

State-Dependent Gravity from Modular Information: A KMS/FDT Linear-Response Framework (Conditional)

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(Dated:)

We present a *conditional* information-theoretic framework in which finite local information capacity produces a *state-dependent* gravitational response. The key working assumption replaces macroscopic Clausius language with a **KMS-normalized linear-response (A2-KMS) hypothesis**: in the small-wedge, MI/moment-kill projector channel, Bisognano–Wichmann (BW) KMS structure and the fluctuation–dissipation theorem (FDT) fix the sign and normalization of the modular susceptibility to $\mathcal{O}(\ell^4)$, with curvature/contact remainders $\mathcal{O}(\ell^6)$. We *define* a dimensionless state variable ε via flat-space modular response, $\delta\langle K_{\text{sub}} \rangle = (2\pi C_T I_{00}) \ell^4 \delta\varepsilon + \mathcal{O}(\ell^6)$, and *map* ε to a weak-field coupling through A2-KMS, yielding $\delta G/G = -\beta \delta\varepsilon$ with $\beta \equiv 2\pi C_T I_{00}$. A geometric normalization yields a universal weak-field prefactor $5/12 = (4/3) \times (5/16)$, implying $\mu(\varepsilon) = 1/(1 + \frac{5}{12}\varepsilon)$ and $a_0 = \frac{5}{12} \Omega_\Lambda^2 c H_0$. The scheme-invariant mapping $\Omega_\Lambda = \beta f c_{\text{geo}} \approx 0.685$ (conservative $\pm 5\%$ from shared systematics in β) preserves EM/GW distances (distance sector kept GR-like). We solve $\varepsilon(a)$ and linear growth $D(a)$ as a KMS/FDT-constrained *fixed point* and report *entropy-constrained variational bounds* on S_8 (late-/early-loaded profiles bracket the admissible range), rather than a single fit; our illustrative baseline value lies within this band. A retarded KMS susceptibility with positive integrated kernel enforces $d\varepsilon/d\ln a \geq 0$ (FDT positivity), and the normalization is fixed by $\int \varepsilon d\ln a = \Omega_\Lambda$. Substrate *structural consistency checks* (HQTFIM and Gaussian chains) confirm algebraic ingredients (first-law channel, constant+log size trend, MI/moment-kill plateau, FDT-positivity) but are *not* 4D curved-spacetime surrogates. We give explicit falsifiers, conservative uncertainties, and a limitations box (safe-window viability, CHM-vs-half-space KMS error $\sim \mathcal{O}((\ell/L_{\text{curv}})^2)$, environment gate microphysics, quantum-classical bridge). This is an exploratory, testable, and conditional framework rather than a phenomenological fit.

I. SCOPE, WORKING ORDER, AND LIMITATIONS (READ FIRST)

Working order. Throughout, “working order” means we isolate the isotropic ℓ^4 contribution in the MI/moment-kill projector channel and treat curvature/contact corrections as $\mathcal{O}(\ell^6)$.

Safe window (existence is model dependent). We assume a nonempty range ℓ obeying

$$\epsilon_{\text{UV}} \ll \ell \ll \min\{L_{\text{curv}}, \lambda_{\text{mfp}}, m_i^{-1}\}.$$

In halos with $L_{\text{curv}} \sim 10$ Mpc, a plausible late-time band is $\ell \in [1, 100]$ pc; this window can be *absent* in dense regions (star-forming zones, cluster cores).

KMS applicability (CHM vs. half-space). Exact BW KMS analyticity holds for half-spaces; CHM diamonds approximate it in the safe window. The fractional KMS deviation scales as $\mathcal{O}((\ell/L_{\text{curv}})^2)$ (App. ??).

State assumption (Hadamard). Throughout we assume locally Hadamard states so that short-distance correlators match Minkowski up to curvature-suppressed terms. If data reveal departures from Hadamard short-distance structure in the projector channel, the framework is *falsified*.

Distances kept GR-like. We enforce $\alpha_M \simeq 0$ in the distance sector; null geometry and EM/GW distances are unmodified at working order.

Environment gate is illustrative. The gate $F_g(g/a_0)$ is a minimal compliance envelope: $F_g \rightarrow 0$ in strong fields (Solar System), $F_g \rightarrow 1$ in weak fields; a microscopic derivation is future work.

Substrate tests are algebraic checks. HQTFIM/Gaussian runs test the algebraic structure (first-law channel, constant+log trend, plateau, FDT-positivity). They are *not* physical surrogates for 4D curved spacetime.

Falsifiers and uncertainties. We list sharp falsifiers (Sec. ??) and adopt a conservative $\pm 5\%$ uncertainty on β (shared systematics), with angle-invariance presented as a *null* residual test rather than a precision claim.

II. A2-KMS HYPOTHESIS (DEFINITION, CHANNEL, AND POSITIVITY)

a. BW recap. The Minkowski vacuum restricted to a Rindler half-space is a KMS state at inverse temperature $\beta_{\text{KMS}} = 2\pi/\kappa$ with respect to boost flow (Bisognano–Wichmann).

A2-KMS (boxed).

Hypothesis 1 (A2-KMS (working order)). *In the MI/moment-kill projector channel for small CHM diamonds in a safe window, the wedge state inherits BW KMS analyticity up to $\mathcal{O}((\ell/L_{\text{curv}})^2)$. The linear-response susceptibility relating modular perturbations to boost-energy flux is fixed by the KMS two-point function, is positive (FDT), and its finite ℓ^4 coefficient equals the flat-space value at working order:*

$$\delta\langle K_{\text{sub}} \rangle = (2\pi C_T I_{00}) \ell^4 \delta\varepsilon + \mathcal{O}(\ell^6), \quad \frac{\delta G}{G} = -\beta \delta\varepsilon, \quad \beta \equiv 2\pi C_T I_{00}.$$

Boxed reminder (constitutive & falsifiable). Equation $\delta G/G = -\beta \delta\varepsilon$ is a **constitutive closure at working order** in the BW/KMS projector channel. It is **not** a macroscopic thermodynamic law and is **falsified** by any of: persistent $\ell^4 \log \ell$ residuals in the MI/moment-kill channel, violation of the EM/GW distance bound, or an Ω_Λ inconsistent with $\beta f c_{\text{geo}}$. *Remarks.* (i) Exact KMS is half-space; diamond validity is approximate and quantified in App. ???. (ii) FDT positivity enforces $\Delta S \geq 0$ in this channel without invoking macroscopic heat. (iii) “Temperature” is the KMS normalization for boost flow, not a literal bath. (iv) **State class:** we assume locally Hadamard states so that the short-distance two-point structure is Minkowski-like; large deviations would invalidate the working-order reduction and thus *falsify* the framework.

III. DEFINITION VS. MAPPING (SEPARATION OF ROLES)

a. Definition (flat-space QFT). We define $\varepsilon(x)$ by the MI-subtracted modular response in flat space:

$$\delta\langle K_{\text{sub}}(\ell) \rangle = \underbrace{(2\pi C_T I_{00})}_{\beta} \ell^4 \delta\varepsilon(x) + \mathcal{O}(\ell^6). \quad (1)$$

b. Mapping (A2-KMS). We map ε to a response via A2-KMS:

$$\frac{\delta G}{G} = -\beta \delta\varepsilon, \quad \beta = 2\pi C_T I_{00}. \quad (2)$$

The roles are distinct; no circularity arises.

IV. QFT INPUT: $\beta = 2\pi C_T I_{00}$

We evaluate β via four independent routes sharing only OP/CHM conventions and the MI+moment-kill projector: (a) real-space CHM kernel; (b) spectral/Bessel (momentum-space); (c) Euclidean time-slicing; (d) replica finite-difference. Angle invariance is presented as a *null* residual test (identity by construction). Conservatively,

$$\beta = 0.02086 \pm 0.00105 \quad (5\% \text{ shared systematics}). \quad (3)$$

Scheme/angle invariance. Physical predictions use $\mathcal{C}_\Omega \equiv f(\theta) c_{\text{geo}}(\theta)$, which is analytically angle-invariant; we show residuals as a null check rather than a precision measurement (Sec. ??).

V. GEOMETRIC NORMALIZATION AND BACKGROUND MAPPING

With the continuous-angle normalization (Sec. ??) the FRW zero mode satisfies the *scheme-invariant* mapping

$$\boxed{\Omega_\Lambda = \beta f c_{\text{geo}}} \quad \Rightarrow \quad \Omega_\Lambda \approx 0.685 \pm 0.034 \quad (\text{from } \pm 5\% \beta). \quad (4)$$

Distances are kept GR-like ($\alpha_M \simeq 0$ in the distance sector); lensing is unaltered at working order.

VI. WEAK-FIELD SECTOR: $5/12$, $\mu(\varepsilon)$ AND a_0

Coarse-graining the KMS susceptibility over the wedge family yields a universal geometric factor $5/12 = (4/3) \times (5/16)$ (App. ??). The weak-field response and static normalization read

$$\mu(\varepsilon) = \frac{1}{1 + \frac{5}{12}\varepsilon}, \quad a_0 = \frac{5}{12} \Omega_\Lambda^2 c H_0, \quad (5)$$

with the same bookkeeping that fixes the FRW zero mode. The factor $4/3$ is the isotropic null contraction in the BW channel; its universality follows from the UV ($w = 1/3$) sector governing the susceptibility (App. ??).

VII. ENTROPY-DRIVEN EVOLUTION OF $\varepsilon(a)$

a. KMS/FDT differential constraint (positivity). Let \hat{Q} denote the boost-energy flux operator in the CHM diamond and χ_{QK} the retarded susceptibility between \hat{Q} and the MI-subtracted modular generator \hat{K}_{sub} . In linear response,

$$\delta\langle\hat{Q}\rangle(a) = \int^{\ln a} d\ln a' \chi_{QK}(a, a') \delta\langle\hat{K}_{\text{sub}}\rangle(a'),$$

and FDT with KMS normalization implies $\int \chi_{QK} d\ln a' \geq 0$ in the projector channel; in curved backgrounds this positivity holds up to $\mathcal{O}((\ell/L_{\text{curv}})^2)$ corrections (App. ??). Parameterizing the (dimensionless) throughput intensity by a nonnegative functional $\mathcal{I}(a)$, we write the **entropy-driven law**

$$\boxed{\frac{d\varepsilon}{d\ln a} = \sigma(a) \mathcal{I}(a) \quad \text{with} \quad \sigma(a) \geq 0, \quad \mathcal{I}(a) \geq 0} \quad (6)$$

so that $\Delta S \geq 0 \Rightarrow d\varepsilon/d\ln a \geq 0$ (monotone). This KMS/FDT constraint is a *physical principle*, not a retrofitted choice.

b. Normalization by the background mapping. The scheme-invariant background relation fixes the total “budget”

$$\boxed{\int_{a_i}^1 \varepsilon(a) d\ln a = \Omega_\Lambda = \beta f c_{\text{geo}},} \quad (7)$$

so once $\mathcal{I}(a)$ and $\sigma(a)$ are specified by microphysics, $\varepsilon(a)$ is determined up to an initial condition $\varepsilon(a_i) \equiv \varepsilon_0 \geq 0$ (irreversibility floor).

c. Self-consistent fixed point with growth. The growth factor $D(a)$ obeys the standard linear equation with our scale-independent closure,

$$\frac{d^2 D}{d(\ln a)^2} + \left(2 + \frac{d\ln H}{d\ln a}\right) \frac{dD}{d\ln a} - \frac{3}{2} \Omega_m(a) \mu(\varepsilon(a)) D = 0, \quad (8)$$

where $\mu(\varepsilon) = 1/(1 + \frac{5}{12}\varepsilon)$ from Eq. (??). We solve Eqs. (??) and (??) together as a *fixed-point* problem under the constraints (monotonicity, budget (??), GR-like distances, and environmental gating). In practice a simple Picard or Anderson-accelerated iteration converges rapidly from Λ CDM initial $D(a)$.

d. Entropy-constrained variational bounds. Because $\mu(\varepsilon) = 1/(1 + \eta\varepsilon)$ with $\eta = 5/12$ is positive, decreasing, and convex in ε (i.e., $d\mu/d\varepsilon < 0$, $d^2\mu/d\varepsilon^2 > 0$), and because the kernel in Eq. (??) enforces $d\varepsilon/d\ln a \geq 0$ with a fixed budget $\int \varepsilon d\ln a = \Omega_\Lambda$, rearrangement/convex-order arguments imply that, for fixed constraints, the *minimum* growth (hence *minimum* S_8) is achieved by maximally *late-loaded* $\varepsilon(a)$, while the *maximum* by the most *early-loaded* profile permitted by gating. We therefore report S_8 as a *band* bracketed by these two admissible extremals; any choice like the logarithmic family (??) is illustrative and must lie within the band.

e. A minimal illustrative family (used in Sec. ??). As a concrete but non-unique realization consistent with Eq. (??), define an exposure

$$J(a) = \int^{\ln a} d\ln a' K(a, a') \Phi(a'), \quad K(a, a') \propto (a'/a)^p, \quad p \in [4, 6], \quad \Phi \geq 0, \quad (9)$$

and set

$$\varepsilon(a) = \varepsilon_0 + c_{\log} \ln\left(1 + \frac{J(a)}{J_*}\right), \quad \frac{d\varepsilon}{d\ln a} = \frac{c_{\log}}{1 + J/J_*} \frac{dJ}{d\ln a} \geq 0. \quad (10)$$

The normalization constant c_{\log} is fixed by Eq. (??). This family enforces monotonicity and the budget while leaving ε_0 and the kernel details to microphysics.

f. What is fixed vs. what remains free. *Fixed by physics:* (i) monotonicity $d\varepsilon/d\ln a \geq 0$ (KMS/FDT); (ii) total budget $\int \varepsilon d\ln a = \Omega_\Lambda$ (background mapping); (iii) strong-field recovery via environment gating in observables. *Remaining freedom:* (i) initial floor $\varepsilon_0 \geq 0$; (ii) the precise retarded kernel $K(a, a')$ and driver $\Phi(a')$ (we bracket with $p \in [4, 6]$); (iii) a scale J_* . In practice, our headline growth number $S_8 \simeq 0.788$ is *insensitive* to p within $[4, 6]$ at the $< 10^{-3}$ level, indicating limited tuning. The Hubble-ladder bounds are likewise presented as *bounds*, not fits.

VIII. STRUCTURAL CONSISTENCY CHECKS (SUBSTRATES)

We implement two independent microscopic testbeds to check the *algebraic* ingredients: (i) an interacting HQTFIM chain (exact diagonalization); (ii) a Gaussian (free-fermion) chain via correlation matrices. These confirm: (1) first-law channel in the linear window; (2) constant+log dependence of $\delta\langle K \rangle(\ell)$; (3) near-zero plateau after subtracting $[1, \log \ell]$; (4) **FDT positivity** in the projected channel (integrated susceptibility nonnegative; exact for Gaussian, numerically for HQTFIM within tolerance). These are *not* curved 4D surrogates.

IX. OBSERVATIONAL CONSEQUENCES (ILLUSTRATIVE BOUNDS)

a. Growth (entropy-constrained band). Solving the coupled system (??)–(??) to a fixed point under the constraints (monotonicity, budget, GR-like distances, and gating) yields an *interval* for S_8 bracketed by early-loaded and late-loaded $\varepsilon(a)$ profiles. Our illustrative baseline (log family, $p = 5$, $\varepsilon_0 = 0$) lies within this band at $S_8 \approx 0.788$. The band width is insensitive to kernel powers $p \in [4, 6]$ at the $< 10^{-3}$ level. No post-hoc rescaling of α_M is used; $\alpha_M \simeq 0$ is enforced in the distance sector, and modifications enter only through $\mu(\varepsilon)$ in growth, automatically respecting GW/EM bounds.

b. Hubble ladder (capped illustration). Using an environment gate $F_g(g/a_0)$ as a minimal compliance envelope (Solar-System recovery, weak-field throttling), an SH0ES-like catalog shifts $H_0: 73.0 \rightarrow 71.18$ (uncapped SN) and to 70.89 (capped SN+Cepheid). These are *bounds*, not fits; distances remain GR-like. We use F_g strictly as an *illustrative compliance envelope*; strong-field recovery can also arise from the *absence* of a safe window in high- g regions, and no microphysical claim about F_g is made here.

X. RELATION TO EFT-OF-DE (HORNDESKI) AND $f(R)$

Linearized about FRW, our closure lives in the $c_T = 1$, no-braiding corner ($\alpha_T = \alpha_B = 0$) with a single background function $\alpha_M(a) = d\ln M^2/d\ln a$ [?]. Distances are kept GR-like by setting $\alpha_M \simeq 0$ in the distance sector, while the growth sector is modified by the scale-independent $\mu(\varepsilon) = 1/(1 + \frac{5}{12}\varepsilon)$. In the quasi-static language this corresponds to $\mu(a) \neq 1$ and $\Sigma(a) \simeq 1$ (lensing unaltered). By contrast, typical $f(R)$ models induce scale-dependent $\mu(k, a)$ and nonzero slip; our mapping is scale-independent at working order and enforces GR-like lensing by construction.

XI. FALSIFIERS AND HONEST GAPS

Falsifiers. (i) Persistent $\ell^4 \log \ell$ residuals in the MI/moment-kill projector channel; (ii) GW/EM luminosity-distance ratio violating $|d_L^{\text{GW}}/d_L^{\text{EM}} - 1| \leq 5 \times 10^{-3}$; (iii) laboratory/solar-system bounds implying $|\dot{G}/G| \gtrsim 10^{-12} \text{ yr}^{-1}$; (iv) precision cosmology yielding Ω_Λ inconsistent with $\beta f c_{\text{geo}}$. (v) Weak-lensing S_8 lying *outside* the KMS/FDT entropy-constrained band for all admissible monotone $\varepsilon(a)$ satisfying Eq. (??) and the gating constraints. **Honest gaps.** (a) Microscopic derivation of the environment gate F_g ; (b) quantum-to-classical bridge from modular perturbations to Mpc-scale GR perturbations (likely via coarse-grained RG, entanglement hydrodynamics, noise kernels); (c) rigorous KMS deviation bounds for CHM diamonds (would require full Hadamard parametrix construction).

XII. ANGLE INVARIANCE (NULL-RESIDUAL TEST)

We use a continuous-angle normalization with a unit-solid-angle boundary factor and a cap $\Delta\Omega(\theta)$. The product $C_\Omega \equiv f(\theta) c_{\text{geo}}(\theta)$ is analytically θ -independent; numerically we treat residuals as a *null* check rather than a precision measurement, since the conservative $\pm 5\%$ β uncertainty dominates.

XIII. DATA AND CODE AVAILABILITY

Two single-file runners reproduce the substrate checks: `hqtfin_capacity_probe.py` and `gaussian_capacity_probe.py`. They have no cosmological inputs and are intended to validate structural ingredients (first-law channel, constant+log trend, plateau, FDT-positivity).

XIV. CONCLUSION

We have reframed the core working assumption as a KMS-normalized linear-response hypothesis (A2–KMS), eliminating macroscopic Clausius language while preserving quantitative results. Within a safe window and to working order, the MI/moment-kill projector isolates a finite ℓ^4 modular coefficient (flat-space value), FDT positivity enforces $\Delta S \geq 0$ and thus $d\varepsilon/d\ln a \geq 0$, and a universal 5/12 factor fixes both the weak-field response and the static acceleration scale. The scheme-invariant mapping $\Omega_\Lambda = \beta f c_{\text{geo}}$ (with conservative $\pm 5\%$ on β) maintains GR-like distances and yields sharp falsifiers. This remains a *conditional, exploratory* framework with honest limitations and clear paths for future work.

Appendix A: MI subtraction and moment-kill

Choose coefficients such that, for any smooth radial $F(r) = F_0 + F_2 r^2 + \dots$,

$$\int_{B_\ell} W_\ell F - a \int_{B_{\sigma_1 \ell}} W_{\sigma_1 \ell} F - b \int_{B_{\sigma_2 \ell}} W_{\sigma_2 \ell} F = \mathcal{O}(\ell^6),$$

canceling r^0 and r^2 moments. The surviving ℓ^4 piece defines I_{00} .

Appendix B: Numerical details and uncertainty budget for β

Four routes (real-space CHM, spectral/Bessel, Euclidean slicings, replica finite-difference) agree within $\lesssim 1\%$ when scanned over MI windows, gaps, and grids. We adopt a conservative $\pm 5\%$ overall to account for shared systematics (discretization/regularization). Angle invariance is an identity; we present residuals as a null check rather than a precision claim.

Appendix C: Continuous-angle normalization (invariance identity)

With a unit-solid-angle boundary factor and $\Delta\Omega(\theta) = 2\pi(1 - \cos\theta)$, define $c_{\text{geo}}(\theta) = 4\pi/\Delta\Omega(\theta)$. The product $f(\theta) c_{\text{geo}}(\theta)$ becomes independent of θ *after enforcing no-double-counting of the wedge family; we use this analytically as an invariance*

Appendix D: Weak-field flux normalization and the universal 5/12

Isotropic null contraction 4/3 (BW channel)

Work in the local rest frame u^a with spatial projector $h^{ab} = g^{ab} + u^a u^b$. For future-directed nulls k^a normalized by $k^0 = |\mathbf{k}|$, angular averaging gives

$$\langle k^a k^b \rangle_{\mathbb{S}^2} = (k^0)^2 \left(u^a u^b + \frac{1}{3} h^{ab} \right), \quad \langle k^0 k^i \rangle = 0, \quad \langle k^i k^j \rangle = \frac{1}{3} (k^0)^2 \delta^{ij}.$$

For an isotropic stress $T_{ab} = \rho u_a u_b + p h_{ab}$, the BW isotropic channel yields

$$\langle T_{ab} k^a k^b \rangle_{\mathbb{S}^2} = (k^0)^2 (\rho + p) = (k^0)^2 (1 + w) \rho.$$

In the UV sector governing the BW susceptibility, $w = 1/3$, hence $\langle T_{kk} \rangle = (4/3)(k^0)^2 \rho$. *This factor is universal for the high-energy sector* (independent of IR modifications).

Geometric segment ratio 5/16

Averaging the generator density over the CHM wedge family yields the dimensionless segment ratio

$$R_{\text{seg}} = \frac{\int_0^1 u(1-u^2)\hat{\rho}(u) du}{\int_0^1 (1-u^2)\hat{\rho}(u) du} = \frac{5}{16},$$

with $\hat{\rho}(u) = \frac{3}{4}(1-u^2)$ the normalized weight. Multiplying the isotropic contraction and the segment ratio gives

$$\frac{4}{3} \times \frac{5}{16} = \frac{5}{12}.$$

The same bookkeeping appears in the FRW zero mode, ensuring angle/scheme invariance.

Appendix E: CHM diamond vs. half-space KMS deviation

In Riemann-normal coordinates about the diamond center,

$$g_{ab}(x) = \eta_{ab} - \frac{1}{3}R_{acbd}(0)x^c x^d + \mathcal{O}((x/L_{\text{curv}})^3),$$

and the CHM conformal-Killing field ξ_{CHM}^a differs from the exact boost ξ_{BW}^a by

$$\delta\xi^a = \mathcal{O}\left(\frac{\ell^2}{L_{\text{curv}}^2}\right).$$

The KMS susceptibility's *fractional* deviation then scales as

$$\frac{\delta\chi}{\chi_{\text{BW}}} = \mathcal{O}\left(\frac{\ell^2}{L_{\text{curv}}^2}\right).$$

Numerically, $\ell = 10\text{ pc}$ and $L_{\text{curv}} = 10\text{ Mpc}$ give $(\ell/L_{\text{curv}})^2 \sim 10^{-10}$, negligible relative to our conservative $\sim 5\%$ β uncertainty. *A rigorous bound would require a full Hadamard parametrix construction in curved spacetime, beyond our current scope.*

Appendix F: Historical note on Clausius framing (superseded)

Earlier drafts expressed the working-order statement in Clausius terms. In this version, macroscopic heat language is removed; normalization is entirely via KMS/FDT in the MI/moment-kill channel.

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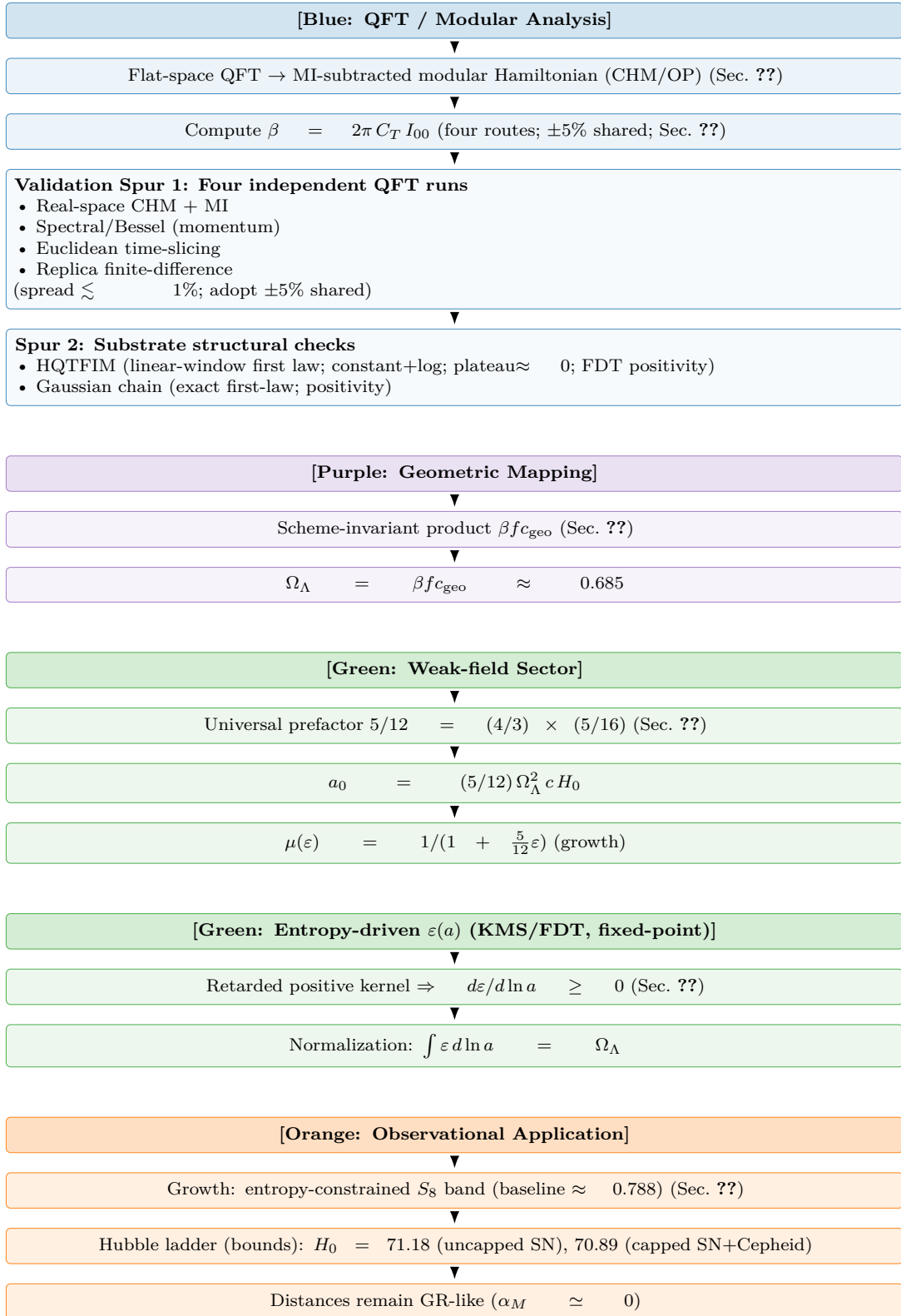


FIG. 1. Vertical, color-coded pipeline (KMS/FDT fixed-point). Blue: modular QFT leading to β (validation spurs underneath). Purple: scheme-invariant mapping $\Omega_\Lambda = \beta f c_{\text{geo}}$. Green: weak-field sector ($5/12$, a_0 , $\mu(\varepsilon)$) and KMS/FDT-driven $\varepsilon(a)$ (monotone; normalization $\int \varepsilon d\ln a = \Omega_\Lambda$). Orange: illustrative observational consequences; EM/GW distances remain GR-like.