

Skyrme Crystals

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Joint work with Derek Harland and Paul Leask

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Skyrme model

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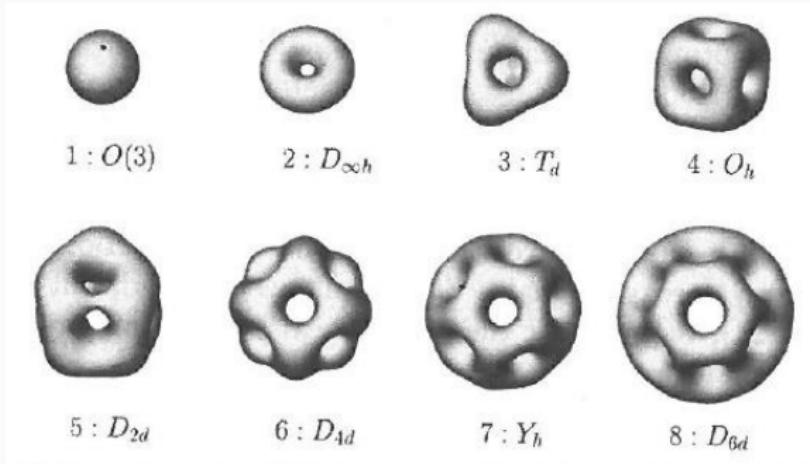
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- Faddeev bound: $E(\varphi) \geq E_0|B|$, unattainable
- Degree B minimizer \leftrightarrow nucleus of atomic weight B

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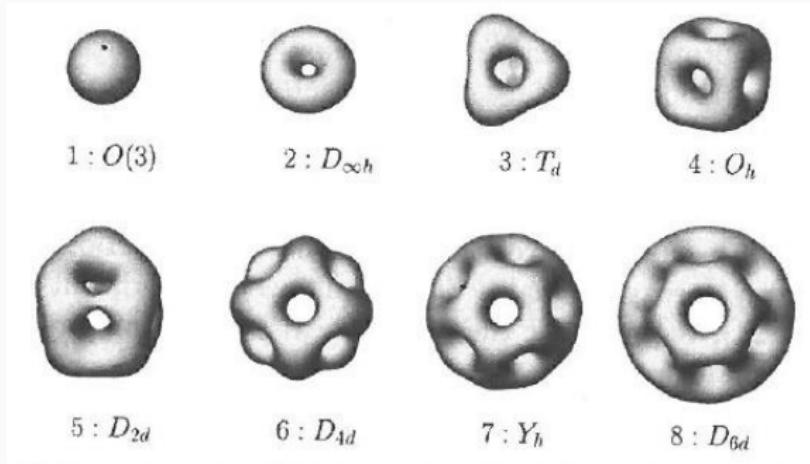
- Numerics



Battye and Sutcliffe

Skyrme model

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Battye and Sutcliffe

- $E/B E_0$ monotonically decreases e.g. 1.232 ($B = 1$), 1.096 ($B = 8$).

Skyrme model

- Suggests Skyrmions may be able to form a **crystal**

$$\varphi : \mathbb{R}^3 / \Lambda \rightarrow G, \quad \Lambda = \{n_1 \mathbf{X}_1 + n_2 \mathbf{X}_2 + n_3 \mathbf{X}_3 : \mathbf{n} \in \mathbb{Z}^3\}$$

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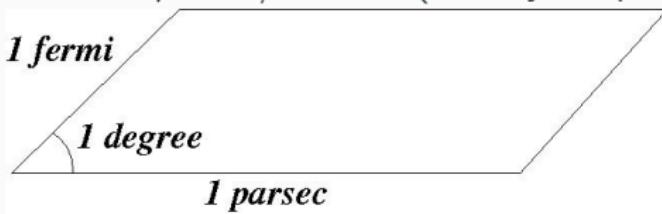
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- But is this really a crystal? Given **any** Λ , B , there exists a degree B minimizer $\varphi : \mathbb{R}^3 / \Lambda \rightarrow G$ (Auckly, Kapitanski).

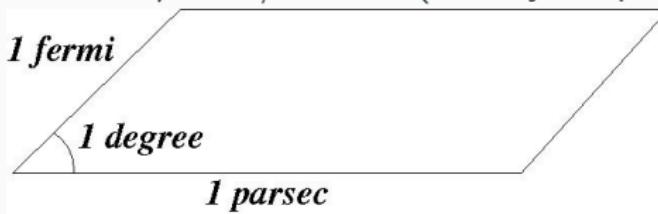


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For most Λ , lifted map $\mathbb{R}^3 \rightarrow G$ clearly isn't a genuine solution: artifact of bc's.

General question

- Given a minimizer $\varphi : \mathbb{R}^k / \Lambda \rightarrow N$ of some energy functional $E(\varphi)$, when is the lifted map $\mathbb{R}^k \rightarrow N$ a genuine crystal?

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- Should be critical (in fact stable) with respect to variations of Λ too.

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- Identify them all with $M = \mathbb{R}^3/\mathbb{Z}^3$, the cubic torus. Now mfd is fixed, but **metric** depends on Λ
$$g_\Lambda = g_{ij} dx_i dx_j, \quad g_{ij} = \mathbf{X}_i \cdot \mathbf{X}_j \text{ const}$$

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- So now

$$E(\varphi, g) = \int_{T^3} (|d\varphi|_g^2 + \frac{1}{4} |\varphi^* \omega|_g^2 + V(\varphi)) \text{vol}_g$$

and we want to minimize w.r.t. both $\varphi \in C_B^2(T^3, G)$ **and** $g \in SPD_3$ (space of symmetric positive definite 3×3 matrices)

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 - What if we fix φ ? Does $E(g)$ attain a min?

Existence and uniqueness of minimizing metrics

- Want to think of E , for a **fixed** $\varphi : T^3 = \mathbb{R}^3 / \mathbb{Z}^3 \rightarrow G$ as a function of the metric g on T^3 :

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- Proof: First note that

$$E_\varphi(g) = \sqrt{|g|} \operatorname{tr}(Hg^{-1}) + \frac{1}{\sqrt{|g|}} \operatorname{tr}(\Omega g) + C\sqrt{|g|}$$

where $H, \Omega \in SPD_3$ and $C \in [0, \infty)$ are fixed.

Existence and uniqueness of minimizing metrics

$$E_2(g) = \int_{T^3} \varphi^* h(\partial_i, \partial_j) g^{ij} \sqrt{|g|} d^3x = \sqrt{|g|} g_{ij}^{-1} H_{ij}$$

$$H_{ij} := \int_{T^3} \varphi^* h(\partial_i, \partial_j) d^3x$$

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$$E_4(g) = \frac{1}{4} \|\varphi^* \omega\|_{L^2(g)}^2 = \frac{1}{4} \int_{T^3} \frac{1}{|g|} g(X_\varphi, X_\varphi) \text{vol}_g = \frac{g_{ij}}{\sqrt{|g|}} \Omega_{ij}$$

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- Surjection $f : (0, \infty)^3 \times O(3) \rightarrow SPD_3$

$$(\boldsymbol{\lambda}, \mathcal{O}) \mapsto \mathcal{O} \begin{pmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_2 & 0 \\ 0 & 0 & \lambda_3 \end{pmatrix} \mathcal{O}^T = \mathcal{O} D_{\boldsymbol{\lambda}} \mathcal{O}^T$$

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J.M. Speight (University of Leeds) • We will show $E \circ f : (0, \infty)^3 \times O(3) \rightarrow \mathbb{R}$ attains a min

Existence

$$\begin{aligned}(E \circ f)(\lambda, \mathcal{O}) &= \text{tr}(H(\mathcal{O}D_\lambda\mathcal{O}^{-1})^{-1}) + \text{tr}(\Omega\mathcal{O}D_\lambda\mathcal{O}^{-1}) + \frac{C}{\lambda_1\lambda_2\lambda_3} \\ &= \text{tr}(\mathcal{O}^{-1}H\mathcal{O}D_\lambda^{-1}) + \text{tr}(\mathcal{O}^{-1}\Omega\mathcal{O}D_\lambda) + \frac{C}{\lambda_1\lambda_2\lambda_3}\end{aligned}$$

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- Consider the smooth functions $O(3) \rightarrow (0, \infty)$

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- Exists $\alpha > 0$ s.t. for all (λ, \mathcal{O}) ,

$$(E \circ f)(\lambda, \mathcal{O}) \geq \alpha \left(\frac{1}{\lambda_1} + \frac{1}{\lambda_2} + \frac{1}{\lambda_3} + \lambda_1 + \lambda_2 + \lambda_3 \right). \quad (*)$$

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- Let $\mathcal{G}_* = \mathcal{O}_* D_{\lambda_*} \mathcal{O}_*^{-1} \in SPD_3$. E_φ attains a min at $g_* = \mathcal{G}_*/\det(\mathcal{G}_*)$.

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- Metric on SPD_3 ?

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- Geodesic through $\mathcal{G}(0)$: $\mathcal{G}(t) = A \exp(t\xi) A^T$ where $AA^T = \mathcal{G}(0)$

$$E_4(\mathcal{G}) = \text{tr}(\Omega \mathcal{G})$$

$$\begin{aligned} E_4(\mathcal{G}(t)) &= \text{tr}(\Omega A \exp(t\xi) A^T) \\ &= \text{tr}(\Omega_A \exp(t\xi)), \quad \Omega_A = A^T \Omega A \end{aligned}$$

$$\left. \frac{d^2}{dt^2} E_4(\mathcal{G}(t)) \right|_{t=0} = \text{tr}(\Omega_A \xi^2) > 0$$

- So E_4 is strictly convex.

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- Hence $E_0 = \det \circ \iota$ is strictly convex

$$E_2(\mathcal{G}) = \text{tr}(H\mathcal{G}^{-1}) = (\widehat{E}_4 \circ \iota)(\mathcal{G})$$

- ι is an isometry, so E_2 is strictly convex
- $\det : SPD_3 \rightarrow \mathbb{R}$ is strictly convex
- Hence $E_0 = \det \circ \iota$ is strictly convex
- So $E = E_2 + E_4 + E_0$ is strictly convex. Hence it has at most one critical point. (Assume $\mathcal{G}_*, \mathcal{G}_{**}$ both cps, apply Rolle's Theorem to $(E \circ \gamma)'$ where γ is the geodesic between them.)

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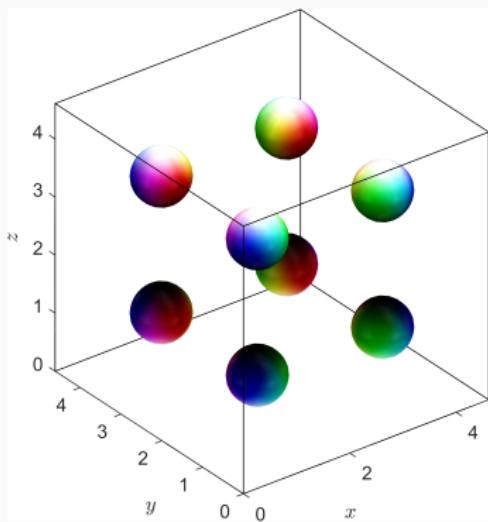
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- Converges much faster than gradient flow.

The Kugler-Shtrikman crystal (massless model)

$$E = \|d\varphi\|^2 + \frac{1}{4}\|\varphi^*\omega\|^2$$



$$(x_1, x_2, x_3) \mapsto (x_2, x_3, x_1)$$

$$(\varphi_0, \varphi_1, \varphi_2, \varphi_3) \mapsto (\varphi_0, \varphi_2, \varphi_3, \varphi_1)$$

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- No reason to expect **degenerate** critical points to survive perturbation

An instructive toy model

$$E_t : \mathbb{R}^2 \rightarrow \mathbb{R}, \quad E_t(x, y) = x^2 + ty^2$$

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The case of the KS crystal

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- Then $M^{\Gamma_R} = M$ and $(R, e) \cdot \varphi_{KS}$ is certainly **not** a nondegenerate cp of $E_0|$

The case of the KS crystal

- For

$$R \in \left\{ \mathbb{I}_4, \underbrace{\begin{pmatrix} (0, 1, 1, 1)/\sqrt{3} \\ * \end{pmatrix}}_{R_\alpha}, \underbrace{\begin{pmatrix} (0, 0, 0, 1) \\ * \end{pmatrix}}_{R_{sheet}}, \underbrace{\begin{pmatrix} (0, 0, 1, 1)/\sqrt{2} \\ * \end{pmatrix}}_{R_{chain}} \right\}$$

Γ_R is nontrivial, and M^{Γ_R} intersects the $SO(4)$ orbit of φ_{KS} transversely: implies $(R, e) \cdot \varphi_{KS}$ is an **isolated** cp of $E_0| : M^{\Gamma_R} \rightarrow \mathbb{R}$

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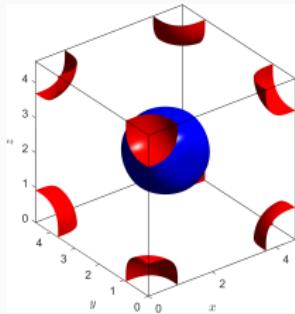
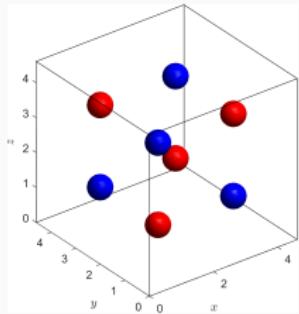
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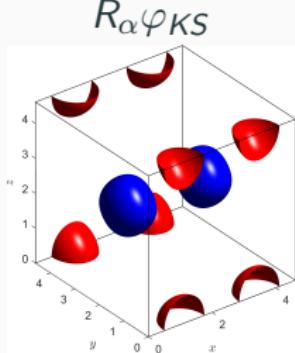
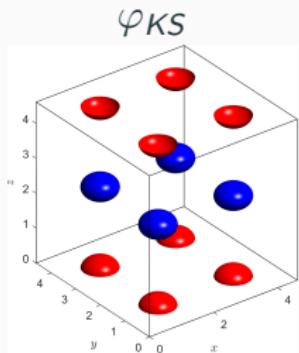
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The KS crystals that (should) survive



$$\varphi_0 = 0.9$$



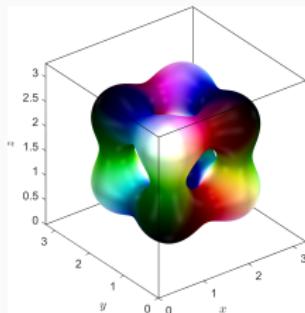
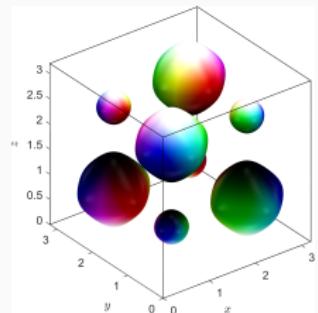
$$R_\alpha \varphi_{KS}$$

$$\varphi_0 = -0.9$$

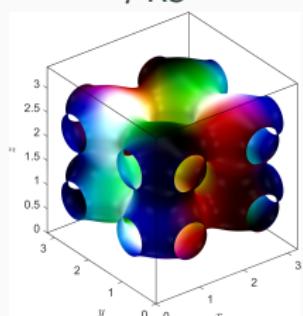
$$R_{sheet} \varphi_{KS}$$

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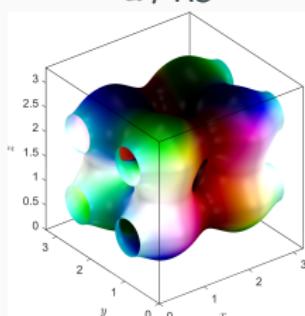
Skyrme crystals at pion mass $t = 1$



φ_{KS}



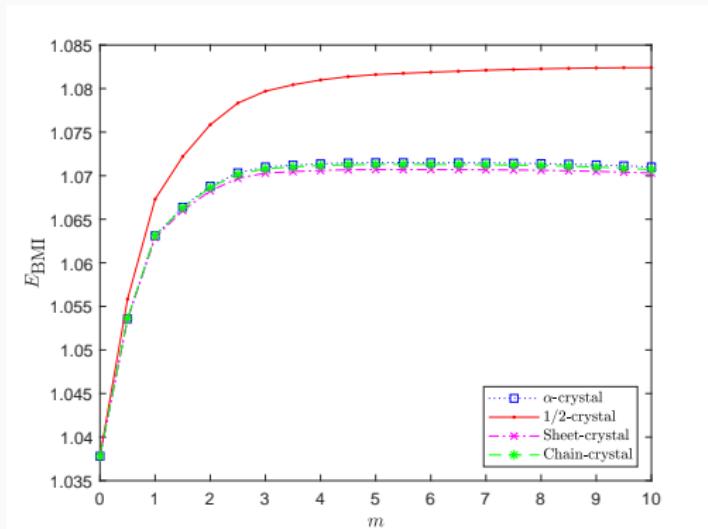
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$R_{sheet} \varphi_{KS}$

$R_{chain} \varphi_{KS}$

Energy ordering: sheet < chain < α < KS



$$g_{\text{sheet}} = \begin{pmatrix} L_1 & 0 & 0 \\ 0 & L_1 & 0 \\ 0 & 0 & L_3 \end{pmatrix}$$
$$L_3 > L_1$$

$$g_{\text{chain}} = \begin{pmatrix} L_1 & 0 & 0 \\ 0 & L_2 & 0 \\ 0 & 0 & L_2 \end{pmatrix}$$
$$L_2 > L_1$$

trigonal, but **not** cubic!

Isospin inertia tensors

$$U_{KS} = \begin{pmatrix} 165.2 & 0 & 0 \\ 0 & 165.2 & 0 \\ 0 & 0 & 165.2 \end{pmatrix}, \quad U_\alpha = \begin{pmatrix} 135.5 & 0 & 0 \\ 0 & 135.5 & 0 \\ 0 & 0 & 167.3 \end{pmatrix},$$

$$U_{sheet} = \begin{pmatrix} 135.8 & 0 & 0 \\ 0 & 135.8 & 0 \\ 0 & 0 & 166.8 \end{pmatrix}, \quad U_{chain} = \begin{pmatrix} 135.6 & 0 & 0 \\ 0 & 135.7 & 0 \\ 0 & 0 & 167.2 \end{pmatrix}.$$

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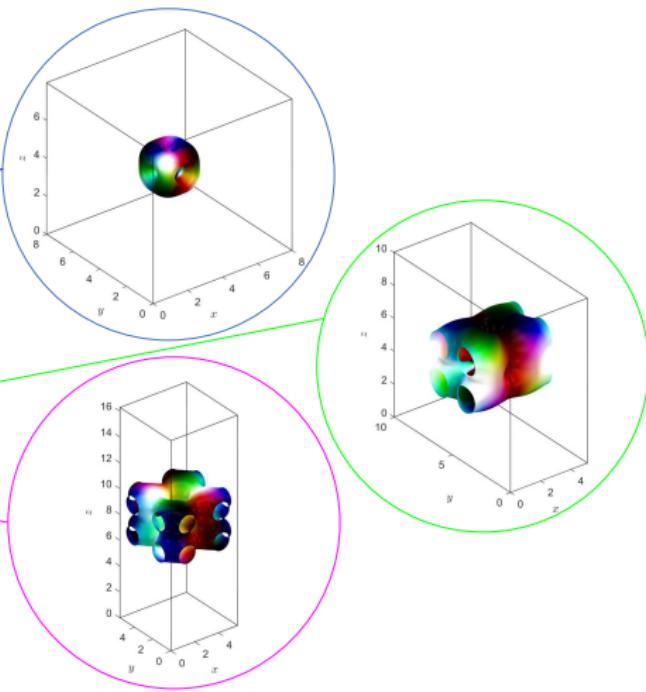
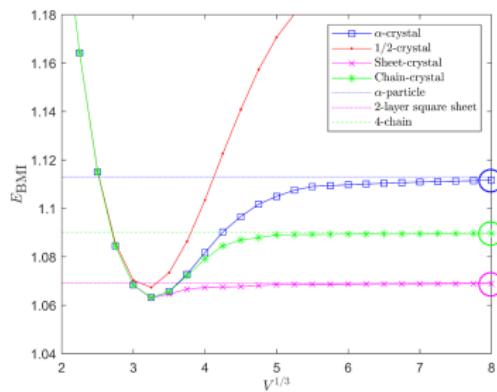
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- Can again solve numerically by ANF

Optimal crystals at fixed baryon density



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- Many examples in condensed matter. True also for nuclear Skyrme model with massive pions
- Extreme case: baby Skyrme model $\varphi : M^2 \rightarrow S^2$

$$E(\varphi) = \int_M \left(\frac{1}{2} |d\varphi|^2 + \frac{1}{2} |\varphi^* \omega|^2 + V(\varphi) \right).$$

Given **any** period lattice $\Lambda \subset \mathbb{R}^2$, can cook up a smooth potential $V : S^2 \rightarrow [0, \infty)$ s.t. $E(\varphi, g)$ has a global min at (φ_*, g_Λ) with φ_* degree 2 and holomorphic.

Concluding remarks

- So this crazy lattice **is** the period lattice of a baby Skyrmion crystal, at least for a (highly contrived) choice of V !

