

Answers to problems

- 1.1. (a) 70 s. (b) 4 years.
- 2.2. (a) Circular. Area π . Normal 0, 0, 1.
 (b) Ellipse. Area $\pi 2\sqrt{2}$, Axis ratio $2^{-\frac{1}{2}}$, Normal $\propto (1, 1, 0)$.
 (c) Ellipse. Area 12π , Axis ratio $\frac{1}{2}$, Normal $\propto (-1, -1, 1)$.
 (d) Linear. $|E| = \frac{1}{2}\sqrt{33}$. Direction $\propto (4, 4, 1)$.
 (e) Circular. Area 2π , Normal 0, 0, 1.
 (f) Ellipse. Area πab , Axis ratio a/b , Normal 0, 1, 0.
- 3.1. Ordered: 0.5 km s^{-1} . Thermal 150 km s^{-1}
- 3.8. (a) $\sim \omega_N^{-1}$. (b) $2 \times 10^{15} \text{ s}$ for temperature 800 K.
- 4.1. $\left\{ \frac{1 + |R^2| - |1 + R^2|}{1 + |R^2| + |1 + R^2|} \right\}^{\frac{1}{2}}, \quad \frac{1}{2}|E_x^2|(1 + |R^2| + |1 + R^2|)$.
- 4.2. $|(\rho - R)/(1 - R\rho)|$.
- 4.4. Smallest $0, 45^\circ$. Greatest $(1 - a)/(1 + a), 90^\circ$.
- 4.5. (a) Elliptical. Axis ratio $\frac{1}{3}$. (b) Partially circularly polarised. Energy flux ratio (circular component)/(unpolarised component) $= \frac{3}{2}$.
- 5.2. (a) $V = c$.
- 5.4. $Gn^2 + J = 0$. See (3.51) for meaning of G and J . If this is true for any Θ , it is true for all Θ .
- 5.5. $\tan \Theta = \{1 \pm (1 - 8 \tan^2 \beta)^{\frac{1}{2}}\}/2 \tan \beta$. See fig. 5.19 for an example.
- 6.3. $q^2 = C^2 - \frac{X(1 - X)}{1 - X - \frac{1}{2}Y^2S^2 \pm Y\{(1 - X)(C^2 - X) + \frac{1}{4}Y^2S^4\}^{\frac{1}{2}}}$.
- 7.1. $d^2y/dx^2 + y\omega^2m/T = 0$.

$$y \sim m^{-\frac{1}{2}} \exp\left(\pm i\omega T^{-\frac{1}{2}} \int^x m^{\frac{1}{2}} dx\right).$$
 Amplitude proportional to $m^{-\frac{1}{2}}$.
- 7.2. (M is the mean molecular mass, K is Boltzmann's constant and γ is the ratio of

the specific heats)

$$\frac{d^2 p}{dz^2} - \frac{T}{P} \frac{d}{dz} \left(\frac{P}{T} \right) \frac{dp}{dz} + \frac{\omega^2 M}{\gamma K T} p = 0,$$

$$\frac{d^2 \zeta}{dz^2} + \frac{1}{P} \frac{dP}{dz} \frac{d\zeta}{dz} + \frac{\omega^2 M}{\gamma K T} \zeta = 0.$$

$$(a) \quad \zeta \sim T^{\frac{1}{2}} \exp \left\{ \pm i \omega \left(\frac{M}{\gamma K} \right)^{\frac{1}{2}} \int^z T^{-\frac{1}{2}} dz \right\},$$

$$p \sim \mp i \omega \gamma P \left(\frac{M}{\gamma K T} \right)^{\frac{1}{2}} \zeta.$$

$$(b) \quad p \sim P^{\frac{1}{2}} \exp \left\{ \pm i \omega \left(\frac{M}{\gamma K T} \right)^{\frac{1}{2}} z \right\},$$

$$\zeta \sim \frac{i}{\omega P} \left(\frac{K T}{\gamma M} \right)^{\frac{1}{2}} p,$$

7.3. $a \propto L^{\frac{1}{2}}, \quad T \propto L^{\frac{1}{2}}, \quad E \propto L^{-\frac{1}{2}}.$

7.5. 0.25 km approx.

8.1. (a) $a_0 a_1$, (b) $a_0 \text{Ai}'(0)$, (c) $-a_1 \text{Ai}(0)$.

8.3. (a) $\pi \{ \text{Bi}(0) \text{Ai}(z) - \text{Ai}(0) \text{Bi}(z) \}.$

(b) $\pi \{ \text{Bi}'(0) \text{Ai}(z) - \text{Ai}'(0) \text{Bi}(z) \}.$

(c) $\{ \text{Bi}(-1) \text{Ai}(z) - \text{Ai}(-1) \text{Bi}(z) \} / \{ \text{Bi}(-1) \text{Ai}(0) - \text{Ai}(-1) \text{Bi}(0) \}.$

(d) No solution; homogeneous boundary conditions.

(e) $\text{Ai}(z)/\text{Ai}(0).$

(f) As for (d)

(g) Any multiple of $\text{Ai}(z).$

(h) $\text{Ai}(z)/\text{Ai}'(-a).$

(i) No solution; conditions incompatible

8.4. (a) a is smallest x that makes $\text{Ai}(-x) = 0$. $y = \text{Ai}(z-a)$. Homogeneous.

(b) a is any solution of $\text{Ai}(-a)/\text{Bi}(-a) = \text{Ai}(1-a)/\text{Bi}(1-a)$. $y = \text{Bi}(-a) \text{Ai}(z-a) - \text{Ai}(-a) \text{Bi}(z-a)$. Homogeneous.

(c) a can have any value provided $\text{Ai}(-a) \neq 0$. $y = \text{Ai}(z-a)/\text{Ai}(-a)$. Inhomogeneous.

(d) a is solution of $\text{Ai}'(-a) = 0$. $y = \text{Ai}(z-a)$. Homogeneous.

(e) a is solution of $\text{Ai}'(-a)/\text{Bi}'(-a) = \text{Ai}(1-a)/\text{Bi}(1-a)$. $y = \text{Bi}(1-a) \text{Ai}(z-a) - \text{Ai}(1-a) \text{Bi}(z-a)$. Homogeneous

(f) a can have any value provided denominator of y is not zero.
 $y = [\text{Bi}(1-a) \text{Ai}(z-a) - \text{Ai}(1-a) \text{Bi}(z-a)] /$
 $[\text{Bi}(1-a) \text{Ai}'(-a) - \text{Ai}(1-a) \text{Bi}'(-a)].$ Inhomogeneous.

8.5. $\frac{d^2 u}{dz^2} - \frac{1}{z} \frac{du}{dz} - zu = 0.$

8.6. (a) $y \sim z^{-\frac{1}{2}} \exp(\pm \frac{1}{4} i z^2).$

- (b) $\arg z = \pm \frac{1}{4}\pi, \pm \frac{3}{4}\pi$.
 (c) $\arg z = 0, \pi, \pm \frac{1}{2}\pi$.
 (d) $i\sqrt{2}$ on all four Stokes lines.
- 9.1. $(2\pi/a)^{\frac{1}{2}} \left(1 - \frac{1}{8a} + \frac{9}{128a^2} - \cdots \right)$.
 $(2\pi/a)^{\frac{1}{2}} \left(1 - \frac{5}{8a} + \frac{129}{128a^2} - \cdots \right)$.
- In each case method (a) gives first term only.
- 9.2. $(\pi/a)^{\frac{1}{2}} \left(1 - \frac{1}{2a} + \frac{3}{4a^2} - \frac{3.5}{8a^3} + \cdots \right)$.
- 11.1. For reflected wave $E_y/\mathcal{H}_y = R_{21}/R_{11} = R_{22}/R_{12}$. Zero reflection when incident wave has $E_y/\mathcal{H}_y = R_{11}/R_{12} = R_{21}/R_{22}$.
- 11.2. $R_0 = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix}$.
- 11.3. $(R_{11} - R_{22})^2 = -4R_{12}R_{21}$.
- 11.4. Incident wave $E_x = 12$, $E_y = 0$. Reflected wave circular $E_y = -iE_x = 5$.
- 11.5. Electric fields of incident and reflected waves have same direction. Amplitude ratio $\frac{1}{3}$.
- 11.6. $R_{11} = (n + Z_0\tau/\rho - 1)/(n + Z_0\tau/\rho + 1)$ where n is refractive index, τ is thickness and ρ is resistivity. $\tau \approx 5 \times 10^{-10}$ m (about two atomic diameters).
- 11.7. $R_{11} = \frac{n^2 \cos \theta - (n^2 - \sin^2 \theta)^{\frac{1}{2}} - \gamma \sin \theta}{n^2 \cos \theta + (n^2 - \sin^2 \theta)^{\frac{1}{2}} + \gamma \sin \theta}$
 where

$$n^2 = \frac{(U - X)^2 - Y^2}{U(U - X) - Y^2}, \quad \gamma = \frac{iXY}{U(U - X) - Y^2}.$$
- 11.10. Angle shift $(z_0 - z_1)D/(D^2 + z_0^2)$, Range shift $2z_0(z_0 - z_1)(D^2 + Z_0^2)^{-\frac{1}{2}}$, for observer on ground at distance D from transmitter.
- 12.4. (a) $f_N^2 = \begin{cases} 0 & 0 \leq z < h_0 \\ f_1^2 + \alpha(z - h_0) & h_0 \leq z \end{cases}$
 (b) $f_N^2 = \begin{cases} 0 & 0 \leq z \leq h_0 \\ \alpha(z - h_0) & h_0 \leq z \leq h_0 + f_1^2/\alpha \\ f_1^2 + \beta(z - h_0 - f_1^2/\alpha) & h_0 + f_1^2/\alpha \leq z \end{cases}$
- 14.4. The problem of the last sentence has no solution.
- 18.2. $\gamma = -(1 + \rho_2\rho_1^*)/(1 + \rho_1\rho_1^*)$.
 (a) $\gamma = 0$. Solution (b) is the penetrating mode.
 (b) $\gamma = -1$. In solution (b) the penetrating mode has been completely swamped. Its fields at the bottom cannot be found from the data.
 (c) $\gamma = -1$. At the bottom the penetrating mode is linearly polarised with $E_y = \rho_2$, $\mathcal{H}_y = -C\rho_2$.