

Recent development in Categorical dynamics

Many people have heard about SDG, carried out by Wraith, AK, Reyes, Bunge, ...

To clarify the objective logic of cont. mechanics

Voll, Gabriel (Gros toposes, Verdier (LaTolla),
taught me

Guth had idea of ABS. There ought to be a non-linear version.

ETCS was the 0-dim part of it.

Halifax - needed: elementary theory of toposes.

"Topological dynamics" (.. & Hedden).

Continuous, smooth, ... so categorical dynamics,
(and many more)

should fit in many different categories.

Use of nilpotents in describing differentiation in algebraic, analytic, ...

Why not C^∞ ? United in: can be embedded in a cat with
tangent bundle is rep'ble, by a unique object.

Fundamental thing was: CCC. From hindsight: Hurewicz' exponentials.

Volterra's functionals, Bernoulli's calc of variations

Hurewicz' campaign for this, communicated it to Fox, who solved
it for sequential spaces. Hurewicz 1948-49 (Princeton) coined the
notion of k -space (compactly generated).

Gale (BAMS Vol. 1, 1950), 1955 Kelley's book, gave an exposition

Many people thought k was k "Kelley". I think I
defined CCC. 1958 K&M adjoint functors; he was aware of
 \exists of k -space in simplicial sets. Since uniqueness,
can be used axiomatically, 1963 Cat of cat's defn using

exp.; Eilenberg-Kelly (LaTolla) Closed monoidal cat.

in particular, since \times is "cartesian", hence the name. They

also noted that highly alg's are CCCs

Axiom ① \mathcal{E} is a cat closed cat.

② TB $T_0 \leftrightarrow T_1 \leftrightarrow T_2 \dots$ amazingly tiny
Extra right adjoint $f^!$ in Groth. notation

$()^!$ means "amazingly!". or ATOM.

$$\begin{array}{ccc} & f_! & \\ \mathcal{E}/r & \xleftrightarrow{\quad} & \mathcal{E} \\ & f_* & \\ & f^! & \end{array}$$

$f_!$, or Σ_r

f_* used notation \prod_r

By composing, get a fractional exponent $()^{1/r}$

$$\frac{X^r \rightarrow Y}{X \rightarrow Y^{1/r}}$$

So T-ary functionals can be viewed (as functions with different codomain). Eilenberg-Mac Lane space.

Differential form. Undreamed of in classical logic

shall use: $T_0 = 1$, $T_1 = T$, & $T_2 = T^2/2!$

(beginning of Taylor series)

'Tiny' by Freyd & Yetter. Relative to the types you are in.
(Every T can become tiny in a [bigger topo].)

Then also have $X^{B/r}$ arbitrary B .

Where do they come from? Imported in math. from Heraclitus
'river twice' Intermediate position: Galileo, Newton, used
tiny motion

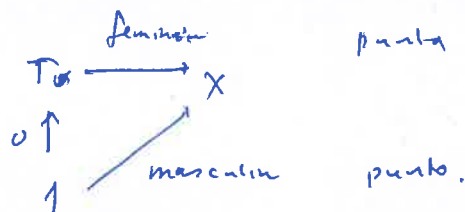
$$T_0 \xrightarrow{\circ} T_1$$

Radical synthetic. (Synthetic: deriv algebraic fun. geometric)

this program is progressing).

$T \xrightarrow{x} X$ is stepping into the river once, at X_0 , but

"Il punto nella punta"



The fact that $1 \xrightarrow{0} T$ is not an iso is shown by: T^T has two points, with empty meet

$$1 \xrightarrow{\tau_0^*} T^T \xrightarrow{\tau_{id}^*} T^T$$

Call X^T "tangent bundle", with $ev_0 : X^T \rightarrow X$,

can get $T_x X$ for $x \in X$. "Tangent space at x ". Has one structure: $T^T \xrightarrow{ev_0} T$ tangent bundle of T

Take tangent space at x , call it R . It is a monoid. This is essentially Euler's def of real no's. It operates on T^T , by "speed-ups".

So X^T has operations by R , so have ratios of velocities. So every high school student can calculate when Achilles overtakes the tortoise (provided $1-x$ is invertible)

Principle in philosophy.

So the tangent space at $x \in X$

$$\text{get } \left[T_x X \rightarrow T_{f(x)} Y \right]$$

$$(X^T)_x \rightarrow (X^T)_{f(x)} \quad \text{compatibly}$$

$$f_x(\lambda \cdot -) = \lambda \cdot f_x$$

for any endomorphism of T fixing 0

If T spectrum of $k[E]$, then

$k[X]$ ~~is~~ spec has $(T^T)_0$ as spec. } in alg. geom

a simple observation in [SDG].

"Infinitesimal generation" ;

Limits of \lim_{\leftarrow} and exponential h (atoms).

Often homogeneity implies linearity (i.e. addition).

A paper I have in Bell volume, concerning addition.

Another synthetic addition,

For much of what I do, T_2 could be anything, not nec $T \times T/2$!

Define $D_n = \{h \in T^T \mid h^{n+1} = 0\} \subset R$

In most books, T is identified with D_1 . To choose an

i.e. $T \cong D_1$ is choosing a unit of time.

T is an infinitesimal time. T^T is 'pure' quantities, whatever

T is ; all the D_i 's are pure.

$$T \longrightarrow D_1 \subseteq T^T$$

By exp adjoint $T \times T \longrightarrow T$, monoid without unit

$D_1 \longrightarrow T^T$ is a monoid homo which has a retraction.

That is the origin of Newton's ().

Parallel to that, have inclusion $D \subset D_2$

$$D_2 = \{t \mid t^3 = 0\}$$

I use this coordinatization for exposition.

Achievement of all my SDG friends, all diff calculus comes out right, for a dense class of spaces

Dynamics. Consider a 'configuration' space C

Consider

$$C^{T_1} \xleftarrow{\alpha^*} C^{T_2}$$

$$\text{where } T_1 \xleftarrow{\alpha} T_2$$

Basic structure : a section ξ of α^* , "prolongation"

\mathcal{E} is a specific information, unlike α^x .

Get a cat. \mathcal{E}^α . If \mathcal{E} , \mathcal{E}^α is a topos $\mathcal{E} \xrightarrow{\quad} \mathcal{E}^\alpha$

Like when you have enriched monoid (concrete).

But $()^D \rightarrow ()^{1/D}$ is not enriched over \mathcal{E} , but over \mathcal{S} (category of \mathcal{S} with $S \rightarrow ST$ is ω_0)

Drawn by Anders and Gonzalo. Think of \mathcal{E}^α as the "cat of laws of motion", 2nd order, as Fibonacci understood in the discrete case \mathcal{E}^α (even when T is not an atom)

is a type of alg structure: a T_2 -tuple of T_1 -ary operation.
An elementary transformation in universal alg.

Assume you have a theory which involves a W -tuple of T -ary op's on C . Can transform into W^T -tuple of unary op's on \mathcal{E}^T

This bunch does not pass underlying set, U is replaced by U^T

Not full bunch $\text{Alg}(A, \mathcal{E}) \rightarrow \mathcal{E}^M$

where M is the monoid, $T^T \subseteq W^T$ modulo eq's a...

In Dynamics, if $W = T_2$, $\text{Alg}(A, \mathcal{E}) = 2^{\text{nd}} \text{ order diff eq's}$

$\mathcal{E}^M = D\text{-action on } X^D$ (Equivalent)

$\text{Alg}(A, \mathcal{E}) \rightarrow \mathcal{E}^M$ full & faithful

implies a vector field on X^D .

$M = \text{monoid of unary op's on a bigger space.}$

Box alg's - the content of $v \wedge$ in terms of unary op's on \mathcal{S} .

$(\mathcal{S} = 2 \times 2 / 2!)$. Similarly, notion of ultrafilter.

A third ex. ...

Also happens to have discrete anchors

A basic transit in univ alg!

There are enough unary ops on $R \times R$ (R a ring) to distinguish ring homs out of $R \rightarrow R'$.

\mathcal{E}^M has a further forgetful to the cat of vector field since a vector field is a prolongation along $1 \rightarrow T$.

When \mathcal{F} is a law of motion, what is a lawful motion?

Is a morphism in a cat.

Clear what a morphism of vector field is.

There is a left adjoint to $\text{Alg}(\mathcal{A}, \mathcal{E}) \rightarrow \mathcal{E}^M$

"Algebra of time", a force moving time forward

Open problem!

Laws of motion are laws of becoming

An ex. would be "infinite anticipation of an infinitesimal law."

This far, we came in 1997 in Montreal.

I have recently come up with a new idea, the def. of force.

A law of motion = prolongation along $T_1 \xrightarrow{\alpha} T_2$

SDG may (by coordinatization) $\mathcal{F}(x, v)(t) \quad t \in D_3$

$$= x + vt + a(x, v) \cdot \frac{t^2}{2}$$

So the information

of the prolongation law is $a(x, v)$

A force law is an endomorph of CT_2 s.t.

$$CT_2 \ni f$$

$$\alpha^x \downarrow$$

$$\alpha^x f \alpha^x = \alpha^x$$

$$CT_1$$

They act on laws of motion

law of motion

implies $f\mathcal{F}$ is also a

force law.

Force laws act on motion laws composition

(So) force laws form a monoid but act
 In engineering - there are 17 different force laws.
 Want one force? Just by composing.

A particular starting point (law), the "geodesic law of motion".
 $\{ \lambda = \lambda \} \quad \{ \lambda \in T^* \}$ "straight line motion"

Unlike vector fields, there are no trivial 2nd order eq's.
 The subset of geodesic law takes the place of trivial laws.

In particular, $\lambda f = f \lambda$ "spray" law.

If you operate ~~on~~ ^{on} a geodesic law ~~on~~ by a spray,
 you again get a geodesic law.

Is composition of force laws commutative? Not just
 addition of vectors, not just addition of the $a(x, v)$ -
 terms (which is commutative).

Typical configuration space in continuum mechanics:

infinite-dimensional spaces. The objective law should be the
 same: body = B E = environment

A typical conf space: all possible placements of B in E,
 i.e. E^B (not nec. monomorphism), $C = E^B$.

Force laws on $C = E^B$ - two special classes; a law might
 be induced from a law on E, just by exponentiation.

"single particle" reasoning: eg geodesic (affine) law on E.

Or $B = E$. "body force" as they call it in cont mech

Internal force: if on B we have an endomorphism of H
 $B \times T_2 \rightarrow B \times T_1 \subset B \times T_2$

Induces a H for 'Hooke', get H^{x_i} force law on E^B .

$$B \times T_2 \Rightarrow H$$

split into two parts

$$B \times T_2 \rightarrow B$$

second order parts
in $B \cdot B$

$$B \times T_2 \rightarrow T_2$$

1st order on $B \times T_2$

$$H_B(t)(b) = b + h_B(t) \frac{t^2}{2}$$

$h_B(t)$ is really 1st derivative since the square of t^2 is 0
some kind of tensor (endomorphism of B)

$$B \rightarrow T_2 T_2$$

$$H_{T_2}(t)(b) = b + k(t) \cdot \frac{t^2}{2}$$

Should have considered first the case when $\|B\|$ (~~is 1~~) is 1.

So these laws are given by "messing with time".

Effect of k : adding $k \cdot v$ to the acceleration $a(x, v)$.

This is viscosity. Law of viscous motion.

If B is any obj., k ... non constant viscosity, varying in B .

$$\text{Constitutive rel'n } H: B \times T_2 \Rightarrow H$$

The h -part - I picked ' h ' elasticity, Hooke tensor,
visco-elastic body.

I haven't mentioned mass, Bernadine^(Engineer) et al they say
"mass is nothing"; they are saying that "mass is just
another force"

Newton writes "mass is a force". We should take it
seriously. Where does it live. Not just on E , nor just on B .
eg Mach's principle. Recent stuff on "Higgs boson",
"Hertzian Mechanics"

Discussion

Torson: Line bundles

I forgot: do force laws commute? All internal forces do commute.
In particular, the 1-dim fam of viscosity laws, how do they compose - the k 's just add.

Endomaps of T_2 which fix T_1 . Composing two of these

They add synthetically addition

Why is addition a module on mult. —

Easy if we assume that R is given a priori.

R is a rig. Have $\text{Mod}_R(\mathcal{E})$, an AB5 cat if \mathcal{E} is a topos.
(comm.)

Therefore have \otimes_R , Hom_R since cat

A line bundle A is just an invertible module, $A \otimes_R B \cong R$,

so B must be A^* , i.e. $A \otimes A^* \xrightarrow{\text{ev}} R$ is invertible

Picard group of \mathcal{E}, R - invertible A 's under \otimes ,

also called H^1 . In physics, we need quantities which

are not pure. Physical dimensions. Only q 's of

same kind can be added, but of different dim can be

multipled. The Pic group may just be $\mathbb{Z} \times \mathbb{Z} \times \mathbb{Z}$

Length is different from mass mass length time

^{constant}
A quantity of type A is just $R \rightarrow A$.

The rules of $\text{Hom} \otimes$ adjointness explain how transform

It is an additive category. Every map in it is a quantity.

Invertibility ^{implies means} locally trivially.

But Picard gp's are usually talked about in petit toposes

Look at grand toposes?

[[Felix Bar: .. non commutativity .. geom. as..]]

I talk about comm/non-comm in the infinitesimal

Stone von Neumann: irred rep's of canon comm relations:

The only non-comm. is due to Leibniz rule.

The opposite of synthetic, i.e. namely algebraic

If A is a comm alg. with a 2nd order diff operator \square

(There is a 2nd order Leibniz rule) e.g. \square Laplacian

There is a binary op'n $*$ ^{on A} that can be defined

$$\frac{1}{2} [f \square g - g \square f] \quad \text{or something}$$

k -bilinear, is a derivation in each variable separately,

and symmetric. This is equiv to a metric on

$\text{Spec } A$, by identifying tangent vectors with derivations

Can define

$$A \xrightarrow[\quad]{\substack{P \\ Q}} \text{Hom}_k(A, A)$$

$$Q_f(g) = f \cdot g$$

$$P_f(g) = f * g$$

$$[P_f, Q_g] = \hbar Q_{f \cdot g} \quad \text{Canonical commutation rel'n.}$$

(So whenever you have e.g. wave eq'n...)