

# **Skyrme Crystals**

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Joint work with Derek Harland and Paul Leask

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University of Leeds

## General question

- When is a soliton on a torus

$$\varphi : \mathbb{R}^k / \Lambda \rightarrow N$$

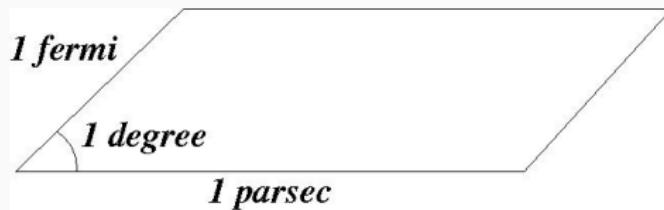
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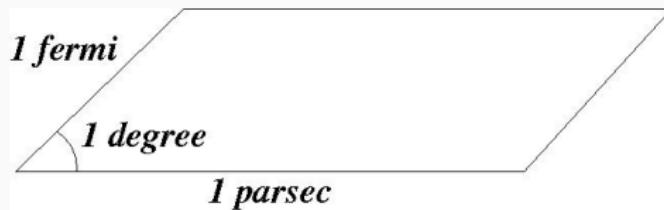


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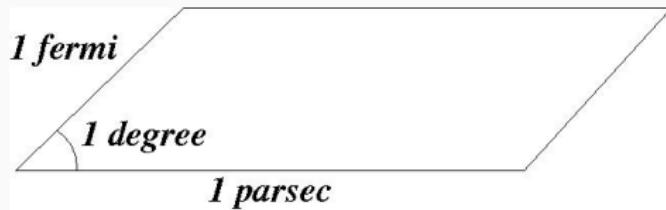
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- Clearly an artifact of b.c.s!
- $\varphi$  should minimize energy  $E$  w.r.t. all variations of field **and period lattice  $\Lambda$**

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$$\Lambda = \{n_1 \mathbf{X}_1 + n_2 \mathbf{X}_2 + \cdots + n_k \mathbf{X}_k : \mathbf{n} \in \mathbb{Z}^k\}$$

$$f : \mathbb{T}^3 \rightarrow \mathbb{R}^k / \Lambda, \quad f(\mathbf{x}) = x_1 \mathbf{X}_1 + x_2 \mathbf{X}_2 + \cdots + x_k \mathbf{X}_k$$

Now mfd is fixed, but **metric** depends on  $\Lambda$

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- Criticality  $\leftrightarrow$  stress tensor
- Manoeuvre works provided  $E(\varphi, g)$  is **geometrically natural**

$$E(\varphi \circ f, g) = E(\varphi, (f^{-1})^* g)$$

$$E : Maps(\mathbb{T}^3, SU(2)) \times SPD_3 \rightarrow \mathbb{R}$$

Two minimization problems:

- Fix  $g$ . Does  $E(\cdot, g) : Maps \rightarrow \mathbb{R}$  attain a min in each homotopy class?

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- Fix  $\varphi$ ? Does  $E(\varphi, \cdot) : SPD_3 \rightarrow \mathbb{R}$  attain a min? YES! And it's a global min, and there are no other critical points!

## The Skyrme energy

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$$E(\varphi, g) = \int_{\mathbb{T}^3} \left( -\frac{1}{2} \operatorname{tr}(L_i L_j) g^{ij} - \frac{1}{16} \operatorname{tr}([L_i, L_j][L_k, L_l]) g^{ik} g^{jl} + V(\varphi) \right) \sqrt{|g|} d^3x$$

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- **Constants:**
  - $H, \Omega$ : symmetric positive **semidefinite** matrices
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- Now  $X_\varphi^g = \sharp_g *_g \varphi^* \omega$ , so

$$|\varphi^* \omega|_g^2 = |X_\varphi^g|_g^2 = \frac{1}{|g|} g(X_\varphi, X_\varphi)$$

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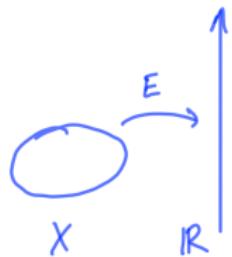
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- Hence

$$E_4(g) = \frac{g_{ij}}{\sqrt{|g|}} \Omega_{ij}$$
$$\Omega_{ij} = \frac{1}{4} \int_{T^3} h(X_i, X_j) d^3x$$

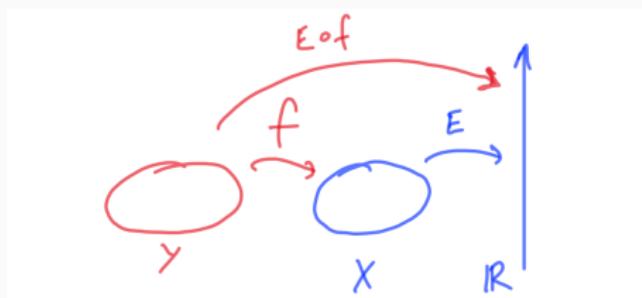
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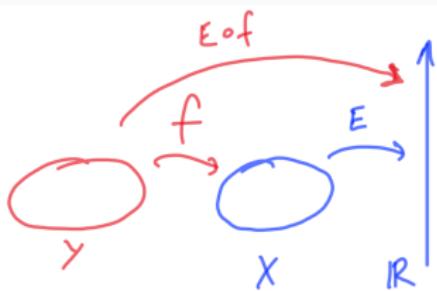
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- $f : (0, \infty)^3 \times O(3) \rightarrow SPD_3$ ,  $f(\lambda, \mathcal{O}) = \mathcal{O} D_\lambda \mathcal{O}^T$
- We will show  $E \circ f : (0, \infty)^3 \times O(3) \rightarrow \mathbb{R}$  attains a min

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$$(E \circ f)(\lambda, \mathcal{O}) = \text{tr}(\mathcal{O}^{-1} H \mathcal{O} D_{\lambda}^{-1}) + \text{tr}(\mathcal{O}^{-1} \Omega \mathcal{O} D_{\lambda}) + \frac{C}{\lambda_1 \lambda_2 \lambda_3}$$

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- Consider the smooth functions  $O(3) \rightarrow (0, \infty)$

$$\mathcal{O} \mapsto (\mathcal{O}^{-1} H \mathcal{O})_{aa}, \quad \mathcal{O} \mapsto (\mathcal{O}^{-1} \Omega \mathcal{O})_{aa}$$

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- Exists  $\alpha > 0$  s.t. for all  $(\boldsymbol{\lambda}, \mathcal{O})$ ,

$$(E \circ f)(\boldsymbol{\lambda}, \mathcal{O}) \geq \alpha \left( \frac{1}{\lambda_1} + \frac{1}{\lambda_2} + \frac{1}{\lambda_3} + \lambda_1 + \lambda_2 + \lambda_3 \right). \quad (*)$$

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- Continuity of  $E$  implies  $(E \circ f)(\lambda_*, \mathcal{O}_*) = E_*$ , i.e.  $E \circ f$  attains a min

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- Metric on  $SPD_3$ ?

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- Geodesic through  $\Sigma(0)$ :  $\Sigma(t) = A \exp(t\xi) A^T$  where  $AA^T = \Sigma(0)$

$$E_4(\Sigma) = \text{tr}(\Omega\Sigma)$$

$$\begin{aligned} E_4(\Sigma(t)) &= \text{tr}(\Omega A \exp(t\xi) A^T) \\ &= \text{tr}(\Omega_A \exp(t\xi)), \quad \Omega_A = A^T \Omega A \end{aligned}$$

$$\frac{d^2}{dt^2} E_4(\Sigma(t)) \Big|_{t=0} = \text{tr}(\Omega_A \xi^2) > 0$$

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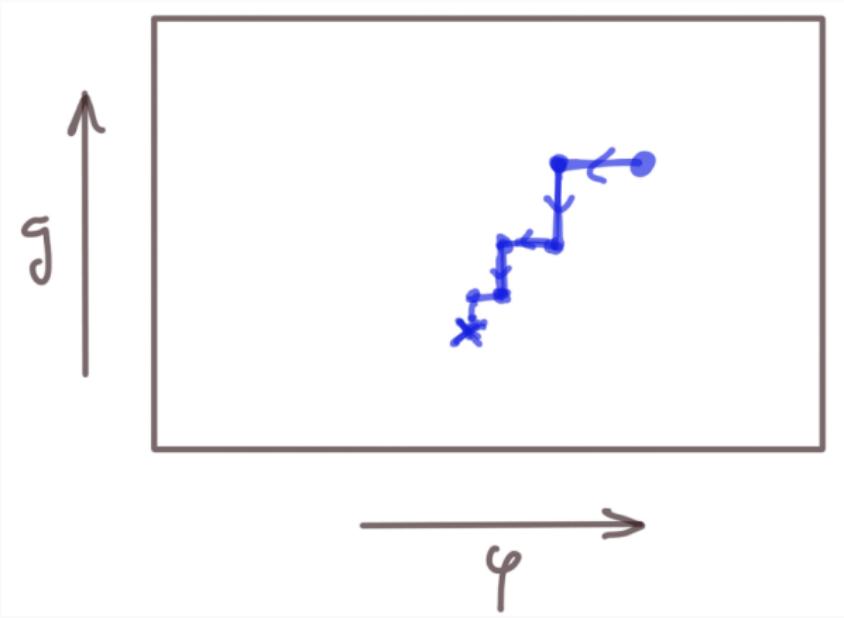
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- $\det : SPD_3 \rightarrow \mathbb{R}$  is convex
- Hence  $E_0 = \det \circ \iota$  is convex
- So  $E = E_2 + E_4 + E_0$  is strictly convex. Hence it has at most one critical point. (Assume  $\Sigma_*$ ,  $\Sigma_{**}$  both cps, apply Rolle's Theorem to  $(E \circ \gamma)'$  where  $\gamma$  is the geodesic between them.)

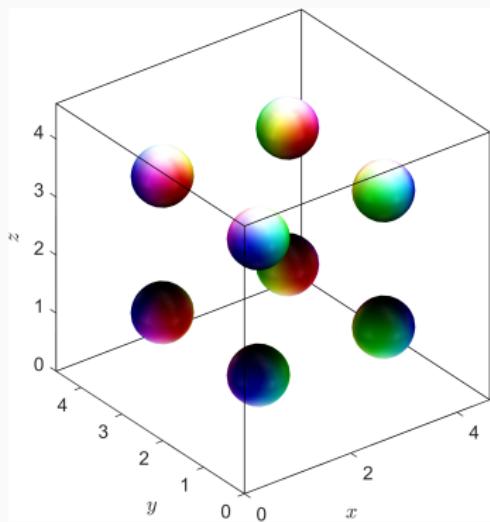
## The numerical problem

$$\ddot{x} = -\operatorname{grad} E(x)$$



# The Kugler-Shtrikman crystal (massless model)

$$E = E_2 + E_4$$



$$(x_1, x_2, x_3) \mapsto (x_2, x_3, x_1)$$

$$(\varphi_0, \varphi_1, \varphi_2, \varphi_3) \mapsto (\varphi_0, \varphi_2, \varphi_3, \varphi_1)$$

$$(x_1, x_2, x_3) \mapsto (x_2, -x_1, x_3)$$

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$$(x_1, x_2, x_3) \mapsto (x_1 + 1/2, x_2, x_3)$$

$$(\varphi_0, \varphi_1, \varphi_2, \varphi_3) \mapsto (-\varphi_0, -\varphi_1, \varphi_2, \varphi_3)$$

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- Turn on pion mass:

$$E_t = E_0 + t \int_{\mathbb{T}^3} (1 - \varphi_0) \sqrt{|g|} d^3x$$

What happens to these critical points?

## The Kugler-Shtrikman crystal: turning on the pion mass

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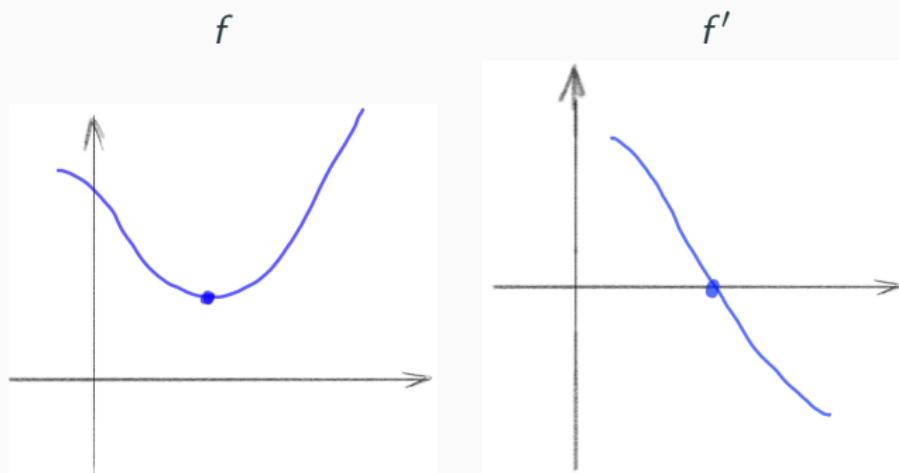
- Massless model has global  $SO(4)$  symmetry: no boundary to break this
- Above solution  $\varphi_{KS}, g_{KS} = L\mathbb{I}_3$  is one point on a  $SO(4)$  orbit of solutions
- Turn on pion mass:

$$E_t = E_0 + t \int_{\mathbb{T}^3} (1 - \varphi_0) \sqrt{|g|} d^3x$$

What happens to these critical points?

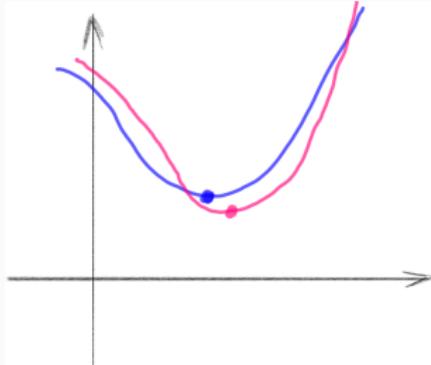
- No reason to expect **degenerate** critical points to survive perturbation

## Degenerate critical points are unstable

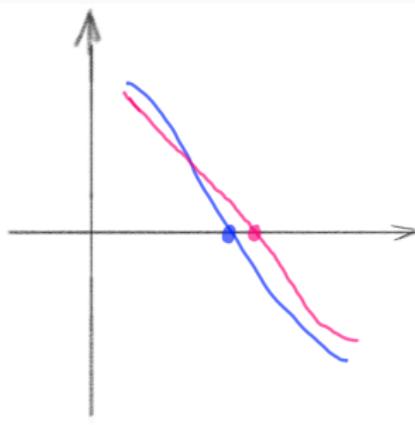


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$f$

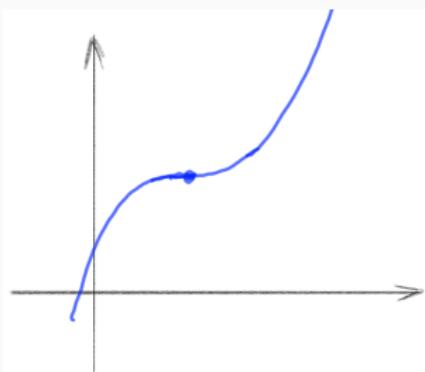


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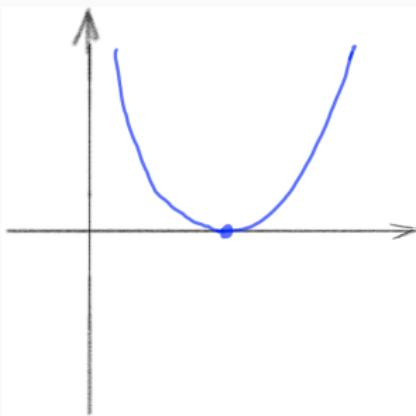


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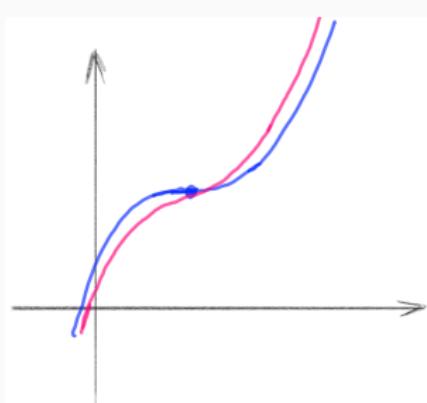


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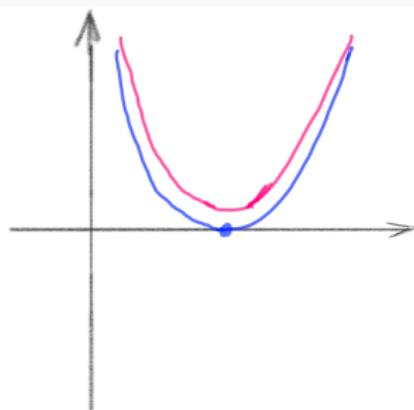


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$$E_t(gx) = E_t(x)$$

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- Nondegenerate  $\Rightarrow$  isolated.

## Symmetry analysis

---

- Apply this to

$$E_t = E_2 + E_4 + t \int_{\mathbb{T}^3} (1 - \varphi_0) \text{vol}_g$$

- $X = C^2(\mathbb{T}^3, SU(2)) \times SPD_3$
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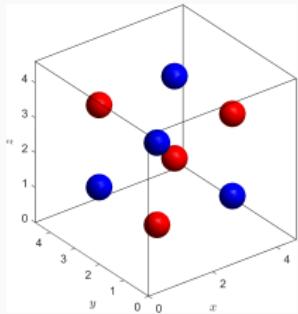
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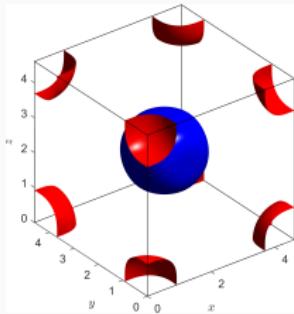
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- Reduces to a problem in representation theory of subgroups of  $O_h$

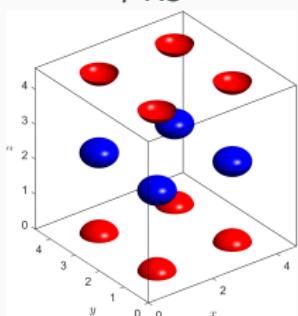
# The KS crystals that (should) survive



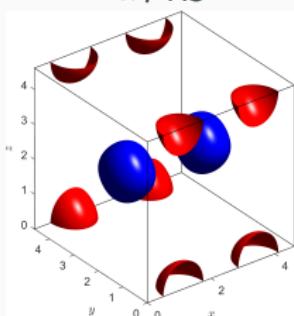
$\varphi_{KS}$



$$\varphi_0 = 0.9$$



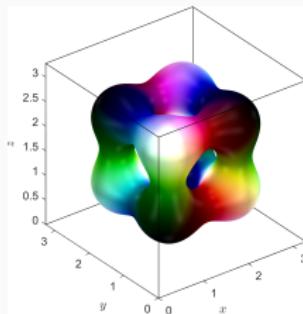
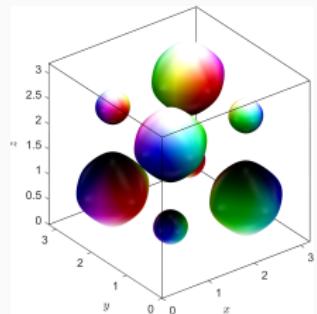
$R_{sheet}\varphi_{KS}$



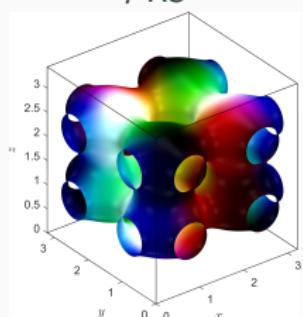
$R_{chain}\varphi_{KS}$

$$\varphi_0 = -0.9$$

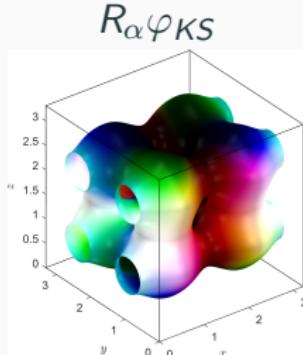
# Skyrme crystals at pion mass $t = 1$



$\varphi_{KS}$

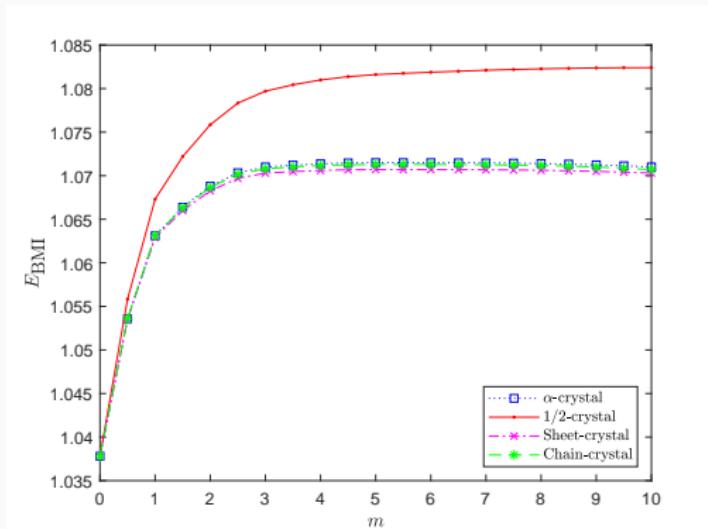


$R_{sheet} \varphi_{KS}$



$R_{chain} \varphi_{KS}$

# Energy ordering: sheet < chain < $\alpha$ < KS



$$g_{\text{sheet}} = \begin{pmatrix} L_1 & 0 & 0 \\ 0 & L_1 & 0 \\ 0 & 0 & L_3 \end{pmatrix}$$
$$L_3 > L_1$$

$$g_{\text{chain}} = \begin{pmatrix} L_1 & 0 & 0 \\ 0 & L_2 & 0 \\ 0 & 0 & L_2 \end{pmatrix}$$
$$L_2 > L_1$$

trigonal, but **not** cubic!

## Isospin inertia tensors

---

$$U_{KS} = \begin{pmatrix} 165.2 & 0 & 0 \\ 0 & 165.2 & 0 \\ 0 & 0 & 165.2 \end{pmatrix}, \quad U_\alpha = \begin{pmatrix} 135.5 & 0 & 0 \\ 0 & 135.5 & 0 \\ 0 & 0 & 167.3 \end{pmatrix},$$

$$U_{sheet} = \begin{pmatrix} 135.8 & 0 & 0 \\ 0 & 135.8 & 0 \\ 0 & 0 & 166.8 \end{pmatrix}, \quad U_{chain} = \begin{pmatrix} 135.6 & 0 & 0 \\ 0 & 135.7 & 0 \\ 0 & 0 & 167.2 \end{pmatrix}.$$

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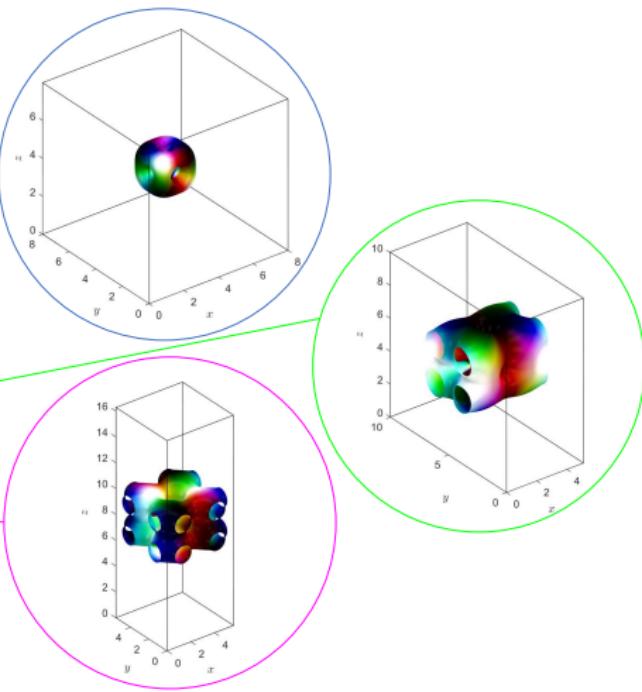
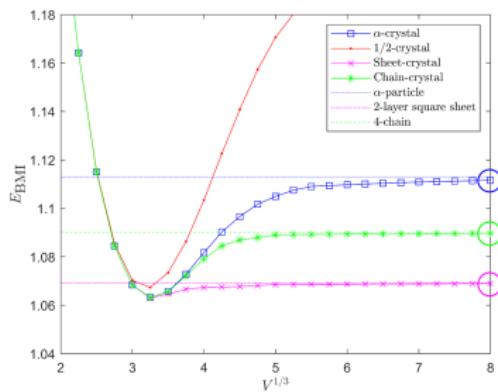
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- Global min is the only c.p.

# Optimal crystals at fixed baryon density



## Concluding remarks

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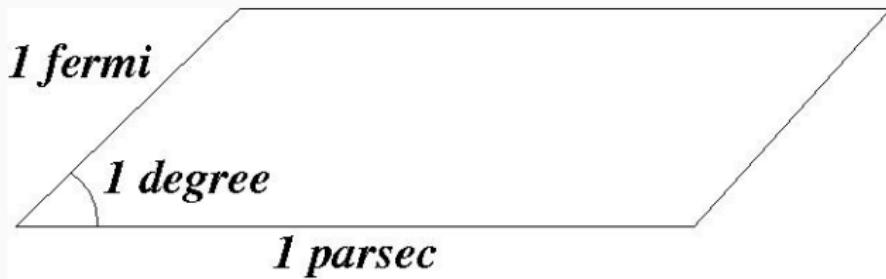
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- Many examples in condensed matter (cf work with Tom Winyard et al). True also for nuclear Skyrme model with massive pions
- Extreme case: baby Skyrme model  $\varphi : M^2 \rightarrow S^2$

$$E(\varphi) = \int_M \left( \frac{1}{2} |d\varphi|^2 + \frac{1}{2} |\varphi^* \omega|^2 + V(\varphi) \right).$$

Given **any** period lattice  $\Lambda \subset \mathbb{R}^2$ , can cook up a smooth potential  $V : S^2 \rightarrow [0, \infty)$  s.t.  $E(\varphi, g)$  has a global min at  $(\varphi_*, g_\Lambda)$  with  $\varphi_*$  degree 2 and holomorphic.

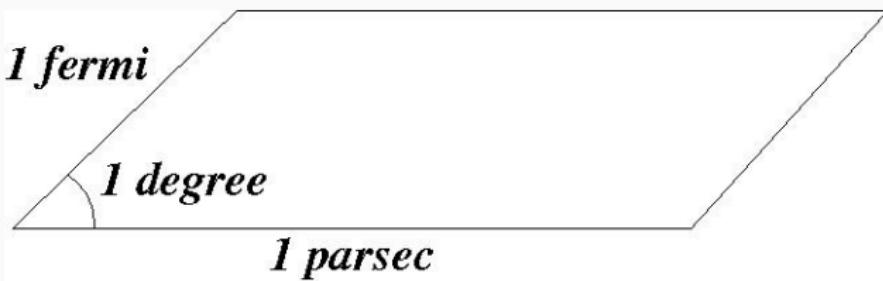
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- Existence result **at fixed volume** very generic

$$E(\varphi, g) = E_2(\varphi, g) + \text{positive, geom nat}$$

any dimension. Compactness argument works. E.g.  $\omega$ -meson Skyrme model