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Nomenclature

\hat{R}	Universal gas constant $\approx 8.314 \text{ J/kg-K}$	F	Helmholtz energy
λ	Lagrange multiplier	G	Gibbs free-energy
\mathcal{L}	The Lagrangian	h'	Sum of contributions from each species towards enthalpy constraint, i.g. $\sum_{j=1}^{\text{NS}} \mathcal{N}_j H_j^\circ$
\mathcal{N}	Total number of moles in the system	h'_0	User-specified enthalpy
\mathcal{N}_j	Number of moles of species j in the system	H_j°	Standard state enthalpy of species j
μ_j	Chemical potential of species j	p_{ref}	Reference temperature used to compute chemical potential
μ_j°	Standard state chemical potential of species j	q_j	Charge of species j (0 for neutrals, 1 for ions, -1 for electrons)
ρ	Density (kg/m^3)	T	Temperature (K)
a_{ij}	Stoichiometric coefficients, e.g. number of atoms of element i in species j	u'	Sum of contributions from each species towards internal energy constraint, i.g. $\sum_{j=1}^{\text{NS}} \mathcal{N}_j U_j^\circ$
b_i	Sum of contributions from each species towards constraint for element i , i.g. $\sum_{j=1}^{\text{NS}} a_{ij} \mathcal{N}_j$	u'_0	User-specified internal energy
b_i°	Specified number of moles of element i	U_j°	Standard state internal energy of species j
$c_{p,j}^\circ$	Standard state specific heat of species j	V	Volume (m^3)
$c_{v,j}^\circ$	Standard state specific heat of species j	MW_j	Molecular weight of species j

1 Minimization Procedures

1.1 Constraints for Both Procedures

There are two constraints that are used the most in both energy minimization procedures. The first, equation [1.1], is always used while the charge constraints in equation [1.2] is only used when ions are present:

$$\sum_{j=1}^{\text{NS}} a_{ij} \mathcal{N}_j - b_i^\circ = 0 \quad (1.1)$$

$$\sum_{j=1}^{\text{NS}} q_j \mathcal{N}_j = 0 \quad (1.2)$$

Since these are used in abundance for both procedures, their constraints are placed in the Lagrangian term. Other constraints (such specified volume, temperature, internal energy, and entropy) are given their rows in the Newton solve without addition into the Lagrangian.

Equation [1.1] says that the number of moles of element i must remain constant during the minimization process, as atoms can not be created or destroyed. b_i° is the number of moles of element i before minimizing, and a_{ij} is the stoichiometric coefficient, e.g. how many atoms of element i are in species j .

Equation [1.2] says that the charge of the system must remain neutral, i.e., if cations are formed, than there must be an equal number of free electrons in the gas to balance it out.

1.2 The Lagrangian

The Lagrangian is defined as;

$$\mathcal{L} = f + \sum_j \lambda_j g_j \quad (1.3)$$

where f is the function we want to minimize, λ_j are Lagrange multipliers, and g_j are constraint functions. Instead of minimizing just f , we aim to minimize the Lagrangian \mathcal{L} .

1.3 Newton-Raphson Iterative Method

A Newton-Raphson iterative procedure is used to minimize the energies. Since this is essentially a root finding system of equations, and we are seeking the solution update, the method looks like the following:

$$\sum_{i=1}^N \frac{\partial f}{\partial x_i} \Delta x_i = -f(x) \quad (1.4)$$

Here, $f(x)$ is the function that is being minimized. In the case of Gibbs and Helmholtz energy minimization, f takes the form of $\partial \mathcal{L} / \partial x$, with x being non-linear variables used to guide the convergence of the system. The non-linear correction variables are of the form $\Delta \ln \mathcal{N}_j$, $\Delta \ln \mathcal{N}$, $\Delta \ln T$, and $\pi_{i,q} = \Delta \pi_{i,q}$. For the last correction variable, it is argued that setting π to 0 at the beginning of each update does not influence the solution for the non-reduced equations, so the Δ is dropped for that term. This will influence the form of the update equations, which will be talked about in due time.

1.4 Gibbs Minimization

The Gibbs energy is a function of temperature, pressure, and composition of a gas:

$$G = G(T, p, \mathcal{N}_1, \mathcal{N}_2, \dots, \mathcal{N}_{NS}) = \sum_{j=1}^{NS} \mathcal{N}_j \mu_j \quad (1.5)$$

Where NS is the number of gas species, \mathcal{N}_j is the number of moles of species j , and μ_j is the chemical potential of species j defined by:

$$\mu_j = \frac{\partial G}{\partial \mathcal{N}_j} = \mu_j^\circ + \hat{R}T \left[\ln \left(\frac{p}{p_{ref}} \right) + \ln \left(\frac{\mathcal{N}_j}{\mathcal{N}} \right) \right] \quad (1.6)$$

Here, \hat{R} is the universal gas constant, T and p are the temperature and pressure of the system, p_{ref} is the reference pressure (usually takes as 101,325 Pa, or 1 bar depending on the literature), and \mathcal{N} is the total number of moles in the system. In order to find chemical equilibrium, the derivative of G wrt \mathcal{N}_j is set to 0 for all species. However, a number of constraints need to be added in.

1.4.1 Constraints

Firstly, the elemental constraint from equation [1.1] is added into the Lagrangian for our Gibbs function. If ions are present, the charge constraint is also put in. Therefore, our Lagrangian is:

$$\mathcal{L} = \sum_{j=1}^{NS} \mu_j \mathcal{N}_j + \sum_{i=1}^{NE} \lambda_i \left[\sum_{j=1}^{NS} a_{ij} \mathcal{N}_j - b_i^\circ \right] + \lambda_q \left[\sum_{j=0}^{NS} q_j \mathcal{N}_j \right] \quad (1.7)$$

The constraints available for the Gibbs energy are:

- Hold temperature (T) and pressure (P) constant.
- Hold enthalpy (H) and pressure (P) constant.
- Hold entropy (S) and pressure (P) constant.

For (TP) minimization, only the elemental and charge (if present) constraints are necessary. For (HP) minimization, the constraint is defined by:

$$h' - h'_0 = 0 \quad (1.8)$$

Where h'_0 is the specified enthalpy, and h' is solved for during the minimization process as:

$$h' = \sum_{j=1}^{NS} \mathcal{N}_j H_j^\circ \quad (1.9)$$

For (SP) minimization, the constraint is defined by:

$$s' - s'_0 = 0 \quad (1.10)$$

Where s'_0 is the specified entropy and s' is solved for in the minimization process as:

$$s' = \sum_{j=1}^{NS} \mathcal{N}_j S_j \quad (1.11)$$

where:

$$S_j = S_j^\circ - \hat{R} \ln \frac{\mathcal{N}_j}{\mathcal{N}} - \hat{R} \ln p \quad (1.12)$$

1.4.2 Derivatives

We now take the derivatives of \mathcal{L} wrt \mathcal{N}_j , and each λ and denote them as f :

$$f_1 = \frac{\partial \mathcal{L}}{\partial \mathcal{N}_j} = \mu_j + \sum_{i=1}^{NE} \lambda_i a_{ij} + \lambda_q q_j = 0 \quad (1.13)$$

$$f_2 = \frac{\partial \mathcal{L}}{\partial \lambda_i} = \sum_{j=1}^{NS} a_{ij} \mathcal{N}_j - b_i = 0 \quad (1.14)$$

$$f_3 = \frac{\partial \mathcal{L}}{\partial \lambda_q} = \sum_{j=0}^{NS} q_j \mathcal{N}_j = 0 \quad (1.15)$$

These three equations are our minimization functions. Through use of the aforementioned Newton-Raphson method, we can solve these equations. Constraint equations for enthalpy and entropy are:

$$f_4 = \frac{h' - h'_0}{\hat{R}T} = \frac{\sum_{j=1}^{NS} \mathcal{N}_j H_j^\circ - h'_0}{\hat{R}T} \quad (1.16)$$

$$f_5 = \frac{s' - s'_0}{\hat{R}} = \frac{\sum_{j=1}^{NS} \mathcal{N}_j S_j - s'_0}{\hat{R}} \quad (1.17)$$

1.4.3 Derivatives of f_1

First, we must do some algebraic manipulation. We start by expanding equation 1.13 with equation 1.6:

$$\frac{\partial \mathcal{L}}{\partial \mathcal{N}_j} = \mu_j^\circ + \hat{R}T \left[\ln \left(\frac{p}{p_{ref}} \right) + \ln \left(\frac{\mathcal{N}_j}{\mathcal{N}} \right) \right] + \sum_{i=1}^{NE} \lambda_i a_{ij} + \lambda_q q_j = 0 \quad (1.18)$$

Dividing through by $\hat{R}T$ and setting $\pi = -\lambda/\hat{R}T$ gives:

$$f_1 = \frac{\mu_j^\circ}{\hat{R}T} + \ln \frac{p}{p_{ref}} + \ln \mathcal{N}_j - \ln \mathcal{N} - \sum_{i=1}^{NE} \pi_i a_{ij} - \pi_q q_j = 0 \quad (1.19)$$

We now need to take the partial derivatives of our functions wrt our non-linear variables as well as each π , and then multiple it by the correction variables. We start with equation 1.19:

$$\frac{\partial f_1}{\partial [\ln \mathcal{N}_j]} \Delta \ln \mathcal{N}_j = \Delta \ln \mathcal{N}_j \quad (1.20)$$

$$\frac{\partial f_1}{\partial [\ln \mathcal{N}]} \Delta \ln \mathcal{N} = -\Delta \ln \mathcal{N} \quad (1.21)$$

$$\frac{\partial f_1}{\partial [\ln T]} \Delta \ln T = -\frac{H_j^\circ}{\hat{R}T} \Delta \ln T \quad (1.22)$$

$$\frac{\partial f_1}{\partial [\pi_i]} \pi_i = -a_{ij} \pi_i \quad (1.23)$$

$$\frac{\partial f_1}{\partial [\pi_q]} \pi_q = -q_j \pi_q \quad (1.24)$$

Combining these into the form of equation 1.4 gives

$$\Delta \ln \mathcal{N}_j - \Delta \ln \mathcal{N} - \sum_{i=1}^{NE} a_{ij} \pi_i - q_j \pi_q - \frac{H_j^\circ}{\hat{R}T} \Delta \ln T = \mu_j + \sum_{i=1}^{NE} \pi_i a_{ij} + \pi_q q_j \quad (1.25)$$

Arguments are made that $\pi_i = 0$ at the start of every iteration, so this equation becomes:

$$\Delta \ln \mathcal{N}_j - \Delta \ln \mathcal{N} - \sum_{i=1}^{NE} a_{ij} \pi_i - q_j \pi_q - \frac{H_j^\circ}{\hat{R}T} \Delta \ln T = -\frac{\mu_j}{\hat{R}T}$$

(1.26)

There will be one of these equation for each species.

1.4.4 Derivatives of f_2

$$\frac{\partial f_2}{\partial [\ln \mathcal{N}_j]} \Delta \ln \mathcal{N}_j = a_{ij} \mathcal{N}_j \Delta \ln \mathcal{N}_j \quad (1.27)$$

$$\frac{\partial f_2}{\partial [\ln \mathcal{N}]} \Delta \ln \mathcal{N} = 0 \quad (1.28)$$

$$\frac{\partial f_2}{\partial [\ln T]} = 0 \quad (1.29)$$

$$\frac{\partial f_2}{\partial [\pi_i]} \pi_i = 0 \quad (1.30)$$

$$\frac{\partial f_2}{\partial [\pi_q]} \pi_q = 0 \quad (1.31)$$

Which gives:

$$\sum_{j=1}^{NS} a_{ij} \mathcal{N}_j \Delta \ln \mathcal{N}_j = b_i^o - \sum_{j=1}^{NS} a_{ij} \mathcal{N}_j$$

(1.32)

1.4.5 Derivatives of f_3

$$\frac{\partial f_3}{\partial [\ln \mathcal{N}_j]} \Delta \ln \mathcal{N}_j = q_j \mathcal{N}_j \Delta \ln \mathcal{N}_j \quad (1.33)$$

$$\frac{\partial f_3}{\partial [\ln \mathcal{N}]} \Delta \ln \mathcal{N}_j = 0 \quad (1.34)$$

$$\frac{\partial f_3}{\partial [\ln T]} = 0 \quad (1.35)$$

$$\frac{\partial f_3}{\partial [\pi_i]} \pi_i = 0 \quad (1.36)$$

$$\frac{\partial f_3}{\partial [\pi_q]} \pi_q = 0 \quad (1.37)$$

Giving:

$$\sum_{j=1}^{NS} q_j \mathcal{N}_j \Delta \ln \mathcal{N}_j = - \sum_{j=1}^{NS} q_i \mathcal{N}_j$$

(1.38)

1.4.6 Derivatives of f_4

$$\frac{\partial f_4}{\partial [\ln \mathcal{N}_j]} \Delta \ln \mathcal{N}_j = \quad (1.39)$$

$$\frac{\partial f_4}{\partial [\ln \mathcal{N}]} \Delta \ln \mathcal{N}_j = 0 \quad (1.40)$$

$$\frac{\partial f_4}{\partial [\ln T]} = \quad (1.41)$$

$$\frac{\partial f_4}{\partial [\pi_i]} \pi_i = 0 \quad (1.42)$$

$$\frac{\partial f_4}{\partial [\pi_q]} \pi_q = 0 \quad (1.43)$$

1.4.7 Derivatives of f_5

$$\frac{\partial f_5}{\partial [\ln \mathcal{N}_j]} \Delta \ln \mathcal{N}_j = \quad (1.44)$$

$$\frac{\partial f_5}{\partial [\ln \mathcal{N}]} \Delta \ln \mathcal{N}_j = 0 \quad (1.45)$$

$$\frac{\partial f_5}{\partial [\ln T]} = \quad (1.46)$$

$$\frac{\partial f_5}{\partial [\pi_i]} \pi_i = 0 \quad (1.47)$$

$$\frac{\partial f_5}{\partial [\pi_q]} \pi_q = 0 \quad (1.48)$$

1.5 Helmholtz Minimization

The Helmholtz energy is defined as:

$$F = G - pV \quad (1.49)$$

Where G is the Gibbs free energy, p is the pressure, and V is the volume. Substitution of G yields:

$$F = \sum_{j=1}^{NS} \mu_j \mathcal{N}_j - pV \quad (1.50)$$

Where the chemical potential is redefined as:

$$\mu_j = \mu_j^\circ + \hat{R}T \ln \left(\frac{\mathcal{N}_j R' T}{V} \right) \quad (1.51)$$

and $R' = \hat{R} \cdot 10^{-5}$ is the same as dividing \hat{R} by the reference pressure of 1 bar. The constraints available for the Helmholtz energy are:

- Hold temperature (T) and volume (V) constant.
- Hold internal energy (U) and volume (V) constant.
- Hold entropy (S) and volume (V) constant.

For (TV) minimization, only the elemental and charge (if present) constraints are necessary. For (UV) minimization, the constraint is defined by:

$$u' - u'_0 = 0 \quad (1.52)$$

Where u'_0 is the specified internal energy, and u' is solved for during the minimization process as:

$$u' = \sum_{j=1}^{NS} \mathcal{N}_j U_j^\circ = u'_0 \quad (1.53)$$

For (SV) minimization, the constraint is defined by:

$$s' - s'_0 = 0 \quad (1.54)$$

Where s'_0 is the specified entropy and s' is solved for during the minimization process as:

$$s' = \sum_{j=1}^{NS} \mathcal{N}_j S_j \quad (1.55)$$

where:

$$S_j = S_j^\circ - \hat{R} \ln \mathcal{N}_j - \hat{R} \ln \left(\frac{R'T}{V} \right) \quad (1.56)$$

While you may be able to include the constraints for internal energy and entropy inside the Lagrangian, we will exclude them from the Lagrangian and use our own residual equation for them. The Lagrangian for the Helmholtz energy minimization process is then defined as:

$$\mathcal{L} = \sum_{j=1}^{NS} \mu_j \mathcal{N}_j - pV + \sum_{i=1}^{NE} \lambda_i \left[\sum_{j=1}^{NS} a_{ij} \mathcal{N}_j - b_i^\circ \right] + \lambda_q \left[\sum_{j=1}^{NS} q_j \mathcal{N}_j \right] \quad (1.57)$$

We first find our function f that are equal to 0 by taking the derivatives wrt \mathcal{N}_j , λ_i , and λ_q :

$$f_1 = \frac{\partial \mathcal{L}}{\partial \mathcal{N}_j} = \mu_j + \sum_{i=1}^{NE} \lambda_i a_{ij} + \lambda_q q_j = 0 \quad (1.58)$$

$$f_2 = \frac{\partial \mathcal{L}}{\partial \lambda_i} = \sum_{j=1}^{NS} a_{ij} \mathcal{N}_j - b_i^\circ = 0 \quad (1.59)$$

$$f_3 = \frac{\partial \mathcal{L}}{\partial \lambda_q} = \sum_{j=1}^{NS} q_j \mathcal{N}_j - 0 = 0 \quad (1.60)$$

The negative of these functions ($-f$) are played on the RHS of the Newton solver. We now form the second derivatives that define our Jacobian in the Newton step. We expand the chemical potential term using equation 1.51 and non-dimensionalize the equation by $\hat{R}T$. We make sure to convert these equations into a form that can be differentiated by our non-linear variables (i.e. $N_j = \exp(\ln(\mathcal{N}_j))$). We also set $\pi = -\lambda/\hat{R}T$. The derivatives for f_1 are then:

$$\frac{\partial f_1}{\partial [\ln \mathcal{N}_j]} \Delta \ln \mathcal{N}_j = \Delta \ln \mathcal{N}_j \quad (1.61)$$

$$\frac{\partial f_1}{\partial [\ln T]} \Delta \ln T = -\frac{U_j^\circ}{\hat{R}T} \ln T \quad (1.62)$$

$$\frac{\partial f_1}{\partial \pi_i} \pi_i = -a_{ij} \pi_i \quad (1.63)$$

$$\frac{\partial f_1}{\partial \pi_q} \pi_q = -q_j \pi_q \quad (1.64)$$

The derivation for equation [1.62] is carried out in more depth in section [2.2]. Therefore, our final result is:

$$\Delta \ln \mathcal{N}_j - \sum_{i=1}^{NE} a_{ij} \pi_i - q_j \pi_q - \frac{U_j^\circ}{\hat{R}T} \Delta \ln T = -\frac{\mu_j}{\hat{R}T}$$

There will be NS number of these equations in our system. Next, we take the derivatives of f_2 :

$$\frac{\partial f_2}{\partial [\ln \mathcal{N}_j]} \Delta \ln \mathcal{N}_j = \mathcal{N}_j a_{i,j} \Delta \ln \mathcal{N}_j \quad (1.65)$$

$$\frac{\partial f_2}{\partial [\ln T]} \Delta \ln T = 0 \quad (1.66)$$

$$\frac{\partial f_2}{\partial \pi_i} \pi_i = 0 \quad (1.67)$$

$$\frac{\partial f_2}{\partial \pi_q} \pi_q = 0 \quad (1.68)$$

Which gives:

$$\boxed{\sum_{j=1}^{\text{NS}} \mathcal{N}_j a_{ij} \Delta \ln \mathcal{N}_j = b_i^\circ - \sum_{j=1}^{\text{NS}} a_{ij} \mathcal{N}_j}$$

There will be NE of these equations. Finally, for f_3 we get:

$$\frac{\partial f_3}{\partial [\ln \mathcal{N}_j]} \Delta \ln \mathcal{N}_j = \mathcal{N}_j q_j \Delta \ln \mathcal{N}_j \quad (1.69)$$

$$\frac{\partial f_3}{\partial [\ln T]} \Delta \ln T = 0 \quad (1.70)$$

$$\frac{\partial f_3}{\partial \pi_i} \pi_i = 0 \quad (1.71)$$

$$\frac{\partial f_3}{\partial \pi_q} \pi_q = 0 \quad (1.72)$$

Which gives only one equation:

$$\boxed{\sum_{j=1}^{\text{NS}} \mathcal{N}_j q_j \Delta \ln \mathcal{N}_j = - \sum_{j=1}^{\text{NS}} q_j \mathcal{N}_j}$$

Finally, we look towards the constraints. For internal energy, we have:

$$f_4(\ln \mathcal{N}_j, \ln T) = \sum_{j=1}^{\text{NS}} \mathcal{N}_j U_j^\circ - u'_0 = u' - u'_0 \quad (1.73)$$

Using identity [2.1] helps with the temperature derivative. The derivatives are then:

$$\frac{\partial f_4}{\partial [\ln \mathcal{N}_j]} \Delta \ln \mathcal{N}_j = \frac{\mathcal{N}_j U_j^\circ}{\hat{R}T} \Delta \ln \mathcal{N}_j \quad (1.74)$$

$$\frac{\partial f_4}{\partial [\ln T]} \Delta \ln T = \sum_{j=1}^{\text{NS}} [T \mathcal{N}_j c_{v,j}^\circ] \Delta \ln T \quad (1.75)$$

$$\frac{\partial f_4}{\partial \pi_i} \pi_i = 0 \quad (1.76)$$

$$\frac{\partial f_4}{\partial \pi_q} \pi_q = 0 \quad (1.77)$$

After non-dimensionalizing, we are left with:

$$\sum_{j=1}^{\text{NS}} \frac{\mathcal{N}_j U_j^\circ}{\hat{R}T} \Delta \ln \mathcal{N}_j + \sum_{j=1}^{\text{NS}} \frac{\mathcal{N}_j c_{v,j}^\circ}{R} \Delta \ln T = \frac{u'_0 - u'}{\hat{R}T}$$

For the entropy constraint we have:

$$f_5(\ln \mathcal{N}_j, \ln T) = \sum_{j=1}^{\text{NS}} \mathcal{N}_j S_j - s'_0 = s' - s'_0 \quad (1.78)$$

Again, identity [2.1] assists us. The derivates are:

$$\frac{\partial f_5}{\partial [\ln \mathcal{N}_j]} \Delta \ln \mathcal{N}_j = \mathcal{N}_j [S_j - \hat{R}] \Delta \ln \mathcal{N}_j \quad (1.79)$$

$$\frac{\partial f_5}{\partial [\ln T]} \Delta \ln T = \mathcal{N}_j c_{v,j}^\circ \Delta \ln T \quad (1.80)$$

$$\frac{\partial f_5}{\partial \pi_i} \pi_i = 0 \quad (1.81)$$

$$\frac{\partial f_5}{\partial \pi_q} \pi_q = 0 \quad (1.82)$$

Which, after non-dimensionalizing gives the singular equation:

$$\sum_{j=1}^{\text{NS}} \mathcal{N}_j \left[\frac{S_j}{\hat{R}} - 1 \right] \Delta \ln \mathcal{N}_j + \mathcal{N}_j \frac{c_{v,j}^\circ}{\hat{R}} \Delta \ln T = \frac{s'_0 - s'}{\hat{R}} \quad (1.83)$$

For convenience, they are all listed here together:

$$\Delta \ln \mathcal{N}_j - \sum_{i=1}^{\text{NE}} a_{ij} \pi_i - q_j \pi_q - \frac{U_j^\circ}{\hat{R}T} \Delta \ln T = -\frac{\mu_j}{\hat{R}T} \quad (1.84)$$

$$\sum_{j=1}^{\text{NS}} \mathcal{N}_j a_{ij} \Delta \ln \mathcal{N}_j = b_i^\circ - \sum_{j=1}^{\text{NS}} a_{ij} \mathcal{N}_j \quad (1.85)$$

$$\sum_{j=1}^{\text{NS}} \mathcal{N}_j q_j \Delta \ln \mathcal{N}_j = - \sum_{j=1}^{\text{NS}} q_j \mathcal{N}_j \quad (1.86)$$

$$\sum_{j=1}^{\text{NS}} \frac{\mathcal{N}_j U_j^\circ}{\hat{R}T} \Delta \ln \mathcal{N}_j + \sum_{j=1}^{\text{NS}} \frac{\mathcal{N}_j c_{v,j}^\circ}{R} \Delta \ln T = \frac{u'_0 - u'}{\hat{R}T} \quad (1.87)$$

$$\sum_{j=1}^{\text{NS}} \mathcal{N}_j \left[\frac{S_j}{\hat{R}} - 1 \right] \Delta \ln \mathcal{N}_j + \mathcal{N}_j \frac{c_{v,j}^\circ}{\hat{R}} \Delta \ln T = \frac{s'_0 - s'}{\hat{R}} \quad (1.88)$$

1.5.1 Reduced Helmholtz Equations

A "simplification" can be made to the above system of equations by solving equations [1.84] for $\Delta \ln \mathcal{N}_j$:

$$\Delta \ln \mathcal{N}_j = \sum_{i=1}^{\text{NE}} a_{i,j} - \frac{\mu_j}{\hat{R}T} \pi_i + q_j \pi_q + \frac{U_j^\circ}{\hat{R}T} \Delta \ln T \quad (1.89)$$

By substituting this into equations [1.85] - [1.87], we can shrink the system of equations dramatically. The resulting set of equations is:

$$\sum_{i=1}^{NE} \left[\sum_{j=1}^{NS} a_{kj} a_{ij} \mathcal{N}_j \right] \pi_i + \left[\sum_{j=1}^{NS} a_{kj} q_j \mathcal{N}_j \right] \pi_q + \left[\sum_{j=1}^{NS} a_{kj} \mathcal{N}_j \frac{U_j^\circ}{\hat{R}T} \right] \Delta \ln T = b_k^\circ - \sum_{j=1}^{NS} a_{kj} \mathcal{N}_j + \sum_{j=1}^{NS} a_{kj} \mathcal{N}_j \frac{\mu_j}{\hat{R}T} \quad (1.90)$$

$$\sum_{i=1}^{NE} \left[\sum_{j=1}^{NS} a_{ij} q_j \mathcal{N}_j \right] \pi_i + \left[\sum_{j=1}^{NS} q_j^2 \mathcal{N}_j \right] \pi_q + \left[\sum_{j=1}^{NS} q_j \mathcal{N}_j \frac{U_j^\circ}{\hat{R}T} \right] \Delta \ln T = \sum_{j=1}^{NS} q_j \mathcal{N}_j \frac{\mu_j}{\hat{R}T} - \sum_{j=1}^{NS} q_j \mathcal{N}_j \quad (1.91)$$

$$\begin{aligned} \sum_{i=1}^{NE} \left[\sum_{j=1}^{NS} a_{ij} \mathcal{N}_j \frac{U_j^\circ}{\hat{R}T} \right] \pi_i + \left[\sum_{j=1}^{NS} q_j \mathcal{N}_j \frac{U_j^\circ}{\hat{R}T} \right] \pi_q + & \left[\sum_{j=1}^{NS} \mathcal{N}_j \left(\frac{U_j^\circ}{\hat{R}T} \right)^2 + \sum_{j=1}^{NS} \mathcal{N}_j \frac{c_{v,j}^\circ}{\hat{R}} \right] \Delta \ln T \\ & = \frac{u'_0 - u'}{\hat{R}T} + \sum_{j=1}^{NS} \mathcal{N}_j \frac{U_j^\circ}{\hat{R}T} \frac{\mu_j}{\hat{R}T} \end{aligned} \quad (1.92)$$

$$\begin{aligned} \sum_{i=1}^{NE} \left[\sum_{j=1}^{NS} a_{ij} \mathcal{N}_j \left(\frac{S_j}{\hat{R}} - 1 \right) \right] \pi_i + \left[\sum_{j=1}^{NS} q_j \mathcal{N}_j \left(\frac{S_j}{\hat{R}} - 1 \right) \right] \pi_q \\ + \left[\sum_{j=1}^{NS} \mathcal{N}_j \left(\frac{S_j}{\hat{R}} - 1 \right) \frac{U_j^\circ}{\hat{R}T} + \sum_{j=1}^{NS} \mathcal{N}_j \frac{c_{v,j}^\circ}{\hat{R}} \right] \Delta \ln T = \frac{s'_0 - s'}{\hat{R}} + \sum_{j=1}^{NS} \mathcal{N}_j \frac{\mu_j}{\hat{R}T} \left(\frac{S_j}{\hat{R}} - 1 \right) \end{aligned} \quad (1.93)$$

There will be NE number of equations 1.90, and one of 1.91 if charge is added, and one of 1.92/1.93 depending on which one you made need.

The equations are colored according to when they are needed. **Red coloring** denotes terms that arise from the elemental constraint condition. These are included no matter what. **Blue** denotes terms that are added when you include the charge constraint. **Violet** terms are added for both UV and SV constraints. **orange** and **teal** are the rows added for either internal energy or entropy constraints.

1.6 Useful Relationships

$$\frac{c_v^\circ}{\hat{R}} = \frac{c_{p,j}^\circ}{\hat{R}} - 1 \quad (1.94)$$

$$\frac{\mu_j^\circ}{\hat{R}T} = \frac{H_j^\circ}{\hat{R}T} - \frac{S_j^\circ}{\hat{R}} \quad (1.95)$$

$$\frac{U_j^\circ}{\hat{R}T} = \frac{H_j^\circ}{\hat{R}T} - 1 \quad (1.96)$$

2 Standard State Chemical Potential Derivatives

Some important derivatives are derived here. First, we state a chain rule identity that is very useful:

$$\frac{\partial}{\partial T} = \frac{\partial[\ln(T)]}{\partial T} \frac{\partial}{\partial[\ln(T)]} \rightarrow \frac{\partial}{\partial[\ln(T)]} = T \frac{\partial}{\partial T} \quad (2.1)$$

Then, taking the derivative of μ_j^0 gives:

$$\frac{\partial}{\partial[\ln(T)]} \left(\frac{\mu_j^0}{RT} \right) = T \frac{\partial}{\partial T} \left(\frac{\mu_j^0}{RT} \right) \quad (2.2)$$

We expand μ_j^0/RT as:

$$\frac{\mu_j^0}{RT} = \frac{H_j^0}{RT} - \frac{S_j^0}{R} \quad (2.3)$$

$$\frac{\mu_j^0}{RT} = \frac{U_j^0}{RT} - \frac{S_j^0}{R} + 1 \quad (2.4)$$

2.1 Gibbs Derivative

First, we take the derivative of [2.3] for Gibbs minimization.

$$T \frac{\partial}{\partial T} \left(\frac{H_j^0}{RT} - \frac{S_j^0}{R} \right) = T \left(\frac{\partial_T(H_j^0) \cdot RT - H_j^0 \cdot R}{(RT)^2} - \frac{\partial_T(S_j^0)}{R} \right) \quad (2.5)$$

Here, ∂_T denotes $(\partial/\partial T)$. If:

$$c_{p,j}^0 = \frac{\partial H_j^0}{\partial T}, \quad \text{and} \quad \frac{c_{p,j}^0}{T} = \frac{\partial S_j^0}{\partial T} \quad (2.6)$$

This reduces to:

$$T \left[\frac{c_{p,j}^0}{RT} - \frac{H_j^0}{RT^2} - \frac{c_{p,j}^0}{RT} \right] = -\frac{H_j^0}{RT} \quad (2.7)$$

Therefore:

$$\frac{\partial}{\partial[\ln(T)]} \left(\frac{\mu_j^0}{RT} \right) = -\frac{H_j^0}{RT} \quad (2.8)$$

2.2 Helmholtz Derivative

Using equation [2.4] gives us:

$$T \frac{\partial}{\partial T} \left(\frac{U_j^0}{RT} + 1 - \frac{S_j^0}{R} \right) = T \left(\frac{\partial_T(U_j^0) \cdot RT - U_j^0 \cdot R}{(RT)^2} - \frac{\partial_T(S_j^0)}{R} \right) \quad (2.9)$$

Recalling $\partial_T(S_j^0)$ from equation [2.6] and using:

$$c_{v,j}^0 = \frac{\partial U_j^0}{\partial T} \quad (2.10)$$

Equation [2.9] reduced to:

$$T \left[\frac{c_{v,j}^0}{RT} - \frac{U_j^0}{RT^2} - \frac{c_{p,j}^0}{RT} \right] \quad (2.11)$$

Since:

$$\frac{c_{v,j}^0}{R} - \frac{c_{p,j}^0}{R} = -1 \quad (2.12)$$

The final derivative is:

$$\boxed{\frac{\partial}{\partial[\ln(T)]} \left(\frac{\mu_j^0}{RT} \right) = -\frac{U_j^0}{RT} - 1} \quad (2.13)$$

3 Coding Tricks