

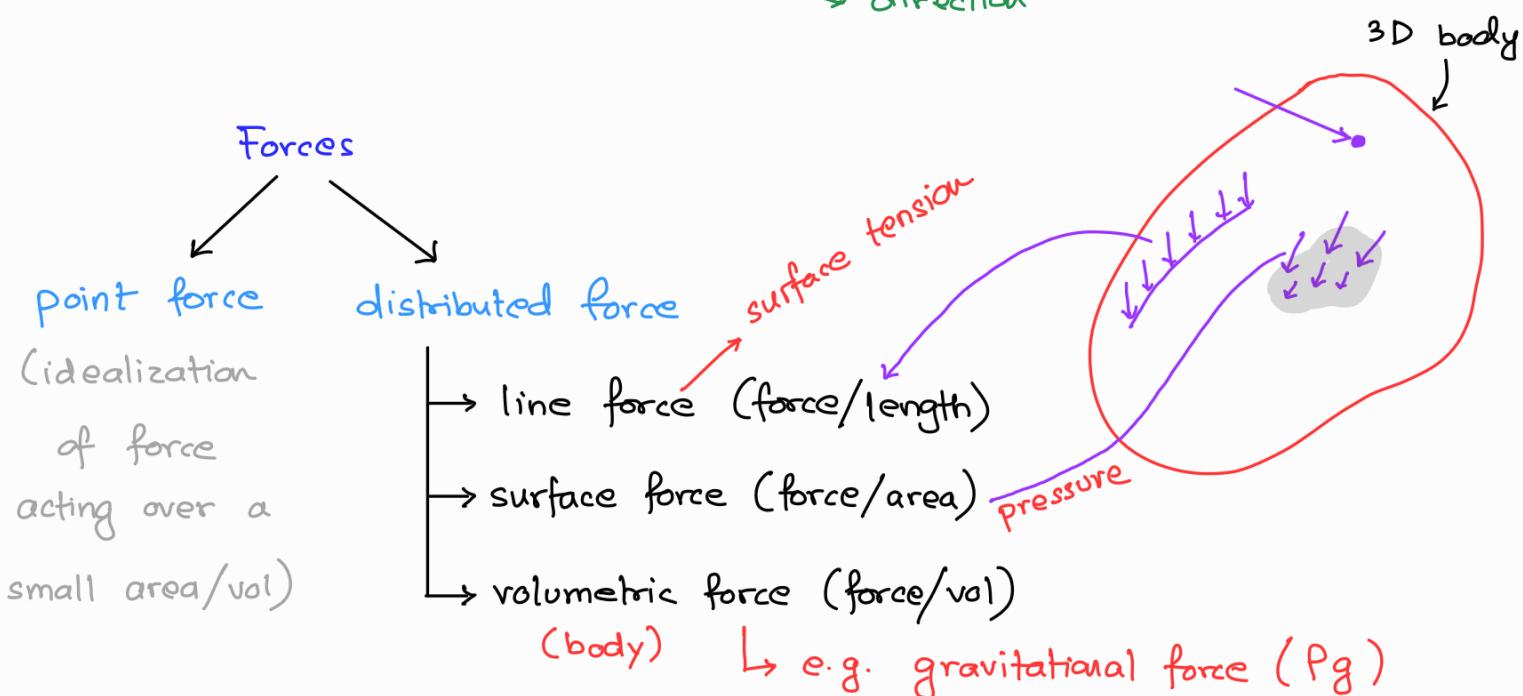
## Chapter 2

Until the last lecture we have been discussing kinematics that was to do with geometry of motion. Now we move to the next chapter of Dumir's book which lays out the precepts of dynamics that deals with analysis of the cause of motion — forces, moments

Let's review some primitive terms and concepts

FORCE : "Quantitative measure of mechanical interaction of material points/bodies in contact or at a distance"

→ it is a vector quantity → magnitude  
→ direction



Contact Force

vs

Non-contact force

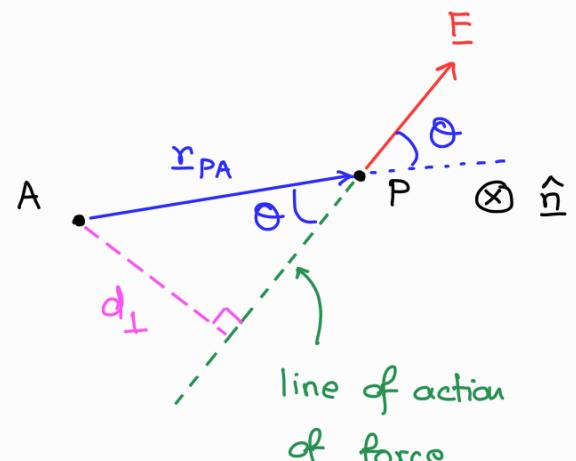
Pressure, surface tension, friction,

Gravity, Magnetic force

## MOMENT of a point force about pt 'A'

Another primitive term

A force tends to push or pull, if the line of action of force passes through A. If not, then we get a rotational effect of force



$$M_A \text{ (Moment about pt A)} = r_{PA} \times F$$

$$= |F| \underbrace{|r_{PA}| \sin \theta}_{d_{\perp}} \hat{n}$$

$d_{\perp}$   
perpendicular  
distance betn pt A  
and line of action  
of force  $F$

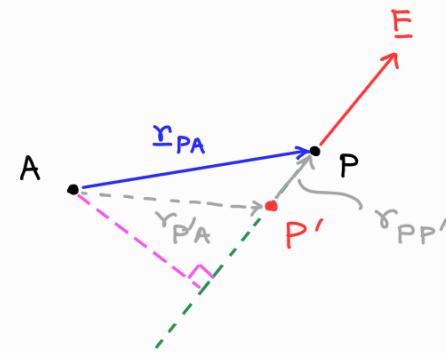
$$\Rightarrow M_A = |F| d_{\perp} \hat{n}$$

$\hat{n}$  is the unit normal vector to the plane formed by pt A and the line of action of force  $F$

Also, note that the moment is the same for any point  $P'$  along the line of action of the force

$$\begin{aligned} M_A &= r_{PA} \times F = (r_{P'A} + r_{PP'}) \times F \\ &= r_{P'A} \times F \end{aligned}$$

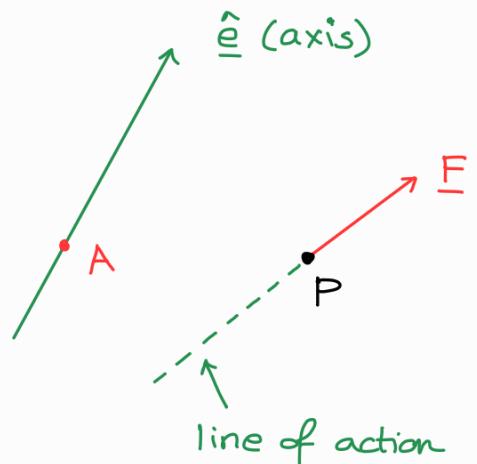
$$(\because r_{PP'} \parallel F \Rightarrow r_{PP'} \times F = 0)$$



## Moment of forces about (along) an axis ( $\hat{e}$ )

Consider the following:

(a) Choose 'P', any point on the line of action of force



(b) Choose any point 'A' on the axis along unit vector  $\hat{e}$

(c) Get  $\underline{M}_A$  = moment of force  $F$  acting at 'P' abt 'A'

$$= \underline{r}_{PA} \times \underline{F}$$

(d) Obtain  $\underline{M}_e$ , the moment of force  $F$  acting about an axis along  $\hat{e}$

$$\underline{M}_e = (\underline{M}_A \cdot \hat{e}) \hat{e}$$

component of  $\underline{M}_A$  along  $\hat{e}$

→ We have already seen that  $\underline{M}_A$  does not depend on 'P' as long as 'P' is anywhere along the line of action of force

Does  $\underline{M}_e$  depend on point 'A's location along the axis ( $\hat{e}$ )?

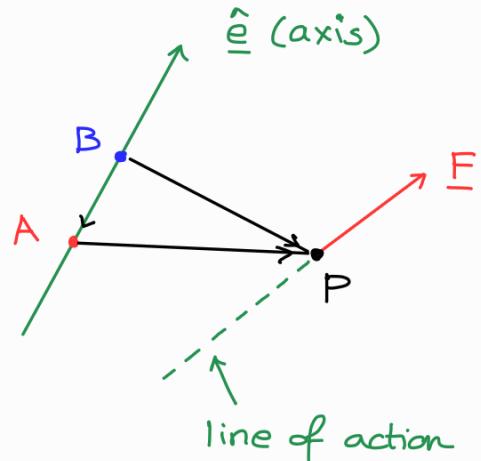
Choose another arbitrary point B on the axis ( $\hat{e}$ )

$$|M'_e| = (\underline{r}_{PB} \times \underline{F}) \cdot \hat{\underline{e}}$$

abt pt. B

$$\begin{aligned} \text{pt. B along } e &= [(\underline{r}_{AB} + \underline{r}_{PA}) \times \underline{F}] \cdot \hat{\underline{e}} \\ &= [(\underline{r}_{AB} \times \underline{F}) + (\underline{r}_{PA} \times \underline{F})] \cdot \hat{\underline{e}} \end{aligned}$$

$$\begin{aligned} &\uparrow \quad // \text{ to } \hat{\underline{e}} \quad \therefore (\underline{r}_{AB} \times \underline{F}) \cdot \hat{\underline{e}} = 0 \\ &\therefore (\underline{r}_{PA} \times \underline{F}) \cdot \hat{\underline{e}} = |M_e| \end{aligned}$$



$$= (\underline{r}_{PA} \times \underline{F}) \cdot \hat{\underline{e}} = |M_e|$$

abt pt A  
along axis  $\hat{\underline{e}}$

$\therefore M_e$  stays the same about any point 'A' as long as 'A' lies along the axis ( $\hat{\underline{e}}$ )

### Special cases

① Moment  $M_e = 0$  if  $\underline{F} \parallel \hat{\underline{e}}$

② "  $M_e = 0$  if  $\underline{F}$  intersects the axis  $\hat{\underline{e}}$   
 $(\because \underline{r}_{PA} \parallel \underline{F})$

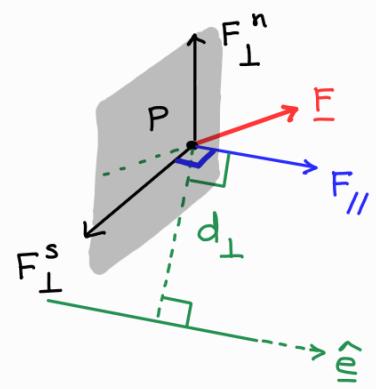
③  $M_e \neq 0$  ONLY IF  $\underline{F}$  and  $\hat{\underline{e}}$  are skew-lines

$\uparrow$   
 directed lines that are neither  $\parallel$   
 nor intersecting

$$\underline{F} = \underline{F}_{\parallel} + \underline{F}_{\perp}^n + \underline{F}_{\perp}^s$$

component      intersects      components of  $\underline{F}$

of  $\underline{F} \parallel$  to  $\hat{e}$     on a plane  $\perp$  to  $\hat{e}$



$$|\underline{M}_e| = |\underline{F}_{\perp}^s| d_{\perp} \quad \text{and} \quad \underline{M}_e = |\underline{M}_e| \hat{e}$$

Couple: set of forces  $\underline{F}$  and  $-\underline{F}$  acting along different lines of action

$\subseteq$  ← symbol of couple vector

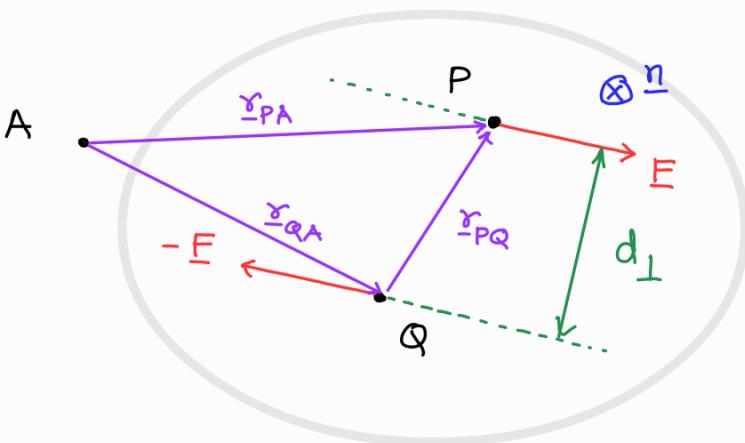
$$\begin{aligned} \underline{M}_A &= \underline{r}_{PA} \times \underline{F} \\ &\quad + \underline{r}_{QA} \times (-\underline{F}) \end{aligned}$$

$$= (\underline{r}_{PA} - \underline{r}_{QA}) \times \underline{F}$$

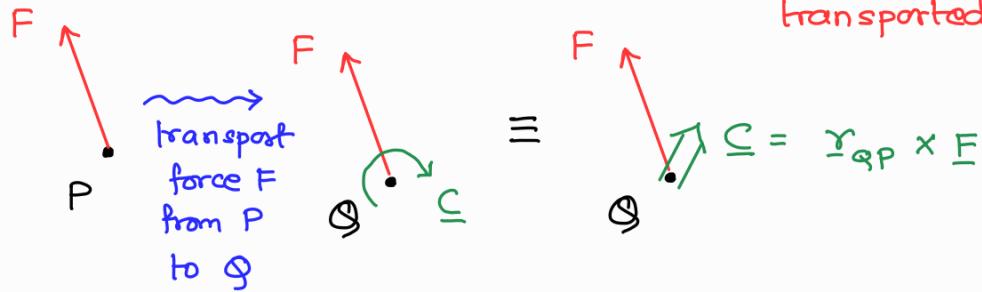
$$= \underline{r}_{PQ} \times \underline{F}$$

$$= |\underline{F}| d_{\perp} \hat{n} \stackrel{\text{def}}{\equiv} \underline{C} \rightarrow \text{does NOT depend upon pt A}$$

(and stays the same abt ANY point in space)



Equivalent force-couple system  $\rightarrow$  what happens when the pt of application of a force  $\underline{F}$  gets transported to another point?



A force  $\underline{F}$  acting at a point  $P$  is equivalent to a force  $\underline{F}$  acting at  $Q$  and a couple  $C = r_{QP} \times \underline{F}$

### Centre of Mass of a set of particles

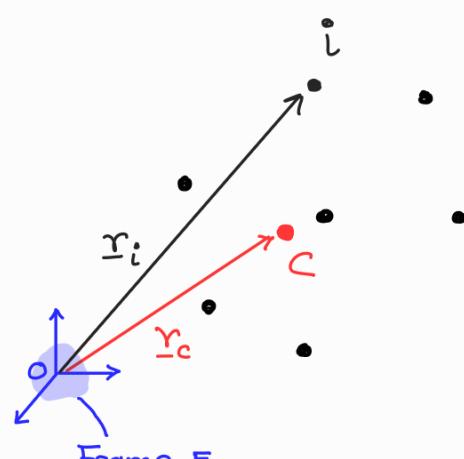
Mass  $\rightarrow$  measure of inertia

property of body due to which an external force is required to accelerate the body

### Centre of mass for a set of particles

If there are  $N$  particles, each having mass  $m_i$  and located  $\underline{r}_i$  w.r.t. the csys of a ref. frame

$$\underline{r}_c = \frac{\sum_{i=1}^N m_i \underline{r}_i}{\sum_{i=1}^N m_i}$$



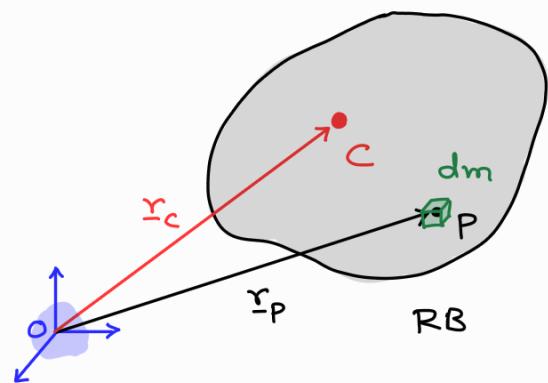
## Centre of mass (COM) of a rigid body

A rigid body is a continuous distribution of point masses

Consider an infinitesimal mass  $dm$   
at any point  $P$  in RB

Suppose  $\underline{r}_P = \overrightarrow{OP}$

Then point  $C$  is the centre  
of mass of the rigid body (RB)  
and its position vector is  $\underline{r}_c$  is:



$$\underline{r}_c = \frac{\int \underline{r}_P dm}{\int dm} = \frac{\int \underline{r}_P \rho dv}{\int \rho dv}$$

density

Total mass:  $m = \int dm$

integration over the volume of the RB

If the mass of RB is approximately distributed along a curve

$$\underline{r}_c = \frac{\int \underline{r}_P dm}{\int dm} = \frac{\int \underline{r}_P \lambda ds}{\int \lambda ds}$$

mass per unit length

integration along the curve

[e.g. thin rod, wire, cable]

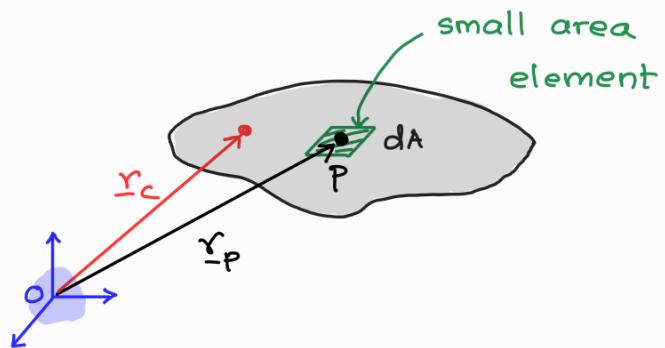
small length element

Similarly, if the mass is distributed over an area

$$\underline{r}_c = \frac{\int \underline{r}_p \sigma dA}{\int \sigma dA}$$

mass per unit area

integration  
over the area



[e.g. thin disc/shell]

Centroid of RB: denoted by  $C^*$

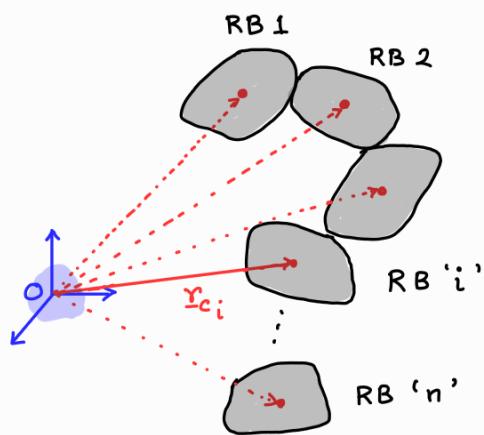
$\underline{r}_{C^*}$  (position vector of centroid) =

$\frac{\int \underline{r}_p dV}{\int dV}$ (Volume element)
$\frac{\int \underline{r}_p dA}{\int dA}$ (Area element)
$\frac{\int \underline{r}_p ds}{\int ds}$ (Line element)

\* Centroid ' $C^*$ ' and COM ' $C$ ' are the same if RB has a uniform density throughout the RB

## Centre of mass of a collection of RB's

For a system of rigid bodies composed of 'n' different parts:



e.g.



Robot arm

the COM is given by  $\underline{r}_{cm}$

(as an assembly of 5 RBs)

$$\underline{r}_c = \frac{\sum_{i=1}^n m_i \underline{r}_{ci}}{\sum_{i=1}^n m_i}$$

position vector of COM  $C_i$  of RB ' $i$ ' measured from  $O$

mass of RB ' $i$ '

# Momentum and Moment of Momentum of RB wrt F

## ① Momentum of RB wrt frame F

also called linear momentum

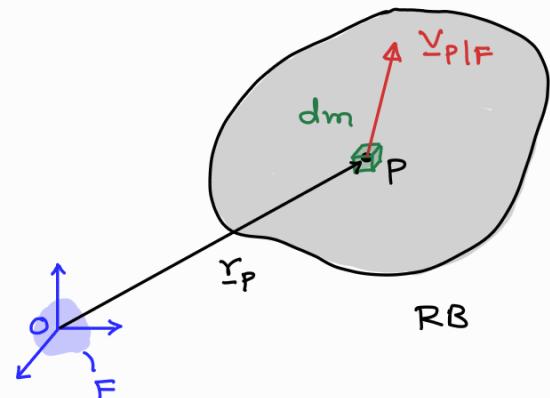
Consider an elementary mass  $dm$

at any point P in RB

Suppose  $\underline{v}_{P/F}$  is the velocity

of point P w.r.t. the frame F

and  $\underline{r}_P = \overrightarrow{OP}$



The momentum of the rigid body is denoted by  $\underline{p}|_F$

and defined as:

$$\underline{p}|_F = \int \underline{v}_{P/F} dm = \int \frac{d\underline{r}_P}{dt} |_F dm$$

integration  
over entire RB

Both velocity and

mass distributions

must be known for integ.

$$= \frac{d}{dt} \left( \int \underline{r}_P dm \right) |_F$$

[Using  
Leibnitz  
integral  
formula]

total mass

$$= \frac{d}{dt} (m \underline{r}_c) |_F$$

$$= m \frac{d}{dt} \underline{r}_c |_F$$

$= m \underline{v}_{c/F}$  = momentum of RB's  
Centre of mass (COM)

$$\underline{r}_c = \frac{\int \underline{r}_P dm}{\int dm}$$

$$\Rightarrow m \underline{r}_c = \int \underline{r}_P dm$$

For a system of RBs, the momentum would be

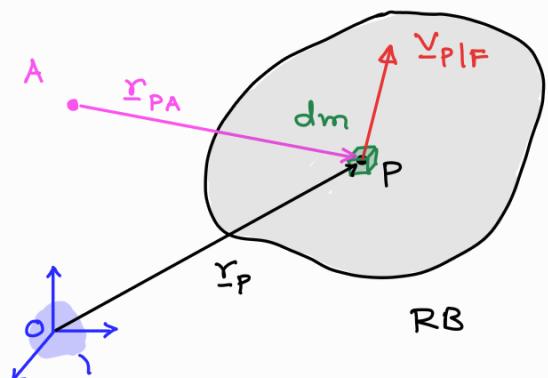
$$\underline{P}|_F = \sum_{i=1}^n m_i \underline{v}_{c_i|F}$$

② Moment of Momentum of a RB about a pt A and  
popular as ANGULAR MOMENTUM      wrt frame 'F'

The angular momentum of a RB  
about a point A is denoted by

$\underline{H}_A|_F$  and defined as:

$$\underline{H}_A|_F = \int \left( \underline{r}_{PA} \times \underline{v}_{PA|F} dm \right)$$



$$\underline{v}_{PA|F} = \underline{v}_{P|F} - \underline{v}_{A|F}$$

↑  
relative velocity of point P  
w.r.t point A compared in F

'A' can be any point in space

$$'A' = \begin{cases} \text{material point of RB (or)} \\ \text{point on fixed frame of ref F (or)} \\ \text{point on a RB that is moving w.r.t F} \end{cases}$$

## EULER's AXIOMS governing dynamics of RBs

- ① foundational statements or principles that are accepted without proofs
- ② they are building blocks of a bigger logical reasoning

There exists a reference frame 'I', with a point 'O' fixed in 'I' (serving as the origin of a csys attached to 'I'), such that, for any system, if the following two axioms : (a) linear momentum balance, and (b) angular momentum balance are satisfied then the frame of reference is called an **Inertial frame of reference**!

### Axiom ① : Linear momentum balance

The time rate of change of the momentum of the COM of RB in inertial frame is equal to the net external force acting on RB

$$\Rightarrow \frac{d}{dt} \left\{ \underbrace{\underline{P}|_I}_{\text{momentum of the RB wrt frame 'I'}} \right\}_I = \dot{\underline{P}}|_I = \underbrace{\underline{F}^{\text{ext}}}_I \quad \text{--- (I)}$$

net external force acting on the RB

$$\Rightarrow \frac{d}{dt} (m \underline{v}_{c/I})|_I = \underline{F}^{\text{ext}}$$

$$\Rightarrow m \underline{\alpha}_{c/I} = \underline{F}^{\text{ext}}$$

Equivalently, the first Euler's axiom is the product of RB's mass times the acceleration of its COM in inertial frame 'I'

$$\underline{F}_{\text{net}}^{\text{ext}} = m \underline{\alpha}_{C|I} = m \dot{\underline{v}}_{C|I} = m \ddot{\underline{x}}_{C|I}$$

Thus, using the first axiom, we can describe the motion of centre of mass of a RB.

Recall that motion of a material point of an RB may be described using different coordinate systems:

(a) Cartesian csys

$$\underline{F}_{\text{net}}^{\text{ext}} = F_1 \hat{\underline{e}}_1 + F_2 \hat{\underline{e}}_2 + F_3 \hat{\underline{e}}_3 = m (\ddot{x}_1 \hat{\underline{e}}_1 + \ddot{x}_2 \hat{\underline{e}}_2 + \ddot{x}_3 \hat{\underline{e}}_3)$$

(b) Cylindrical csys

$$\underline{F}_{\text{net}}^{\text{ext}} = F_r \hat{\underline{e}}_r + F_\phi \hat{\underline{e}}_\phi + F_z \hat{\underline{e}}_z$$

$$= m \left[ (\ddot{r} - r \dot{\phi}^2) \hat{\underline{e}}_r + \frac{1}{r} \frac{d}{dt} (r^2 \dot{\phi}) \hat{\underline{e}}_\phi + \ddot{z} \hat{\underline{e}}_z \right]$$

(c) Path csys

$$\underline{F}_{\text{net}}^{\text{ext}} = F_t \hat{\underline{e}}_t + F_n \hat{\underline{e}}_n = m \left( \ddot{s} \hat{\underline{e}}_t + \frac{\dot{s}^2}{P} \hat{\underline{e}}_n \right)$$

Note that there can be no force component  $F_b$  normal to the osculating plane.

## Axiom ② : Angular momentum balance

The time rate of change of angular momentum of RB about a specified reference point 'O' in an inertial frame 'I' is equal to the sum of all moments/torques caused by external loads

$$\textcircled{II} \quad \frac{d}{dt} \left\{ \underline{H}_{O|I} \right\}_I = \dot{\underline{H}}_{O|I} = \underline{\underline{M}}_{O,\text{net}}^{\text{ext}}$$

/

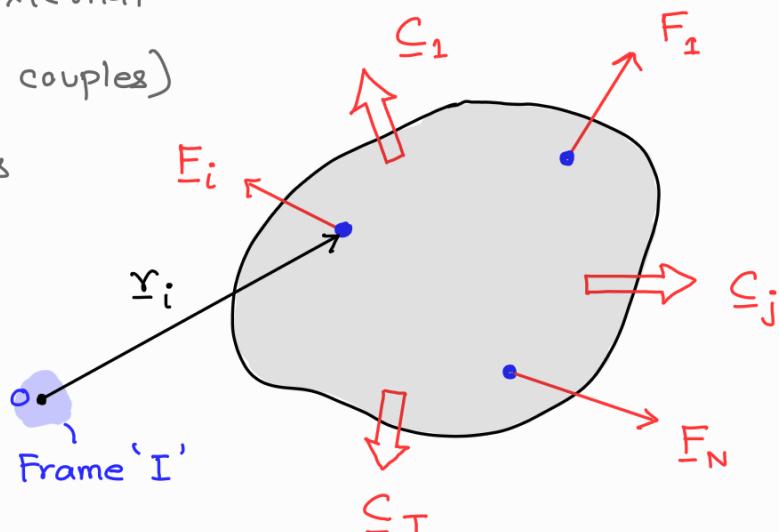
angular momentum  
of RB wrt frame 'I'

net moment  
due to all ext.  
loads about 'O'

Eqns  $\textcircled{I}$  and  $\textcircled{II}$  leads to six scalar equations.

What do we mean by net external forces and moments/torques?

In general, there may be  $N$  external forces (excluding forces causing couples) and  $T$  torques/couples of forces acting on a RB



Net external force

$$\underline{F}_i = \sum_{i=1}^N \underline{F}_i$$

Net external moment abt pt 'O'

$$\underline{\underline{M}}_O = \sum_{i=1}^N \underline{r}_i \times \underline{F}_i + \sum_{j=1}^T \underline{C}_j$$

↑

position vector of the point of application of  $\underline{F}_i$  on RB

## Inertial frames of reference

Frames of reference which satisfies the two Euler's axioms are called inertial frame of reference.

In simple words, an inertial frame is a reference frame relative to which a uniform motion ( $v_{PIF} = \text{constant}$ ) can be sustained without force. Therefore, any disturbance of a RB from a state of rest or uniform motion w.r.t. inertial frame can occur only in response to unbalanced forces and moments, acting separately or together.

Now the question is what physical ref. frame (or RB) in the real world can be treated as a reference frame?

Plainly, a fixed reference frame/body (or a ref. frame at rest) is an inertial frame. But a frame/body can be identified as fixed in space only relative to other bodies known to be fixed in space. So the idea of an inertial frame being fixed in space is meaningless.

It is impossible to find an exact physical inertial ref. frame!

Our most natural choice appears to be the Earth frame.

However, we know that the Earth rotates and revolves and therefore, can we use Earth frame as an inertial frame?

We cannot use the Earth frame as an inertial frame for dynamics of astronomical bodies, however, we could approximately use it for bodies for which Earth's movement does not affect the calculations very much.