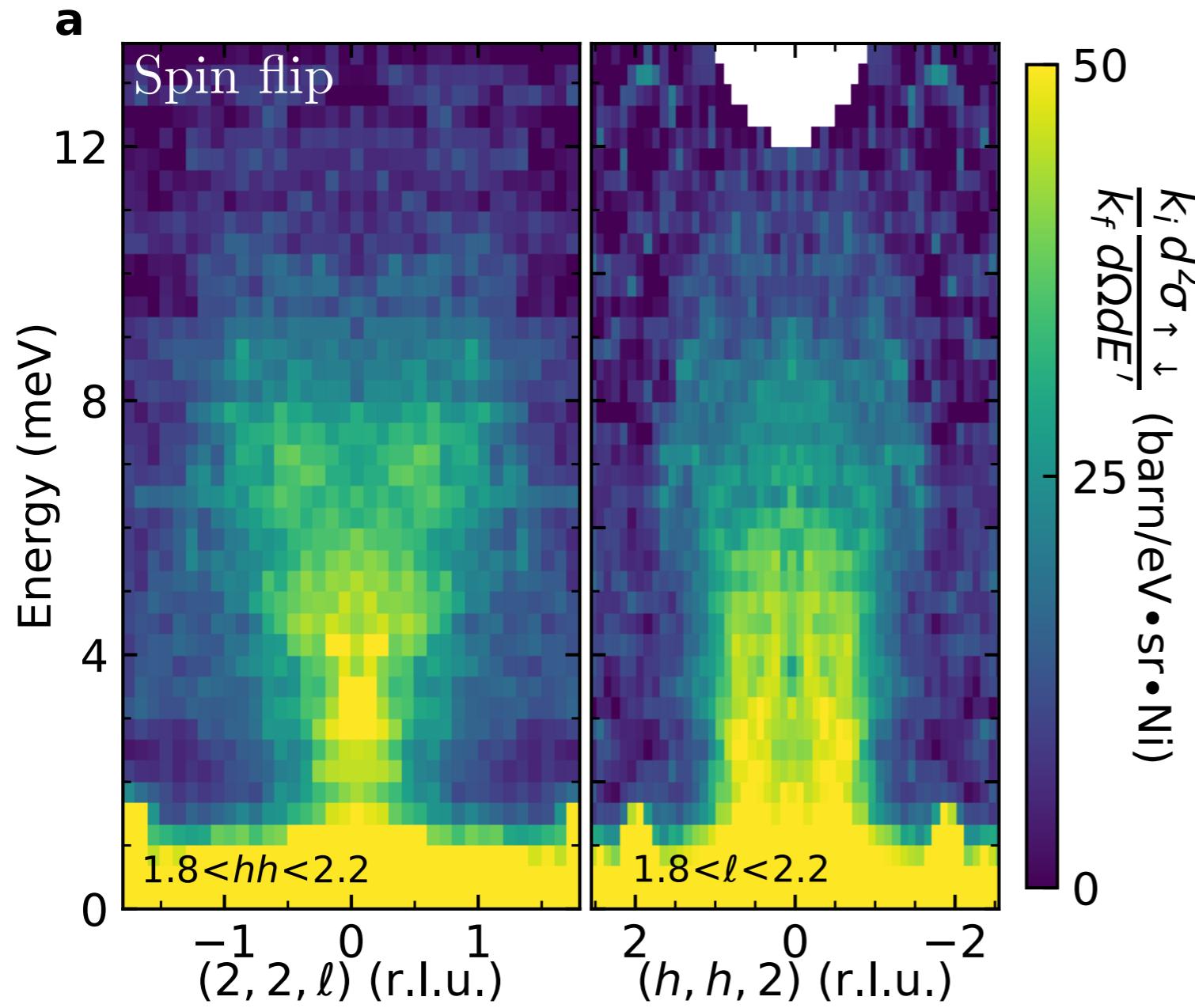


EXCITATIONS FROM A COULOMB PHASE



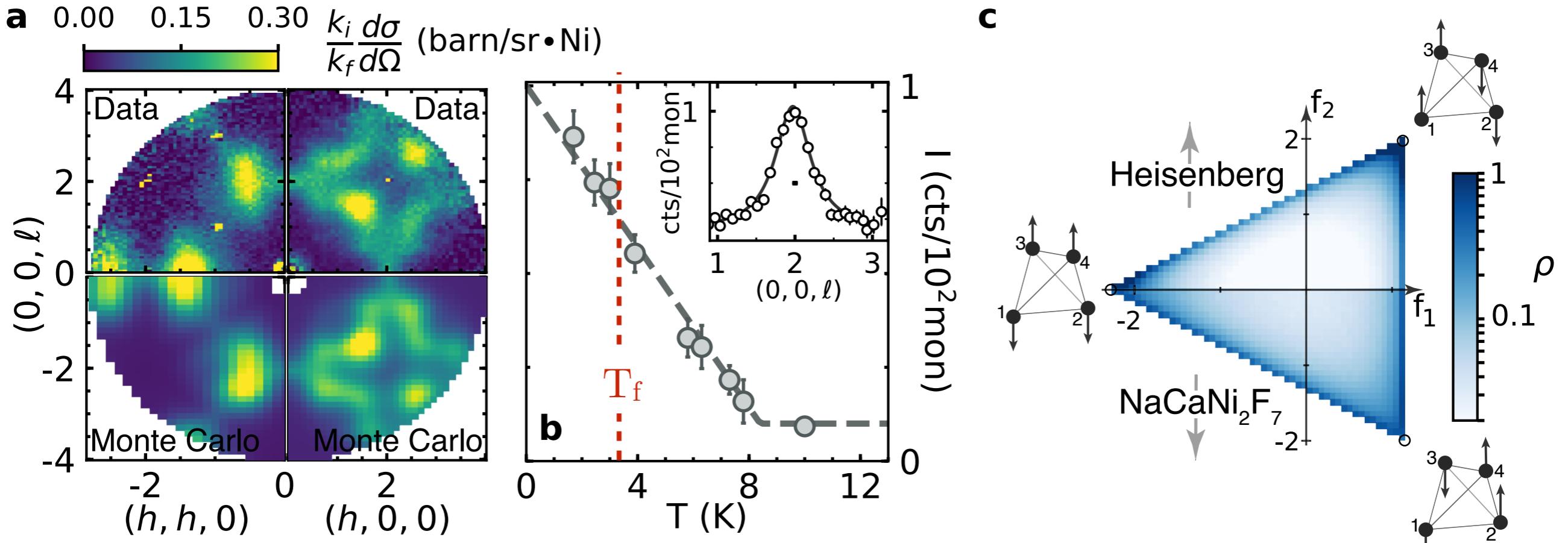
- Continuum of excitations in a three-dimensional magnet.
- Total signal is $S(S + 1)$ and 90% of the spectral weight is inelastic (> 0.4 meV).

EXCITATIONS FROM A COULOMB PHASE



- No resonant modes, but well defined structure.
- Spectral weight build up at $3.5 \text{ meV} \approx J_1$.

UNCONVENTIONAL SPIN FREEZING

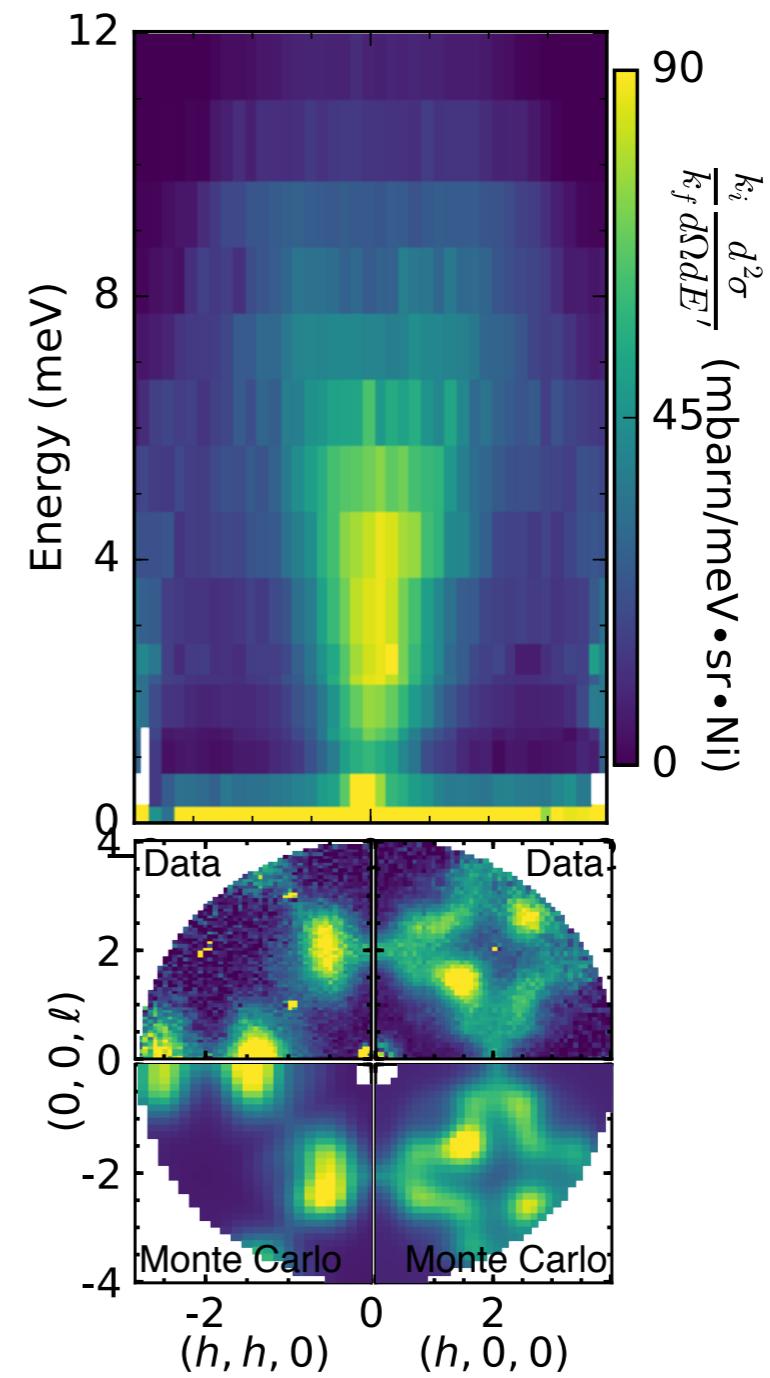
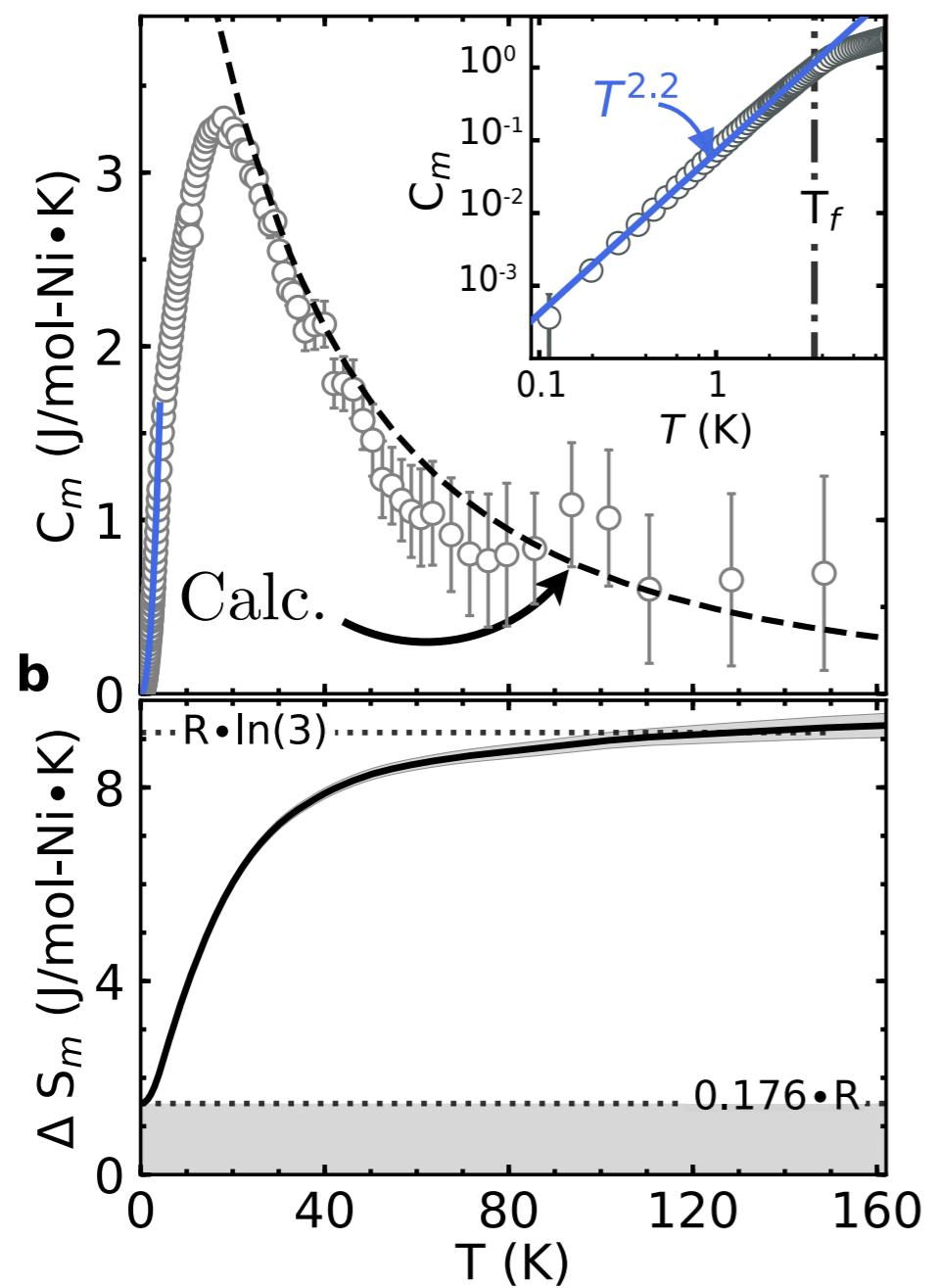


- Unconventional frozen ground state.

$$\langle m_{elastic} \rangle^2 / (g^2 S(S+1)) = 10\%$$

- Study local structure with numerics and knowledge of Hamiltonian.
- Frozen moment configurations have approximately $\mathbf{S}_{tot} = \sum_{i \in \Delta} \mathbf{S}_i = 0$.

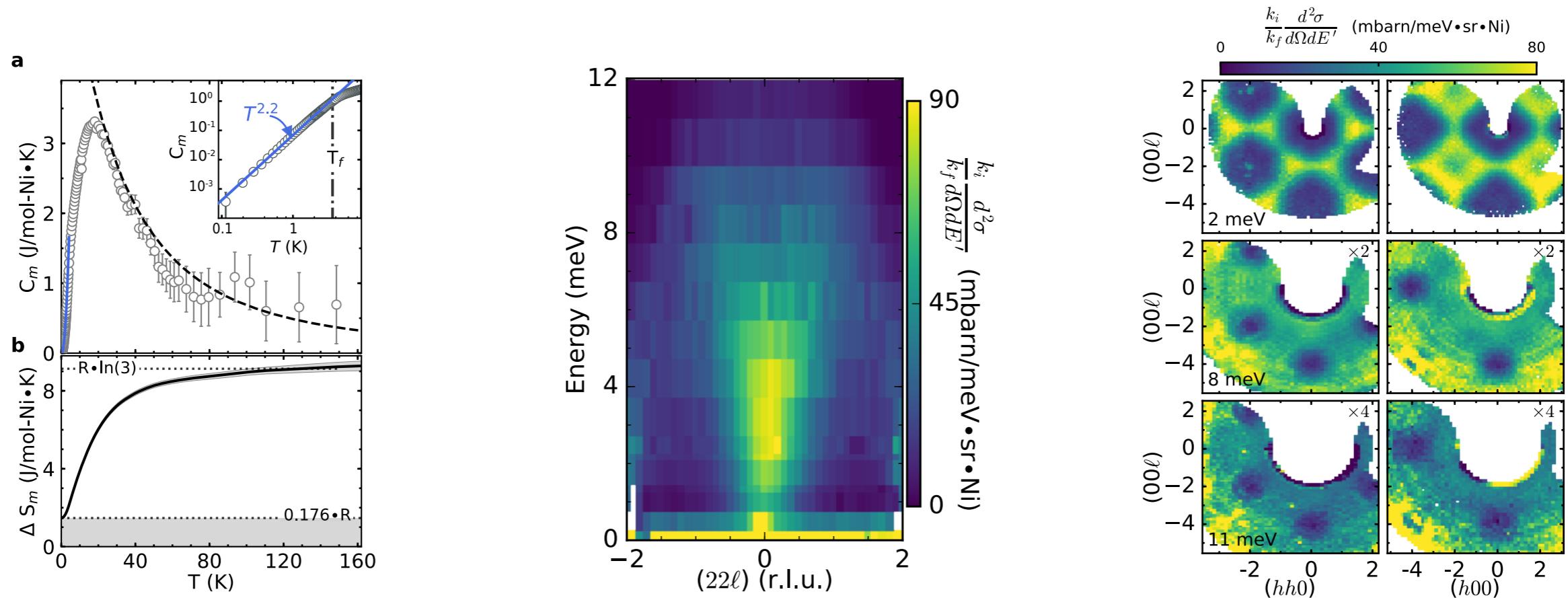
THE OVERALL PICTURE



- Classical spin liquid above 20 K. Villain's cooperative paramagnet.
- Residual entropy is $0.176/\text{spin} \rightarrow$ broken ergodicity, disorder?

$$\Delta S/R \ln(3) \approx \langle m_{elastic} \rangle^2 / (g^2 S(S+1))$$

SUMMARY



- Role of “quantum” → 90% of spectral weight inelastic.
- Existence of quasiparticles? Fractionalization?
- Nature of the frozen state? Kinetically arrested spin liquid?

Thank You!